

The Principle of Least Action and its Geometric Aspects

... What I understood ...

UCLA IPAM, February 2016

Bruno Lévy ALICE Géométrie & Lumière CENTRE INRIA Nancy Grand-Est

Part. 1. The Least Action Principle

Part. 2. Optimal Transport – Elementary intro.

Part. 3. The Semi-Discrete Case

Part. 4. Understanding What's Going On...

Part. 5. Concluding Words





Yann Brenier

The polar factorization theorem (Brenier Transport)

"Each time the Laplace operator is used in a PDE, it can be replaced with the Monte-Ampère operator, and then interesting things occur"

Innía



Yann Brenier

The polar factorization theorem (Brenier Transport)



Cédric Villani Optimal Transport Old & New Topics on Optimal Transport

"Each time the Laplace operator is used in a PDE, it can be replaced with the Monte-Ampère operator, and then interesting things occur"

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A computer programmer's adventure in MathLand



June 2015 Institut Fourier

March 2015 Bonn

Febr 2015 BIRS (Canada)

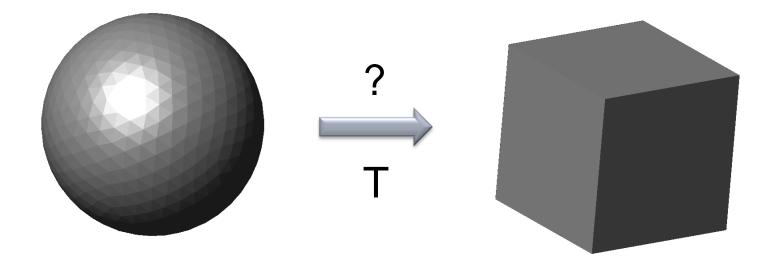
Febr 2015 LJLL

Discuss. with Quentin Mérigot (computational geometry – geometric measure theory)





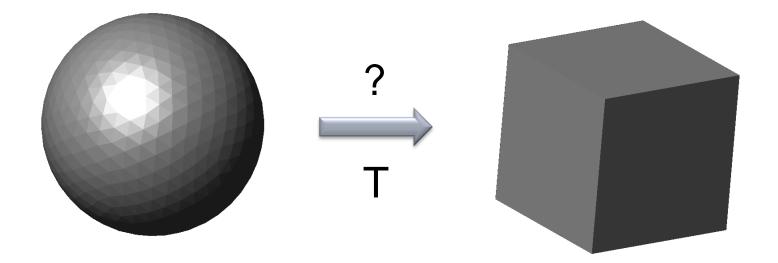
Part. 1 The Least Action Principle



How to "morph" a shape into another one of same mass while minimizing the "effort" ?



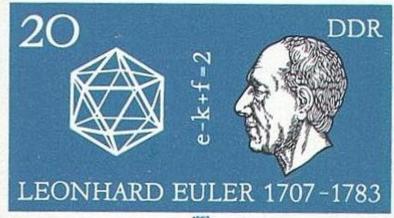
Part. 1 The Least Action Principle



How to "morph" a shape into another one of same mass while minimizing the "effort" ?

The "effort" of the best T defines a distance between the shapes





Hamilton, Legendre, Maupertuis

Lagrange



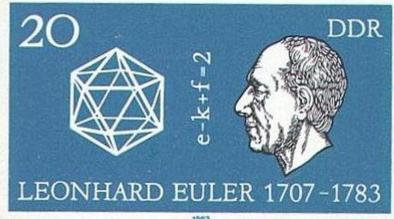
The Least Action Principle

Axiom 1: There exists a function L(x, x, t)

that describes the state of a physical system

Short summary of the 1st chapter of Landau, Lifshitz Course of Theoretical Physics





Hamilton, Legendre, Maupertuis

Lagrange



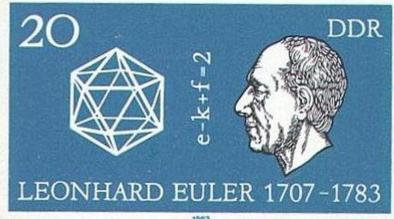
The Least Action Principle

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position

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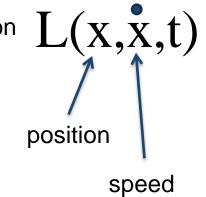
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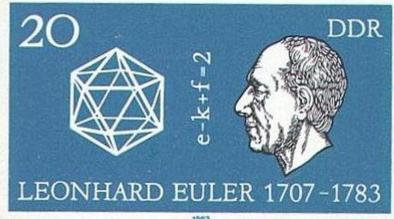
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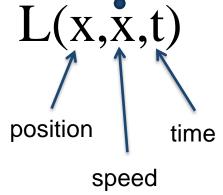
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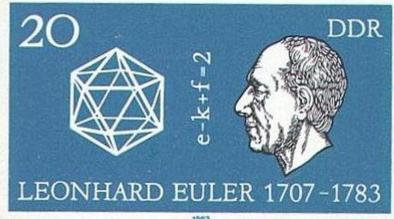
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The Least Action Principle

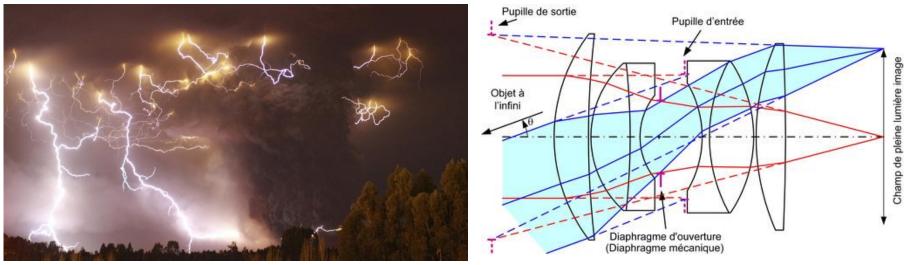
Axiom 1: There exists a function L(x, x, t)

that describes the state of a physical system

Axiom 2: The movement (time evolution) of the physical system minimizes the following integral

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$$\int_{t_1}^{t_2} L(\mathbf{x}, \mathbf{x}, t) dt$$



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Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$



Axiom 1: There exists L Axiom 2: The movement minimizes

$$\int_{t_1}^{t_2} L(x, x, t) dt$$

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t} = \frac{x+vt}{t}$$

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Theorem 2:

$$\mathbf{\hat{x}} \frac{\partial \mathbf{L}}{\partial \mathbf{\hat{x}}} - \mathbf{L} = \mathsf{cte}$$

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Theorem 2:

OX

$$\frac{\partial L}{\partial \cdot}$$
 - L = cte

Homogeneity of time \rightarrow Preservation of **energy**

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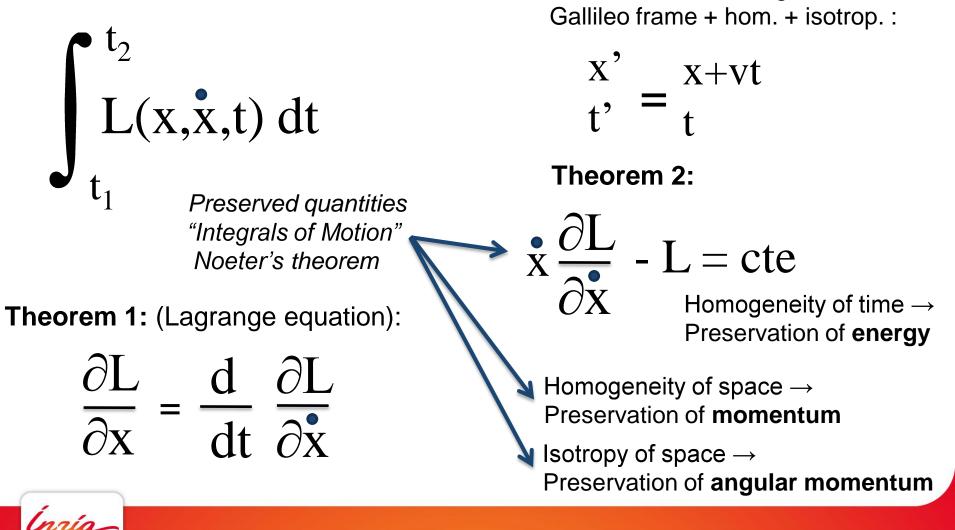
Isotropy of space \rightarrow Preservation of **angular momentum**



Axiom 3:

Invariance w.r.t. change of

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Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

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Free particle:

Theorem 3: v = cte (Newton's law I)

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Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$

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UX

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$$\frac{\partial L}{\partial x} - L = cte$$

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Particle in a field:

Expression of the Lagrangian: $L = \frac{1}{2} m v^2 - U(x)$



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Free particle:

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Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$

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Particle in a field:

Expression of the Lagrangian: $L = \frac{1}{2} m v^2 - U(x)$ Expression of the Energy: $E = \frac{1}{2} m v^2 + U(x)$

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Particle in a field:

Expression of the Lagrangian: $L = \frac{1}{2} m v^{2} - U(x)$ Expression of the Energy: $E = \frac{1}{2} m v^{2} + U(x)$ Theorem 4: $mx = -\nabla U \text{ (Newton's law II)}$

(relativistic setting)

Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$

Axiom 3:

Invariance w.r.t. Lorentz change of frame

$$\frac{x'}{t'} = \frac{(x-vt) \times \gamma}{(t - vx/c^2) \times \gamma}$$

$$\gamma = 1 / \sqrt{(1 - v^2 / c^2)}$$



The Least Action Principle (relativistic setting – just for fun...)

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$$\gamma = 1 / \sqrt{(1 - v^2 / c^2)}$$

Theorem 5:

$$E = \frac{1}{2} \gamma m v^2 + mc^2$$



(quantum physics setting – just for fun...)

In quantum mechanics non just the extreme path contributes to the probability amplitude

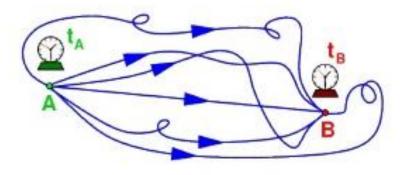
$$K(B, A) = \sum_{over all possible paths} \phi[x(t)]$$

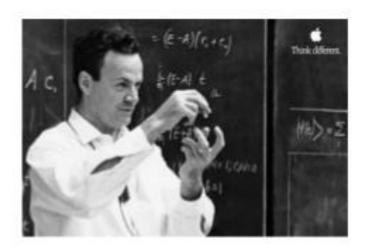
where $\Phi[x(t)] = A \exp\left\{\frac{i}{\hbar}S[x(t)]\right\}$

