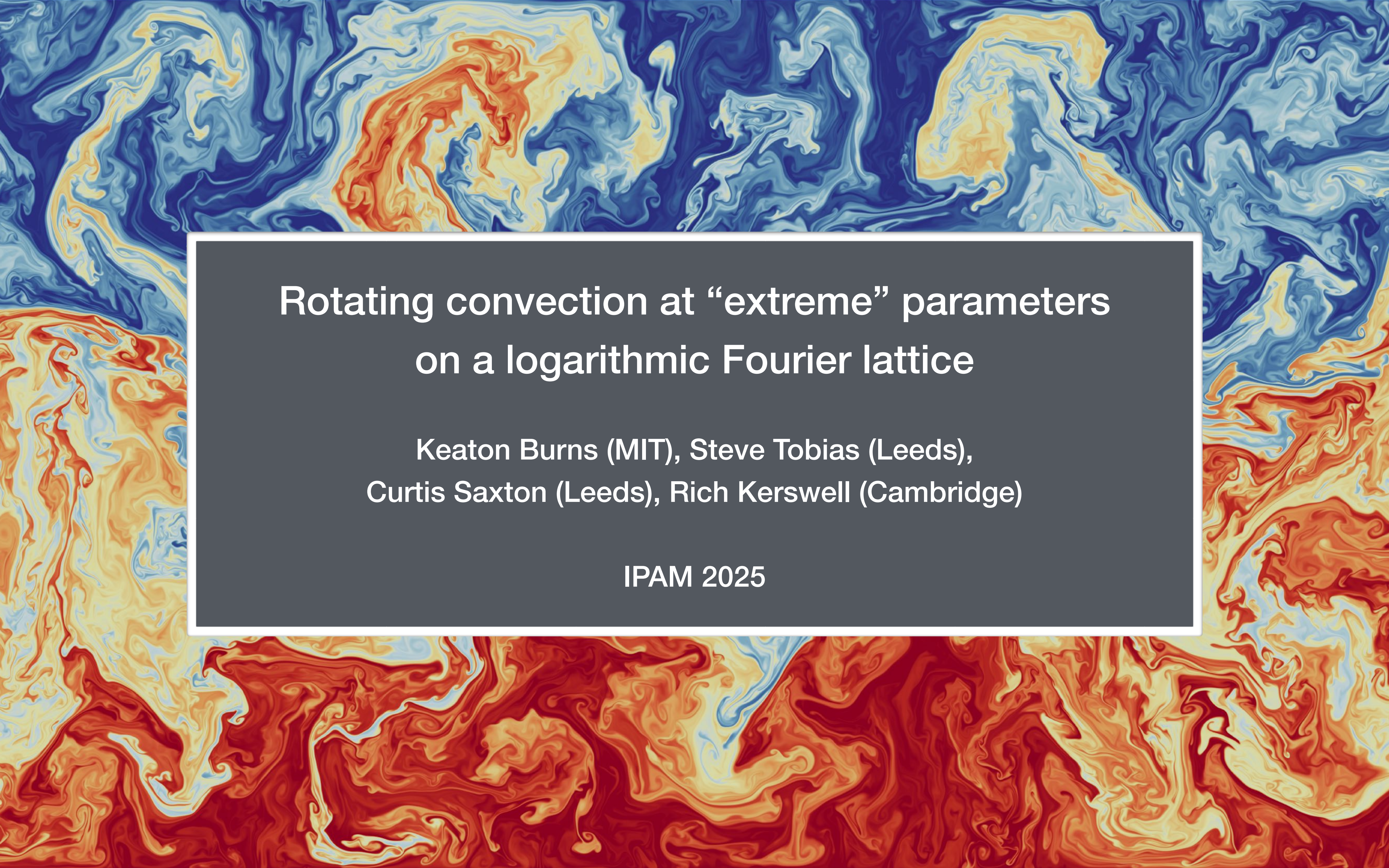




Rotating convection at “extreme” parameters on a logarithmic Fourier lattice

Keaton Burns (MIT), Steve Tobias (Leeds),
Curtis Saxton (Leeds), Rich Kerswell (Cambridge)

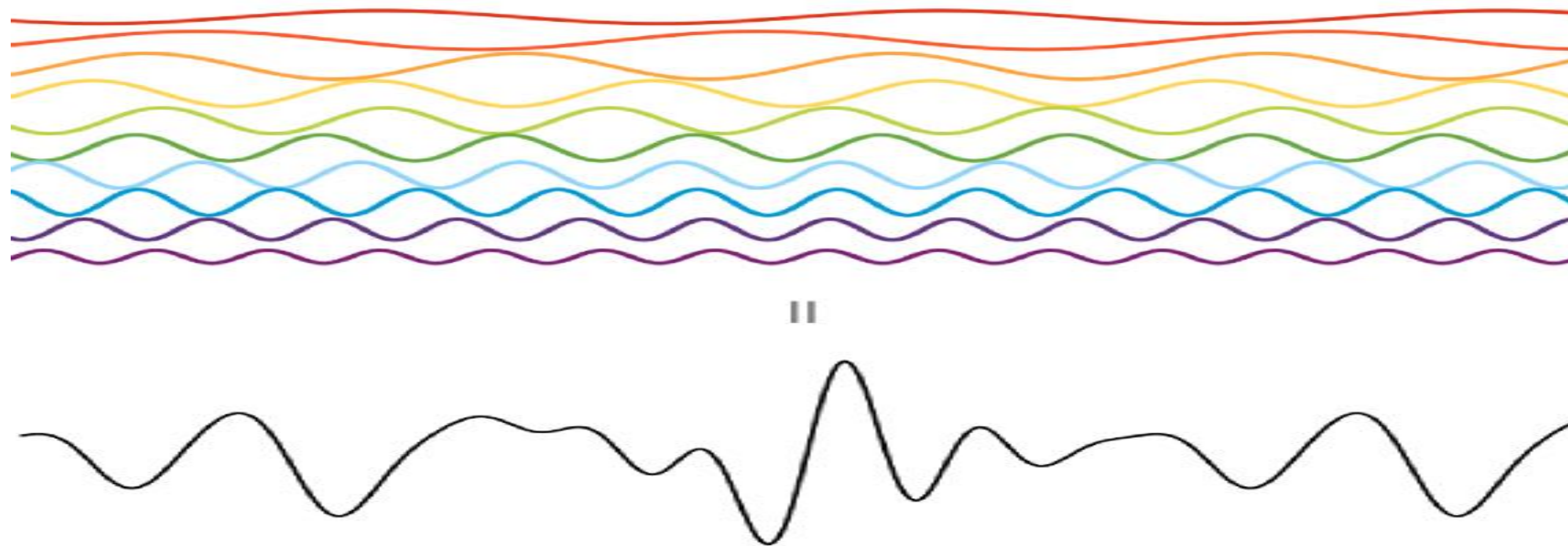
IPAM 2025

Fourier spectral methods

Fourier series: $\phi_n(x) = e^{inx}$

- **Exponential convergence** for smooth **periodic** functions
- **Fast transforms** for computing coefficients
- **Diagonal** differential operators:

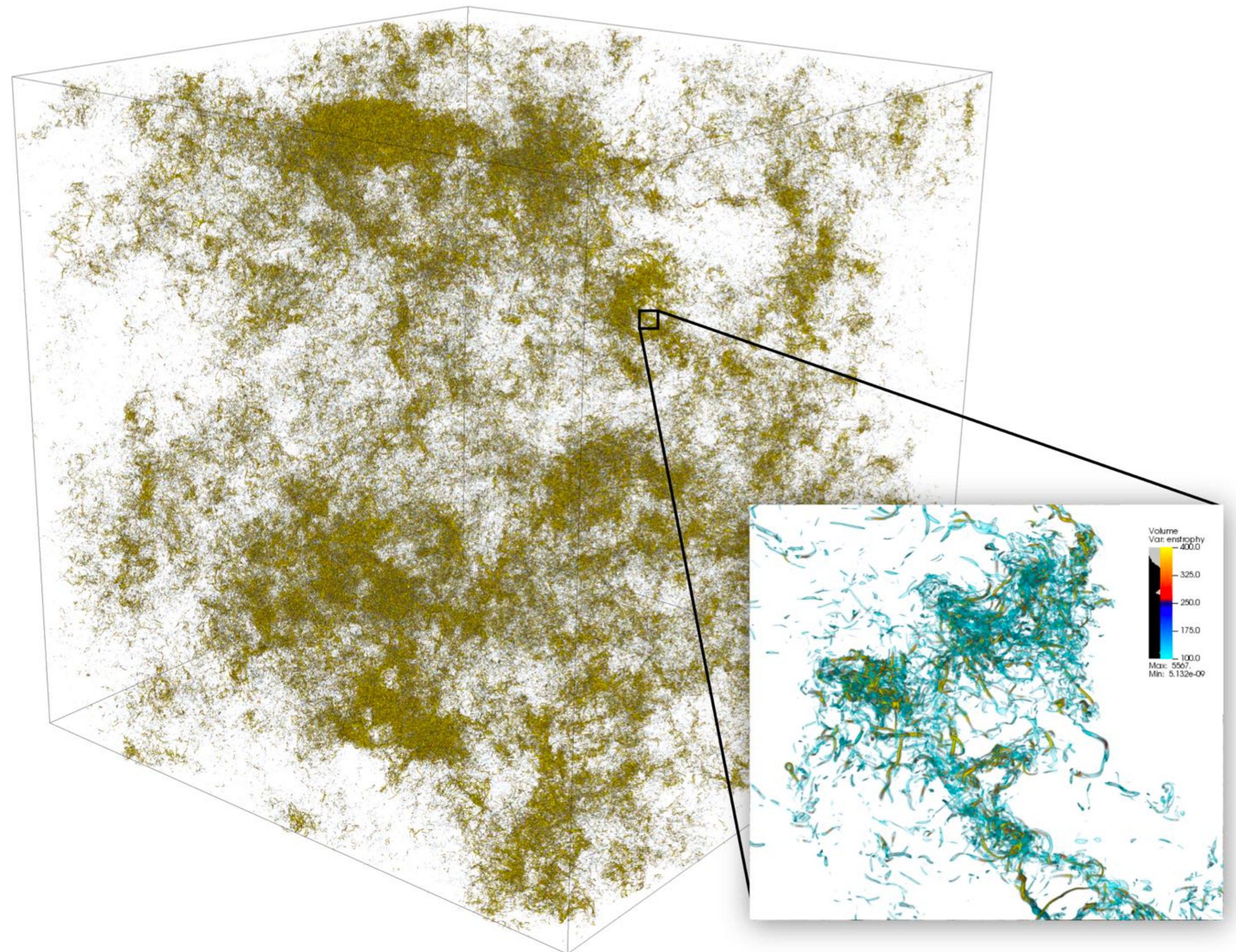
$$\langle \phi_m | \partial_x \phi_n \rangle = in \delta_{m,n}$$



World's largest turbulence simulations

Yeung & Ravikumar (2024)

- Fourier pseudospectral method
- $32,768^3$ grid points
- 32,768 GPUs (Frontier)
- $R_\lambda \approx 2550$



Cost of Direct Numerical Simulations (DNS)

Spatial dynamic range:

$$N_x = \frac{L_x}{\Delta x} \sim \text{Re}^{3/4}$$

Decade of Moore's Law ~ Decade in Reynolds

Temporal dynamic range:

$$N_t = \frac{T}{\Delta t} \sim \frac{L_x}{\Delta x}$$

DNS today:

Re ~ 10⁶

DNS of Earth:

~50 years

DNS of Sun:

~90 years

Simulation degrees of freedom (DOF):

$$N = N_x^3 N_t \sim \text{Re}^3$$

“Log-Lattice” Fourier discretizations

Fourier discretization:

$$u(x) = \sum_k \hat{u}(k) e^{ikx}$$

Standard (linear) wavenumbers:

$$k \in \frac{2\pi}{L} \mathbb{Z} \quad N_k \sim \frac{k}{k_0} \sim \frac{L}{\Delta x}$$

Logarithmic wavenumbers:

$$k \in \frac{2\pi}{L} \{\pm \lambda^n\}_{n \in \mathbb{N}} \quad \text{Spacing factor } \lambda$$

$$N_k \sim \log(k/k_0) \sim \log(L/\Delta x)$$

Properties:

$$\hat{u}(-k) = \hat{u}^*(k)$$

$$\widehat{\partial_x u}(k) = i k \hat{u}(k)$$

$$(u, v) = \sum_k \hat{u}(k) \hat{v}^*(k)$$

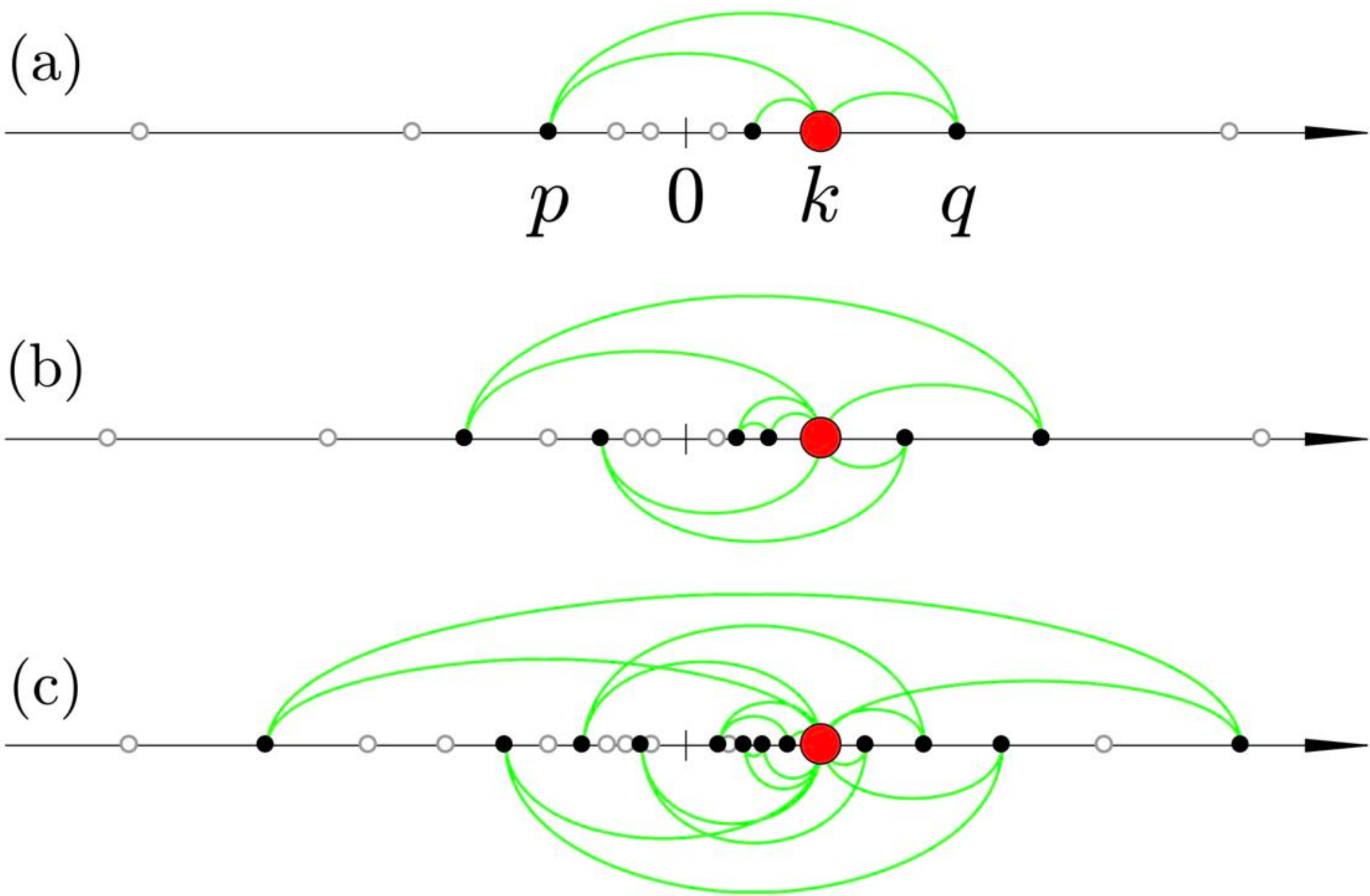
$$(\partial_x u, v) = - (u, \partial_x v)$$

Correct linear dynamics

Log-Lattice triads

Theorem 1. *The following three cases describe all nondegenerate lattices with spacing factors $\lambda \geq 1.05$:*

- (a) $\lambda = 2$, and all triads at the unity are given in table 1(a);
- (b) $\lambda = \sigma$, the plastic number (2), and all triads at the unity are given in table 1(c);
- (c) λ satisfies $1 = \lambda^b - \lambda^a$, where (a, b) are mutually prime integers not larger than 62, excluding also the pairs $(a, b) = (1, 3)$ and $(4, 5)$. All triads at the unity are given in table 1(b).



