

A free boundary problem from Brownian bees in the infinite swarm limit in \mathbb{R}^d

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Including joint work with

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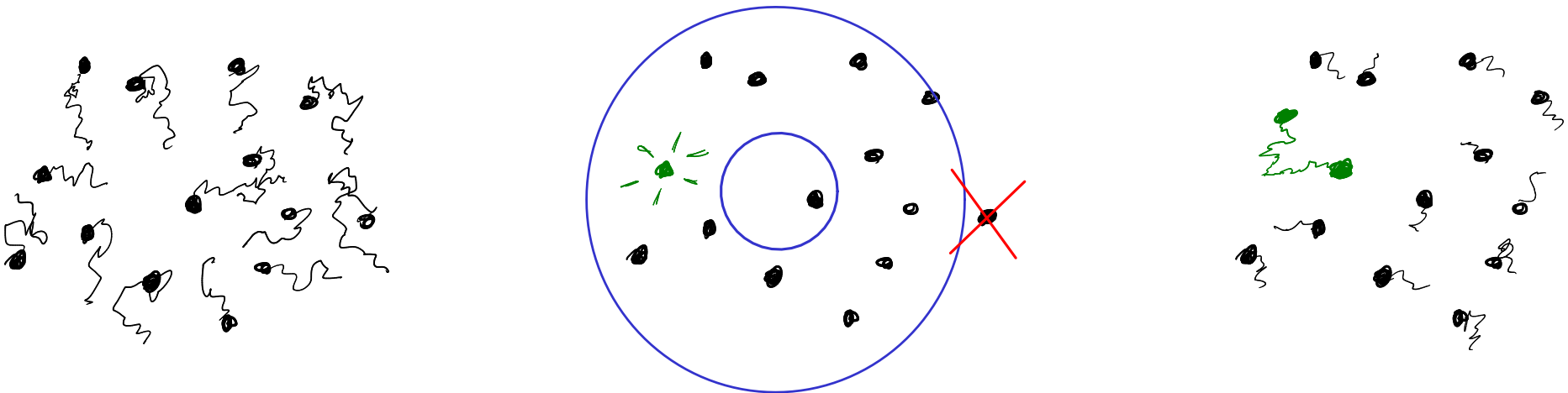
and with Erin Beckman (Concordia)

The **Brownian Bees Process** is a special case of Branching Brownian Motion with selection (**N -BBM**).

There are N particles, moving as independent Brownian motions in \mathbb{R}^d .

Each particle branches independently, at rate one.

At each branch event, the particle farthest from the origin is removed.



General N -**BBM process** with spatial selection:

There are N particles, moving as independent Brownian motions in \mathbb{R}^d .

Each particle branches independently, at rate one.

At each branch event, remove the particle with minimal “fitness”, according to fitness function $V(x) : \mathbb{R}^d \rightarrow \mathbb{R}$.

Thus, the total number of particles is always N .

$$\mathbf{X}(t) = \{X_1(t), \dots, X_N(t)\} \subset \mathbb{R}^d$$

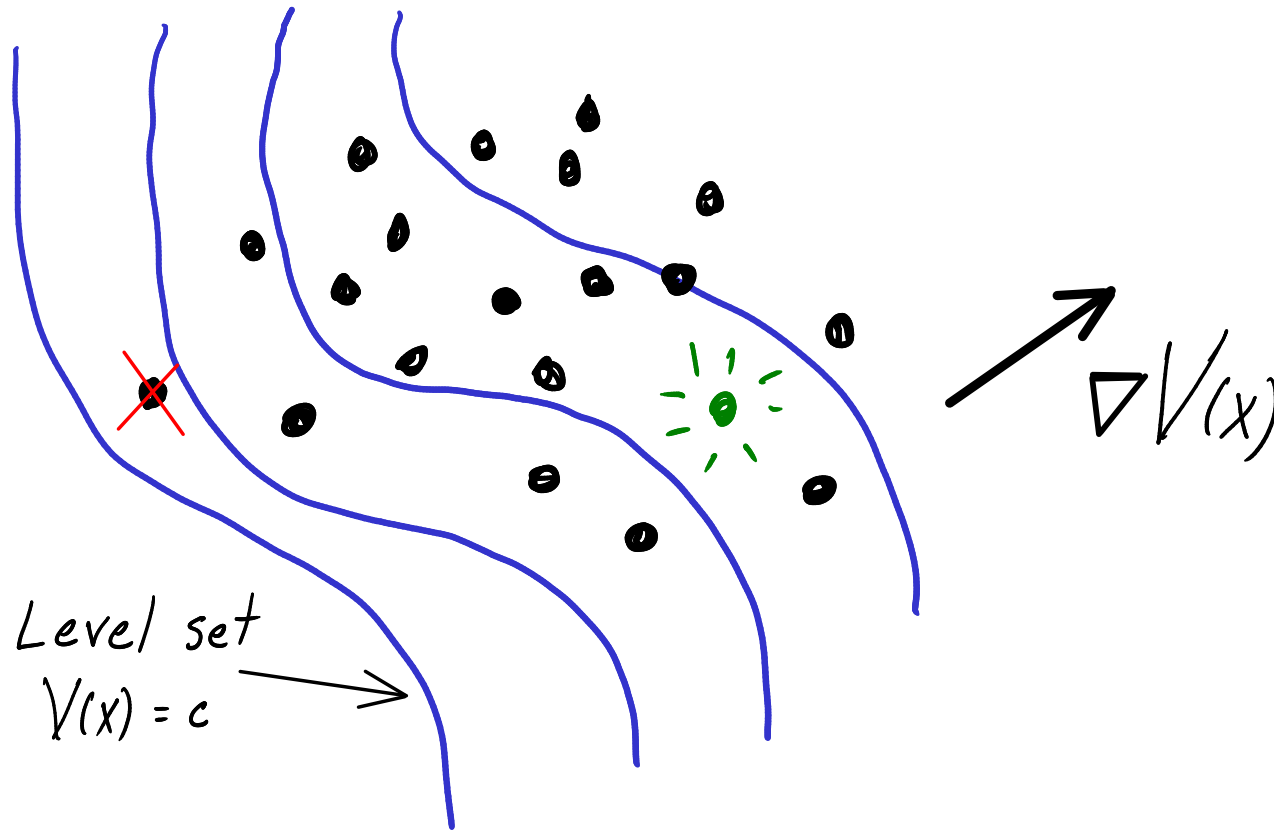
Particle positions: $X_1(t), \dots, X_N(t) \in \mathbb{R}^d$

Fitness values: $V(X_1(t)), \dots, V(X_N(t))$

Brunet, Derrida, Mueller, Munier (2006)

N. Berestycki, L. Zhao (2018).

Selection pressure pushes the ensemble in the direction of ∇V .