Feynman's path integral formula

$$K(B,A) = \int_{A}^{B} \exp\left(\frac{i}{\hbar}S[B,A]Dx(t)\right)$$

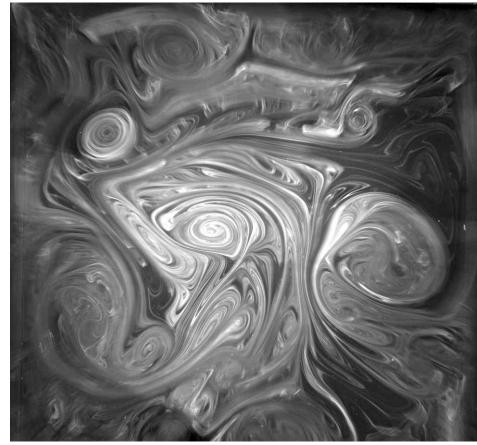
 $P(B, A) = |K(2, 1)|^2$





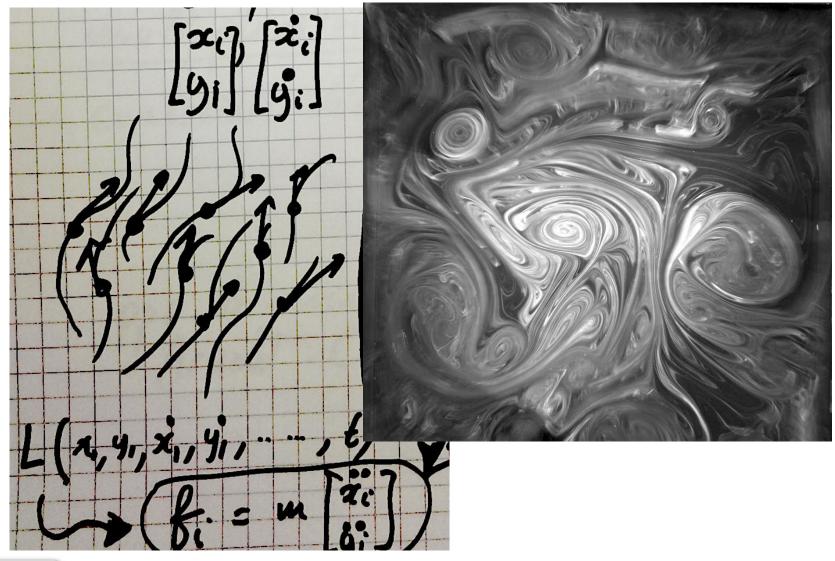


Fluids

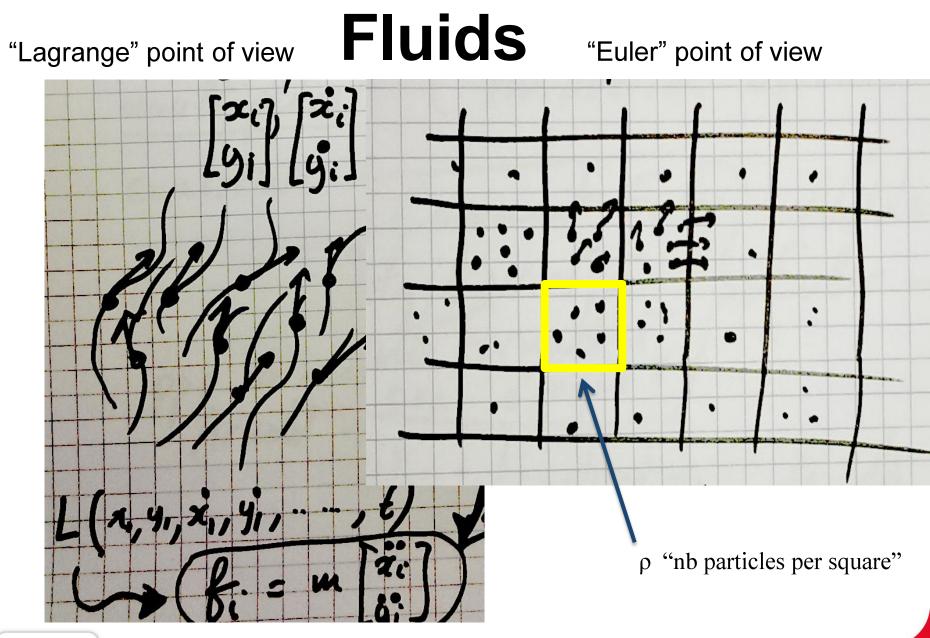




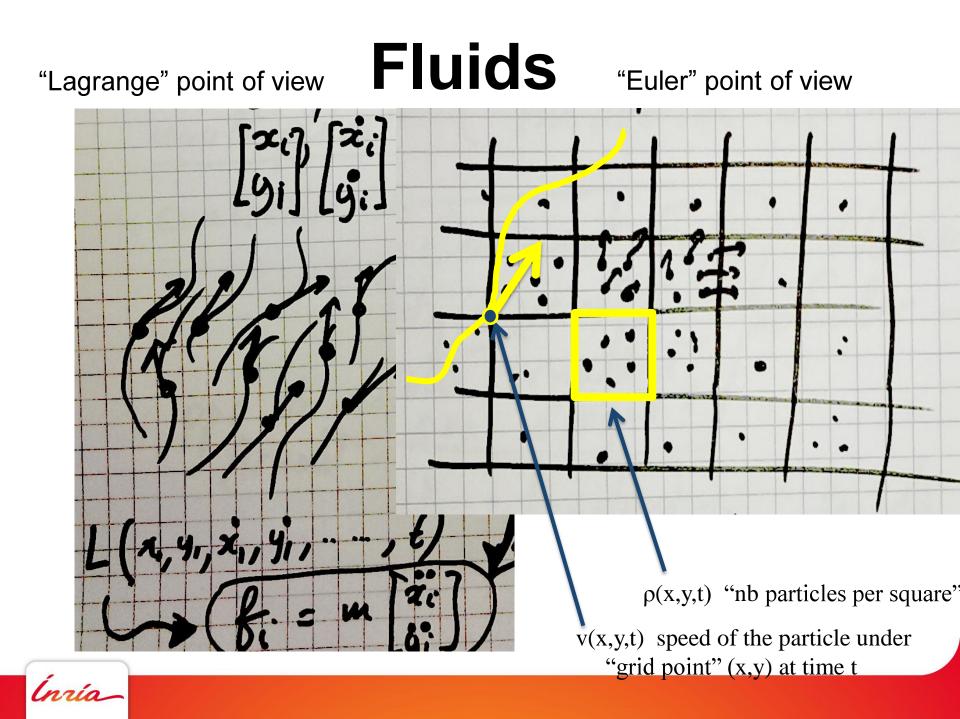
"Lagrange" point of view **Fluids**

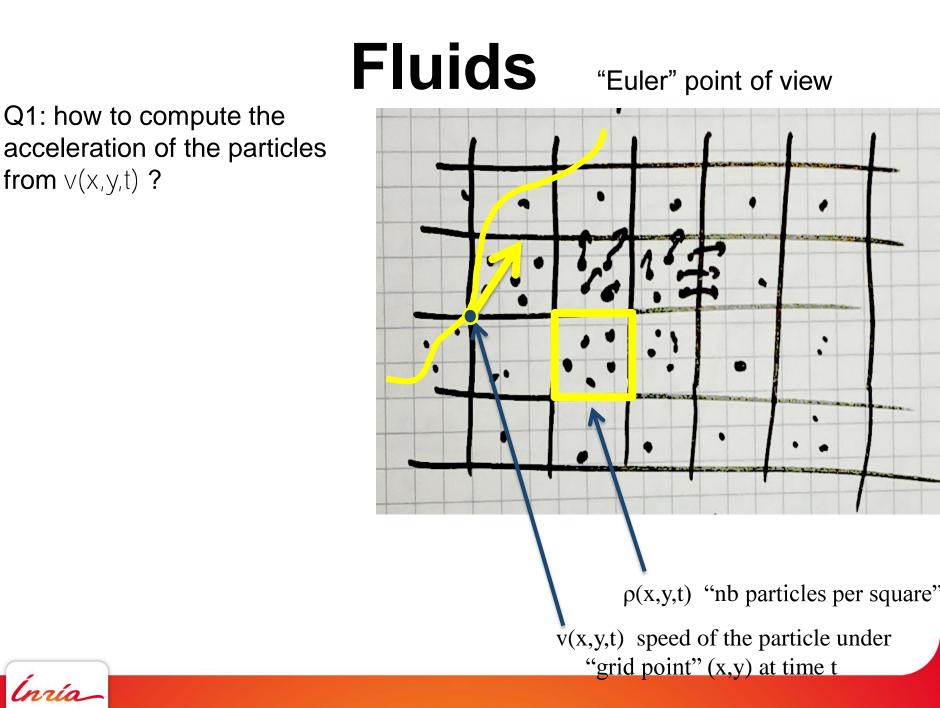


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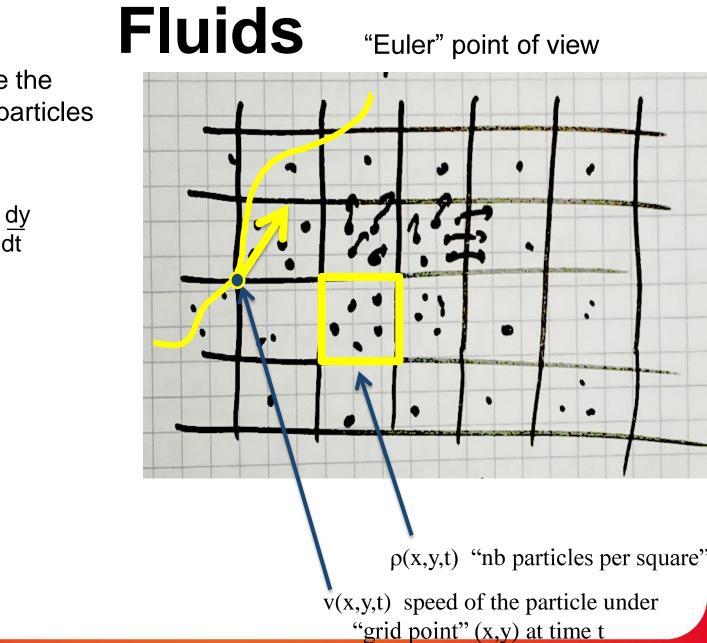


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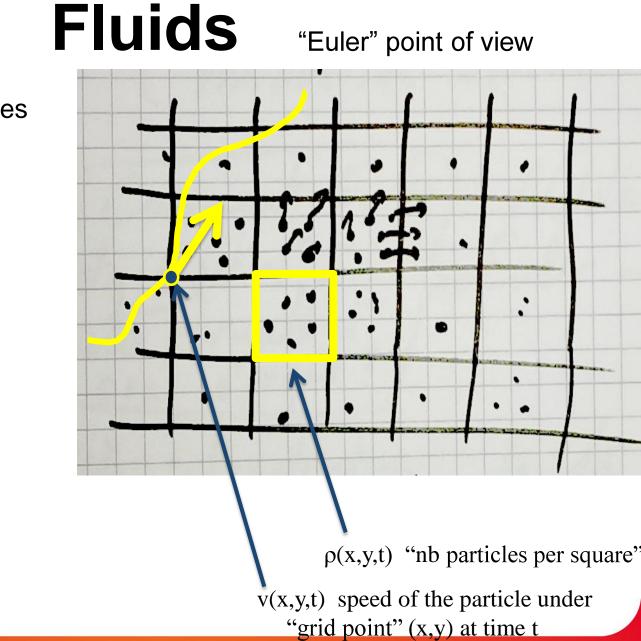




Q1: how to compute the acceleration of the particles from v(x,y,t) ?

 $\frac{\mathrm{d} \vee}{\mathrm{d} t} = \frac{\partial \mathbf{v}}{\partial t} + \frac{\partial \mathbf{v}}{\partial \mathbf{x}} \frac{\mathrm{d} \mathbf{x}}{\mathrm{d} t} + \frac{\partial \mathbf{v}}{\partial \mathbf{y}} \frac{\mathrm{d} \mathbf{y}}{\mathrm{d} t}$

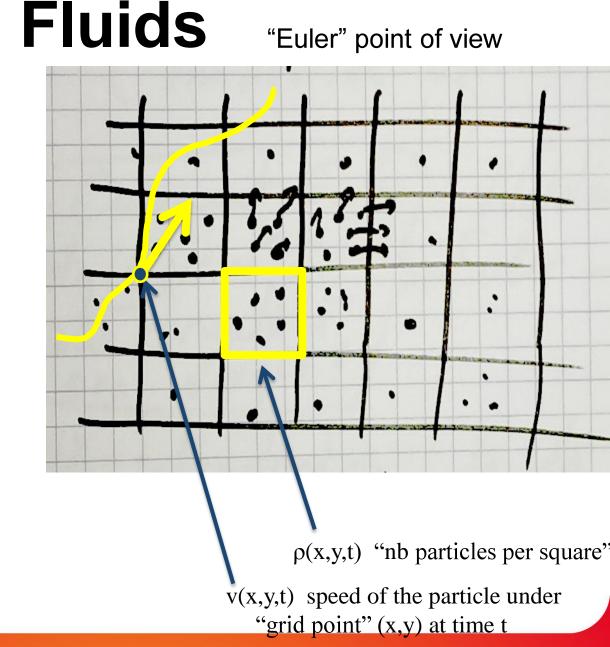
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Q1: how to compute the acceleration of the particles from
$$v(x,y,t)$$
 ?

$$\frac{\mathrm{d}v}{\mathrm{d}t} = \frac{\partial v}{\partial t} + \frac{\partial v}{\partial x}\frac{\mathrm{d}x}{\mathrm{d}t} + \frac{\partial v}{\partial y}\frac{\mathrm{d}y}{\mathrm{d}t}$$
$$= \frac{\partial v}{\partial t} + \frac{\partial v}{\partial x}\frac{v_x}{v_x} + \frac{\partial v}{\partial y}\frac{v_y}{v_y}$$

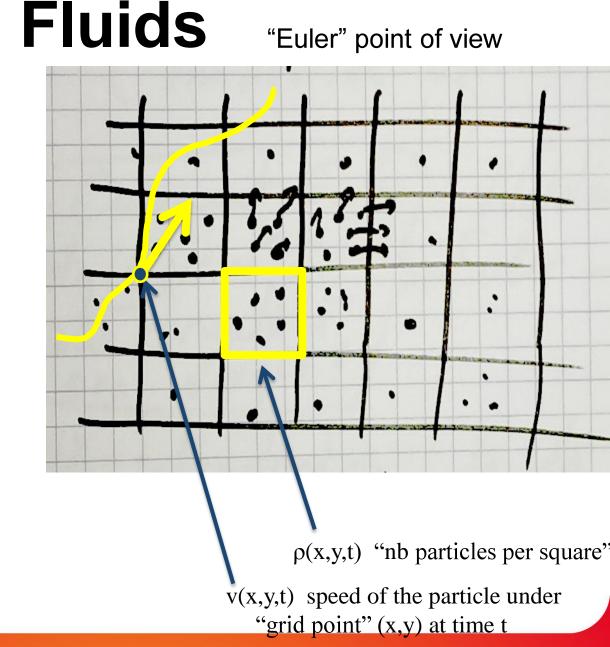
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Q1: how to compute the acceleration of the particles from v(x,y,t) ?

 $\frac{dV}{dt} = \frac{\partial v}{\partial t} + \frac{\partial v}{\partial x}\frac{dx}{dt} + \frac{\partial v}{\partial y}\frac{dy}{dt}$ $= \frac{\partial v}{\partial t} + \frac{\partial v}{\partial x}\frac{v_x}{v_x} + \frac{\partial v}{\partial y}\frac{v_y}{v_y}$ $= \frac{\partial v}{\partial t} + \frac{v.\nabla v}{v}$

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Q1: how to compute the acceleration of the particles from v(x,y,t) ?

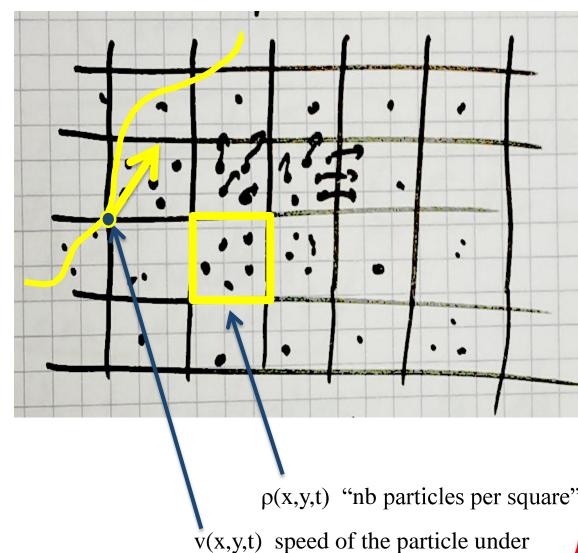
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Fluids "Euler" point of view

Q1: how to compute the acceleration of the particles from v(x,y,t) ?

- $\frac{\mathrm{d} v}{\mathrm{d} t} = \frac{\partial v}{\partial t} + \frac{v \cdot \nabla v}{v}$
- Q2: incompressible fluids ?



"grid point" (x,y) at time t

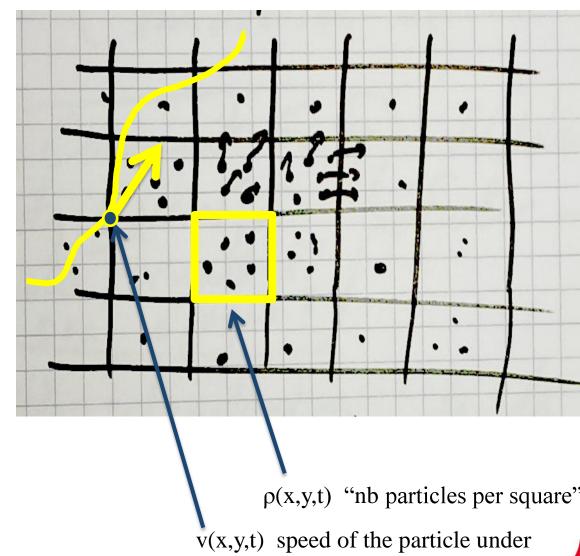




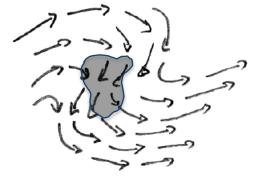
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"grid point" (x,y) at time t





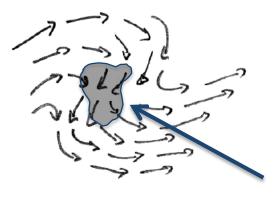
Fluids "Euler" point of view Q1: how to compute the acceleration of the particles from v(x, y, t) ? $\frac{\mathrm{d} \mathrm{v}}{\mathrm{d} \mathrm{t}} = \frac{\partial \mathrm{v}}{\partial \mathrm{t}} + \mathrm{v} \cdot \nabla \mathrm{v}$ Q2: incompressible fluids ? $\rho(x,y,t)$ "nb particles per square" what goes in = v(x,y,t) speed of the particle under what goes out "grid point" (x,y) at time t Ínría

Fluids "Euler" point of view

Q1: how to compute the acceleration of the particles from v(x,y,t) ?

- $\frac{\mathrm{d} v}{\mathrm{d} t} = \frac{\partial v}{\partial t} + \frac{v \cdot \nabla v}{v}$
 - Q2: incompressible fluids ?

div(v) = 0



what goes in = what goes out



 $\rho(x,y,t)$ "nb particles per square"

v(x,y,t) speed of the particle under "grid point" (x,y) at time t

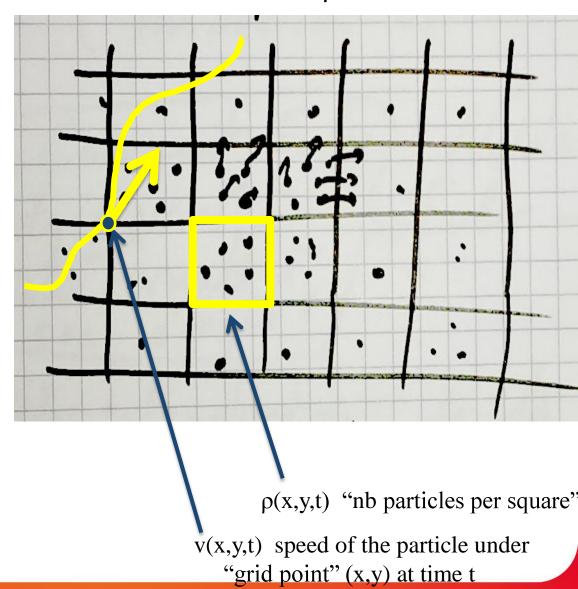


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- Q2: incompressible fluids ? $di \lor (\lor) = 0$

Q3: mass preservation ?





Fluids "Euler" point of view $\rho(x,y,t)$ "nb particles per square" v(x,y,t) speed of the particle under

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Q2: incompressible fluids ?

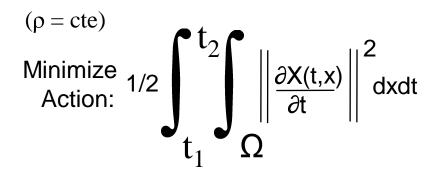
div(v) = 0

Q3: mass preservation ?

$$\frac{d \rho}{dt} = - di \vee (\rho \vee)$$

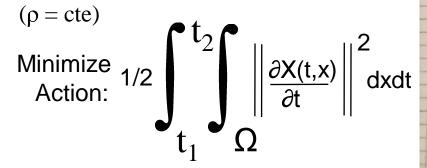
(Continuity equation)

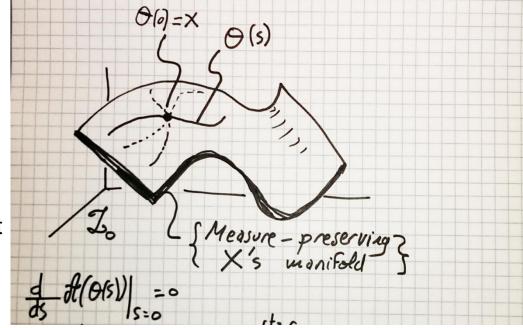
Start with Lagrange coordinates: "particle trajectories": X(t, X)



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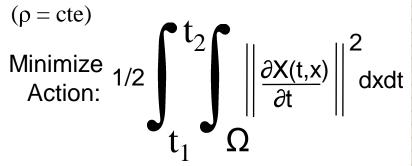
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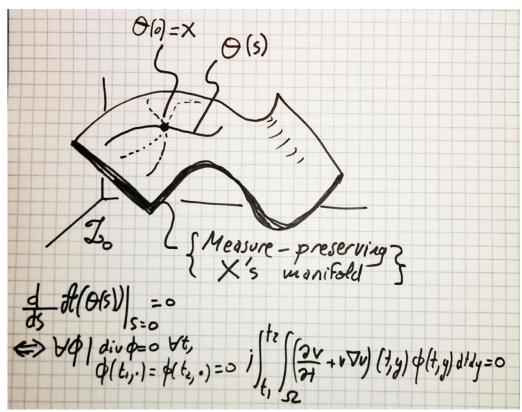




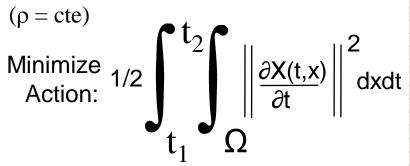


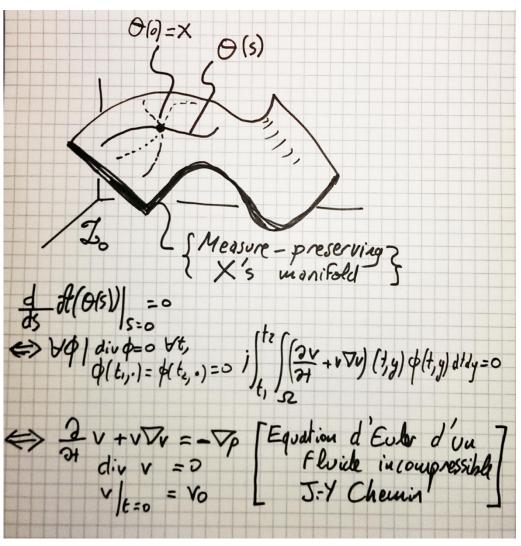
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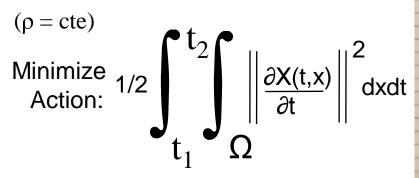
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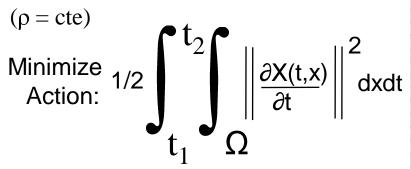


s.t. X satisfies mass preservation (X is measure-preserving, more on this later...)