(a)

i	1	2	3
p_i	2	-1	1/2
q_i	-1	2	1/2

(b)

i	1	2	3	4	5	6
p_i	λ^b	$-\lambda^a$	λ^{b-a}	$-\lambda^{-a}$	λ^{-b}	λ^{a-b}
q_i	$-\lambda^a$	λ^b	$-\lambda^{-a}$	λ^{b-a}	λ^{a-b}	λ^{-b}

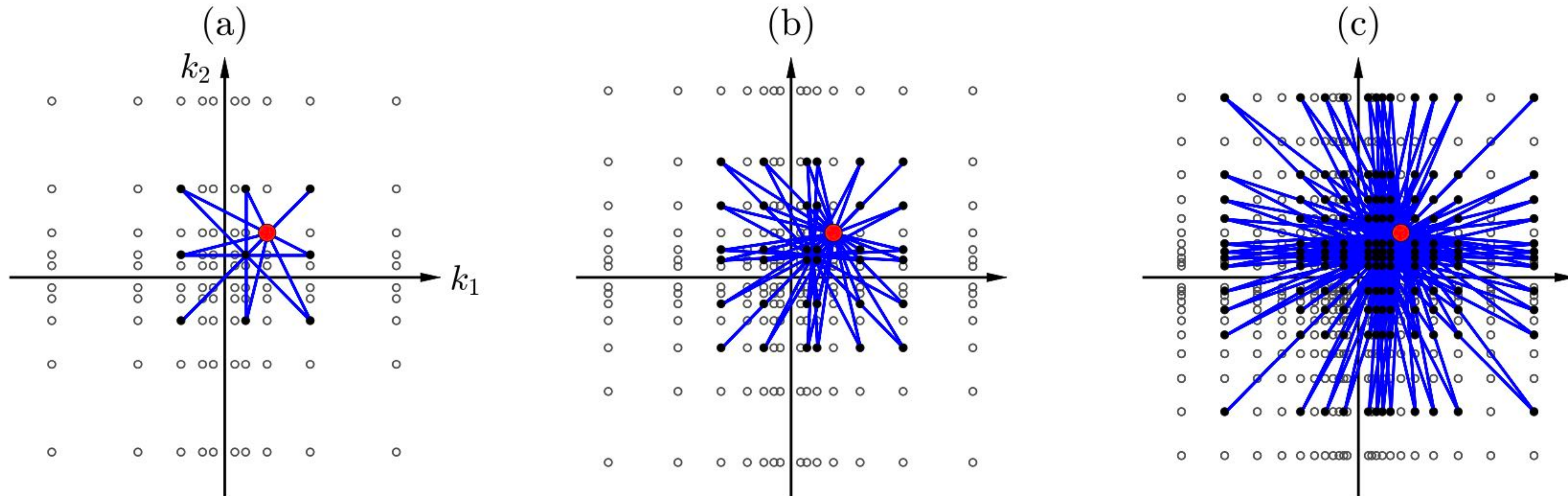
(c)

i	1	2	3	4	5	6	7	8	9	10	11	12
p_i	σ^3	$-\sigma$	σ^2	$-\sigma^{-1}$	σ^{-3}	σ^{-2}	σ^5	$-\sigma^4$	σ	$-\sigma^{-4}$	σ^{-5}	σ^{-1}
q_i	$-\sigma$	σ^3	$-\sigma^{-1}$	σ^2	σ^{-2}	σ^{-3}	$-\sigma^4$	σ^5	$-\sigma^{-4}$	σ	σ^{-1}	σ^{-5}

Log-Lattice triads

Theorem 1. *The following three cases describe all nondegenerate lattices with spacing factors $\lambda \geq 1.05$:*

- (a) $\lambda = 2$, and all triads at the unity are given in table 1(a);
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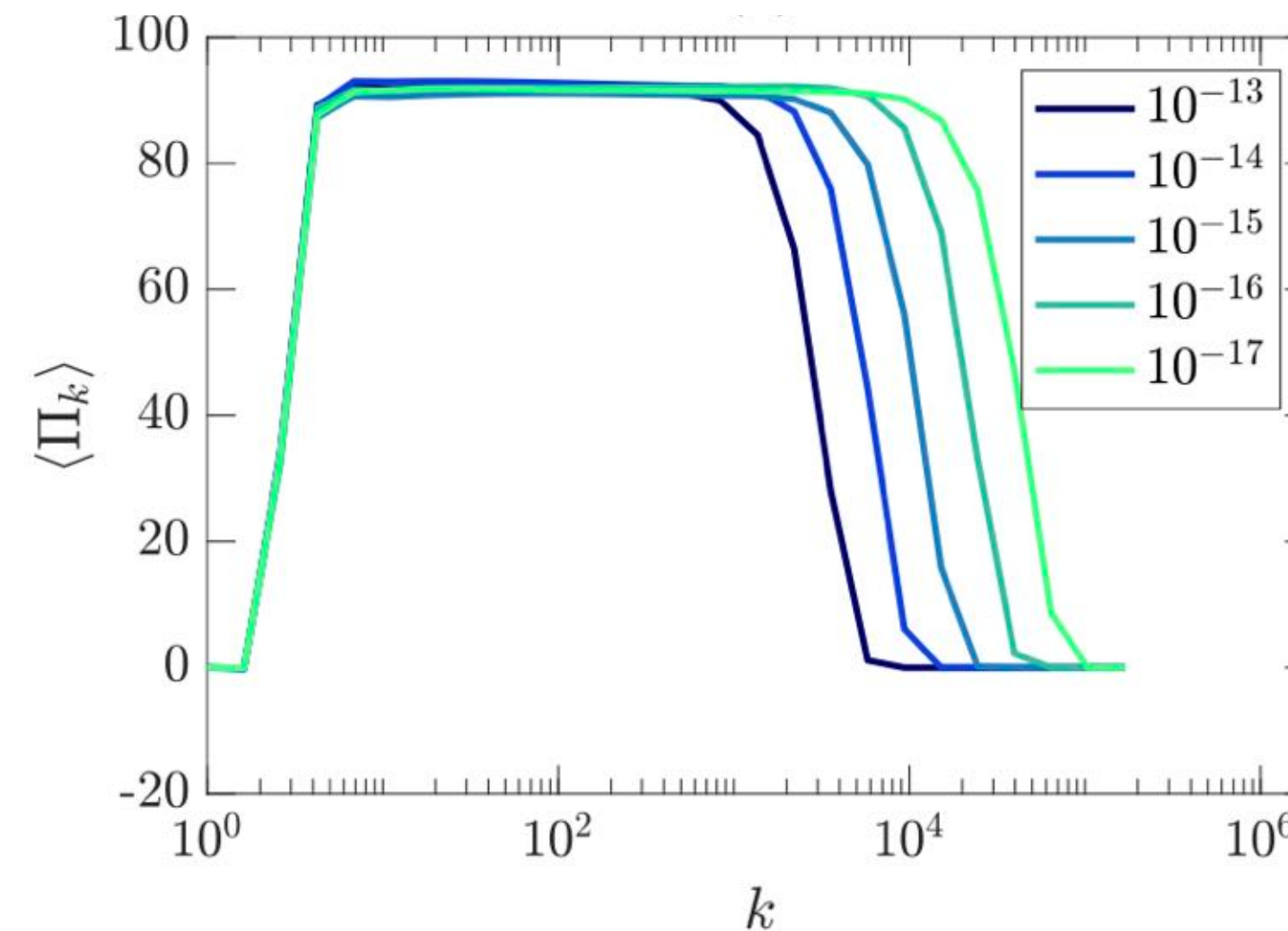


Log-Lattice isotropic turbulence

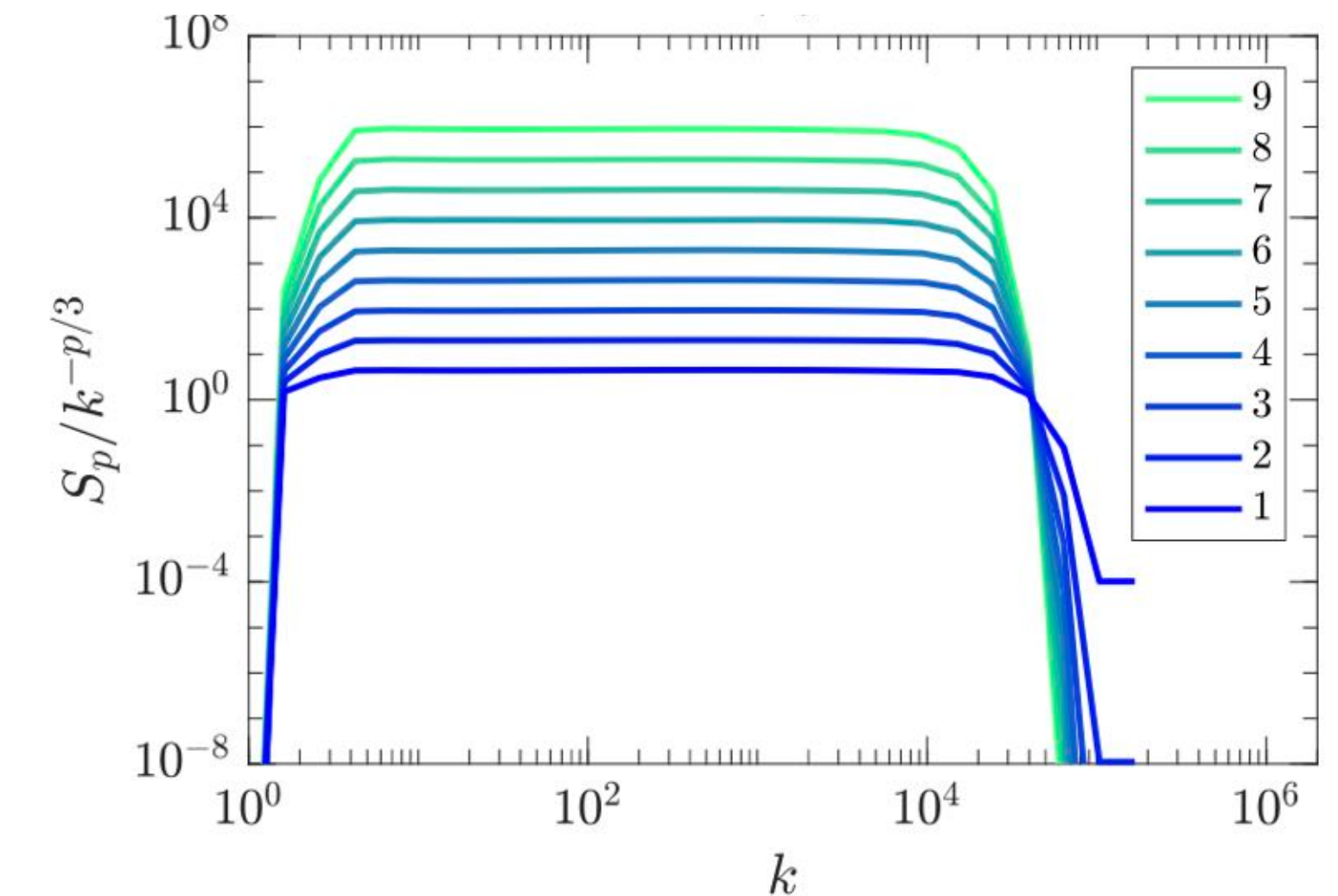
Campolina & Mailybaev (2021)

- 3D log-lattice
- Forward cascade, wrong slope
- Good structure function scaling
- Good velocity PDFs