The process is unchanged by replacing V with $F \circ V$, F increasing. For the **Brownian Bees** process, $V(x) = -|x|$; the selection has a confining effect on the population.

Some interesting questions:

How does the structure of $V(x)$ effect the evolution of the population?

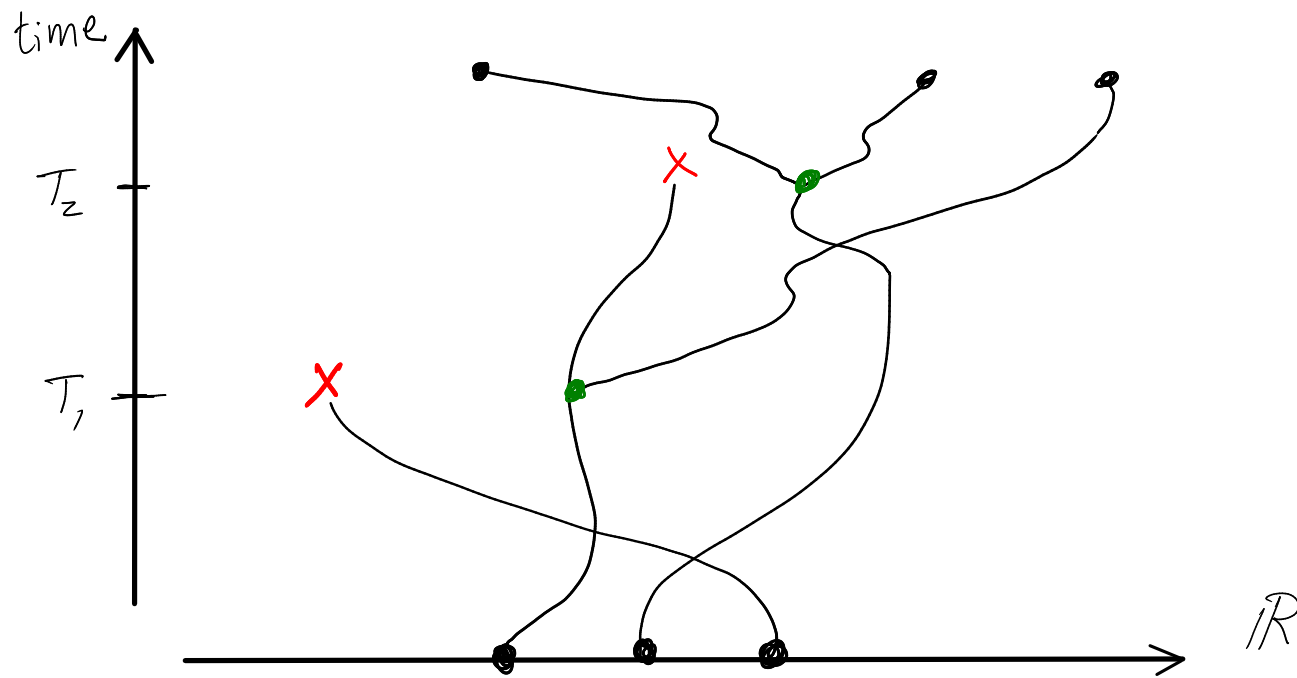
What is the behavior as $t \rightarrow \infty$?

Is there a hydrodynamic limit as $N \rightarrow \infty$?

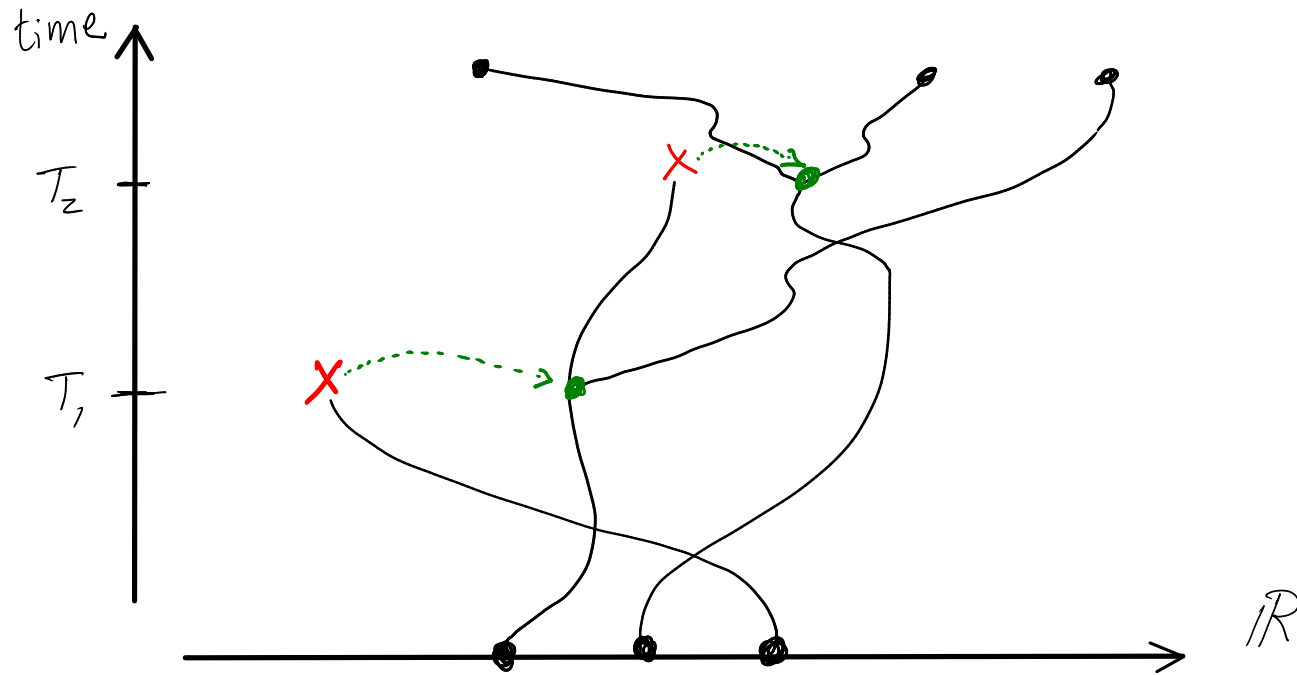
Do these limits commute?

How does V effect the genealogy of the process?