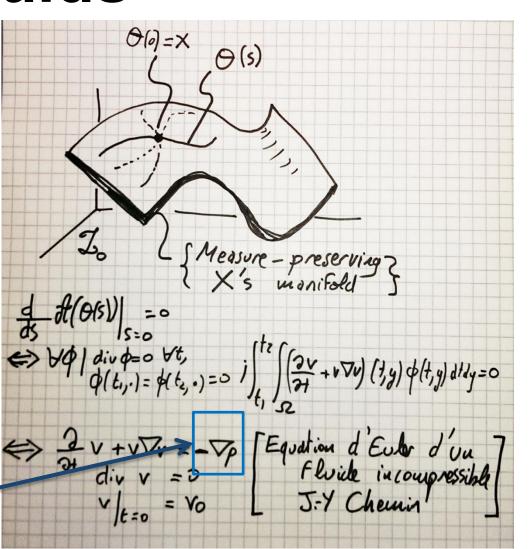
Acceleration of the **final** particle "under the grid"

00=X $\Theta(s)$ S Measure - preserving ? वमु $\Rightarrow \forall \phi \mid div \phi = 0 \forall t, \\ \phi(t_{ij}) = \phi(t_{ij}, j) = 0 \quad j \int_{t_{ij}}^{t_{ij}} \int_{t_{ij}}^{t_{ij}} \int_{t_{ij}}^{t_{ij}} \left(\frac{2v}{2t} + v \nabla v \right) (t_{ij}) \phi(t_{ij}) dt dy = 0$ -Vp [Equation d'Euler d'un Fluide incompressible To J=Y Chemin

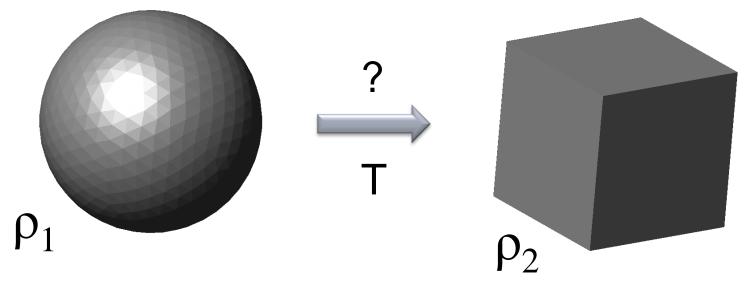
Start with Lagrange coordinates: "particle trajectories": X(t, X)



s.t. X satisfies mass preservation (X is measure-preserving, more on this later...)

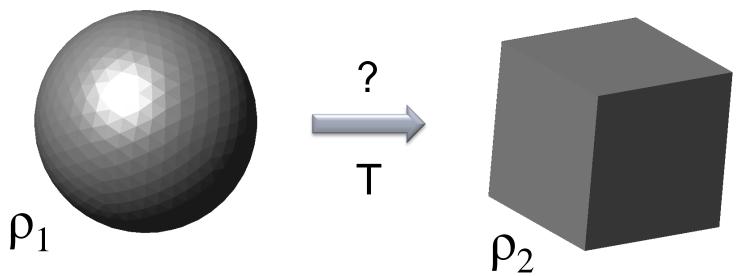


Fluids – Benamou Brenier





Fluids – Benamou Brenier



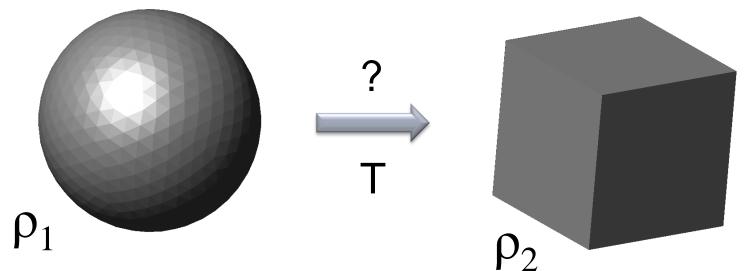
Minimize

$$A(\rho, v) = (t_2 - t_1) \int_{t_1}^{t_2} \int_{\Omega} \rho(x, t) ||v(t, x)||^2 dx dt$$

$$s.t. \ \rho(t_1, .) = \rho_1 \quad ; \quad \rho(t_2, .) = \rho_2 \quad ; \quad \frac{d \rho}{dt} = - \operatorname{div}(\rho v)$$



Fluids – Benamou Brenier



Minimize

$$A(\rho,v) = (t_2-t_1) \int_{t_1}^{t_2} \rho(x,t) ||v(t,x)||^2 dx dt$$

$$t_1 \Omega$$
S.t. $\rho(t_1,.) = \rho_1$; $\rho(t_2,.) = \rho_2$; $\frac{d \rho}{dt} = - \operatorname{div}(\rho \vee)$
Minimize $C(T) = \int_{\Omega}^{t_1} \rho(x) ||x - T(x)||^2 dx$
S.t. T is measure-preserving

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Optimal Transport an elementary introduction







(X;µ)

(Y;v)

Two measures
$$\mu$$
, v such that $\int_X d\mu(x) = \int_Y dv(x)$







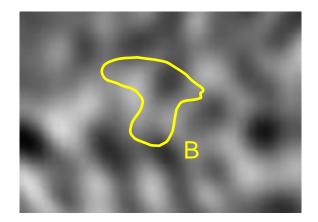
(X;µ)

(Y;v)

A map T is a *transport map* between μ and ν if $\mu(T^{-1}(B)) = \nu(B)$ for any Borel subset B of Y





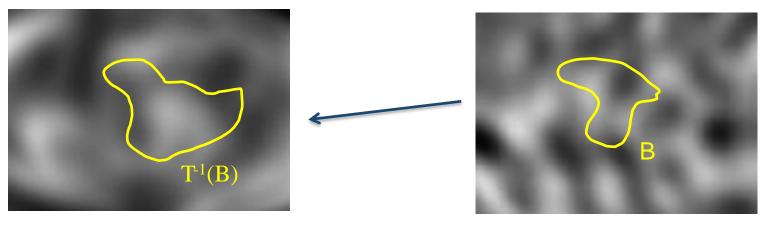


(X;µ)

(Y;v)

A map T is a *transport map* between μ and \vee if $\mu(T^{-1}(B)) = \vee(B)$ for any Borel subset B









A map T is a *transport map* between μ and ν if $\mu(T^{-1}(B)) = \nu(B)$ for any Borel subset B







(X;µ)

(Y;v)

A map T is a *transport map* between μ and ν if $\mu(T^{-1}(B)) = \nu(B)$ for any Borel subset B (or $\nu = T \# \mu$ the *pushforward* of μ)









(Y;v)

Monge's problem (1787): Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$



Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

- Difficult to study
- If μ has an atom (isolated Dirac), it can only be mapped to another Dirac (T needs to be a map)



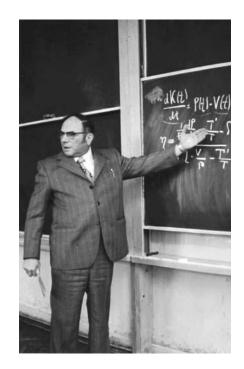
Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem (1942):

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\nu(y)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$

that minimizes
$$\iint_{X \times Y} || x - y ||^2 d_{Y(x,y)}$$





Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\nu(y)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$

" $\gamma(x,y)$ " : How much sand goes from x to y

that minimizes $\iint_{X \times Y} || x - y ||^2 d_{\gamma(x,y)}$



<u>Monge's problem:</u>

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem:

Find a measure
$$\gamma$$
 defined on X x Y
such that $\int_{X \text{ in } X} d\gamma(x,y) = dv(y)$
and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$
Everything that is
transported from x sums to " $\mu(x)$ "

that minimizes $\iint_{X \times Y} || x - y ||^2 d_{\gamma(x,y)}$



<u>Monge's problem:</u>

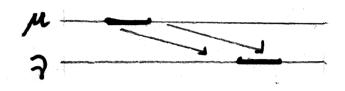
Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem:

Find a measure
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such that $\int_{X \text{ in } X} d\gamma(x,y) = dv(y)$
and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$
Everything that is
transported to y sums to "v(y)"

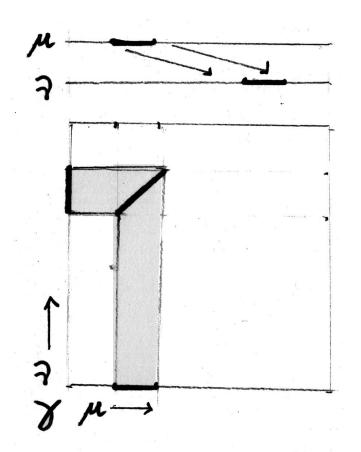
that minimizes $\iint_{X \times Y} || x - y ||^2 d_{Y(x,y)}$





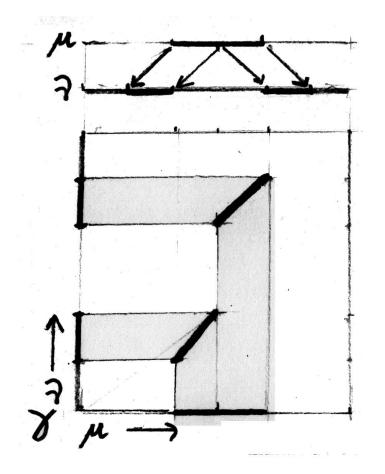
Transport plan – example 1/2 : translation of a segment





Transport plan – example 1/2 : translation of a segment

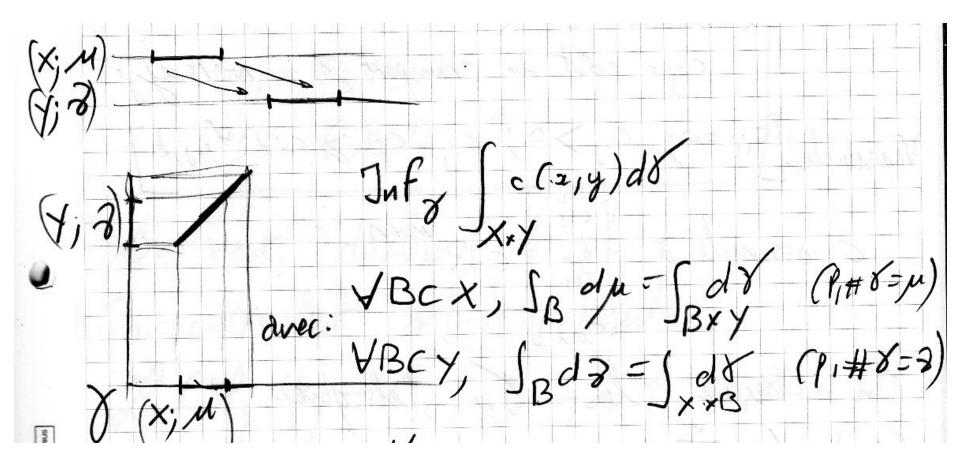




Transport plan – example 2/2 : spitting a segment

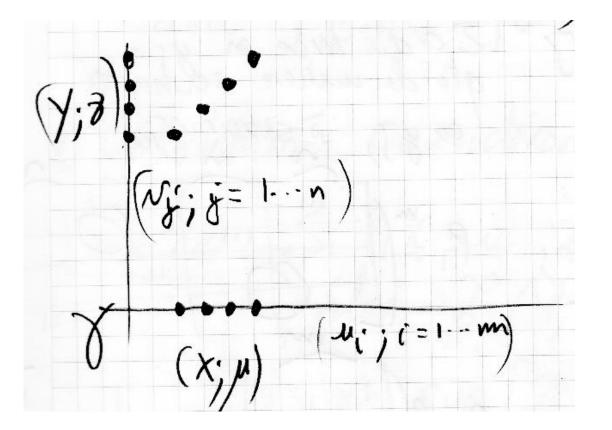


Part. 2 Optimal Transport – Duality



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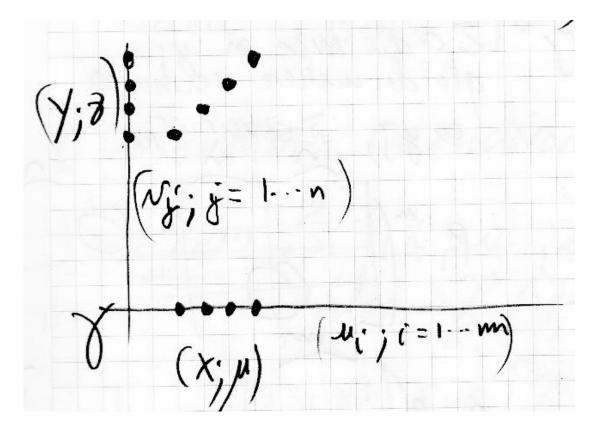
Part. 2 Optimal Transport – Duality



Duality is easier to understand with a discrete version Then we'll go back to the continuous setting.



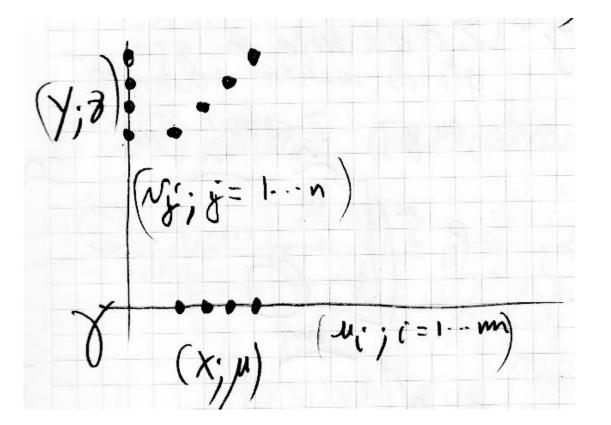
Part. 2 Optimal Transport – Duality



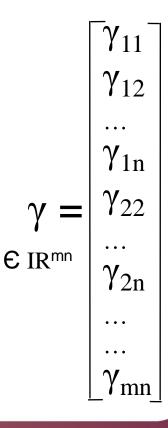
(DMK): Min <C, γ > $s.t. \begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$



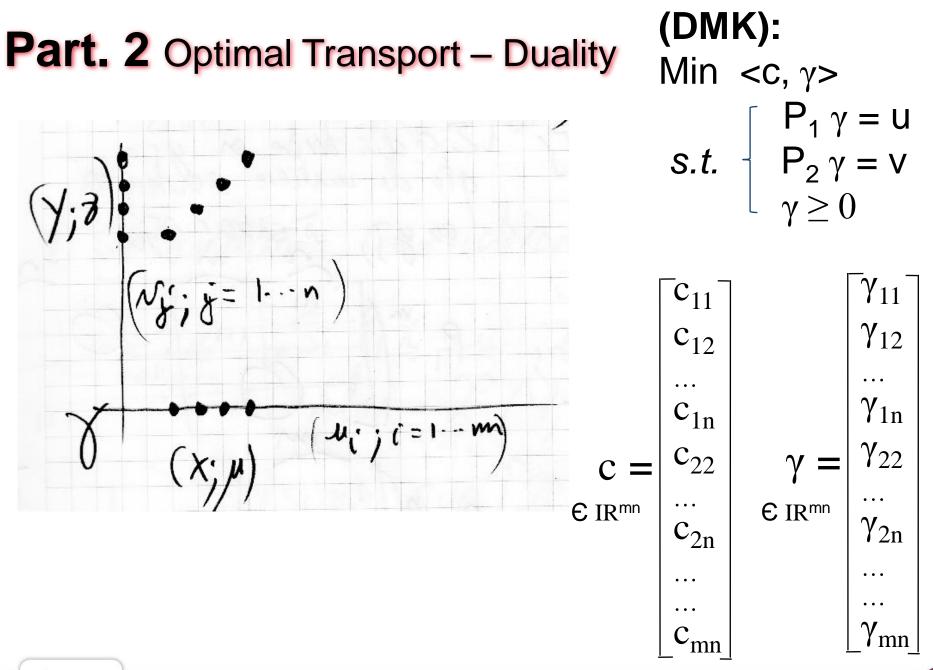
Part. 2 Optimal Transport – Duality



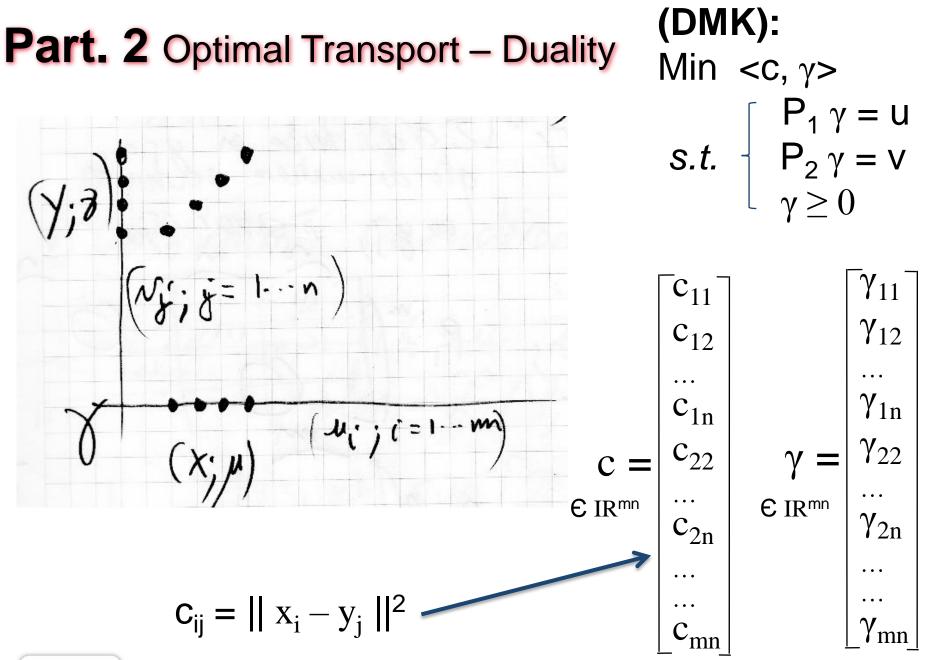
(DMK): Min <C, γ > s.t. $\begin{cases}
P_1 \gamma = u \\
P_2 \gamma = v \\
\gamma \ge 0
\end{cases}$