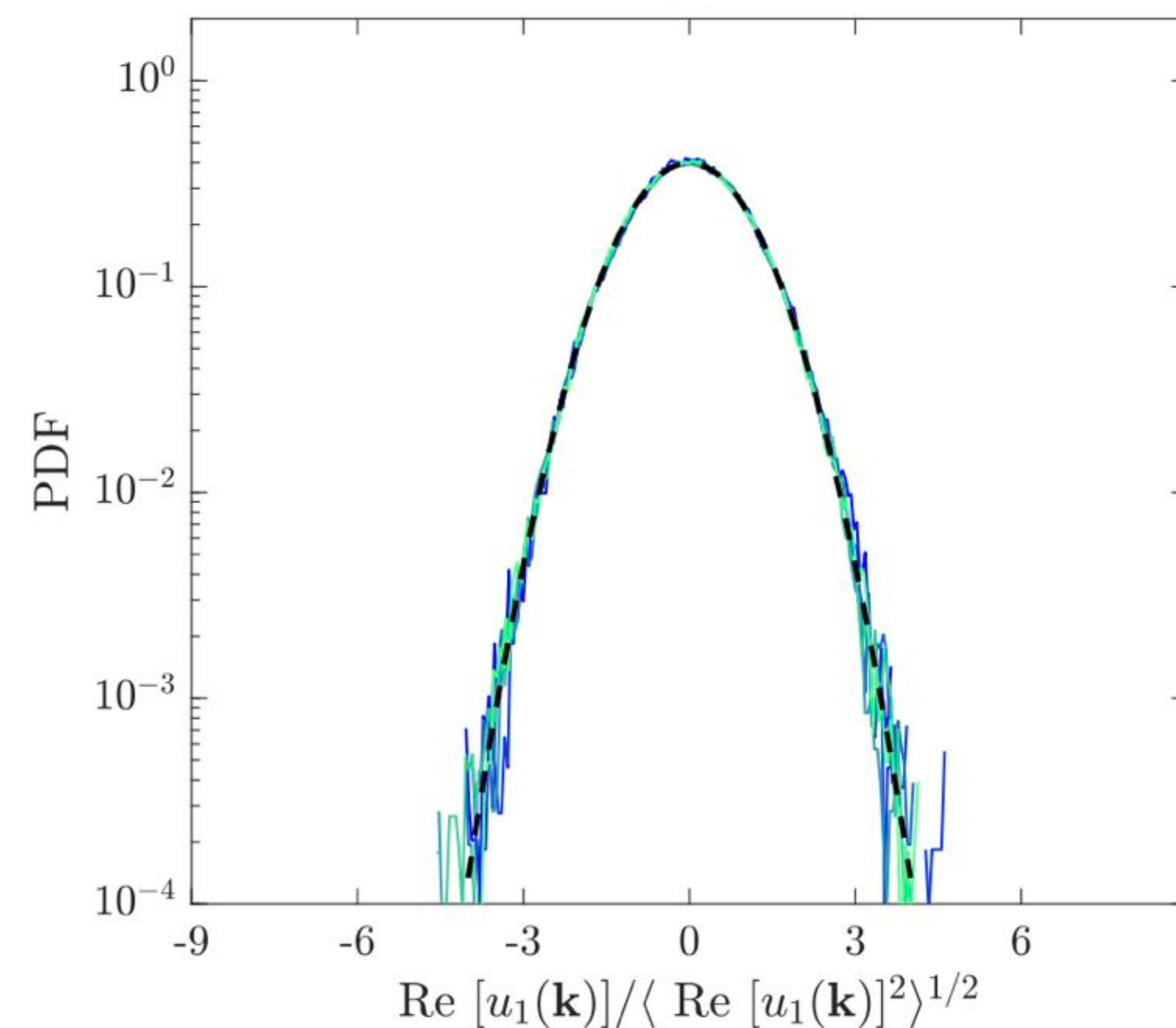
Mean energy flux



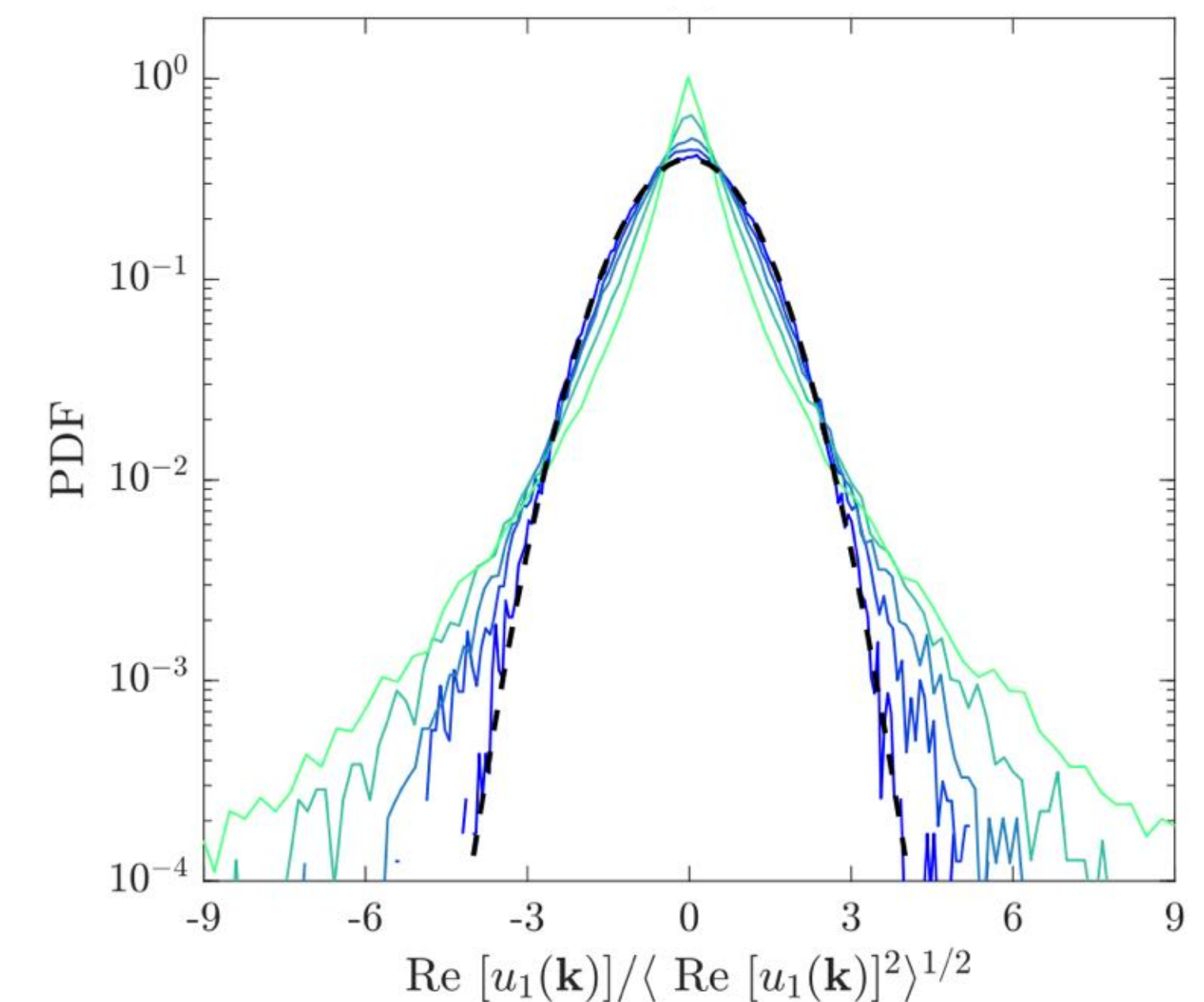
Comp. structure func.



Inertial range velocity PDF



Viscous range velocity PDF

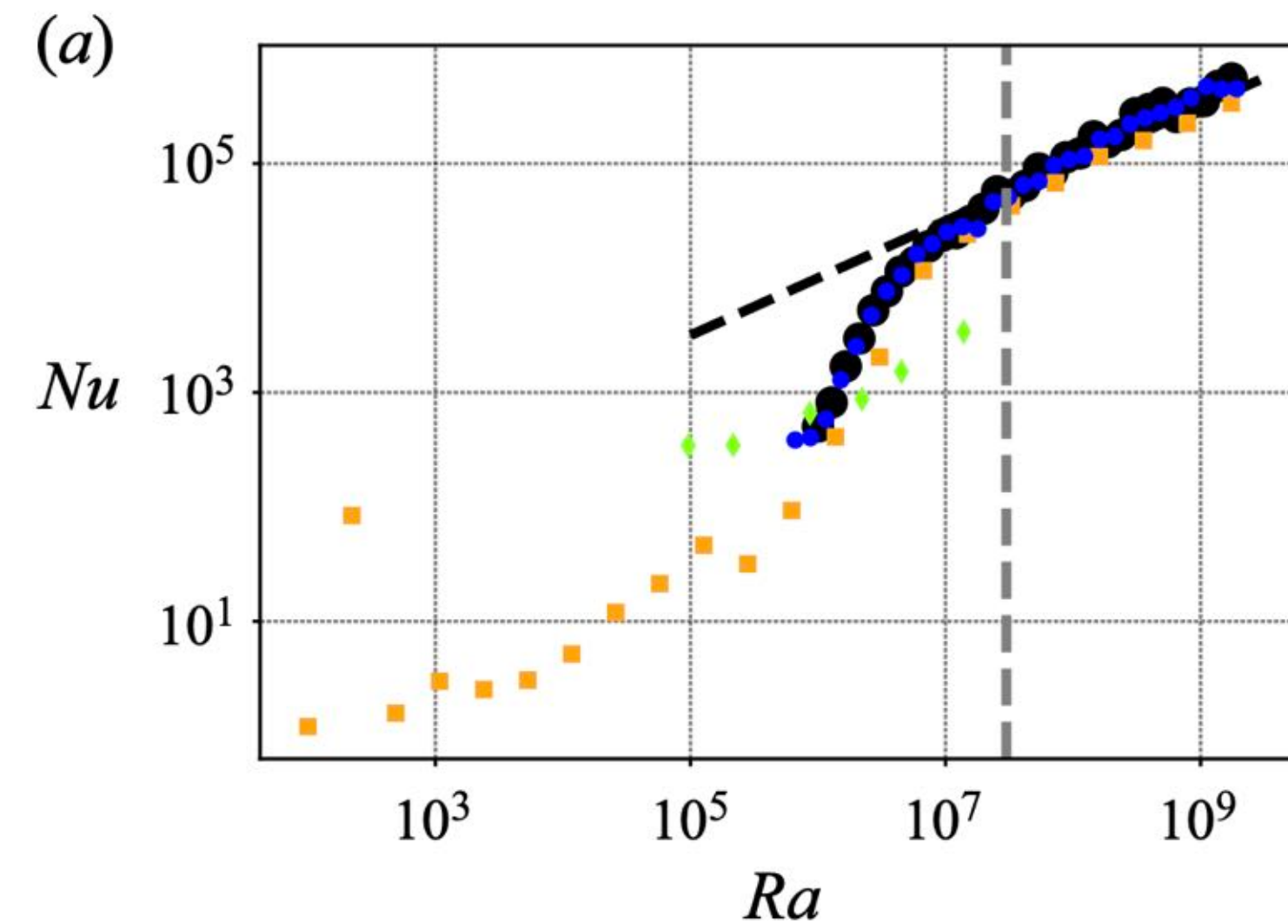


Log-Lattice homogeneous RBC

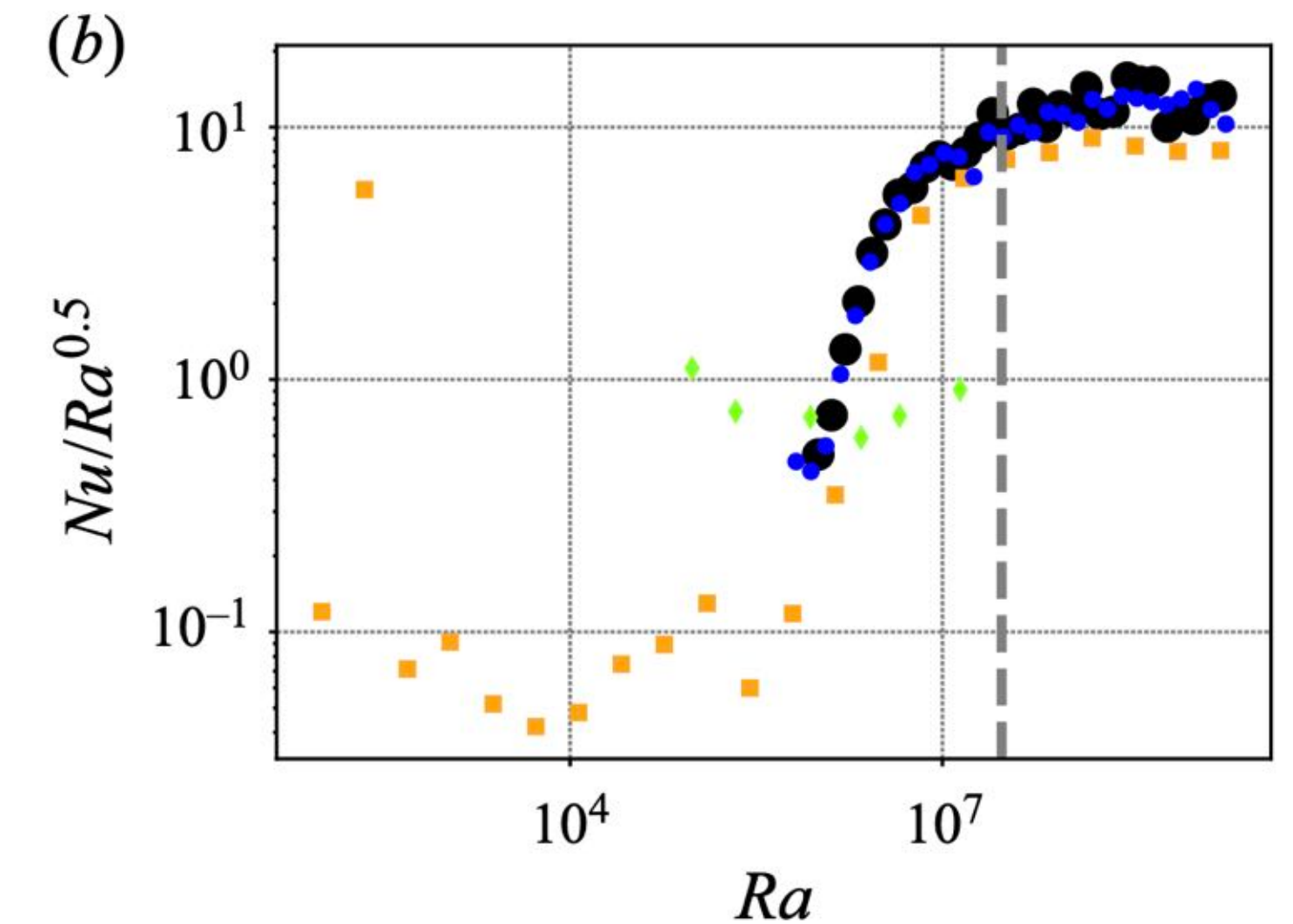
Barral & Dubrulle (2023)

- 3D log-lattice
- Need large-scale drag
- Friction modifies low Pr
- See hom. “ultimate regime” with GL scalings for Nu, Re