There have been several works on the case $d = 1$ with monotone fitness: $V(x) = x$.
The selection pressure produces a “drift” to the right.



There have been several works on the case $d = 1$ with monotone fitness: $V(x) = x$. The selection pressure produces a “drift” to the right.



Alternatively, you might say that the left-most particle jumps to the location of another randomly chosen particle.

For each N , the selection pushes the ensemble to the right with positive speed c_N :

$$\lim_{t \rightarrow \infty} \max_k \frac{X_k(t)}{t} = \lim_{t \rightarrow \infty} \min_k \frac{X_k(t)}{t} = c_N > 0,$$

almost surely. Moreover,

$$c_N \sim 2 - \frac{\pi^2}{(\log N)^2}, \quad \text{as } N \rightarrow \infty.$$

Bérard and Guérou, Comm. Math. Phys, 2010.

N. Berestycki, L. Zhao, Ann. Appl. Prob. 2018.

There are some outstanding conjectures related to the **genealogy**:

It is expected that the ancestral process behaves like the Bolthausen-Sznitman coalescent on a time scale $(\log N)^3$.

This allows for multiple simultaneous mergers in the family ancestry, because a single particle may run far ahead and then quickly branch to make up a significant fraction of the population.

Brunet, Derrida, Mueller, Munier, Phys Rev E, 2006.

P. Maillard, PTRF 2016.

J. Berestycki, N. Berestycki, J. Schweinsberg, Ann. Probability 2013.

Hydrodynamic limit, $d = 1$:

For $d = 1$, with monotone fitness, De Masi, Ferrari, Presutti, Soprano-Lotto (2017) proved a hydrodynamic limit for the empirical measure

$$\mu_t^N = \frac{1}{N} \sum_{k=1}^N \delta_{X_k(t)}$$

As $N \rightarrow \infty$, $\mu_t^N \rightarrow v(t, x)$ which solves

$$v_t = v_{xx} + v, \quad x \geq \gamma(t)$$

$$v(t, \gamma(t)) = 0$$

$$\int_{\gamma(t)}^{\infty} v(t, x) dx = 1$$

$\gamma(t)$ defines a **free boundary**.

Global existence of PDE solution: J. Berestycki, É. Brunet, S. Penington 2018

This problem is equivalent to

$$v_t = v_{xx} + v, \quad x \geq \gamma(t),$$

$$v(t, \gamma(t)) = 0,$$

$$\gamma'(t) = -v_{xx}(t, \gamma(t))$$

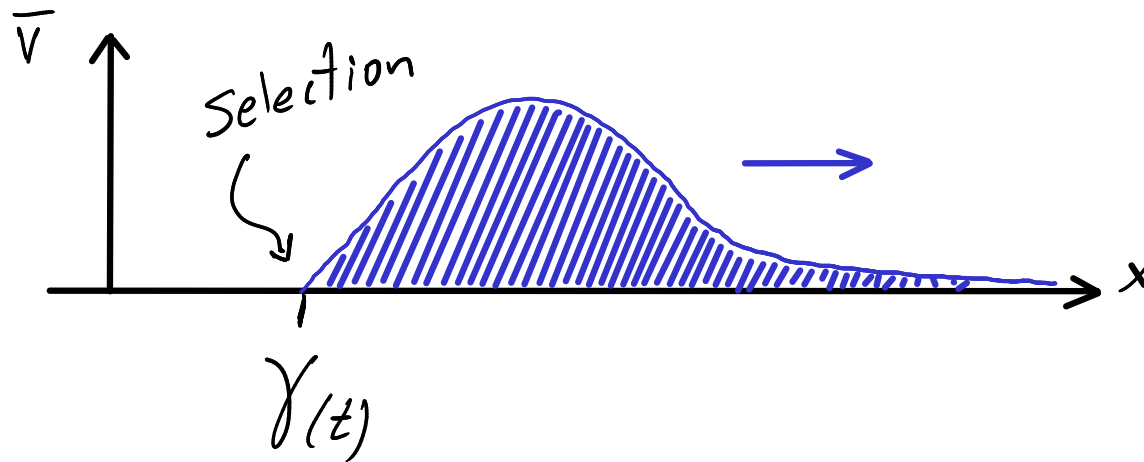
$$v_x(t, \gamma(t)) = 1$$

Recall that the supercooled Stefan problem involves $\gamma'(t) = v_x(t, \gamma(t))$.

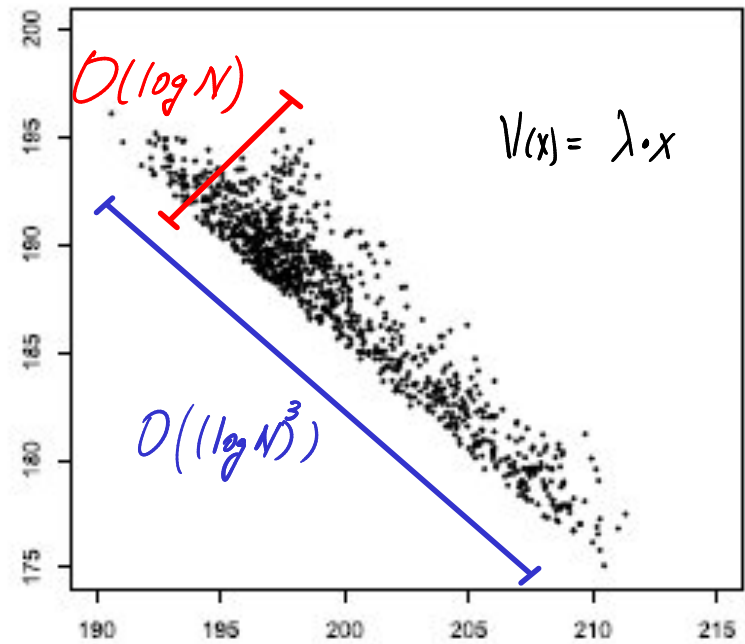
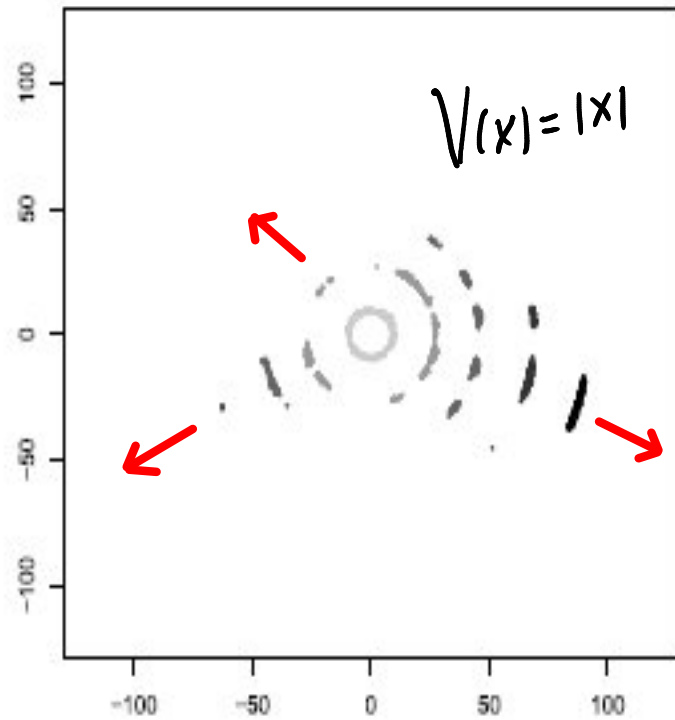
Chayes, Kim (2012), Delarue, Nadtochiy, Shkolnikov (2019)