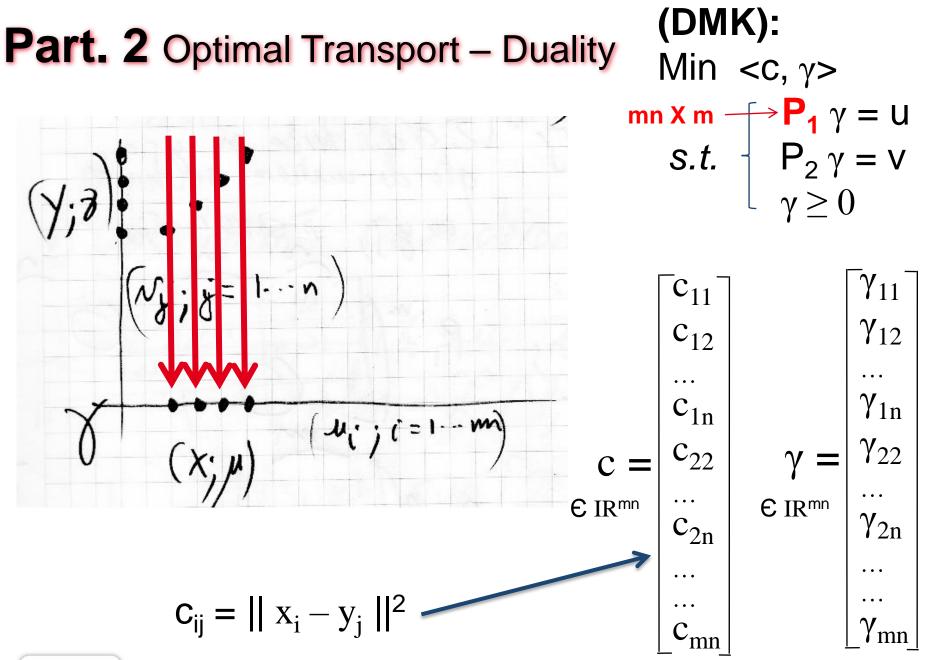




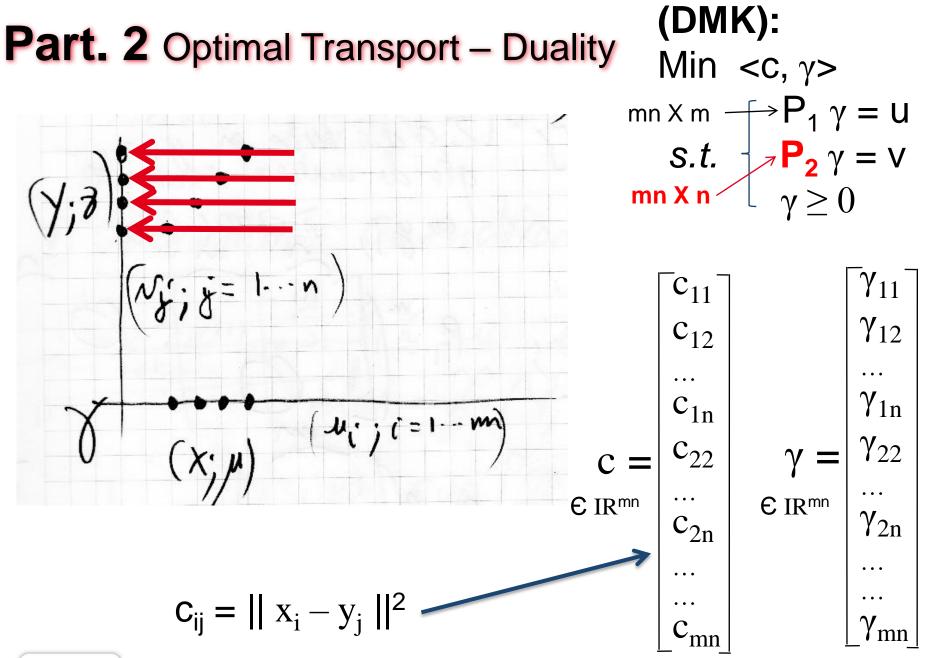


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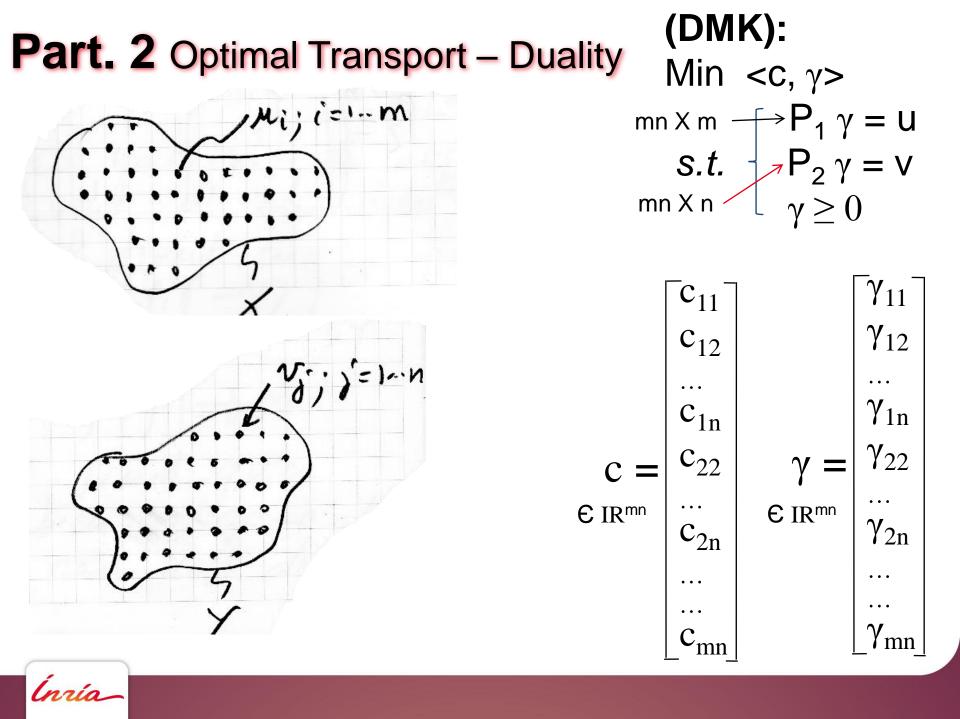




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Part. 2 Optimal Transport – Duality

< u, v > denotes the dot product between u and v

(DMK): Min <C, γ > $\begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$

Consider $\lfloor (\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$



Part. 2 Optimal Transport – Duality(DMK):
Min <C, γ >
S.t.Min <C, γ >
P₁ γ = U
P₂ γ = V
 $\gamma \ge 0$

Consider
$$\lfloor (\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$



Part. 2 Optimal Transport – Duality(DMK):
Min <C, γ >
S.t.Min <C, γ >
P₁ γ = U
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Part. 2 Optimal Transport – Duality(DMK):
Min <C, γ >
S.t.Min <C, γ >
P₂ γ = V
 $\gamma \ge 0$

Consider
$$\lfloor (\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$

 $\begin{array}{l} \text{Consider now: Inf} \left[\begin{array}{c} \text{Sup}[\ \ \Box(\phi,\psi) \end{array} \right] \right] \\ \gamma \geq 0 \quad \begin{array}{c} \phi \ \in \mathrm{IR}^m \\ \psi \in \mathrm{IR}^n \end{array} \end{array}$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma>
S.t.
$$\begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$$

Consider
$$\lfloor (\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma>
S.t.
$$\begin{cases} P_1 \ \gamma = u \\ P_2 \ \gamma = v \\ \gamma \ge 0 \end{cases}$$

Consider
$$\lfloor (\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$

Consider now: Inf $\begin{bmatrix} Sup[\sqcup(\phi, \psi) \end{bmatrix} \end{bmatrix} = Inf [< C, \gamma >]$ (DMK) $\gamma \ge 0 \quad \substack{\phi \in IR^m \\ \psi \in IR^n} \quad \substack{\gamma \ge 0 \\ P_1 \gamma = u \\ P_2 \gamma = v}$



(DMK): **Part. 2** Optimal Transport – Duality Min < c, γ > s.t. $\begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$ Inf [Sup[<c, γ > - < ϕ , P₁ γ - u> - < ψ , P₂ γ - v>]] $\gamma \geq 0 \quad \phi \ \in IR^m$ $\psi \in IR^n$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma >
S.t.
$$\begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$$
Inf $\begin{bmatrix} Sup[- <\phi, P_1 \gamma - u > - <\psi, P_2 \gamma - v >] \end{bmatrix}$ $\gamma \ge 0 \quad \end{v} \in IR^m \quad \end{v} \in IR^m \quad \end{v} \in IR^m \quad \end{v} = 0$ Sup[Inf[\gamma > - <\phi, P_1 \gamma - u > - <\psi, P_2 \gamma - v >] \end{bmatrix} $\varphi \in IR^m \quad \end{v} \neq 0$

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$$\begin{array}{l} \mbox{(DMK):}\\ \mbox{Min } < c, \gamma > \\ \mbox{S.t.} & \left[\begin{array}{c} P_1 \ \gamma = u \\ P_2 \ \gamma = v \\ \gamma \ge 0 \end{array} \right] \\ \mbox{Y } \geq 0 \end{array} \\ \begin{array}{l} \mbox{Inf} \left[\begin{array}{c} \mbox{Sup} \left[< c, \gamma > - < \phi, P_1 \ \gamma - u > - < \psi, P_2 \ \gamma - v > \right] \right] \\ \mbox{Y } \geq 0 \end{array} \\ \mbox{Sup} \left[\begin{array}{c} \mbox{Inf} \left[< c, \gamma > - < \phi, P_1 \ \gamma - u > - < \psi, P_2 \ \gamma - v > \right] \right] \\ \mbox{Exchange Inf Sup} \end{array} \\ \begin{array}{c} \mbox{Sup} \left[\begin{array}{c} \mbox{Inf} \left[< c, \gamma > - < \phi, P_1 \ \gamma - u > - < \psi, P_2 \ \gamma - v > \right] \right] \\ \mbox{Exchange Inf Sup} \end{array} \\ \mbox{Sup} \left[\begin{array}{c} \mbox{Inf} \left[< c, \gamma > - < \phi, P_1 \ \gamma - u > - < \psi, P_2 \ \gamma - v > \right] \right] \\ \mbox{Expand/Reorder/Collect} \end{array} \\ \mbox{Sup} \left[\begin{array}{c} \mbox{Inf} \left[< \gamma, c - P_1^t \ \phi - P_2^t \ \psi > + < \phi, u > + < \psi, v > \right] \right] \\ \mbox{W } \in \mbox{IR}^n \end{array} \right] \end{array}$$

Part. 2 Optimal Transport – Duality(DMK):
Min \gamma >
$$S.t. \begin{bmatrix} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{bmatrix}$$
Inf $\begin{bmatrix} Sup[- <\phi, P_1 \gamma - u > - <\psi, P_2 \gamma - v >] \end{bmatrix}$
 $\gamma \ge 0 \quad \psi \in IR^n$ Exchange Inf SupSup[Inf[\gamma > - <\phi, P_1 \gamma - u > - <\psi, P_2 \gamma - v >] \end{bmatrix}
 $\psi \in IR^n$ Expand/Reorder/CollectSup[Inf[< $\gamma, c - P_1^t \phi - P_2^t \psi > + <\phi, u > + <\psi, v >] \end{bmatrix}$
 $\psi \in IR^n$ Interpret

Part. 2 Optimal Transport – Duality(DMK):
Min
$$<$$
C, $\gamma >$
S.t. $\begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$

$$Sup[Inf[< \gamma, c-P_1^t \phi - P_2^t \psi > + <\phi, u > + <\psi, v >]]$$

$$\varphi \in IR^m \quad \gamma \ge 0$$

$$\psi \in IR^n$$

Interpret

$$Sup[<\phi,u> + <\psi, v>]$$
(DDMK)
$$\varphi \in IR^{m}$$

$$\psi \in IR^{n}$$

$$P_{1}^{t} \varphi + P_{2}^{t} \psi \leq C$$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma >
$$\begin{bmatrix} P_1 \gamma = U \\ P_2 \gamma = V \\ \gamma \ge 0 \end{bmatrix}$$
 $\sup [\inf [< \gamma, c - P_1^t \phi - P_2^t \psi > + <\phi, u > + <\psi, v >]] \\ \psi \in IR^n \\ \psi \in IR^n \end{bmatrix}$ Interpret $\sup [<\phi, u > + <\psi, v >]$
 $\psi \in IR^n \\ \psi \in IR^n \\ P_i^t \phi + P_2^t \psi \le c$ (DDMK)
 $\phi_i + \psi_j \le c_{ij} \quad \forall (i,j)$

Kantorovich's problem:

Find a measure
$$\gamma$$
 defined on X x Y
such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$
and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\nu(x)$
that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Dual formulation of Kantorovich's problem (Continuous):

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi d\mu + \int_Y \psi dv$

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \leq \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi d\mu + \int_Y \psi dv$



Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in $L^1(\mu)$ and ψ in $L^1(v)$ Such that for all x,y, $\varphi(x) + \psi(y) \leq \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi d\mu + \int_Y \psi dv$

Point of view of a "transport company": Try to maximize transport price

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$

What they charge for loading at x

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$ What they charge for loading at x What they charge for unloading at y

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\nu(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v)

Such that for all x,y, $\varphi(x) + \psi(y) \leq \frac{1}{2} ||x - y||^2$

that maximize $\int_X \phi(x) d\mu + \int_Y \psi(y) dv$

Your point of view: Try to minimize transport cost

Price (loading + unloading) cannot be greater than transport cost (else you do the job yourself)

What they charge for loading at x

What they charge for unloading at y

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$



Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$

If we got two functions φ and ψ that satisfy the constraint

Then it is possible to obtain a better solution by replacing ψ with φ^c defined by: For all y, $\varphi^c(y) = \inf_{x \text{ in } X} \frac{1}{2} ||x - y||^2 - \varphi(y)$



Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(V) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dV$

If we got two functions ϕ and ψ that satisfy the constraint

Then it is possible to obtain a better solution by replacing ψ with ϕ^c defined by: For all y, $\phi^c(y) = \inf_{x \text{ in } X} \frac{1}{2} ||x - y||^2 - \phi(y)$

- ϕ^c is called the **c-conjugate** function of ϕ
- If there is a function φ such that $\psi = \varphi^c$ then ψ is said to be **c-concave**
- If ψ is c-concave, then $\psi^{cc} = \psi$

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes $\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$

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Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

 ψ is called a **"Kantorovich potential"**

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Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

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$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

 ψ is called a **"Kantorovich potential"**

What about our initial problem ?



Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

 ψ is called a **"Kantorovich potential"**

What about our initial problem ? (i.e., this is T() that we want to find ...)



Theorem 1.

 $\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$

where $\partial_c \psi = \{(x,y) | \phi(x) + \psi(y) = c(x,y)\}$ denotes the so-called c-subdifferential of ψ .



Theorem 1.

 $\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$

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Proof: see OTON, chap. 10.

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Proof: see OTON, chap. 10.

Heuristic argument (at the beginning of the same chapter):

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Theorem 1.