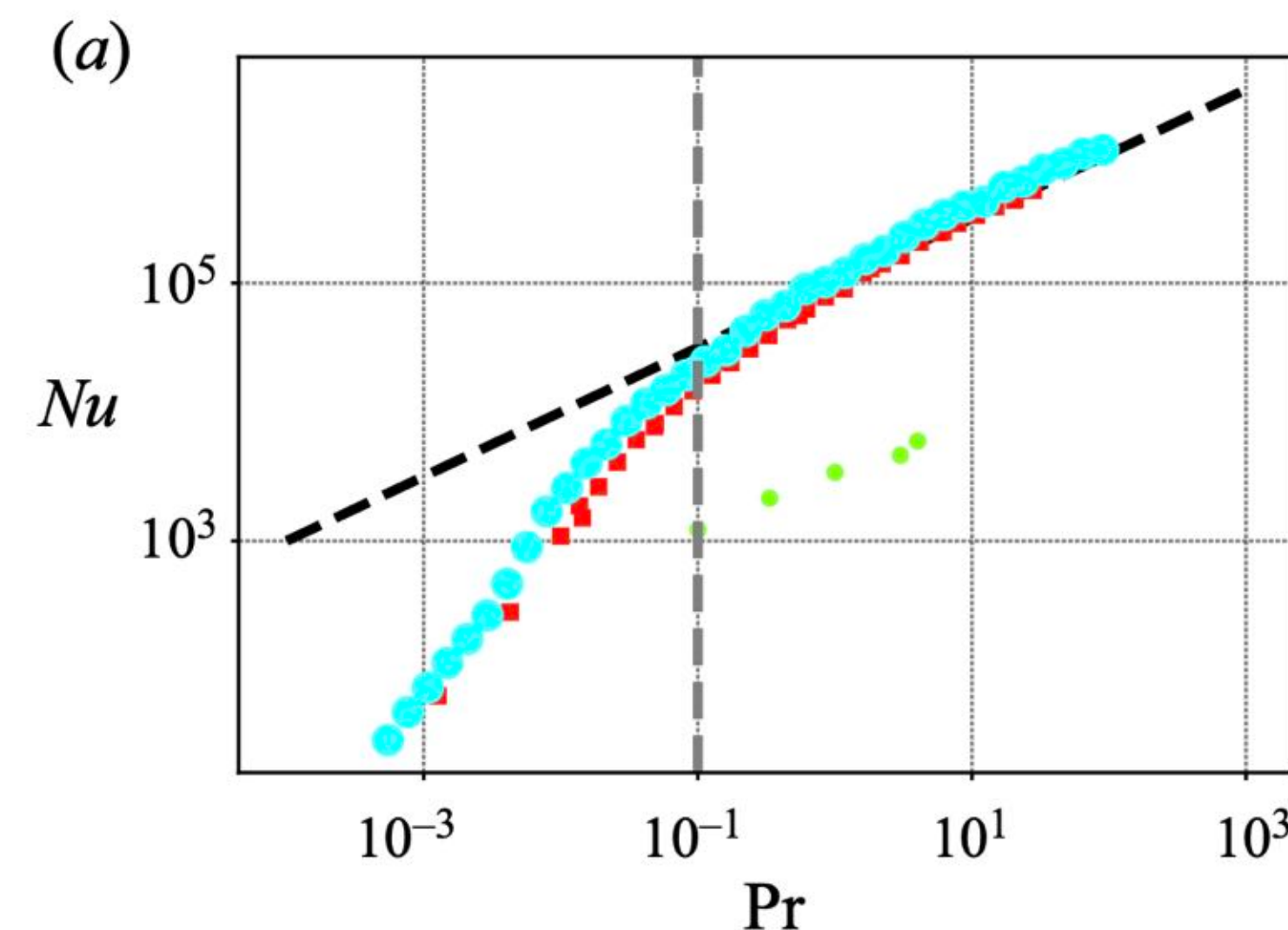
Nusselt at Pr=1



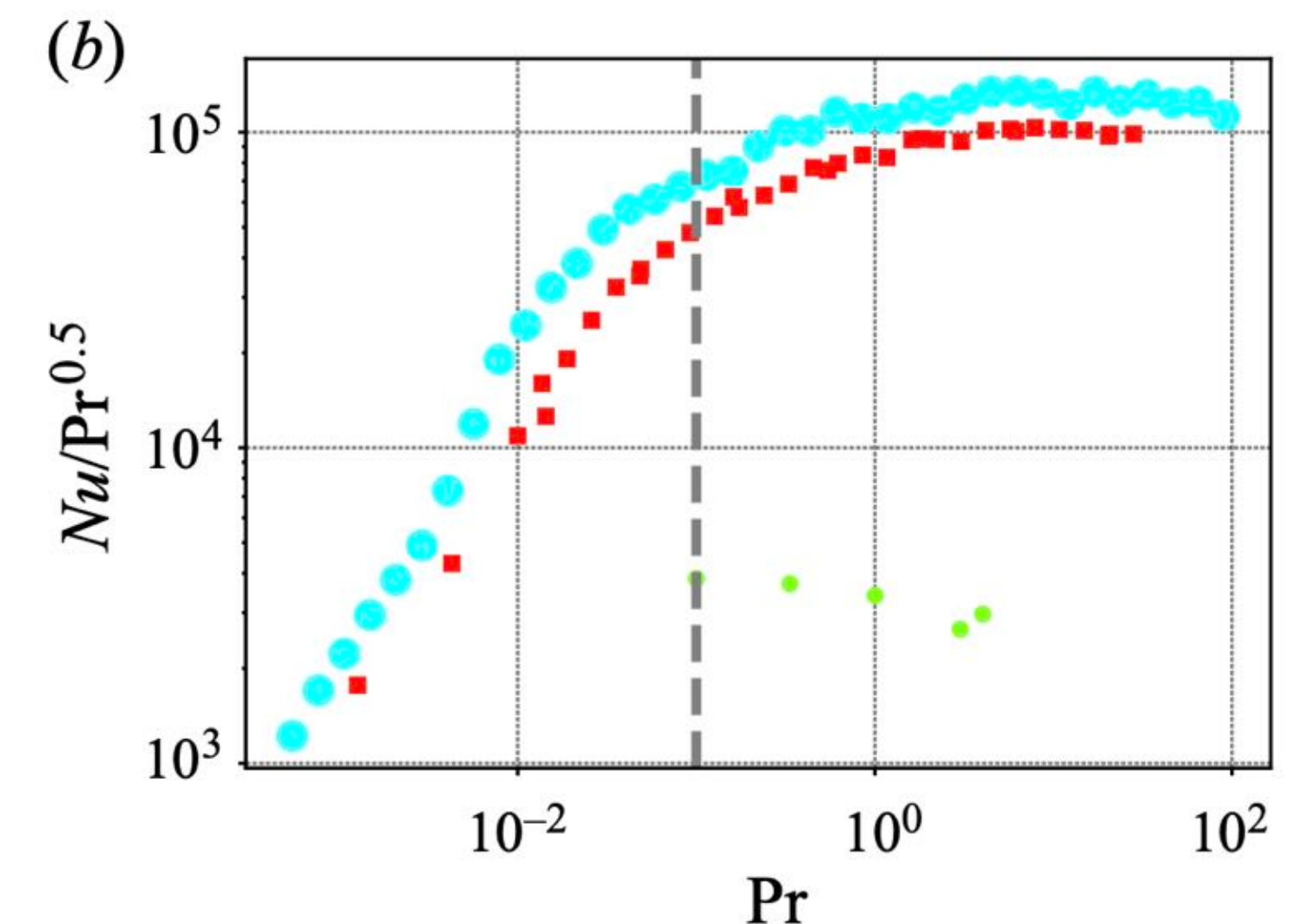
$Nu / Ra^{1/2}$



Nusselt at Ra=10^8



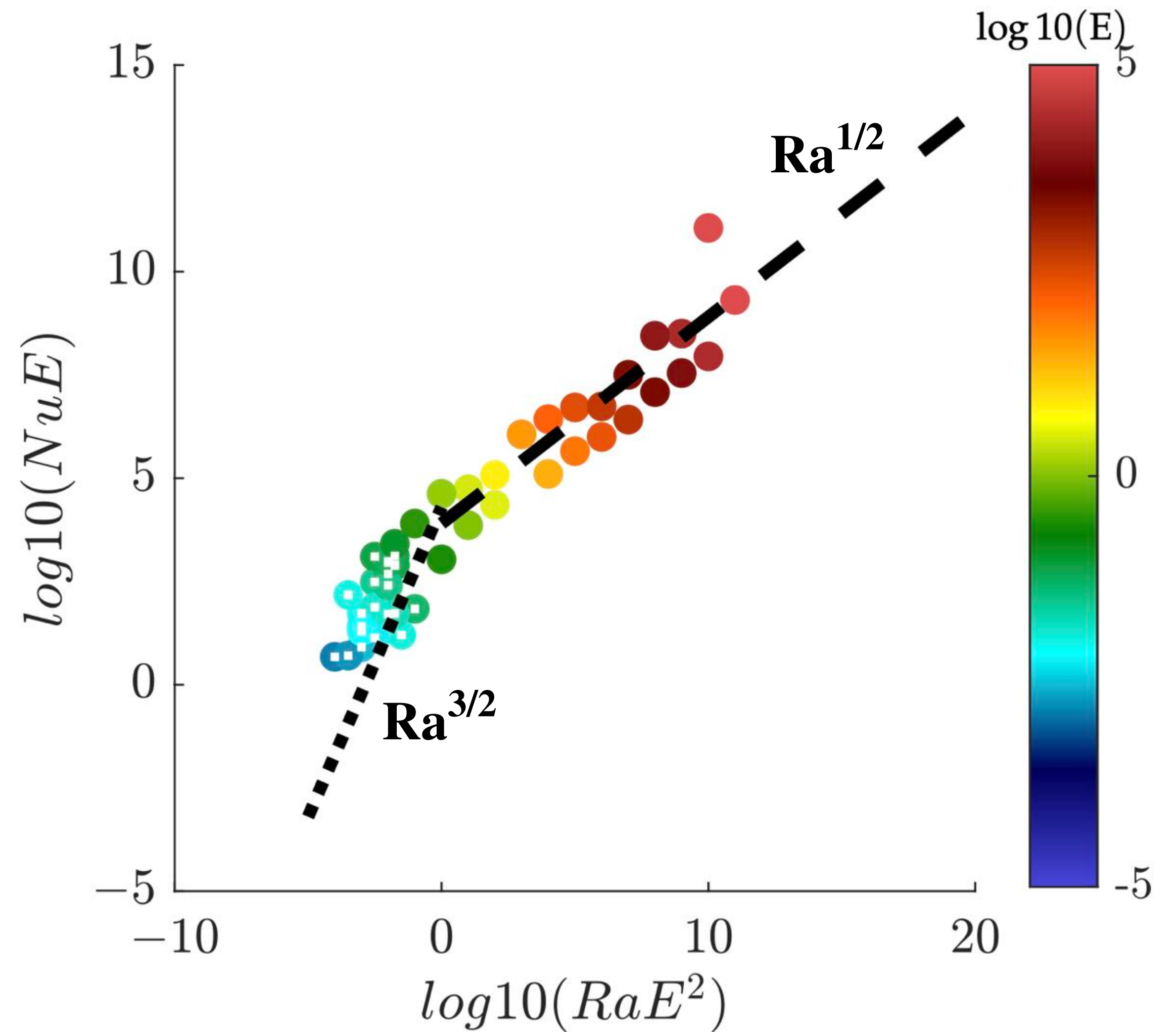
$Nu / Pr^{1/2}$



Log-Lattice homogeneous RRBC

Pikeroen et al. (2023)

- 3D log-lattice
- Need large-scale drag
- See geostrophic and ultimate regime scalings



Extreme RBC requires BL resolution globally

Transition to the Ultimate Regime in Two-Dimensional Rayleigh-Bénard Convection

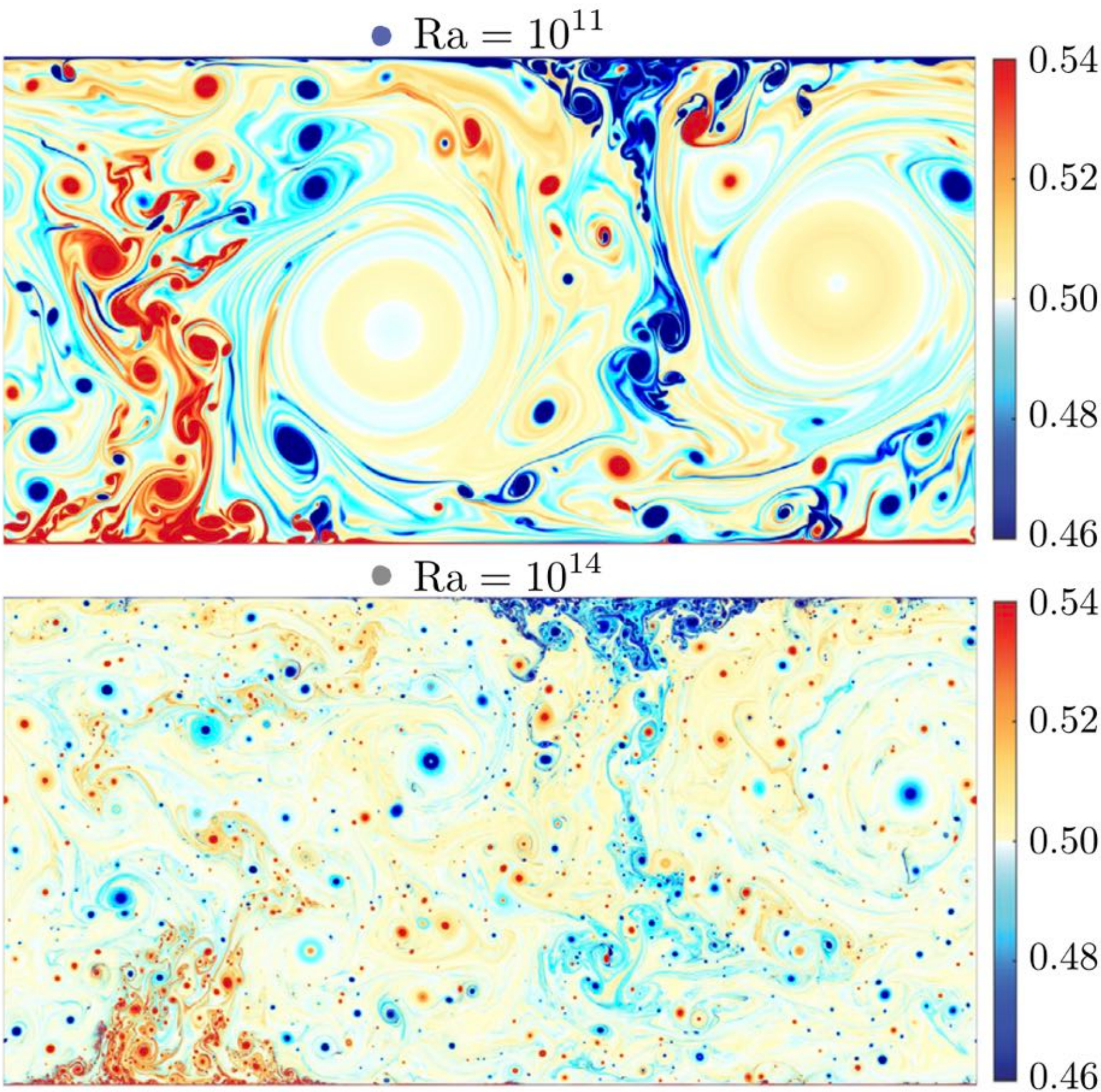
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²*Dipartimento di Ingegneria Industriale, University of Rome “Tor Vergata”, Via del Politecnico 1, Roma 00133, Italy*

³*Max Planck Institute for Dynamics and Self-Organization, 37077 Göttingen, Germany*