There is an explicit traveling wave solution

$$\bar{v}(t, x) = \phi(x - \gamma(t)), \quad \phi(x) = xe^{-x}, \quad \gamma(t) = 2t.$$



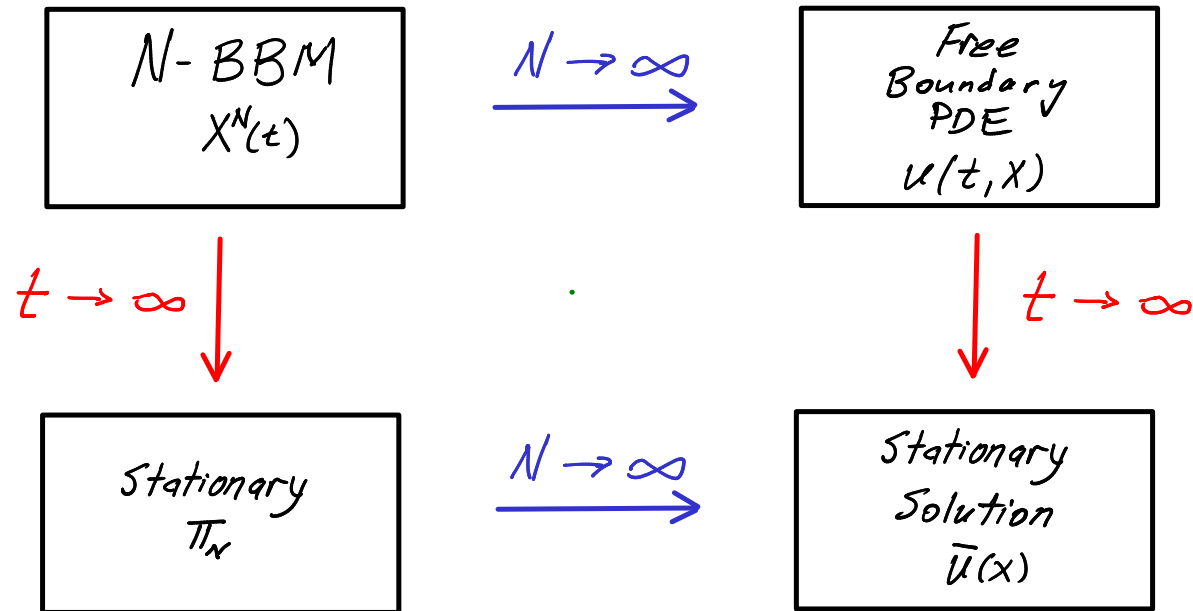
For $d \geq 2$, or with non-monotone fitness, there are relatively few results.



Images from N. Berestycki, L. Zhao, Ann. Appl. Prob. 2018.

Our work: N-BBM in dimension $d \geq 1$, with fitness $V(x) = -|x|$.

Super-level sets of V are compact (Euclidean balls). So, the selection has a confining effect on the ensemble.



Hydrodynamic limit, $d \geq 1$:

Let μ_t^N be the empirical measure for the N-BBM with fitness $V(x) = -|x|$ in \mathbb{R}^d , and define

$$M_t^N = \max_{k \in \{1, \dots, N\}} |X_k(t)|.$$

Suppose $\{X_k(0)\}_{k=1}^N$ are i.i.d, $X_k(0) \sim \mu_0 \in \mathcal{P}(\mathbb{R}^d)$. Then for every $t > 0$,

$$\begin{aligned} \mu_t^N(dx) &\rightarrow u(t, x) dx \text{ weakly} \\ M_t^N &\rightarrow R_t \end{aligned}$$

almost surely, as $N \rightarrow \infty$, where $(u(t, x), R_t)$ solves a free-boundary problem with initial condition μ_0

Free boundary problem in \mathbb{R}^d for the hydrodynamic limit: Given initial measure $\mu_0 \in \mathcal{P}(\mathbb{R}^d)$, find a pair $(u(x, t), R_t)$ such that

$$u_t = \Delta u + u, \quad |x| < R_t, \quad t > 0$$

$$u(t, x) = 0, \quad |x| \geq R_t, \quad t > 0$$

$$\int_{|x| \leq R_t} u(t, x) dx = 1$$

$$u(t, x) \rightarrow \mu_0(dx), \text{ weakly as } t \rightarrow 0$$

Theorem: For any Borel probability measure μ_0 on \mathbb{R}^d , there is a unique, global solution (u, R_t) to this free boundary problem. Moreover, R_t is finite and continuous for all $t > 0$, and $R_t \rightarrow R_0 := \inf\{r > 0 \mid \mu_0(\mathcal{B}(r)) = 1\} \in [0, \infty]$.

Free boundary problem

$$u_t = \Delta u + u, \quad |x| < R_t, \quad t > 0$$

$$u(t, x) = 0, \quad |x| \geq R_t, \quad t > 0$$

$$\int_{|x| \leq R_t} u(t, x) dx = 1$$

$$u(t, x) \rightarrow \mu_0(dx), \text{ weakly as } t \rightarrow 0$$

$u(x, t)$ may not be spherically symmetric, even though the free boundary is a sphere: $\{|x| = R_t\}$.