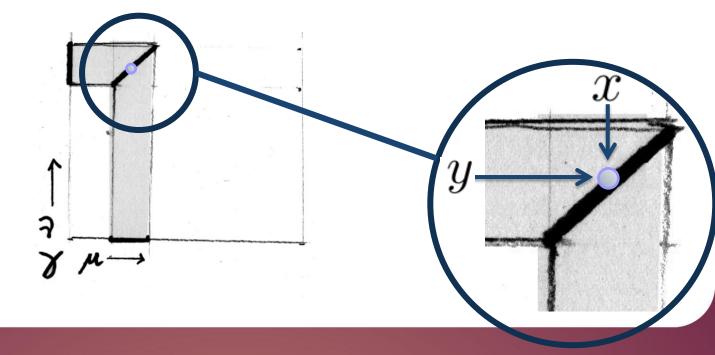
$$\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$$

where $\partial_c \psi = \{(x,y) | \phi(x) + \psi(y) = c(x,y)\}$ denotes the so-called c-subdifferential of ψ .

Proof: see OTON, chap. 10.

Heuristic argument (at the beginning of the same chapter):

Consider a point (x, y) on the c-subdifferential $\partial_c \psi$, that satisfies $\phi(y) + \psi(x) = c(x, y)$ (1).





Theorem 1.

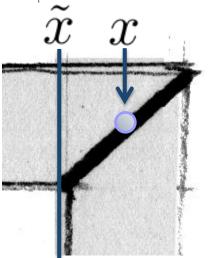
$$\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$$

where $\partial_c \psi = \{(x,y) | \phi(x) + \psi(y) = c(x,y)\}$ denotes the so-called c-subdifferential of ψ .

Proof: see OTON, chap. 10.

Heuristic argument (at the beginning of the same chapter):

Consider a point (x, y) on the c-subdifferential $\partial_c \psi$, that satisfies $\phi(y) + \psi(x) = c(x, y)$ (1). By definition, $\phi(y) = \psi^c(y) = \inf_x c(x, y) - \psi(x)$, thus $\forall \tilde{x}, \phi(y) \leq c(\tilde{x}, y) - \psi(\tilde{x})$, or $\phi(y) + \psi(\tilde{x}) \leq c(\tilde{x}, y)$ (2). By substituting (1) into (2), one gets $\psi(\tilde{x}) - \psi(x) \leq c(\tilde{x}, y) - c(x, y)$ for all \tilde{x} .





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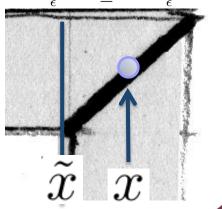
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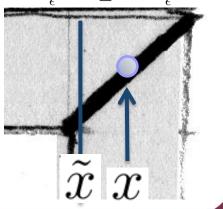
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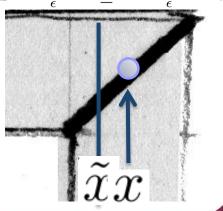
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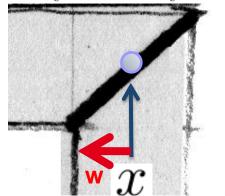
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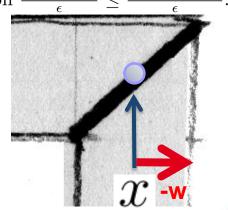
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The same derivation can be done with -w instead of w, and one gets:

$$\forall w, \nabla \psi(x) \cdot w = \nabla_x c(x, y) \cdot w$$
, thus $\forall (x, y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x, y) = 0$.





Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

In the L_2 case, i.e. $c(x,y) = 1/2||x-y||^2$, we have $\forall (x,y) \in \partial_c \psi, \nabla \psi(x) + y - x = 0$, thus, whenever the optimal transport map T exists, we have $T(x) = x - \nabla \psi(x) = \nabla (||x||^2/2 - \psi(x))$.



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Part. 2 Optimal Transport – convexity

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Proof.

$$\begin{split} \psi(x) &= \inf_{y} \frac{|x-y|^2}{2} - \phi(y) \\ &= \inf_{y} \frac{||x||^2}{2} - x \cdot y + \frac{||y||^2}{2} - \phi(y) \\ -\bar{\psi}(x) &= \phi(x) - \frac{||x||^2}{2} = \inf_{y} -x \cdot y + \left(\frac{||y||^2}{2} - \phi(y)\right) \\ \bar{\psi}(x) &= \sup_{y} x \cdot y - \left(\frac{||y||^2}{2} - \phi(y)\right) \end{split}$$



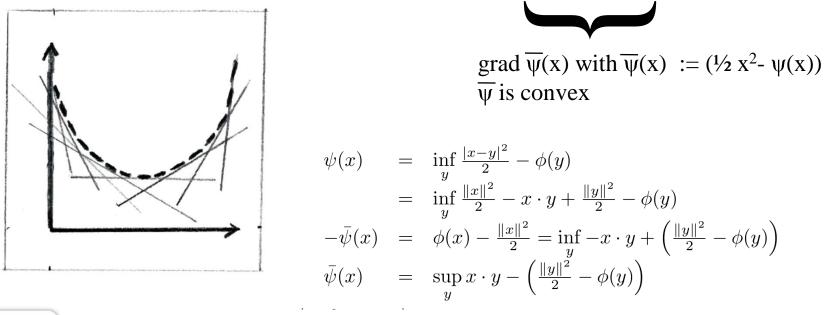
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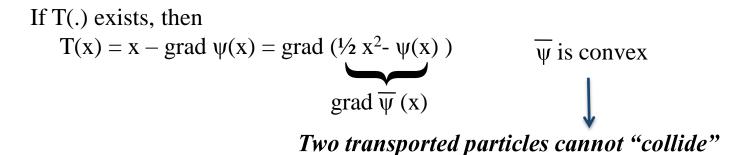


Part. 2 Optimal Transport – no collision

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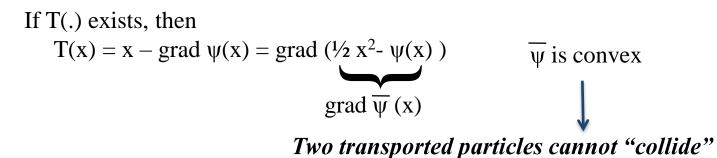


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Proof. By contradiction, suppose that you have $t \in (0, 1)$ and $x_1 \neq x_2$ such that:

$$(1-t)x_1 + tT(x_1) = (1-t)x_2 + tT(x_2)$$

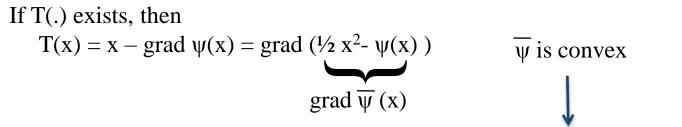


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Two transported particles cannot "collide"

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Part. 2 Optimal Transport – Monge-Ampere

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

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What about our initial problem ? If T(.) exists, then one can show that: T(x) = x - grad $\psi(x)$ = grad (¹/₂ x²- $\psi(x)$)

grad $\overline{\psi}(x)$ with $\overline{\psi}(x) := (\frac{1}{2} x^2 - \psi(x))$

for all borel set A, $\int_A d\mu = \int_{T(A)} (|JT|) dv$ (change of variable)

Jacobian of T (1st order derivatives)



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Det. of the Hessian of $\overline{\psi}$ (2nd order derivatives)



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When μ and ν have a density u and v , $(H \overline{\psi}(x))$. $v(\text{grad } \overline{\psi}(x)) = u(x)$ Monge-Ampère equation

Part. 2 Optimal Transport – summary

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$



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Brenier, Mc Cann, Trudinger: *The optimal transport map is then given by:* $T(x) = \text{ grad } \overline{\psi}(x)$

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For a given volume, ball is the shape that minimizes border area



L₁ **Sobolev inegality:** Given f: $IR^n \rightarrow IR$ sufficiently regular

$$\int |\operatorname{grad} f| \ge n \operatorname{Vol}(B_2^n)^{1/n} \left(\int f^{n/(n-1)} \right) (n-1)/n$$

Explanation in [Dario Cordero Erauquin] course notes



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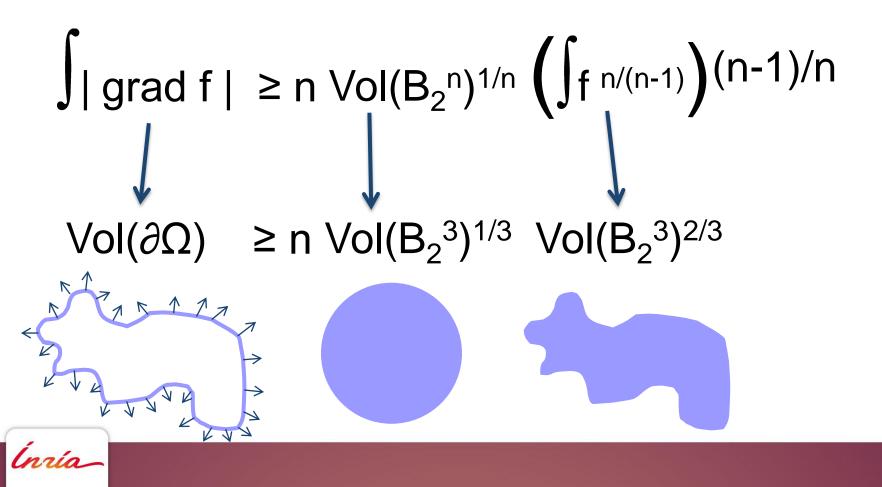
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L₁ Sobolev inegality: a proof with OT [Gromov]

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We suppose w.l.o.g. that
$$\int f^{n/(n-1)} = 1$$

det (Hess $\overline{\Psi}$) $^{1/n} \leq (\Delta \overline{\Psi})/n$
 $Vol(B_2^n) f^{n/(n-1)}(x) = det Hess \overline{\Psi}$
 $Vol(B_2^n) = Vol(B_2^n) \int f^{n/(n-1)} = \int f Vol(B_2^n) f^{1/(n-1)}$



L₁ Sobolev inegality: a proof with OT [Gromov]

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 $\int |\operatorname{grad} f| \ge n \operatorname{Vol}(B_2^n)^{1/n}$

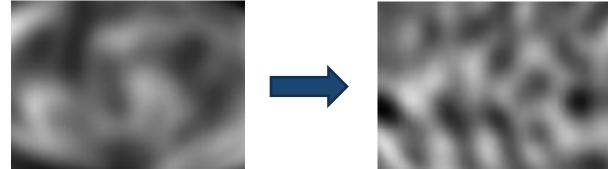




Semi-Discrete Optimal Transport

Ínría

Part. 3 Optimal Transport – how to program ? $(X;\mu)$ (Y;v)



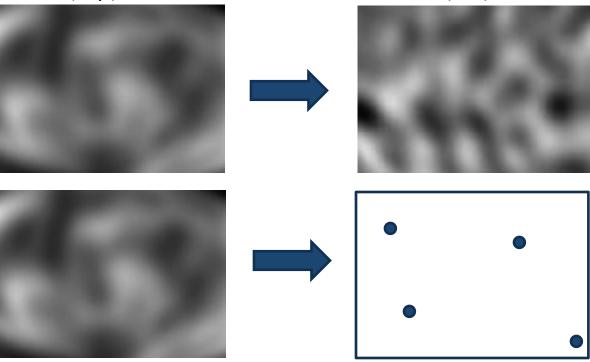
Continuous



Part. 3 Optimal Transport – how to program ? (X;µ) (Y;v)

Continuous

Semi-discrete





Part. 3 Optimal Transport – how to program ? (X;µ) (Y;v)

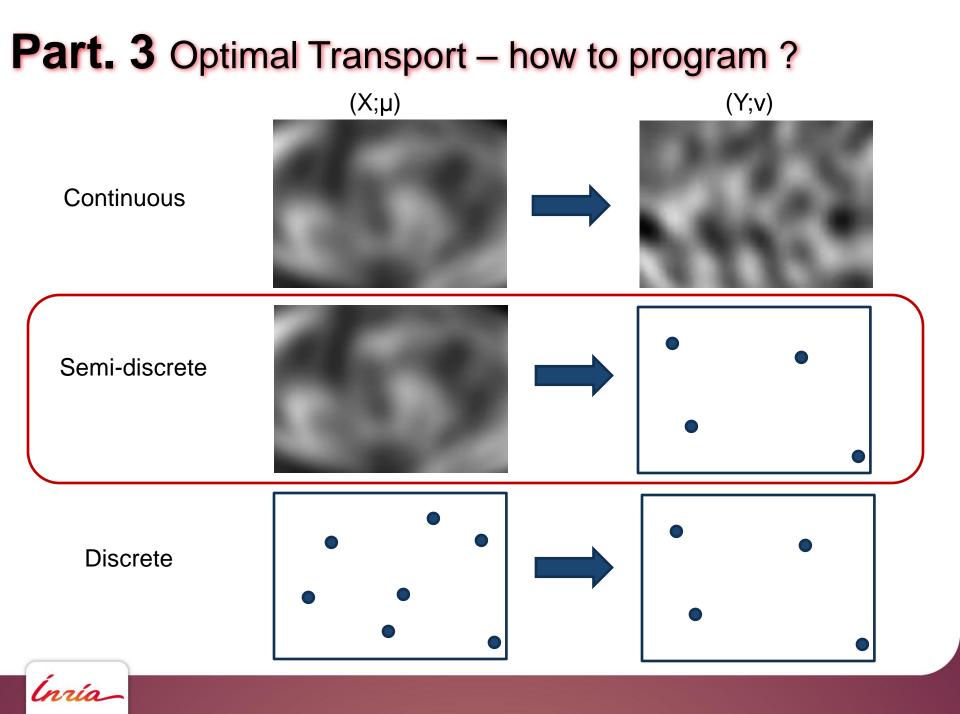
Continuous

Semi-discrete

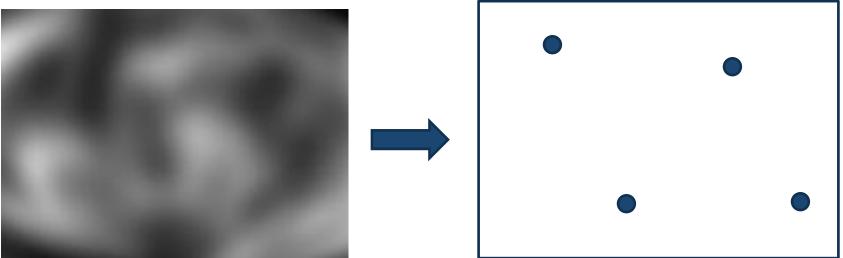


Discrete





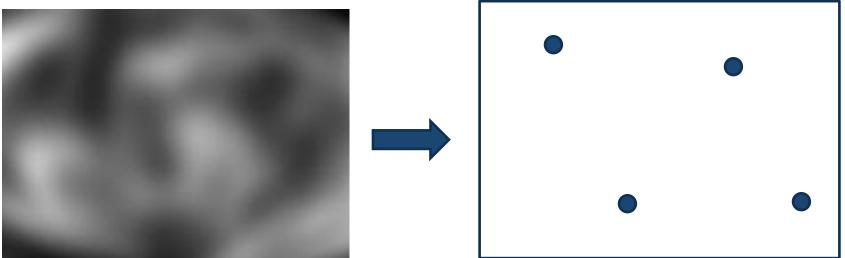
Part. 3 Optimal Transport – semi-discrete (X;µ) (Y;v)



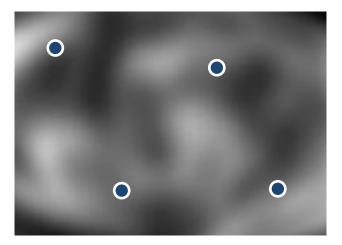
(DMK)
$$\sup_{\psi \in \psi^c} \int_X \psi^c(x) d\mu + \int_Y \psi(y) d\nu$$



Part. 3 Optimal Transport – semi-discrete (X;µ) (Y;v)

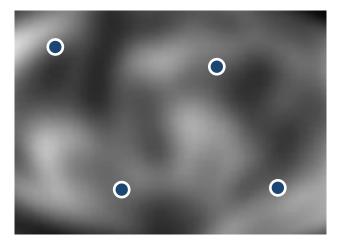






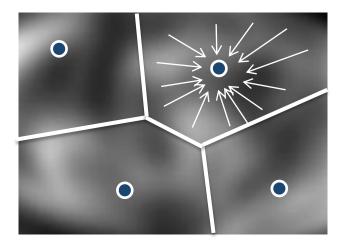
$$(\text{DMK}) \quad \begin{array}{l} \underset{\psi \in \psi^c}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{C}}{\overset{\text{W}}{\overset{\text{C}}{\overset{\text{C}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}}}}}}}}}}}}}}}}}}}}}} } } \\ \\ \overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}}}}}} {\overset{\text{Sup}}{\overset{\text{Sup}}}}{\overset{Sup}}}}}} } \\} \overset{\text{Sup}}{\overset{Sup}}}}} \\} \overset{\text{Sup}}}{\overset{Sup}}}}} } \\} \overset{\text{Sup}}}{\overset{Sup}}}} \overset{\text{Sup}}}{\overset{Sup}}}}} } \\} \overset{\text{Sup}}}{\overset{Sup}}}} \overset{\text{Sup}}}{\overset{Sup}}}} } \overset{Sup}}} } \overset{Sup}}} \\} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}} \overset{Sup}}} \overset{Sup}} \overset{Sup}} \overset{Sup}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup$$





$$\begin{array}{ll} \text{(DMK)} & \underset{\psi \in \psi^c}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{V}^c}{\overset{\text{W}^c}{(x)d\mu}}}} \int_X \psi^c(x)d\mu + \int_Y \psi(y)d\nu \\ \\ \int_X \inf_{y_j \in Y} \left[\| x - y_j \|^2 - \psi(y_j) \right] d\mu \\ & \sum_j \psi(y_j) \ v_j \end{array}$$