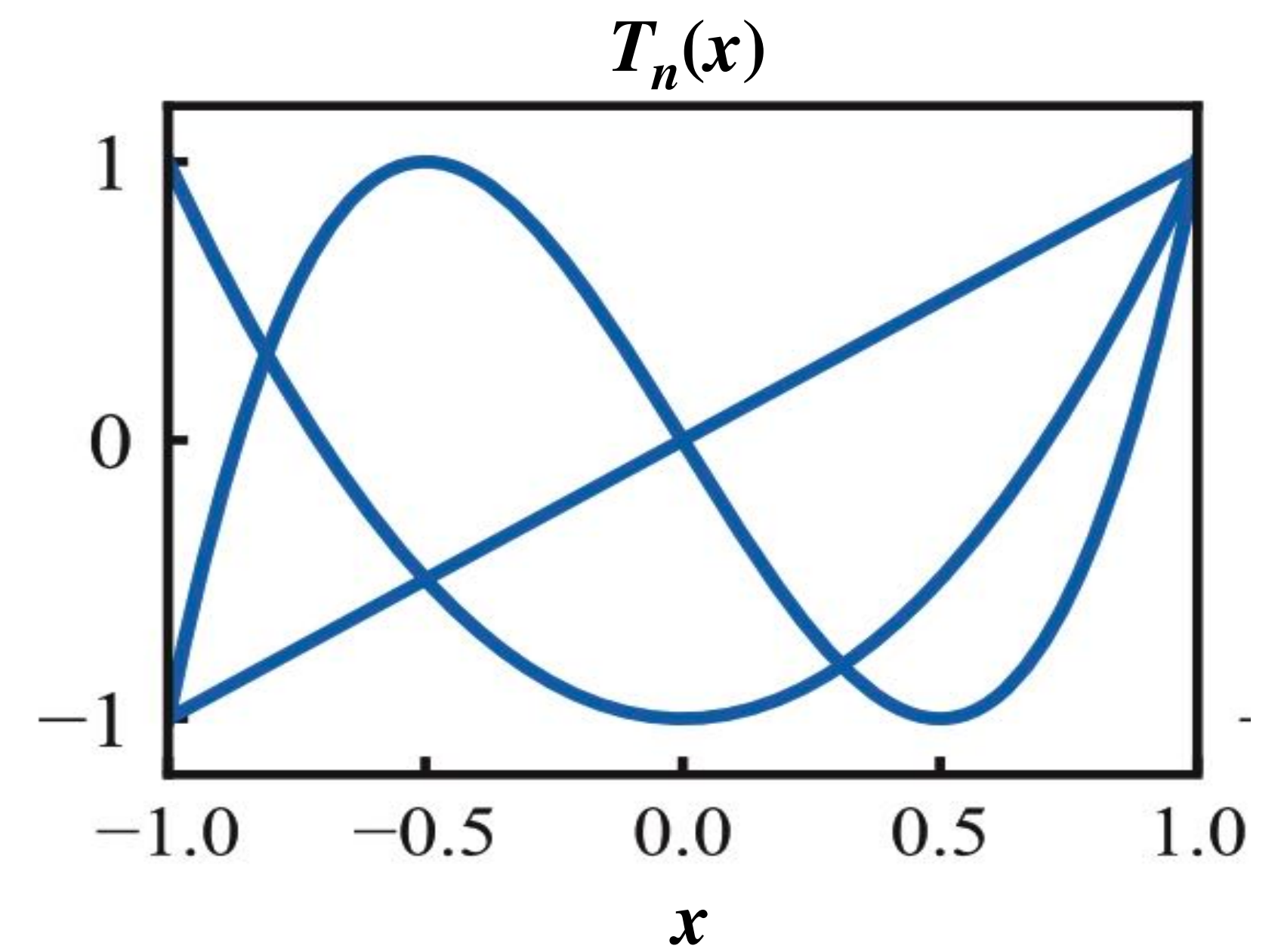
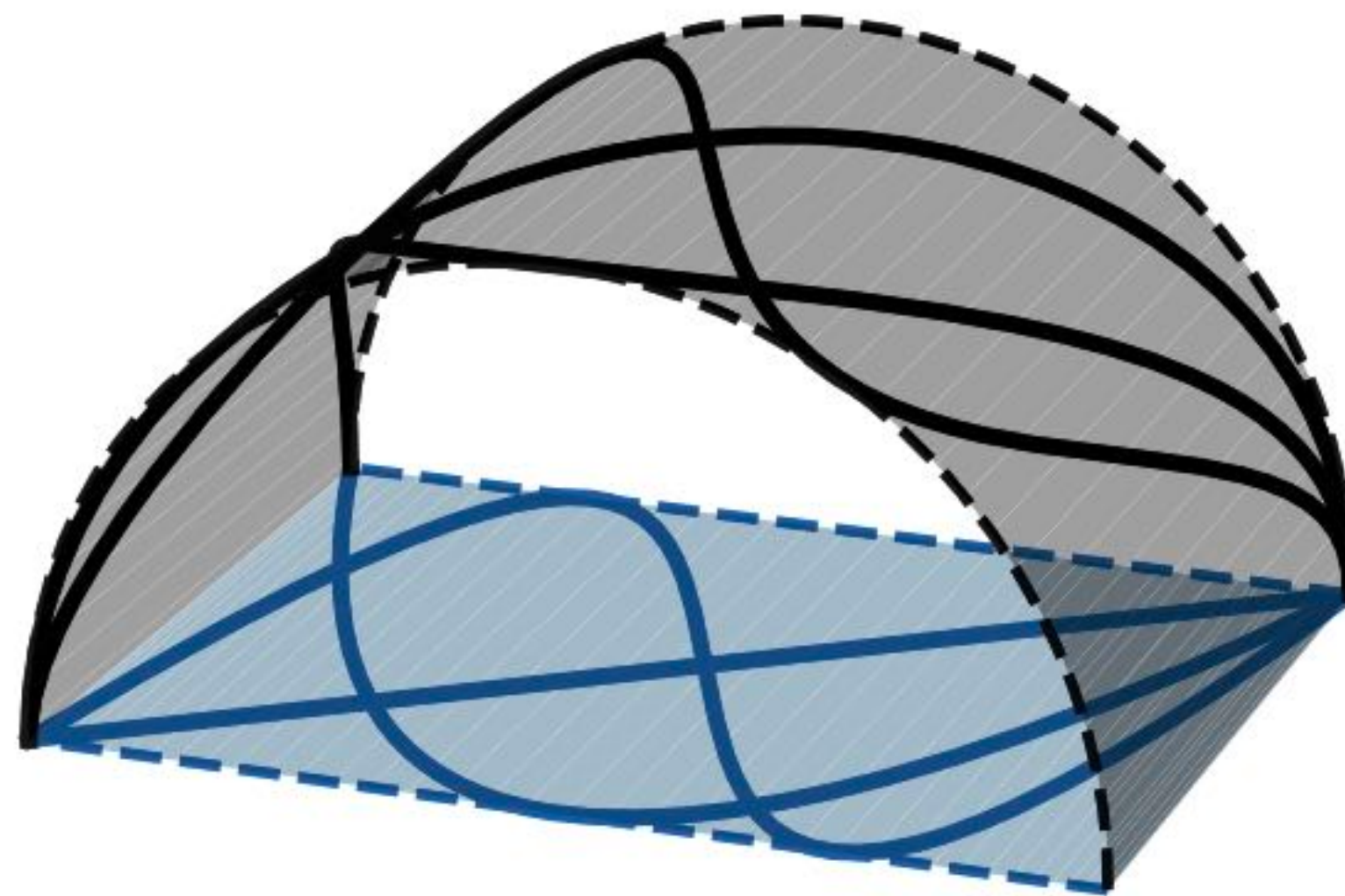
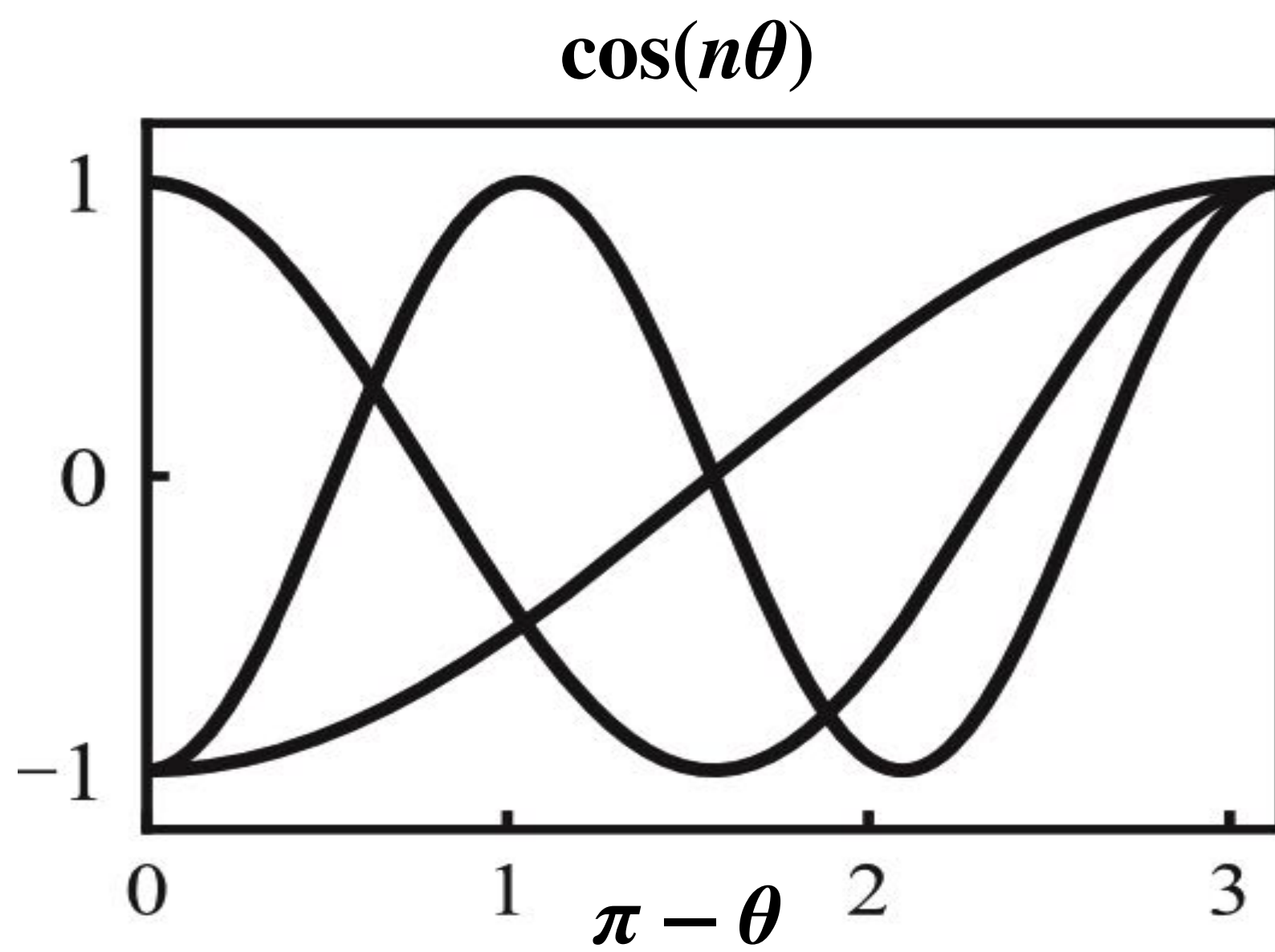
 (Received 4 January 2018; published 6 April 2018)



Ra	Pr	Γ	$N_x \times N_z$
10^8	1	2	512×256
2.15×10^8	1	2	768×384
4.64×10^8	1	2	768×384
10^9	1	2	1024×512
2.15×10^9	1	2	1280×640
4.64×10^9	1	2	1536×768
10^{10}	1	2	2048×1024
2.15×10^{10}	1	2	3072×1536
4.64×10^{10}	1	2	3072×1536
10^{11}	1	2	4096×2048
2.15×10^{11}	1	2	5120×2560
4.64×10^{11}	1	2	7680×3840
10^{12}	1	2	8192×4096
2.15×10^{12}	1	2	10240×5120
4.64×10^{12}	1	2	10240×5120
10^{13}	1	2	12288×6144
2.15×10^{13}	1	2	14336×7168
4.64×10^{13}	1	2	16384×8192
10^{14}	1	2	20480×10240