If μ_0 has an atom at the edge of its support, then R_t very rapidly moves outward, $R_t \gg R_0 + \sqrt{t}$ as $t \rightarrow 0$. However, for any $\alpha \in (0, 1/2)$,

$$R_t - R_s \leq C_\alpha (t - s)^\alpha, \quad 0 \leq s < t.$$

Large time limit for the free boundary problem:

Let R_∞ be chosen so that the principal Dirichlet eigenvalue of $(-\Delta)$ is $\lambda = 1$:

$$\begin{aligned} -\Delta \bar{u} &= \bar{u}, & |x| &\leq R_\infty, \\ \bar{u}(x) &= 0, & |x| &= R_\infty. \end{aligned}$$

Then,

$$\|u(t, \cdot) - \bar{u}(\cdot)\|_\infty \leq Ct^{-1}, \quad \text{and} \quad |R_t - R_\infty| \leq Ct^{-1}$$

In particular, for $d = 1$: $R_\infty = \pi/2$, $\bar{u}(x) = \alpha \cos(x)$.

Stationary distribution:

For each N , the N-BBM problem for $d \geq 1$ with $V(x) = -|x|$ has a stationary distribution π^N on $(\mathbb{R}^d)^N$. ($\mathbf{X}(t) = (X_1(t), \dots, X_N(t))$ is positive Harris recurrent)

The density of particles under π^N converges to \bar{u} as $N \rightarrow \infty$. For a configuration $\mathbf{X} \in (\mathbb{R}^d)^N$, let $\mu_{\mathbf{X}}^N$ be its empirical measure. For any $\epsilon > 0$, and any $A \subset \mathbb{R}^d$,

$$\pi^N \left(\mathbf{X} \in (\mathbb{R}^d)^N \mid \left| \mu_{\mathbf{X}}^N(A) - \int_A \bar{u}(x) dx \right| > \epsilon \right) \rightarrow 0$$

as $N \rightarrow \infty$. Also,

$$\pi^N \left(\mathbf{X} \in (\mathbb{R}^d)^N \mid \left| \max_{k=1, \dots, N} |X_k| - R_\infty \right| > \epsilon \right) \rightarrow 0$$

N -BBM
 $X^N(t)$

$N \rightarrow \infty$
→

Free
Boundary
PDE
 $u(t, x)$

$t \rightarrow \infty$
↓

Stationary
 π_x

$N \rightarrow \infty$
→

Stationary
Solution
 $\bar{u}(x)$

↓ $t \rightarrow \infty$

Observations about the free boundary problem

$$u_t = \Delta u + u, \quad |x| < R_t, \quad t > 0$$

$$u(t, x) = 0, \quad |x| \geq R_t, \quad t > 0$$

$$\int_{|x| \leq R_t} u(t, x) dx = 1$$

$$u(t, x) \rightarrow \mu_0(dx), \text{ weakly as } t \rightarrow 0$$

Assuming u is C^2 at the boundary, we formally derive

$$\int_{|x|=R_t} (\nu \cdot \nabla u) dS(x) = -1, \quad \text{and} \quad \frac{d}{dt} R_t = \int_{|x|=R_t} \Delta u(x, t) dS(x)$$

So, this has some similarities to a problem of Caffarelli and Vázquez (1992):

$$u_t = \Delta u, \quad u > 0, \quad x \in \Omega_t$$

with $u = 0$ and $|\nabla u| = c$ on the free boundary $\partial\Omega_t$.

Hence speed of free boundary is $\frac{u_t}{|\nabla u|} = c^{-1} \Delta u$.

The free boundary problem is closely related to a **parabolic obstacle problem**.

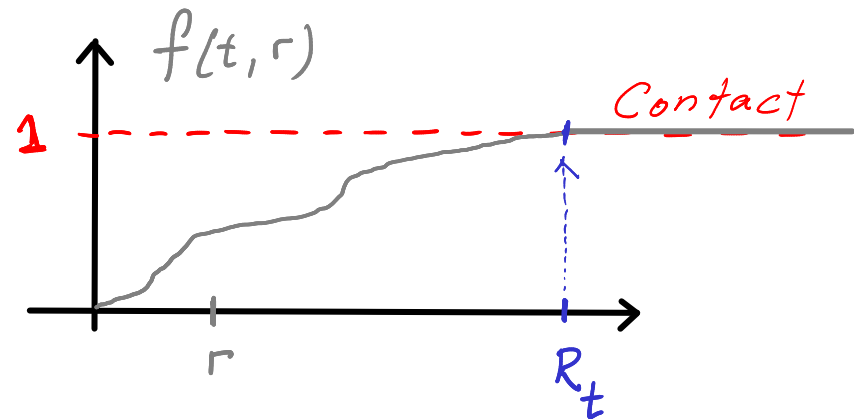
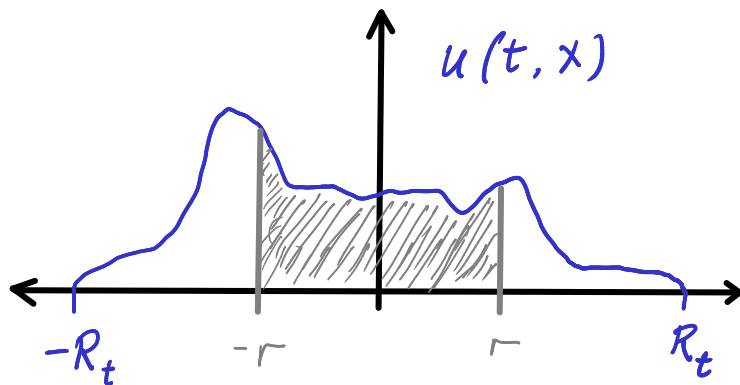
The function $f(t, r) = \int_{|x| \leq r} u(t, x) dx, \quad r \geq 0$

satisfies $f \leq 1$ and

$$f_t = f_{rr} - \frac{d-1}{r} f_r + f, \quad \text{if } f(t, r) < 1$$

$r \mapsto f_r(t, r)$ is continuous

R_t is the boundary of the contact set $\{f(t, \cdot) = 1\}$. At R_t , $f_r(t, R_t) = 0$.



Obstacle problems: see A. Friedman 1982.