(DMK) Sup
$$\psi \in \psi^c$$
 $G(\psi) = \sum_j \int_{\text{Lag } \psi(yj)} ||x - y_j||^2 - \psi(y_j) d\mu + \sum_j \psi(y_j) v_j$

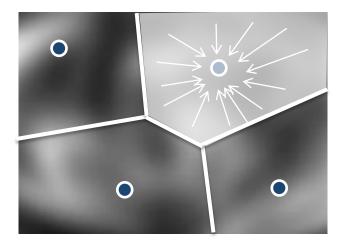
Where: Lag $\psi(yj) = \left\{ \begin{array}{cc} x & | & || x - y_j ||^2 - \psi(y_j) \\ & < || x - y_j ||^2 - \psi(y_{j'}) \end{array} \right\}$ for all j' $\neq j$

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 for all $j' \neq j$

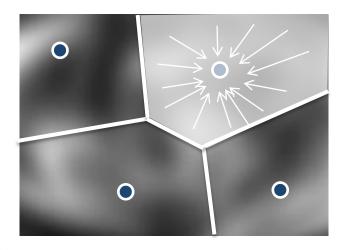
Laguerre diagram of the y_j 's (with the L₂ cost || x - y ||² used here, Power diagram)



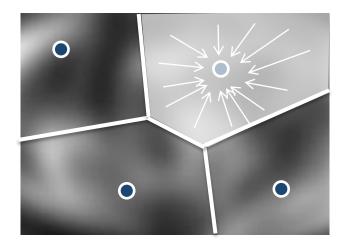
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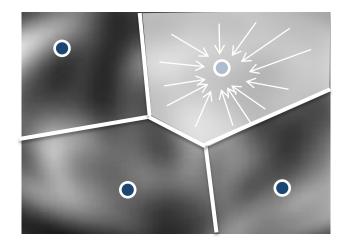
(DMK)
$$\begin{split} & \underset{\psi \in \psi^{c}}{\operatorname{Sup}} \quad G(\psi) = \sum_{j} \int_{\operatorname{Lag} \psi(yj)} \| x - y_{j} \|^{2} - \psi(y_{j}) \, d\mu + \sum_{j} \psi(y_{j}) \, v_{j} \\ & \text{Where: } \operatorname{Lag} \psi(yj) = \left\{ \begin{array}{c} x \mid \| x - y_{j} \|^{2} - \psi(y_{j}) < \| x - y_{j} \|^{2} - \psi(y_{j'}) \end{array} \right\} \text{ for all } j' \neq j \\ & & & \\ & \underset{(\text{with the } L_{2} \operatorname{cost} \| x - y \|^{2} \operatorname{used here}, \operatorname{Power diagram}) \\ \end{split}$$



 ψ is determined by the weight vector $[\psi(y_1) \ \psi(y_2) \ \dots \ \psi(y_m)]$



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For all weight vector, $\boldsymbol{\psi}$ is c-concave



Part. 3 Power Diagrams

Voronoi diagram: $Vor(x_i) = \{ x \mid d^2(x, x_i) < d^2(x, x_j) \}$

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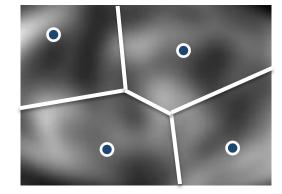
Power diagram: $Pow(x_i) = \{ x \mid d^2(x,x_i) - \psi_i < d^2(x,x_j) - \psi_j \}$



Part. 3 Power Diagrams



Part. 3 Optimal Transport

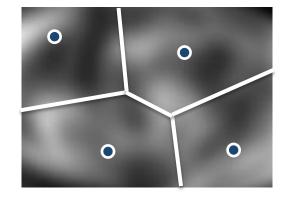


<u>Theorem</u>: (direct consequence of MK duality alternative proof in [Aurenhammer, Hoffmann, Aronov 98]):

Given a measure μ with density, a set of points (y_j) , a set of positive coefficients v_j such that $\sum v_j = \int d\mu(x)$, it is possible to find the weights $W = [\psi(y_1) \ \psi(y_2) \ \dots \ \psi(y_m)]$ such that the map $T_S{}^W$ is the unique optimal transport map between μ and $v = \sum v_j \ \delta(y_j)$



Part. 3 Optimal Transport



<u>Theorem:</u> (direct consequence of MK duality alternative proof in [Aurenhammer, Hoffmann, Aronov 98]):

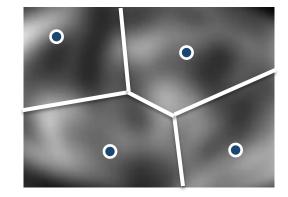
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Proof:
$$G(\psi) = \sum_{j} \int_{\text{Lag } \psi(yj)} ||x - y_j||^2 - \psi(y_j) d\mu + \sum_{j} \psi(y_j) v_j$$

Is a concave function of the weight vector $[\psi(y_1) \ \psi(y_2) \ \dots \ \psi(y_m)]$



Part. 3 Optimal Transport



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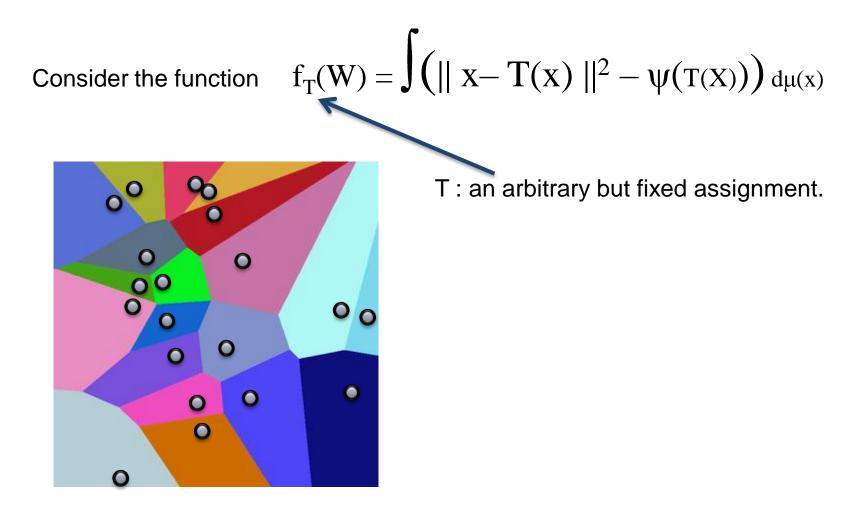


Idea of the proof

Consider the function
$$f_T(W) = \int (||x - T(x)||^2 - \psi(T(X))) d\mu(x)$$

The (unknown) weights $W = [\Psi(y_1) \Psi(y_2) \dots \Psi(y_m)]$





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Idea of the proof

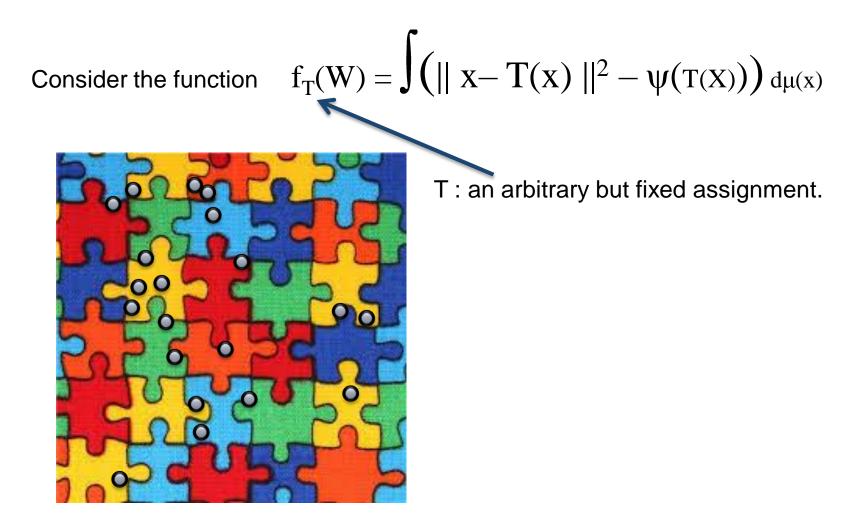
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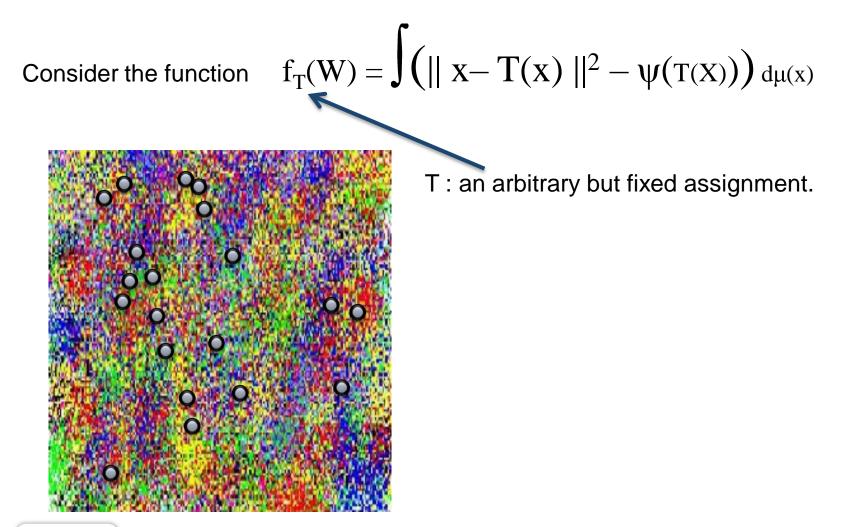


T : an arbitrary but fixed assignment.

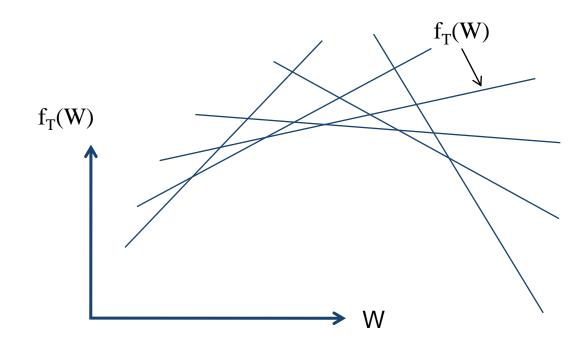




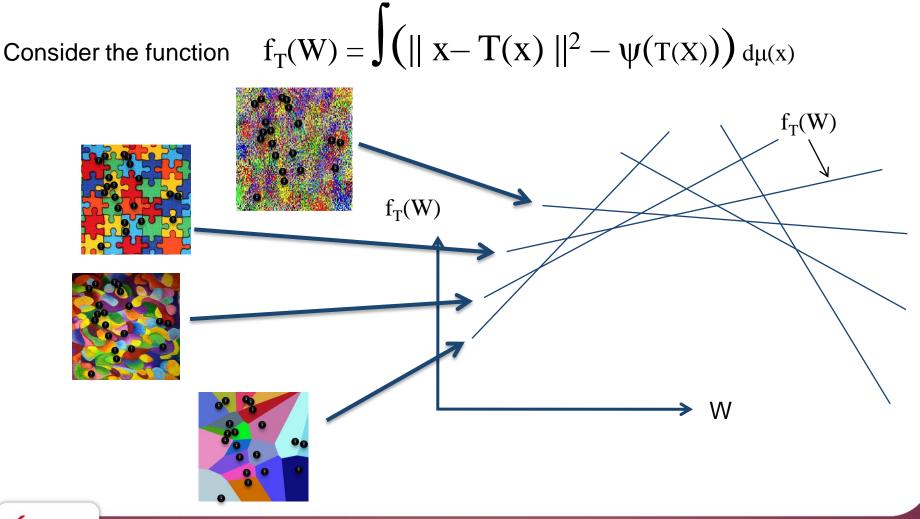
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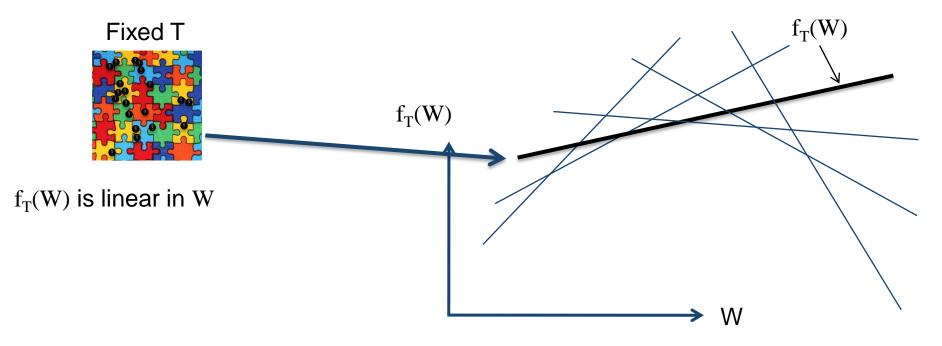






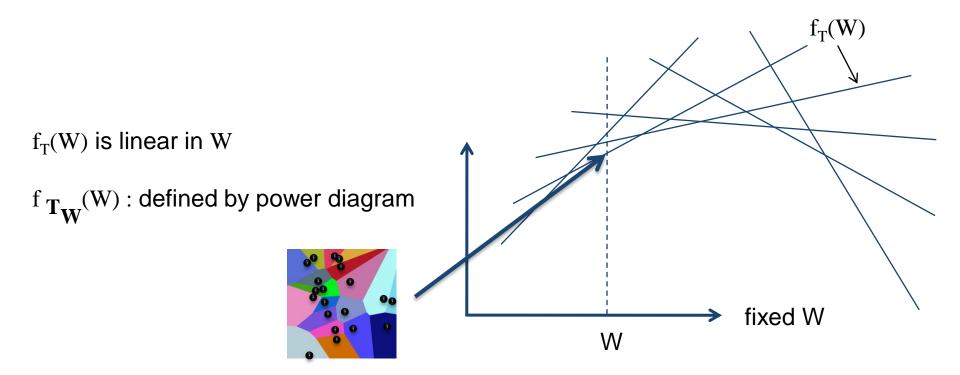


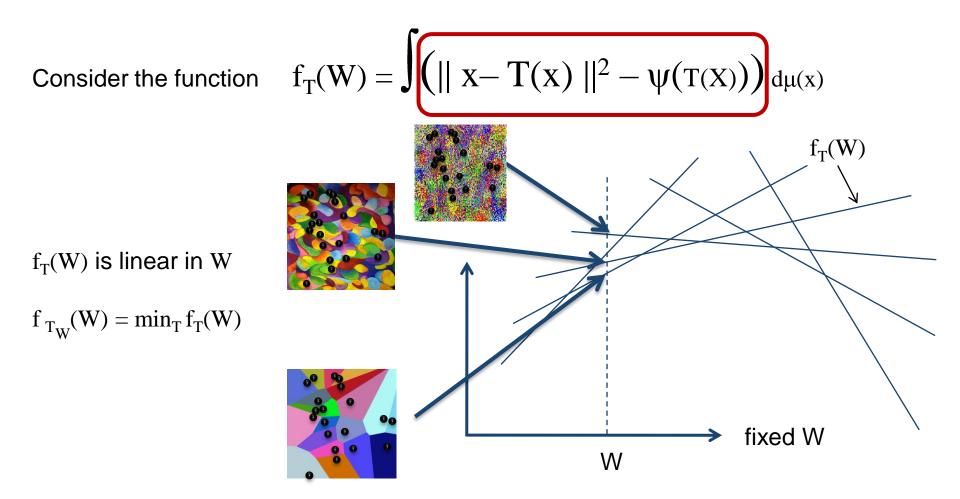
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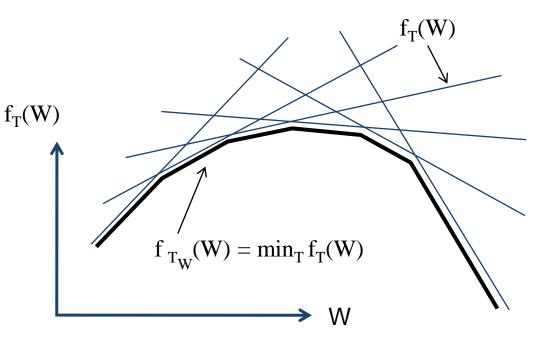




Idea of the proof

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 $f_T(W)$ is linear in W f: W $\rightarrow f_{T_W}(W)$ is **CONCAVE !!** (because its graph is the lower enveloppe of linear functions)





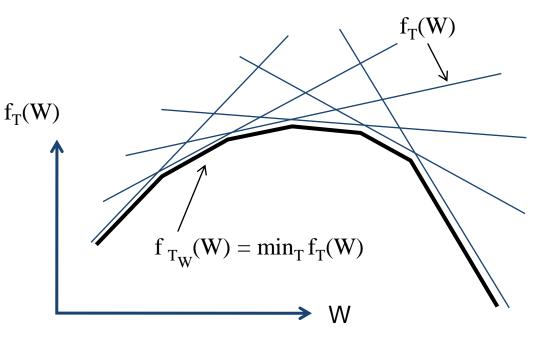
Part. 3 Optimal Transport – the AHA paper

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Consider
$$g(W) = f_{T_W}(W) + \sum v_j \psi_j$$





Part. 3 Optimal Transport – the AHA paper

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$$f_{T}(W)$$

 $f_{T_{W}}(W) = \min_{T} f_{T}(W)$
 W

(111)

Consider
$$g(W) = f_{T_W}(W) + \sum v_j \psi_j$$

 $\partial g / \partial \Psi_{j} = V_{j} - \int_{pow(yj)} ||x - y_{j}||^{2} d\mu(x)$ and g is concave.