General strategy for well-posedness of original free boundary problem:

- (i) Well-posedness of the obstacle problem with initial condition $f_0(r) = \mu_0(\{|x| < r\})$. This defines $R_t = \inf\{r > 0 \mid f(r, t) = 1\}$.
- (ii) R_t is continuous for $t > 0$.
- (iii) Given R_t , prove that

$$u(x, t) = e^t \int_{\mathbb{R}^d} \mu_0(dx) \rho^R(x, y, t)$$

solves the free boundary problem, where $y \mapsto \rho^R(x, y, t)$ is density of Brownian motion that avoids the boundary before time t :

$$\int_A \rho^R(x, y, t) dy = \mathbb{P}_x(B_t \in A, \quad |B_s| < R_s \quad \forall s \in (0, t]).$$

The solution f to the obstacle problem is approximated by $f = \lim_{n \rightarrow \infty} f_n$,

$$\partial_t f_n = \partial_r^2 f_n - \frac{d-1}{r} \partial_r f_n + f_n - (f_n)^n, \quad r > 0$$

Which means that u is the limit of an approximating sequence $u_n(t, x)$ satisfying

$$\partial_t u_n = \Delta u_n + u_n - u_n A_n([u_n], x, t)$$

where the nonlinear, nonlocal term A_n is

$$A_n([u_n], x, t) = n \left(\int_{B(|x|)} u_n(t, y) dy \right)^{n-1}$$

which strongly penalizes mass at points $|x| > R = R_{n,t}$ such that

$$\int_{B(R)} u_n(t, y) dy = 1,$$

but is negligible at points $|x| < R$.

The obstacle problem provides monotonicity that has an analogue for the stochastic system. The radial distribution function

$$f(t, r) = \int_{|x| \leq r} u(t, x) dx, \quad r \geq 0$$

satisfies the obstacle problem: $f \leq 1$ with

$$f_t = f_{rr} - \frac{d-1}{r} f_r + f, \quad \text{where } f < 1$$

$$f(t, 0) = 0, \quad t > 0,$$

$$r \mapsto f_r(t, r) \text{ is continuous.}$$

Given $f(t, r)$, the solution at time $t + \delta$ may be approximated by operator splitting

$$e^\delta G_\delta * \mathcal{C}_{e^{-\delta}} f(t, r) \leq f(r, t + \delta) \leq \mathcal{C}_1(e^\delta G_\delta * f(t, r))$$

where operator \mathcal{C}_m cuts mass to level m :

$$\mathcal{C}_m g(r) = \min(g(r), m)$$

For the stochastic particle system, this idea suggests a coupling of the N-BBM with unconstrained BBM.

Let $\mathbf{X}^{(N)}(t)$ be the N-BBM. Let $\mathbf{X}^+(t)$ be a regular BBM, unconstrained.

If the particles are labeled by *rank*: $|X_1^{(N)}| \leq |X_2^{(N)}| \leq \dots \leq |X_N^{(N)}|$, then

$$|X_k^+| \preceq |X_k^{(N)}|, \quad k = 1, \dots, N.$$

That is, the CDF of $|\mathbf{X}^+|$ stochastically dominates the CDF of $|\mathbf{X}^N|$:

$$F^{(N)}(t, \cdot) \preceq F^+(t, \cdot), \quad F^{(N)}(t, \cdot) \preceq \mathcal{C}_1 F^+(t, \cdot)$$

where

$$F^{(N)}(t, r) = \frac{1}{N} \sum_{k=1}^N \mathbb{1}_{\{|X_k(t)| \leq r\}} \quad F^+(t, r) = \frac{1}{N} \sum_{k=1}^{N_t^+} \mathbb{1}_{\{|X_k^+(t)| \leq r\}}.$$

Compare this to: $f(r, t + \delta) \leq \mathcal{C}_1 (e^\delta G_\delta * f(t, r))$

This coupling gives an upper bound on $F^{(N)}$.

A lower bound can be obtained by cutting first, then evolving the free BBM until $N_t^+ = N$.

Define $\mathbf{X}^+(0) \subset \mathbf{X}^{(N)}(0)$ by removing particles from $\mathbf{X}^{(N)}(0)$ with highest ranks, then

$$|X_k^{(N)}(t)| \preceq |X_k^+(t)|, \quad k = 1, \dots, N_t$$

as long as $N_t^+ \leq N$. Which implies

$$\mathbb{1}_{\{N_t^+ < N\}} F^+(t, \cdot) \preceq F^{(N)}(t, \cdot)$$

Compare this to: $e^\delta G_\delta * \mathcal{C}_{e^{-\delta}} f(t, r) \leq f(r, t + \delta)$

The first version of this coupling comes from De Masi, Ferrari, Presutti, Soprano-Lotto (2017).

For **general fitness function** $V(x) : \mathbb{R}^d \rightarrow \mathbb{R}$, one expects the following limit problem:

For $\ell \in \mathbb{R}$, let Ω_ℓ be the super-level set of fitness function V :

$$\Omega_\ell = \{x \in \mathbb{R}^d \mid V(x) > \ell\}$$

Find $(u(t, x), \ell_t)$ such that

$$u_t = \Delta u + u, \quad x \in \Omega_{\ell_t},$$

$$u(t, x) = 0, \quad x \in \partial\Omega_{\ell_t}$$

$$\int_{\Omega_{\ell_t}} u(t, x) dx = 1$$

So, the free boundary is $\partial\Omega_{\ell_t}$, the ℓ_t -level set of V .

ℓ_t is the fitness-level at which selection happens.

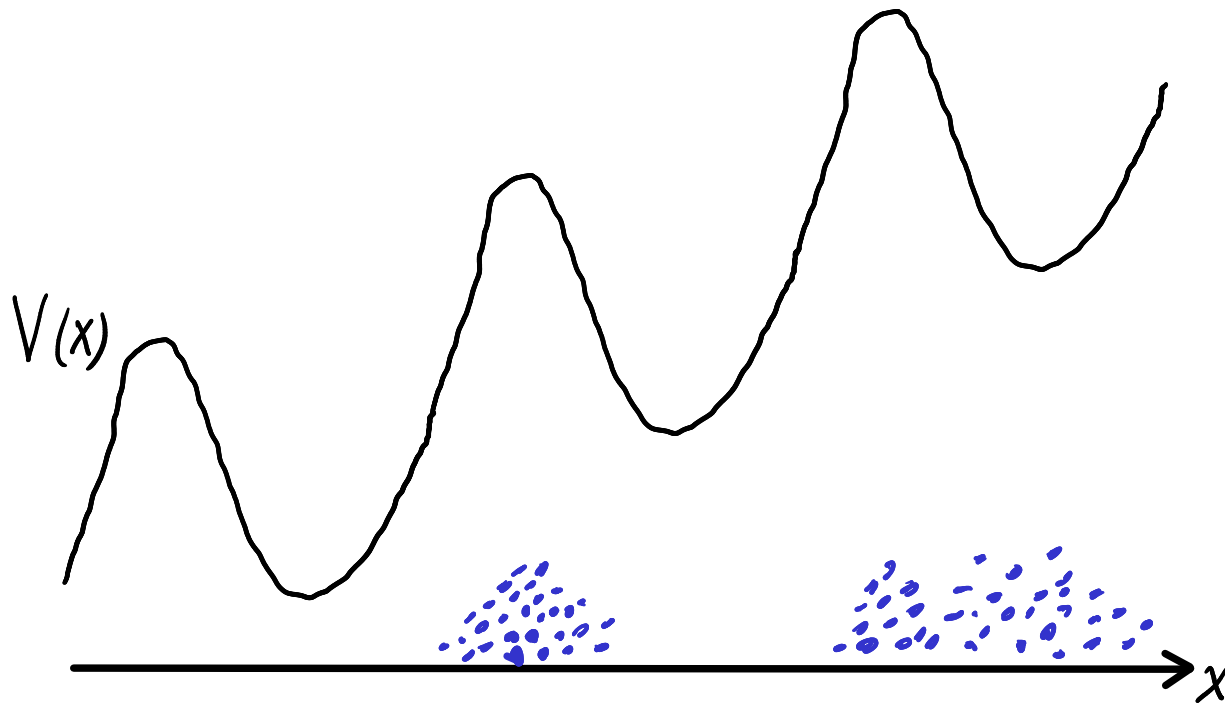
The topology of Ω_{ℓ_t} may change as ℓ_t varies.