Semi-discrete OT Summary:

(DMK) $\begin{array}{c} \text{Sup} \\ \psi \in \psi^{c} \end{array} \quad G(\psi) = \int_{X} \psi^{c}(x) d\mu + \int_{Y} \psi(y) dv \end{array}$

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$$G(\psi) = g(W) = \sum_{j} \int_{\text{Lag } \psi(yj)} \| x - y_j \|^2 - \psi(y_j) \, d\mu + \sum_{j} \psi(y_j) \, v_j \text{ is concave}$$



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 (= 0 at the maximum)



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 (= 0 at the maximum)

Desired mass at y_j

Mass transported to y_i



The [AHA] paper summary:

- The optimal weights minimize a convex function
- The gradient of this convex function is easy to compute

Note: the weight w(s) correspond to the Kantorovich potential $\psi(x)$ (solves a "discrete Monge-Ampere" equation)

The algorithm:

Input: two tetrahedral meshes M_1 and M_2 **Output:** a morphing between M_1 and M_2



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The [AHA] paper summary:

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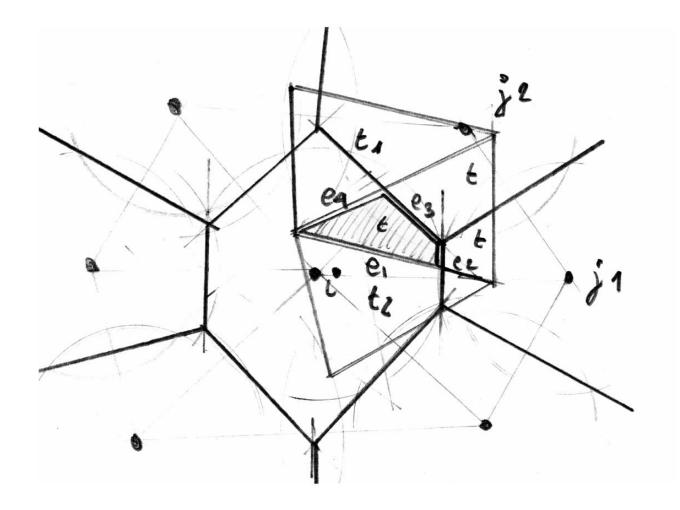
The algorithm:

Input: two tetrahedral meshes M_1 and M_2 **Output:** a morphing between M_1 and M_2

Step 1: sample M₂ with N points $(s_1 \dots s_N)$ **Step 2:** initialize the weights $(w_1 \dots w_N) = (0 \dots 0)$ **Step 3:** minimize $g(w_1 \dots w_N)$ with a quasi-Newton algorithm: For each iterate $(s_1 \dots s_N)^{(k)}$: Compute Pow((w_i, s_i)) \cap M₁ [Nivoliers, L 2014, Curves and Surfaces]



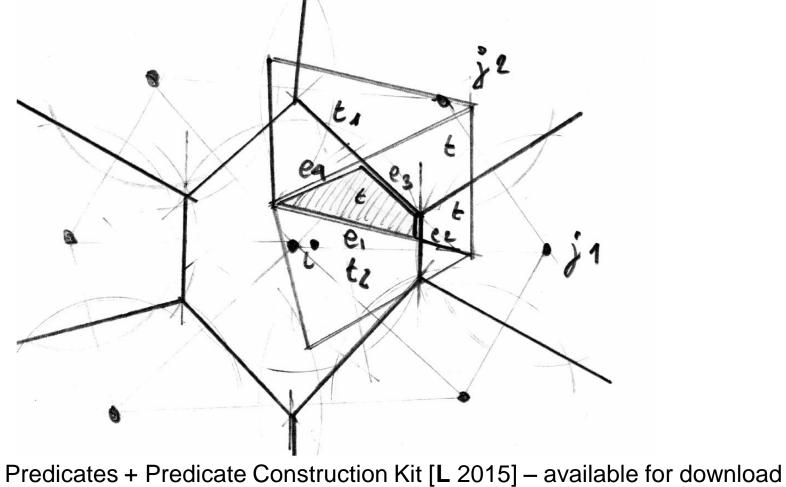
Compute Pow((w_i, s_i)) \cap M₁ [Nivoliers, L 2014, Curves and Surfaces]





Compute Pow((w_i , s_i)) \cap M₁ [Nivoliers, L 2014, Curves and Surfaces]

Implementation in **GEOGRAM** (http://alice.loria.fr/software/geogram



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The [AHA] paper summary:

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Compute g and grad g

+ Multilevel version [Merigot 2011] (2D),

[L 2014 arXiv, M2AN 2015] (3D & relation with Centroidal Voronoi Tesselation)

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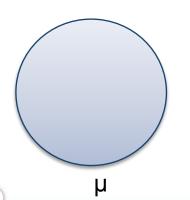
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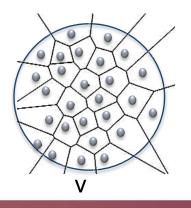
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The algorithm:

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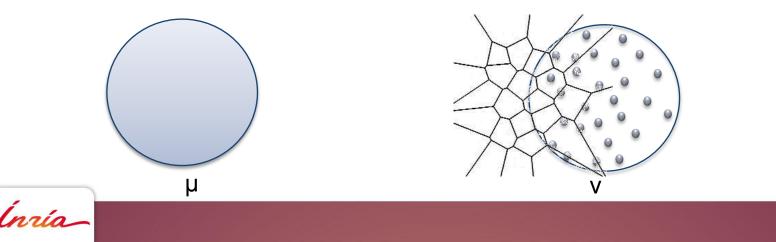
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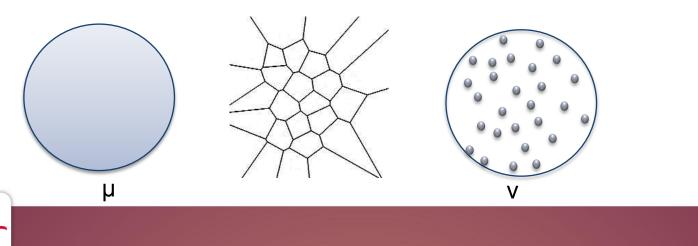
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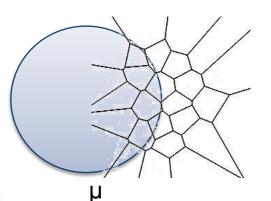
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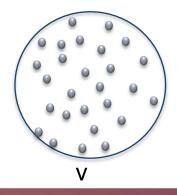
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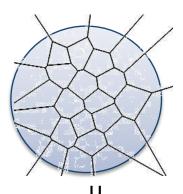
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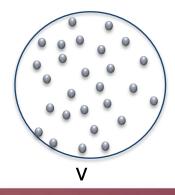
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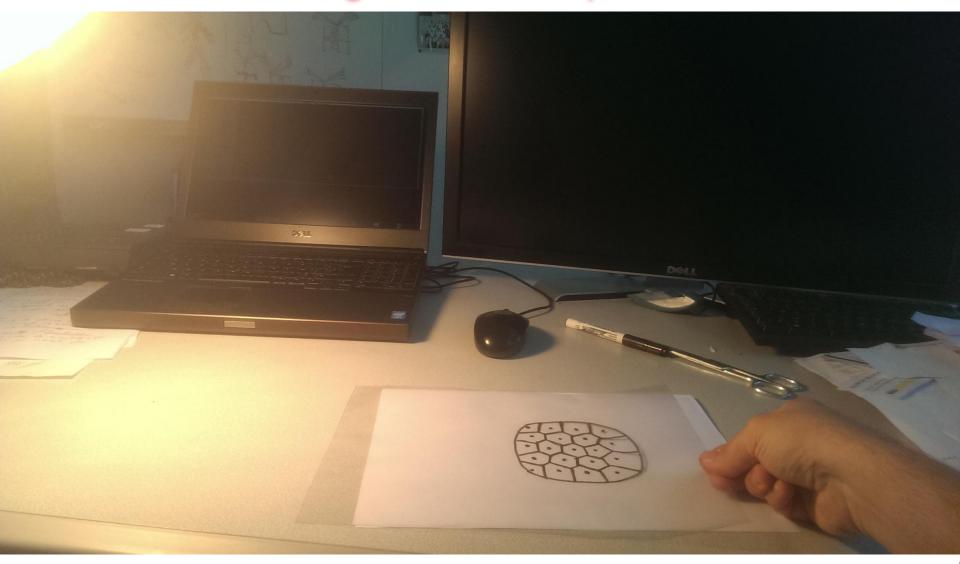
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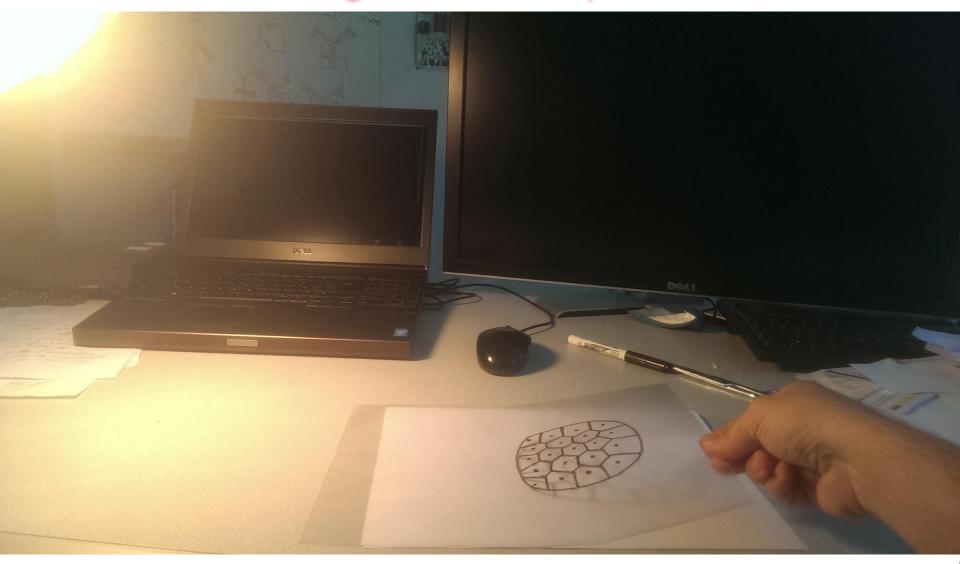




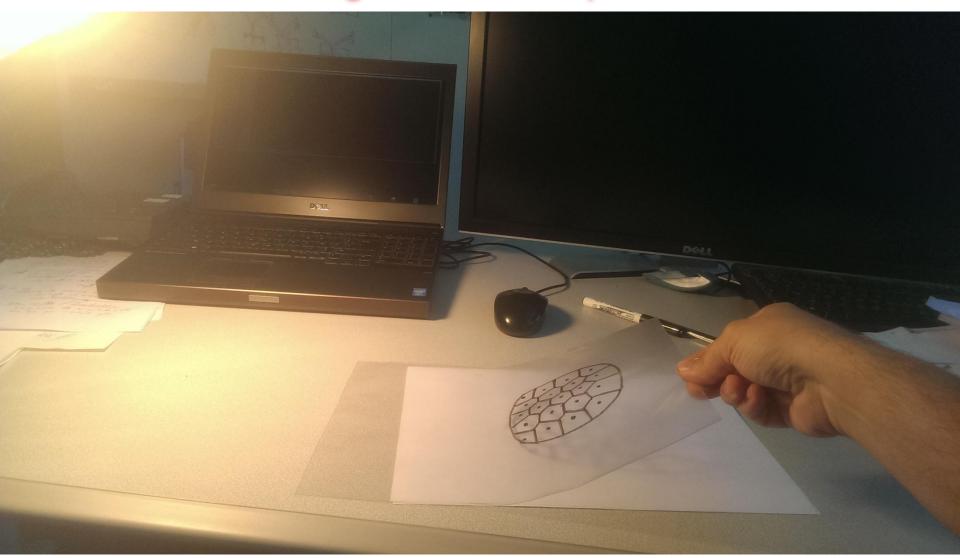




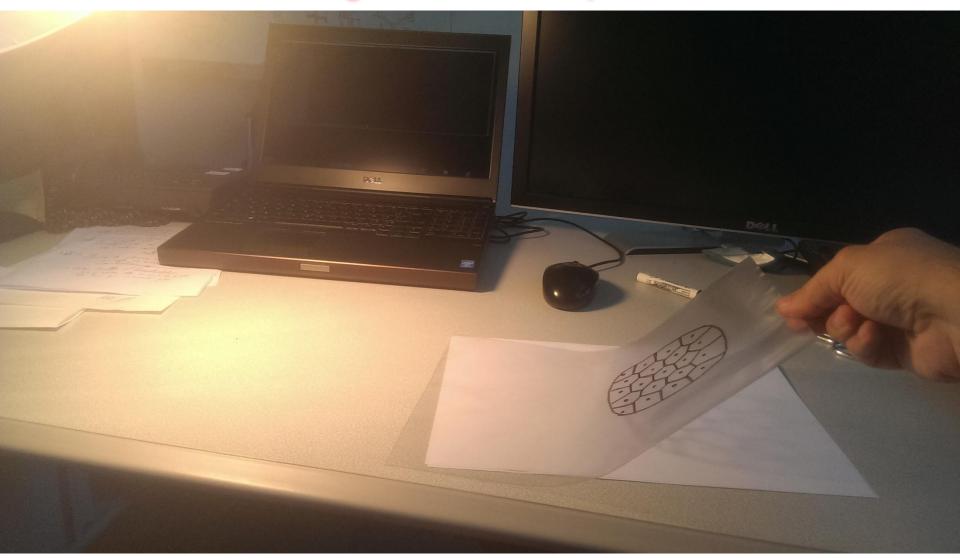
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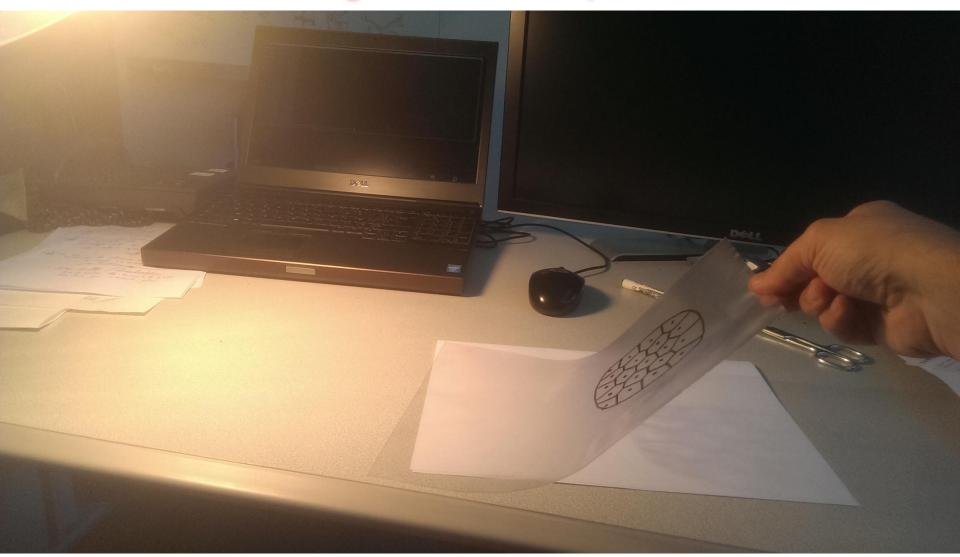
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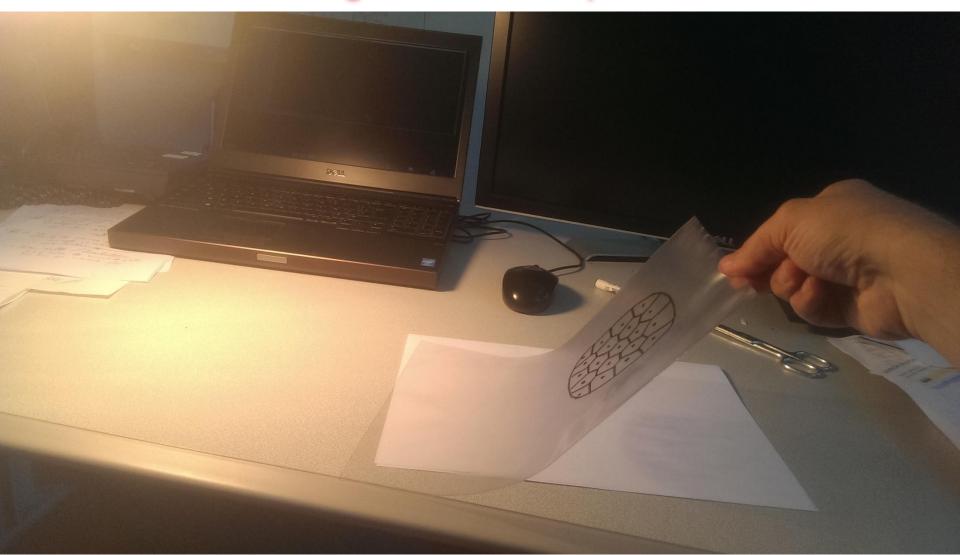
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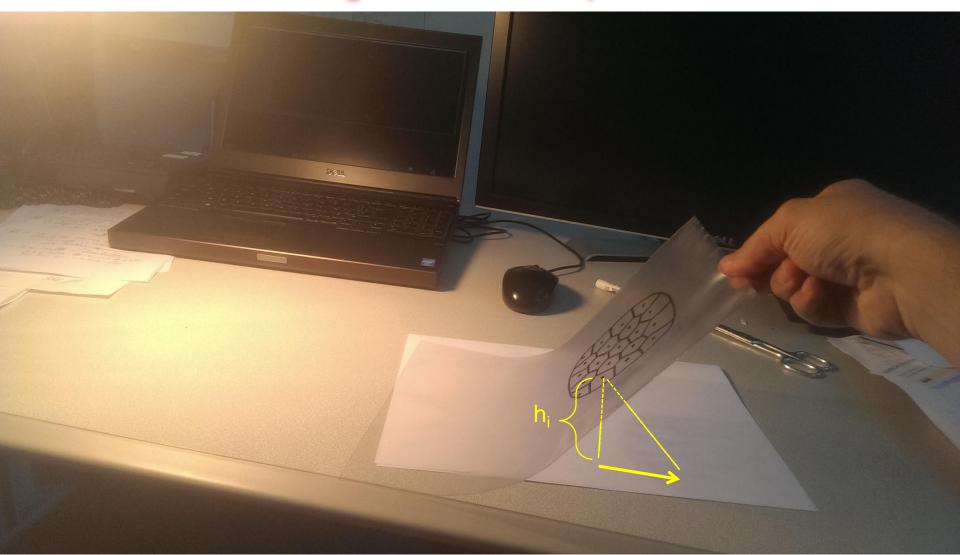
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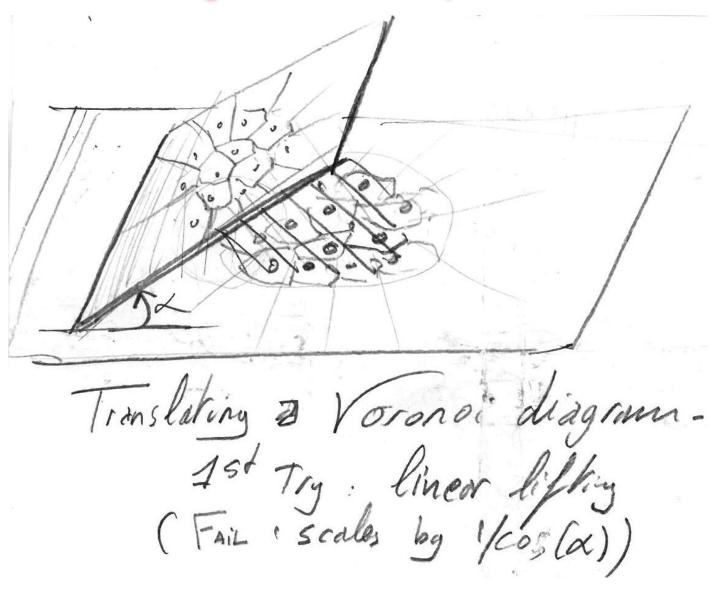