Back to N-BBM:

What is the effect of fitness barriers? Suppose $d = 1$, but non-monotone fitness.

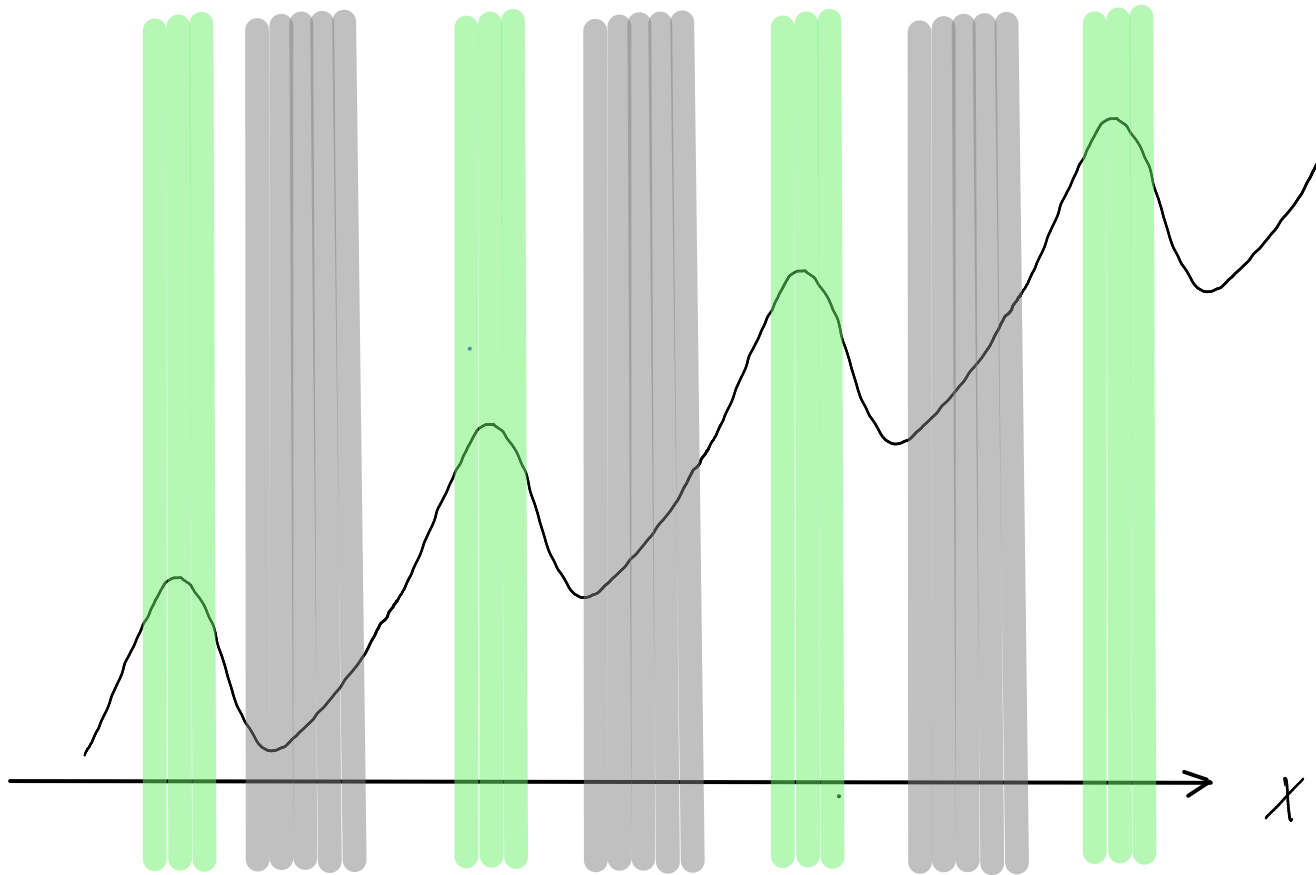
$$V(x) = x + \psi(x), \quad \psi(x) \text{ is } L\text{-periodic.}$$

Particles must pass through a fitness valley before moving to better territory.



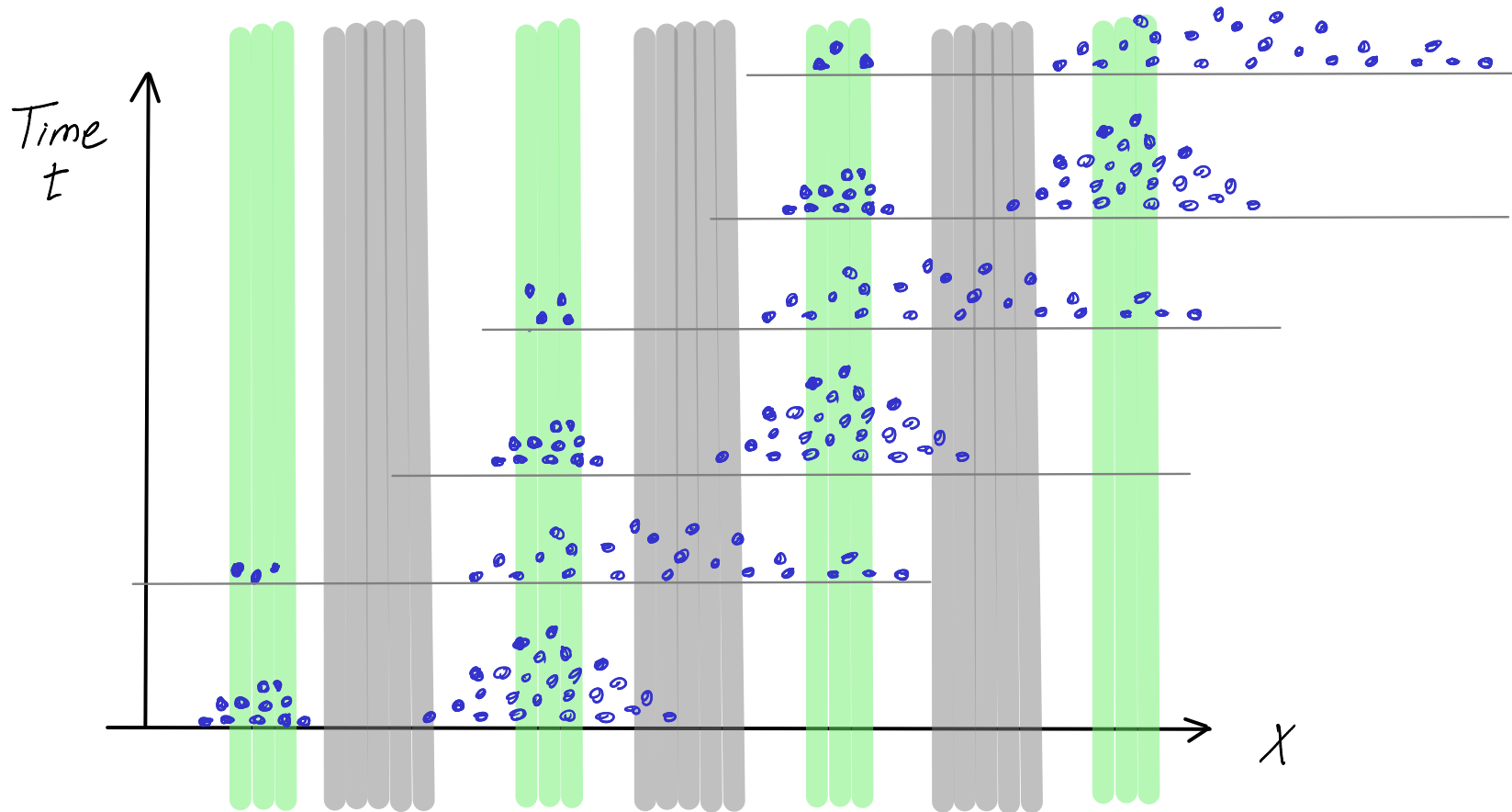
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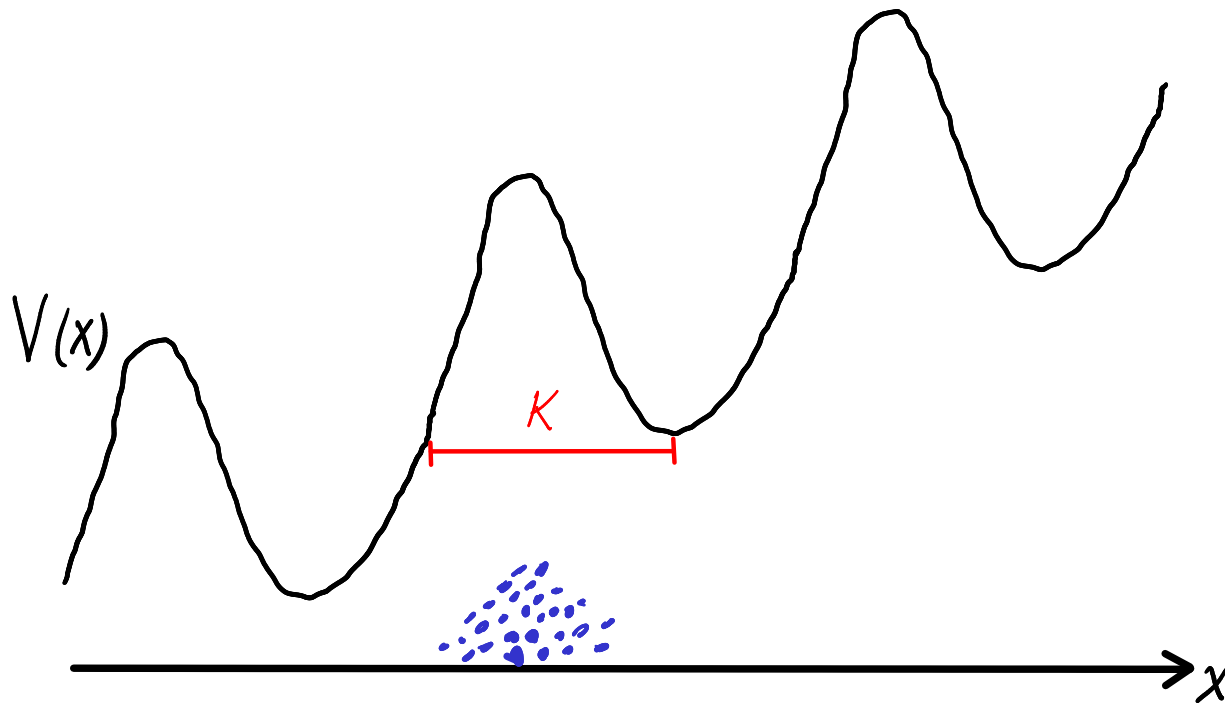
Theorem: for each N there is a unique stationary distribution, within a moving frame. Moreover, there is c_N such that

$$\lim_{t \rightarrow \infty} \max_k \frac{X_k(t)}{t} = \lim_{t \rightarrow \infty} \min_k \frac{X_k(t)}{t} = c_N$$

holds almost surely and in L^1 .

Erin Beckman, N. (2019)

Unlike the case of monotone fitness, the hydrodynamic limit may be a standing wave (eigenfunction) inside a local fitness valley, if $K > 2R_\infty = \pi$.



Numerical simulations show a metastable behavior when $K > 2R_\infty$.

Is $c_N > 0$?

Conclusion/summary:

- What aspects of this extend to general fitness V ?
- What is effect of V on the population?
- Traveling waves? metastability? rate of adaptation?
- Strong selection principle?
- Variants on graphs?

For references and further reading:

- J. Berestycki, Brunet, N., Penington: *A free boundary problem arising from branching Brownian motion with selection*, arXiv:2005.09384
- J. Berestycki, Brunet, N., Penington: *Brownian bees in the infinite swarm limit*, to appear soon...