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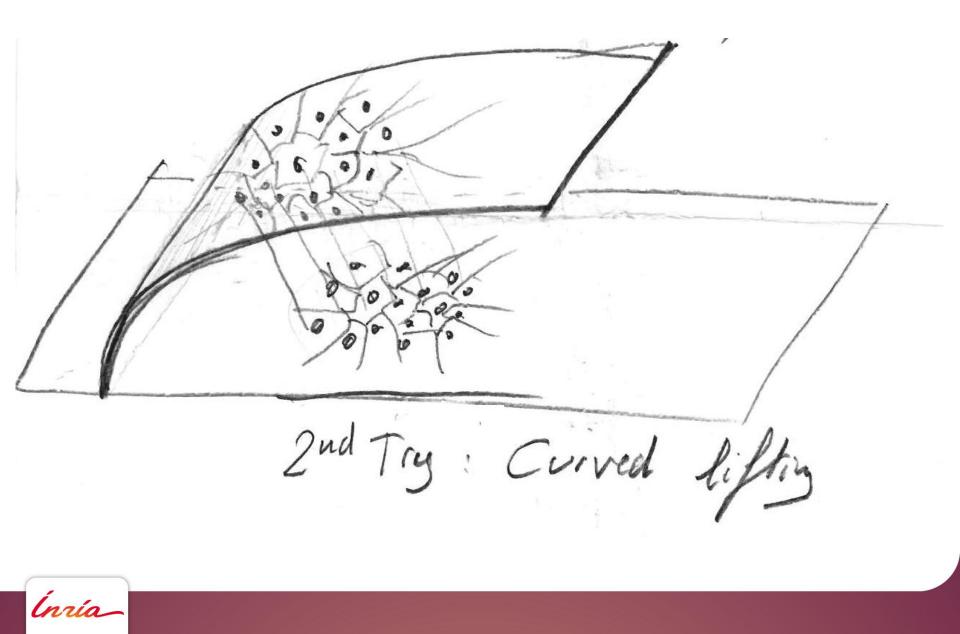


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Part. 4 Power Diagrams & Transport





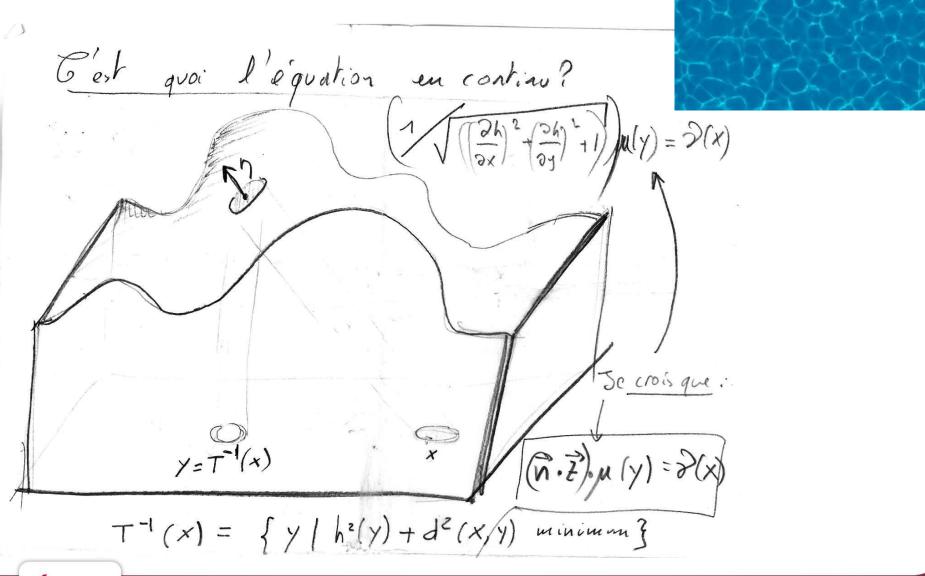


"converging beams" con compensate the cos(x) expansion by "re-concentrating" the paints

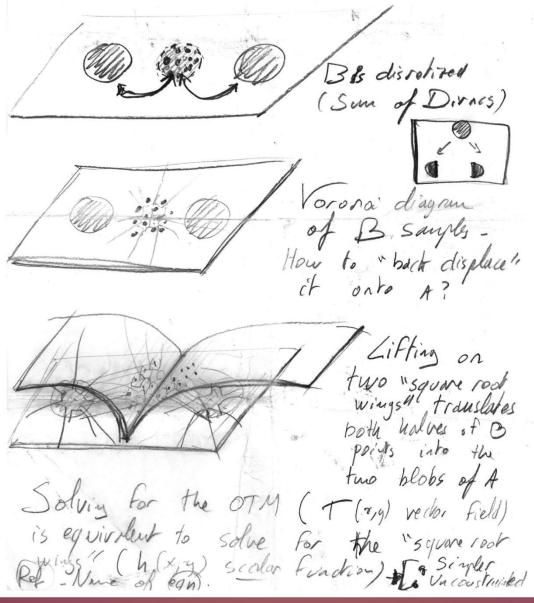


Part. 4 Power Diagrams & Transport $d^2(p_{i,q}) \stackrel{+h_i^2}{-w_i} \langle d^2(p_{i,q}) \stackrel{+h_i^2}{-w_j} \lor_j$ $d^{2}(p_{i}, q-T) < d^{2}(j, q-T)$ V, $(p_i - q + T)^2 \leq (p_j - q + T)^2 \qquad \forall i$ $d^{2}(p_{i},q) + 2T.(p_{i}-q) + T^{2} \leq d^{2}(p_{j},q) + 2T.(p_{j}-q) + T^{2} \vee_{j}^{-}$ d²(pi,q) + 2T.pi <d²(pj,q) +2T.pj $W_i^2 = -2T \cdot p_i'$ + che hi?: (2ripir Che); hi= VZ(T-pi - min(T-p)) Granstation d'un diagramme de Uronoi sectionnel-Delevement en racine cané -

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Part. 4 Optimal Transport – 2D examples Numerical Experiment: *A disk becomes two disks*

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Part. 4 Optimal Transport – 3D examples Numerical Experiment: *A sphere becomes a cube*

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Part. 4 Optimal Transport – 3D examples Numerical Experiment: *A sphere becomes two spheres*

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Part. 4 Optimal Transport – 3D examples Numerical Experiment: *Armadillo to sphere*

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Part. 4 Optimal Transport – 3D examples Numerical Experiment: *Other examples*

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Part. 4 Optimal Transport – 3D examples Numerical Experiment: *Varying density*

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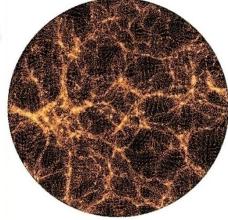
Part. 4 Optimal Transport – 3D examples Numerical Experiment: *Performances*

nb masses										
time (s)	1.45	3.2	7.3	17.3	55	154	187	671	1262	2649

TABLE 4. Statistics for the Armadillo \rightarrow sphere optimal transport with varying number of masses (see third row of Figure 12). Timings are given in seconds. The multi-level algorithm with BRIO pre-ordering and degree 2 regressions is used.



Part. 4 Optimal Transport – 3D examples Numerical Experiment: *Performances*



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Note that a few years ago, several hours of supercomputer time were needed for computing OT with a few thousand Dirac masses, with a combinatorial algorithm in $O(n^3)$



Part. 4 Optimal Transport – 3D examples Numerical Experiment: *Performances*

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nb masses	1000	2000	5000	10000	30000	50000	10^{5}	3×10^5	5×10^5	10^{6}
time (s)	1.45	3.2	7.3	17.3	55	154	187	671	1262	2649

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With the semi-discrete algorithm, it takes less than 10 seconds on my laptop



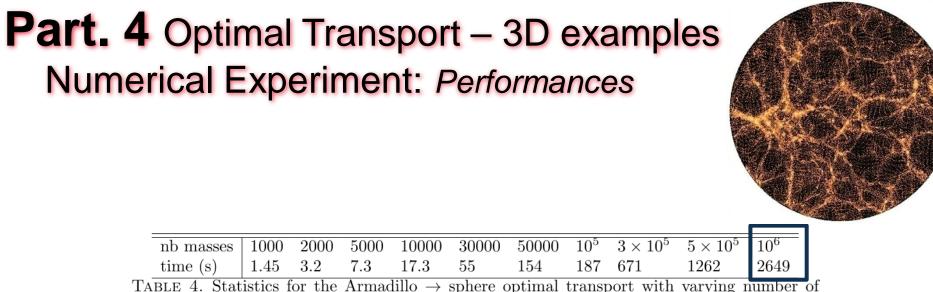


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In the semi-discrete setting, my 3D version of Merigot's multigrid algorithm computes OT for **1 million Dirac masses** in less than 1 hour on a laptop PC

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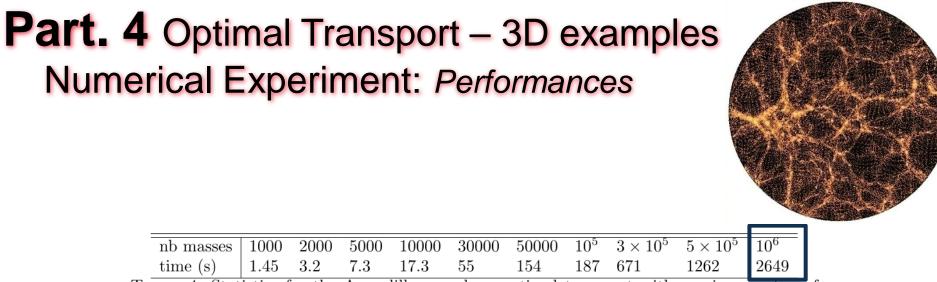


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In the semi-discrete setting, my 3D version of Merigot's multigrid algorithm computes OT for **1 million Dirac masses** in less than 1 hour on a laptop PC

Even much faster convergence can probably be reached with a true Newton solver (and several acceleration tricks), still investigating...

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Other topics

- •Euler equation in more complicated setting: [Merigot & Mirebeau]
- •Using semi-discrete OT to solve other PDEs [Benamou, Carlier, Merigot, Oudet]
- •New fluid simulation methods "power particles" [DeGoes et.al]



Conclusions – Open questions

* Connections with physics, Legendre transform and entropy ?

[Cuturi & Peyré] – regularized discrete optimal transport – why does it work ? Hint 1: Minimum action principle subject to conservation laws Hint 2: Entropy = dual of temperature ; Legendre = Fourier[(+,*) → (Max,+)]...

* More continuous numerical algorithms ? [Benamou & Brenier] fluid dynamics point of view – very elegant, but 4D problem !! FEM-type adaptive discretization of the subdifferential (graph of T) ?

* Can we characterize OT in other semi-discrete settings ? measures supported on unions of spheres piecewise linear densities

* Connections with computational geometry ?

Singularity set **[Figalli]** = set of points where T is discontinuous Looks like a "mutual power diagram", anisotropic Voronoi diagrams

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Conclusions - References

Some references (that this presentation is based on)

A Multiscale Approach to Optimal Transport, **Quentin Mérigot**, Computer Graphics Forum, 2011

Variational Principles for Minkowski Type Problems, Discrete Optimal Transport, and Discrete Monge-Ampere Equations Xianfeng Gu, Feng Luo, Jian Sun, S.-T. Yau, ArXiv 2013

Minkowski-type theorems and least-squares clustering **AHA! (Aurenhammer, Hoffmann, and Aronov),** SIAM J. on math. ana. 1998

Topics on Optimal Transportation, 2003 Optimal Transport Old and New, 2008 **Cédric Villani**

Jean-David Benamou & Yann Brenier

a fluid formulation of Optimal Transportation, 2000

Laudau and Lifschitz – Course of Theoretical Physics – Volumes I and III



Conclusions - References Other references

Polar factorization and monotone rearrangement of vector-valued functions **Yann Brenier**, Comm. On Pure and Applied Mathematics, June 1991

A computational fluid mechanics solution of the Monge-Kantorovich mass transfer problem, **J.-D. Benamou**, **Y. Brenier**, Numer. Math. 84 (2000), pp. 375-393

Pogorelov, Alexandrov – Gradient maps, Minkovsky problem (older than AHA paper, some overlap, in slightly different context, formalism used by Gu & Yau)

Rockafeller – Convex optimization – Theorem to switch inf(sup()) – sup(inf()) with convex functions (used to justify Kantorovich duality)

New textbook: **Filippo Santambrogio** – Optimal Transport for Applied Mathematician, Calculus of Variations, PDEs and Modeling – Jan 15, 2015

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Online resources

All the sourcecode/documentation available from: alice.loria.fr/software/geogram

Computes semi-discrete OT in 3D Scales up to millions Dirac masses on a laptop

L., A numerical algorithm for semi-discrete L2 OT in 3D, ESAIM Math. Modeling and Analysis, accepted (draft: <u>http://arxiv.org/abs/1409.1279</u> <= to be fixed: bug in MA equation in this version, fixed in M2AN journal version)

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Downloads: alice.loria.fr/software GEOGRAM & GRAPHITE Video of course on OT: www.loria.fr/~levy

Acknowledgements

Funding: European Research Council & ANR ("french NSF")



GOODSHAPE ERC-StG-205693

VORPALINE ERC-PoC-334829

ANR MORPHO, ANR BECASIM



New project: EXPLORAGRAM (Inria) with Q. Mérigot and J.-D. Benamou

Quentin Merigot, Yann Brenier, Boris Thibert, Emmanuel Maitre, Jean-David Benamou, Filippo Santambrogio, Edouard Oudet, Hervé Pajot. ANR TOMMI, ANR GEOMETRYA