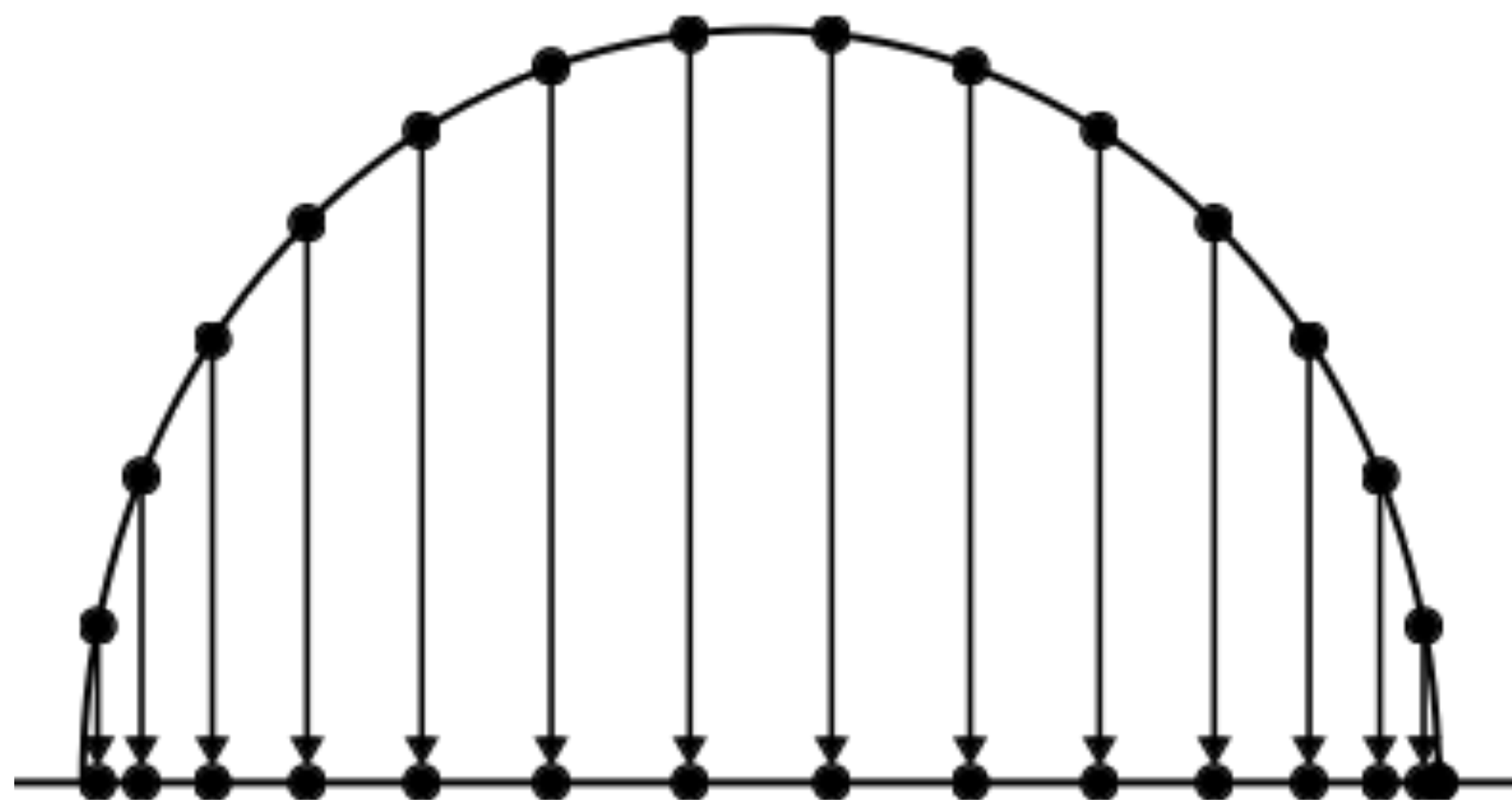
Sparse Chebyshev Methods

Chebyshev polynomials: Fourier in disguise

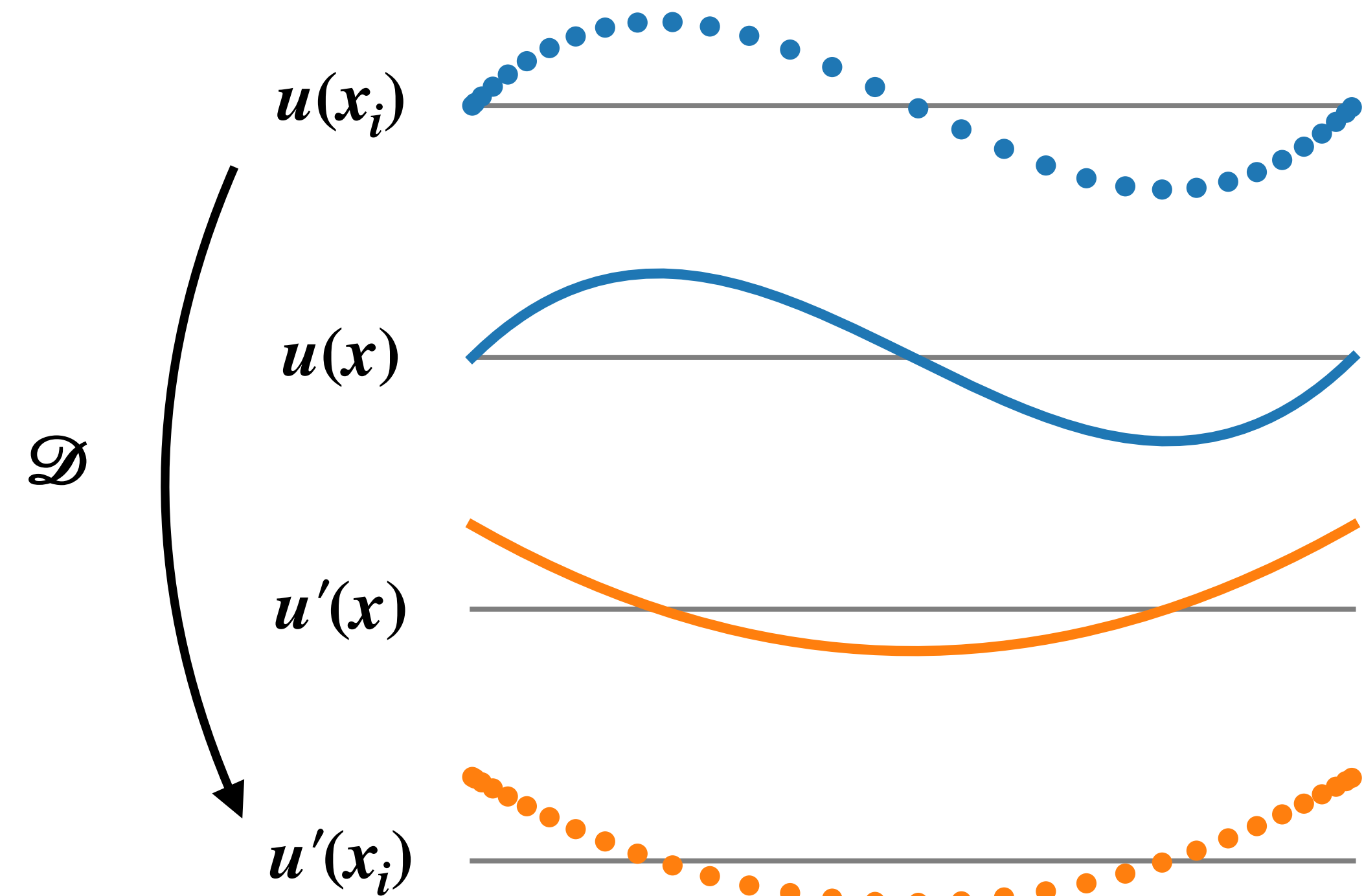


Chebyshev Collocation

Chebyshev nodes



Differentiation

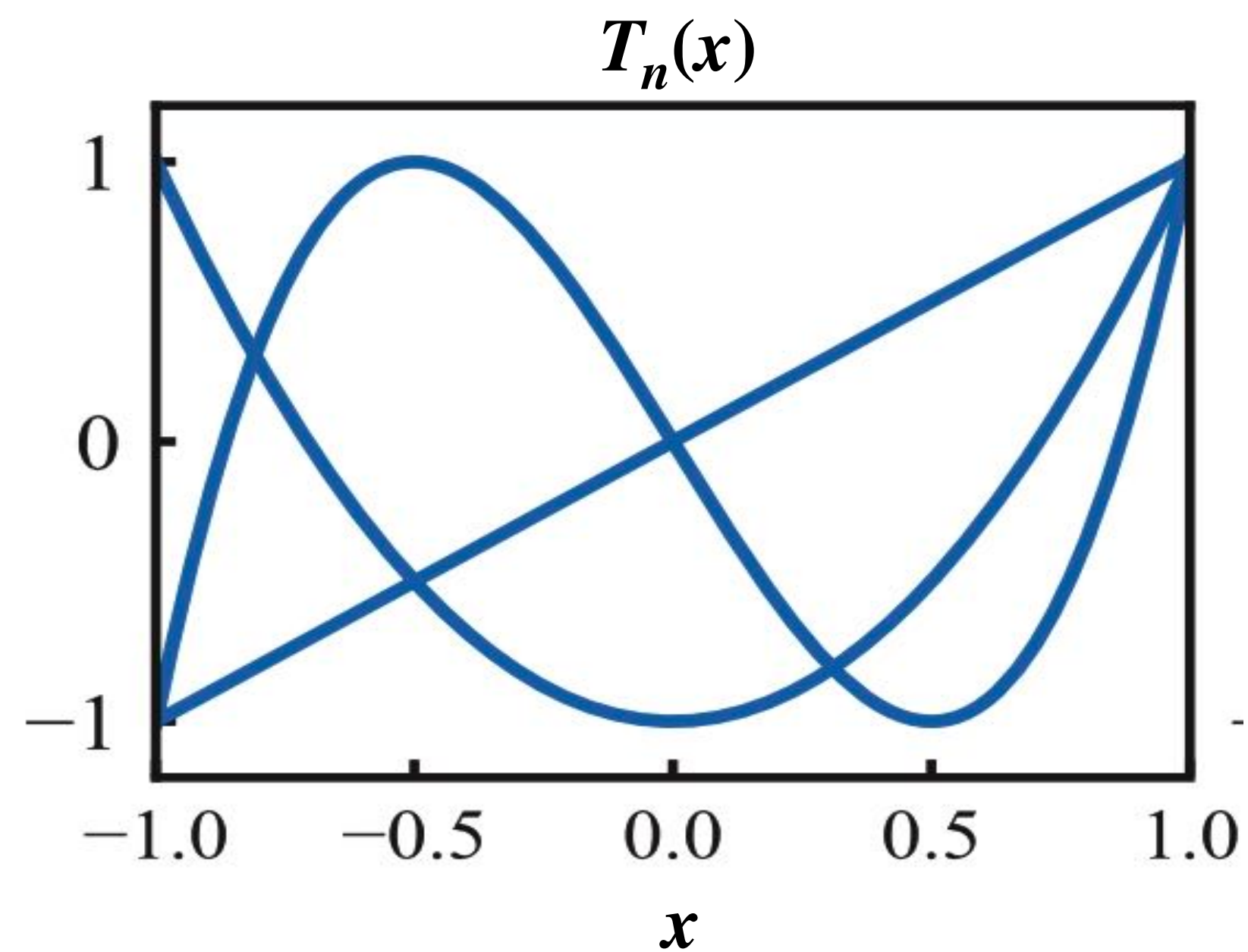


- Dense matrices
- Poor conditioning

Classical/Lanczos Tau Method

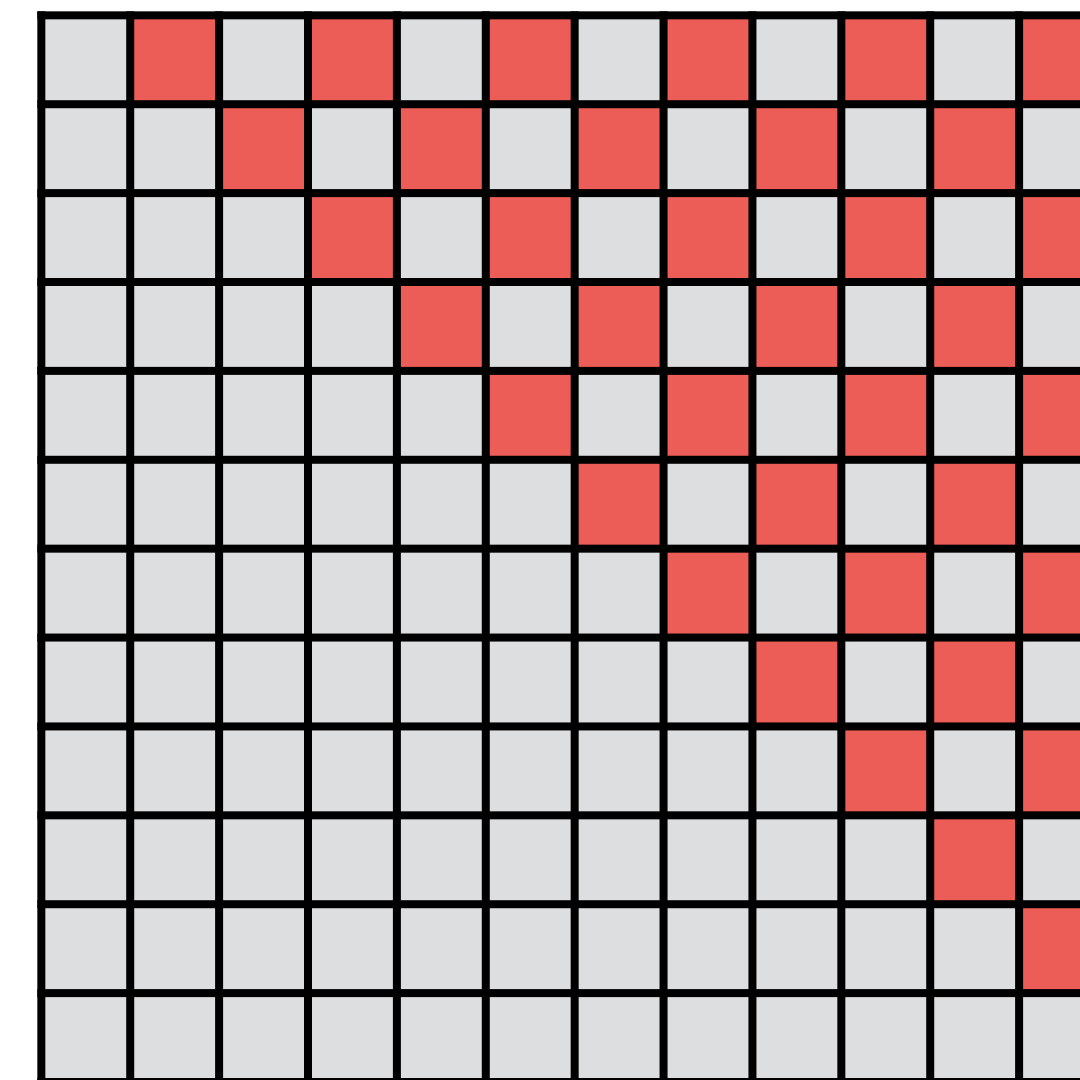
Chebyshev basis functions

$$u(x) = \sum_{n=0}^N u_n T_n(x)$$



Differentiation

$$\mathcal{D}_{m,n} = \langle T_m | \partial_x T_n \rangle$$



- Dense matrices
- Poor conditioning

Lanczos, Liu, Ortiz



Quasi-Inverse Technique

$$u(x) = \sum_{n=0}^N u_n T_n(x), \quad u'(x) = \sum_{n=0}^{N-1} u'_n T_n(x)$$

Derivative recursion:

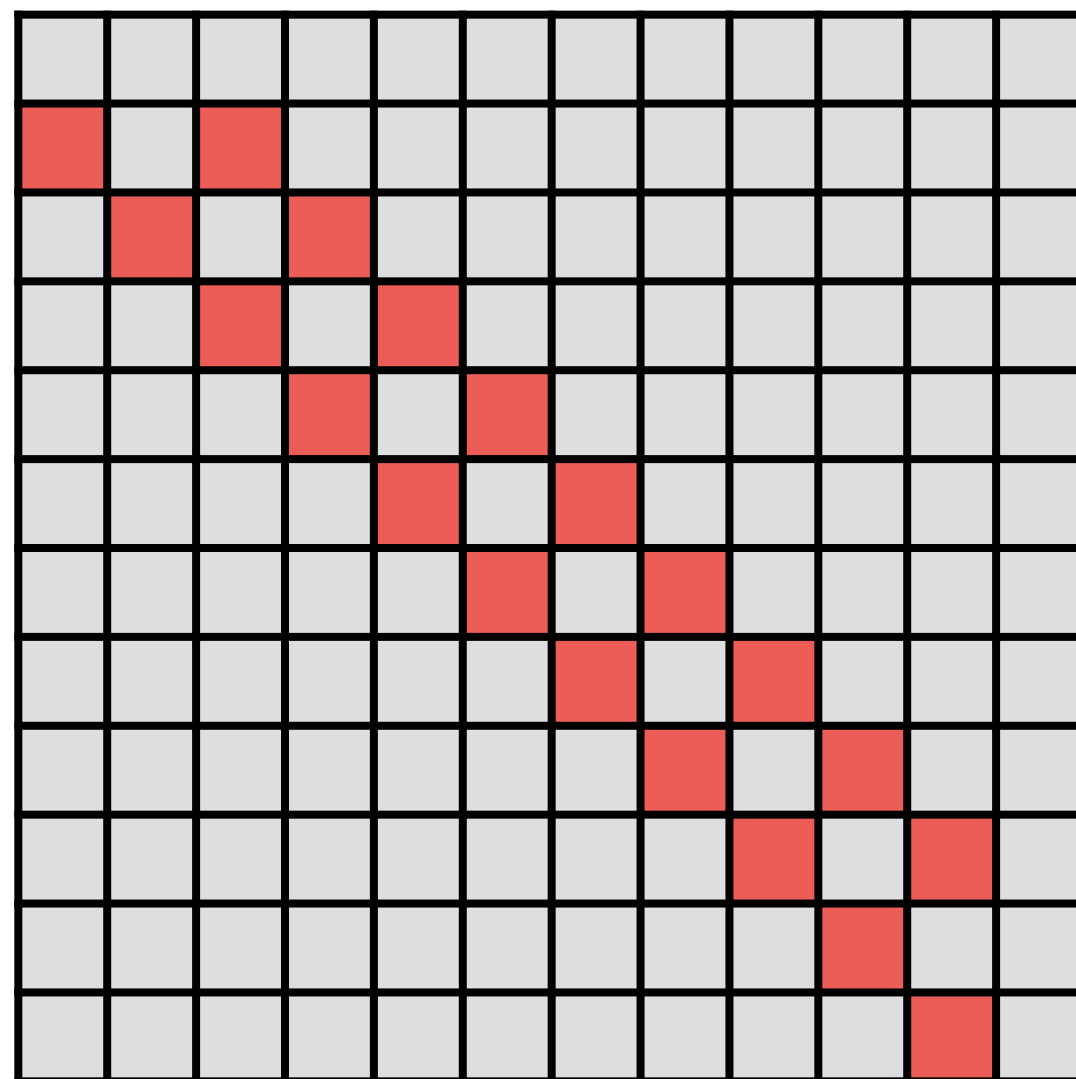
$$T_n(x) = \frac{T'_{n+1}(x)}{n+1} - \frac{T'_{n-1}(x)}{n-1} \implies c_{n-1}u'_{n-1} - u'_{n+1} = 2nu_n$$

Pseudoinverse

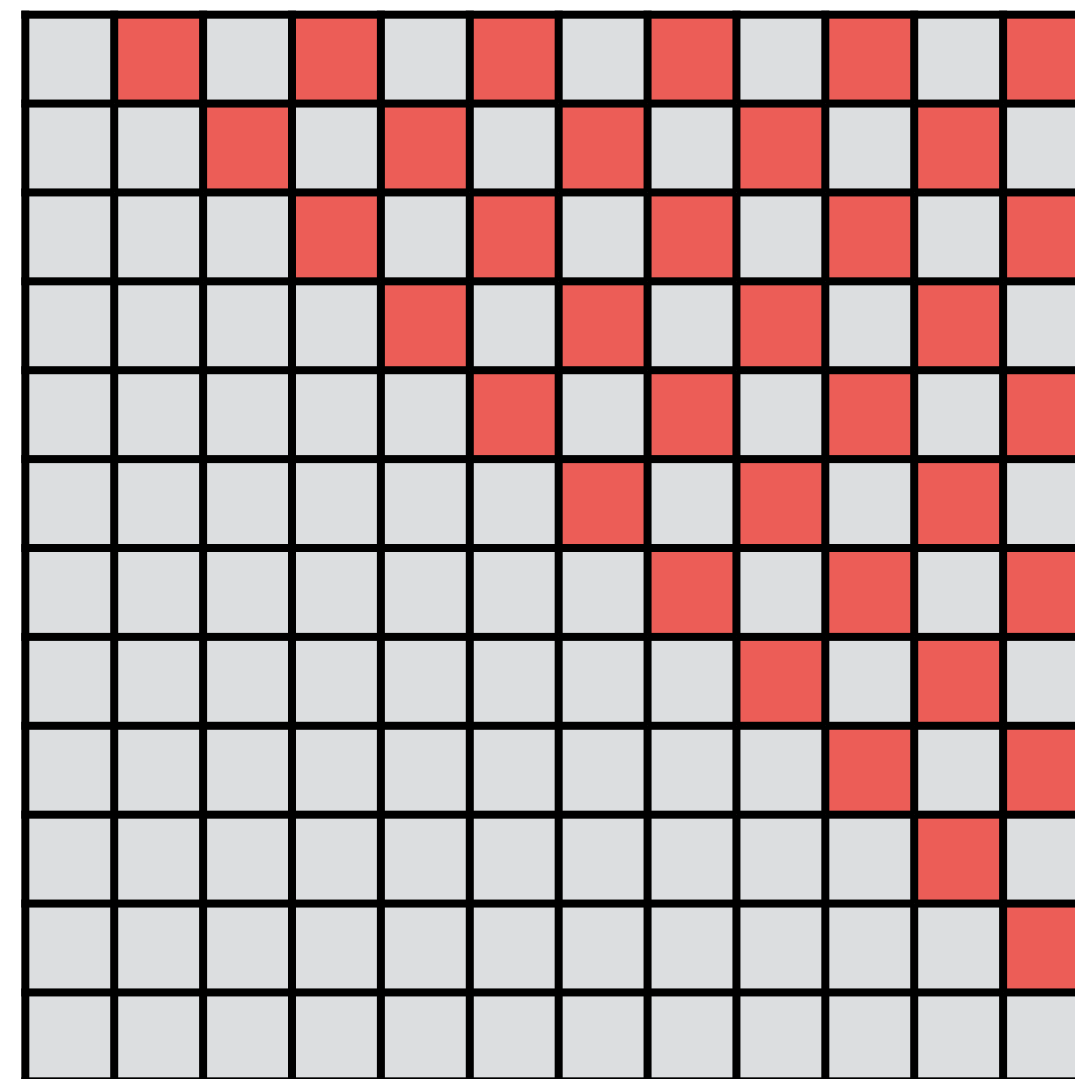
$\mathcal{B}$

$\mathcal{D}$

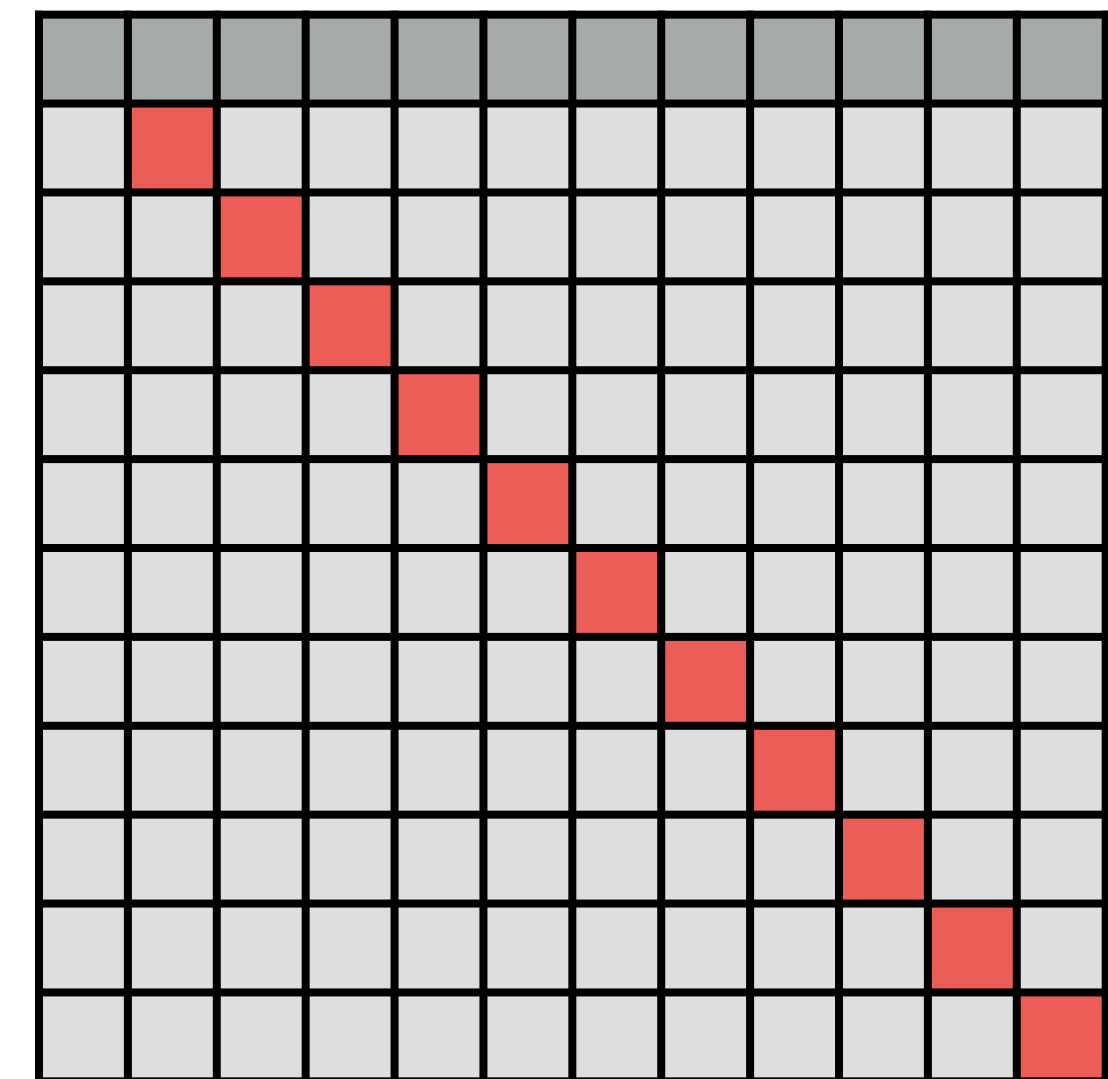
$I^{(1)}$



•



=



Gottlieb, Orszag, Greengard, Dang-Vu, Julien & Watson + friends

Ultraspherical Method

Chebyshev “trial” functions:

$$u(x) = \sum_{n=0}^N u_n T_n(x)$$

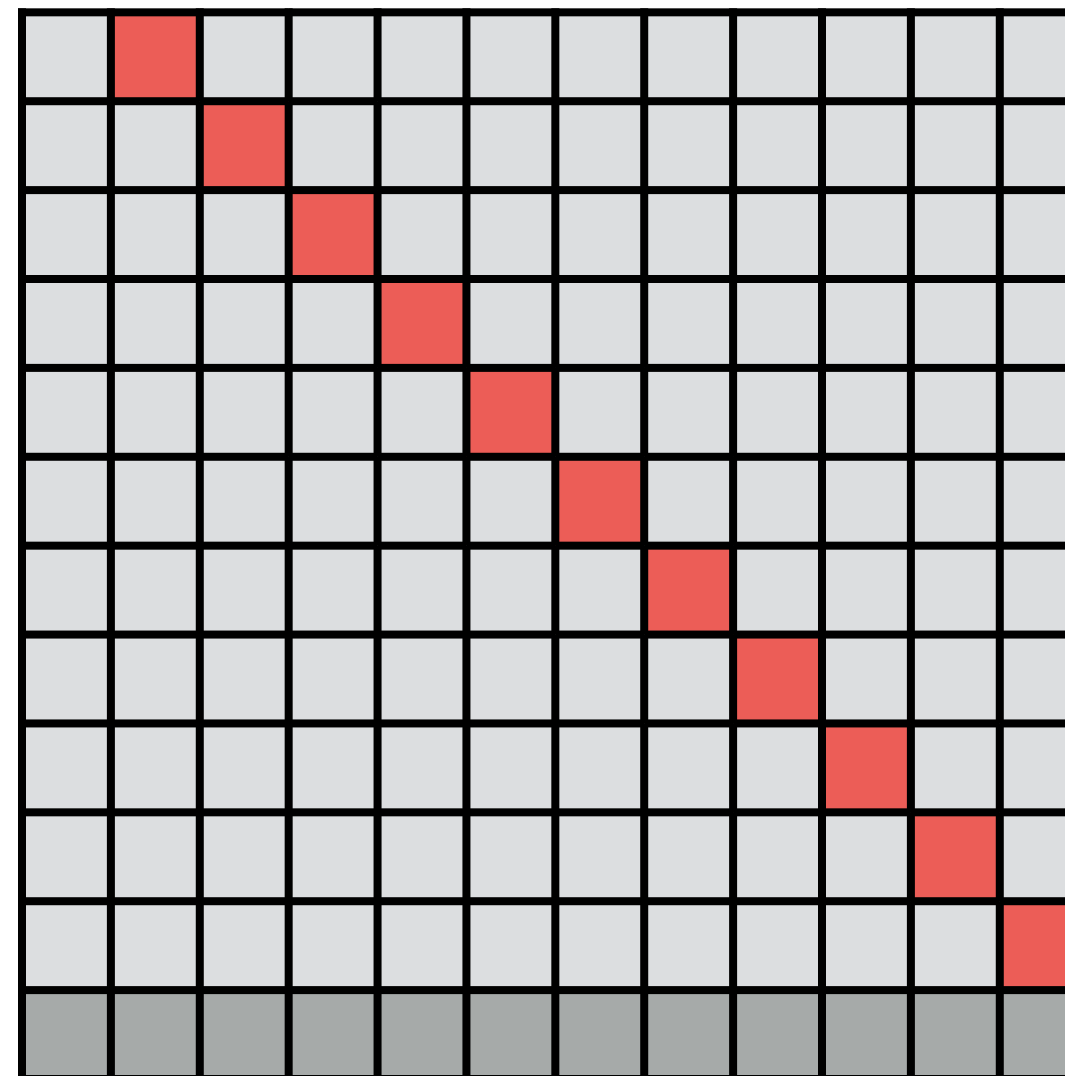
Ultraspherical “test” functions:

$$C_n^{(k)}(x) \propto \partial_x^k T_{n+k}(x)$$

$$w(x) = (1 - x^2)^{k-1/2}$$

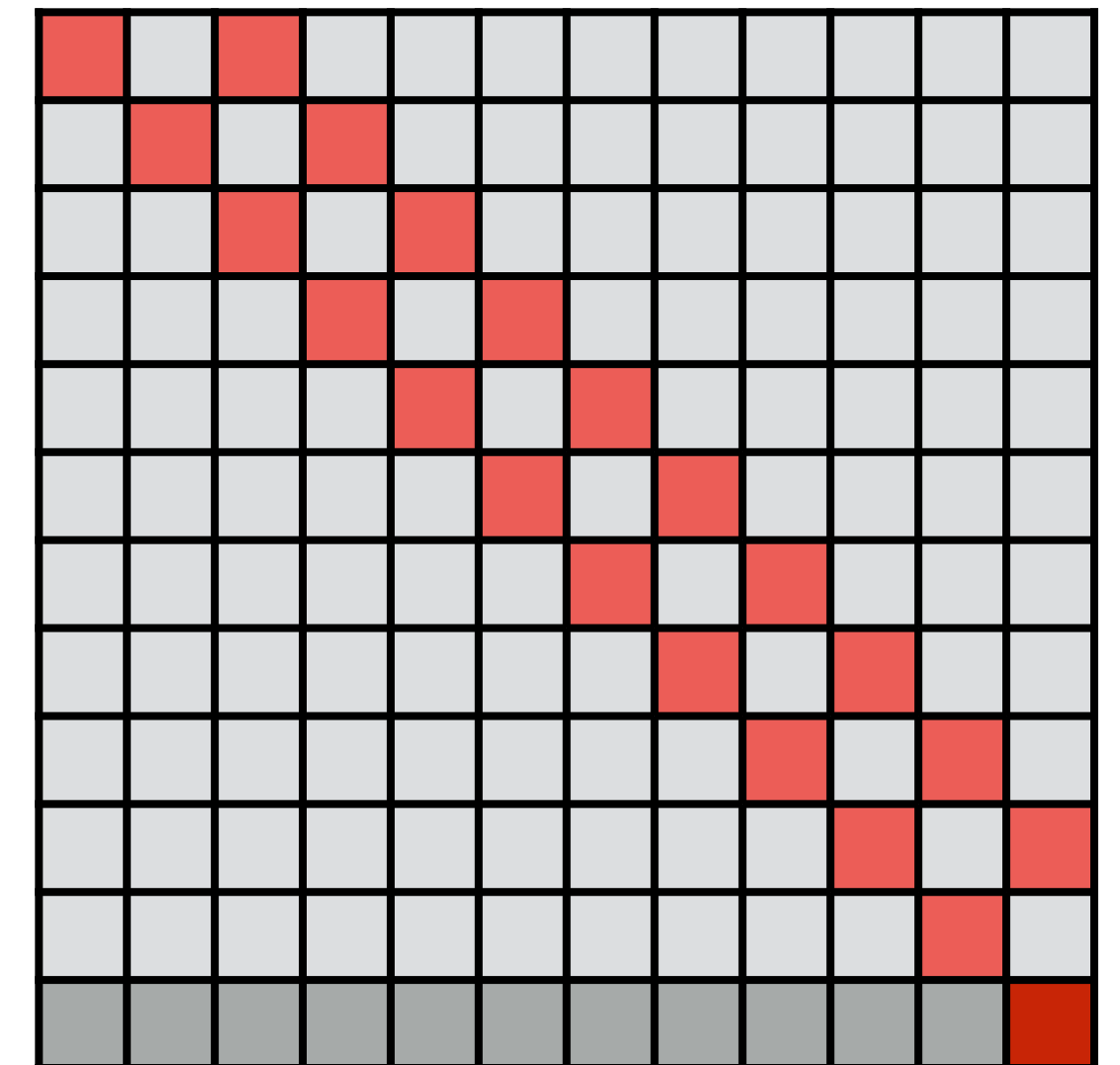
Differentiation

$$\mathcal{D}_{m,n} = \langle C_m^{(1)} | \partial_x T_n \rangle$$



Conversion

$$\tilde{I}_{m,n} = \langle C_m^{(1)} | T_n \rangle$$

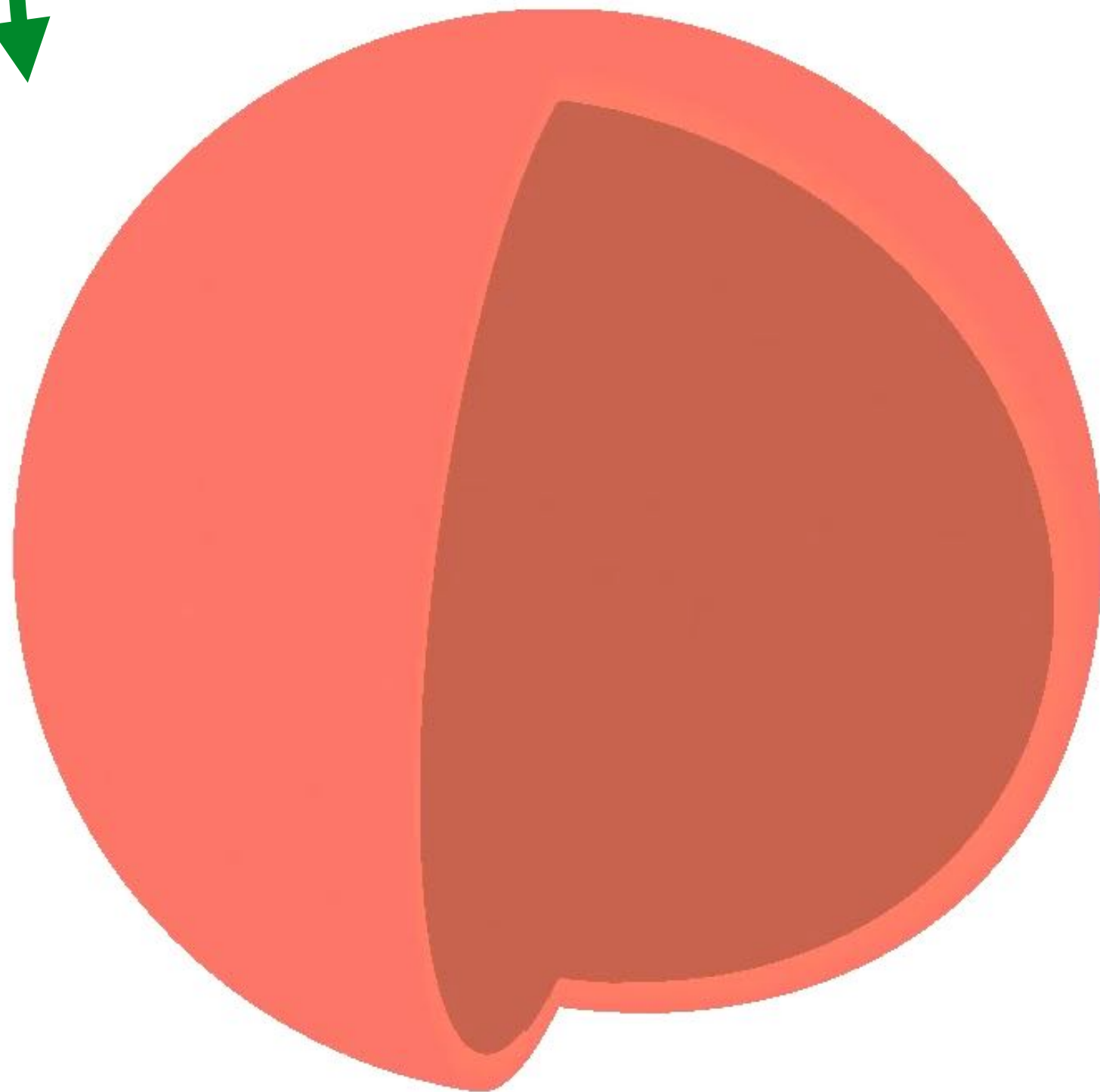


$$\tilde{I} \sim \mathcal{B}$$

Dedalus Project

```
problem.add_equation("div(u) = 0")
problem.add_equation("dt(u) - \nu*Lap(u) + grad(p) + b*g = - u@grad(u)")
problem.add_equation("dt(b) - \kappa*Lap(b) = - u@grad(b)")
```

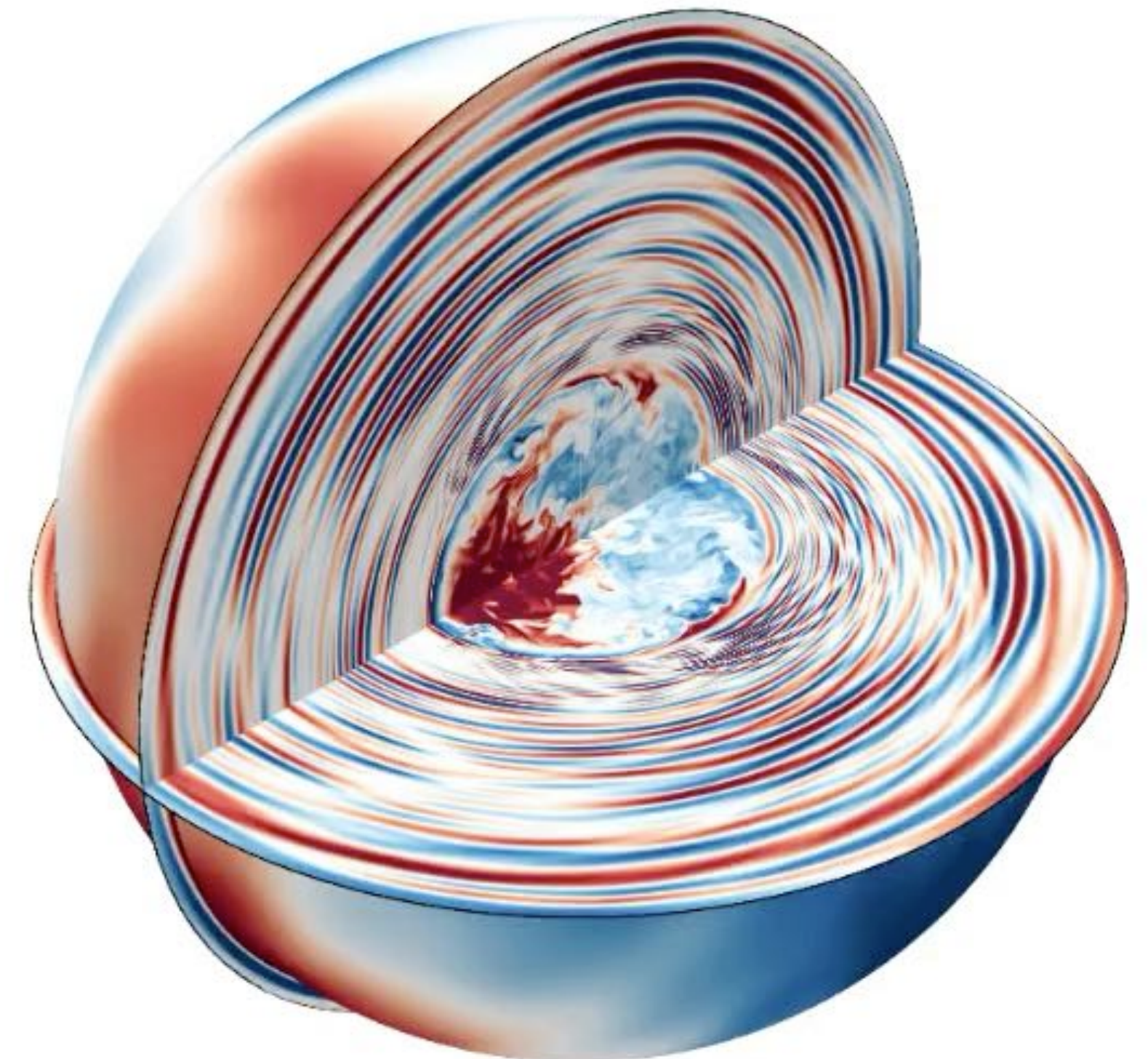
Rapid solver development
Spiral-defect chaos



Flexible equations
NLS quantum graphs

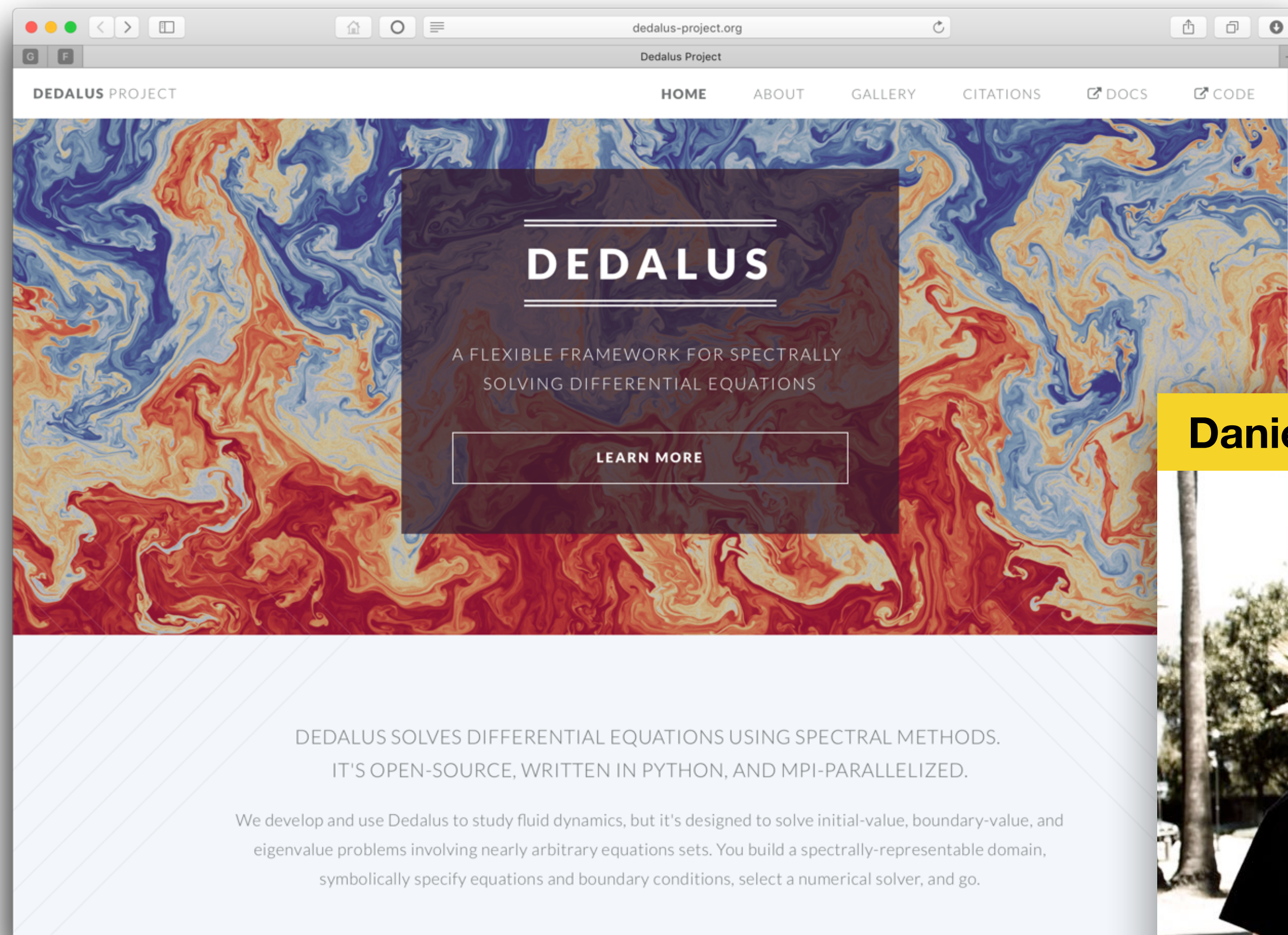


High performance
Turbulent wave excitation



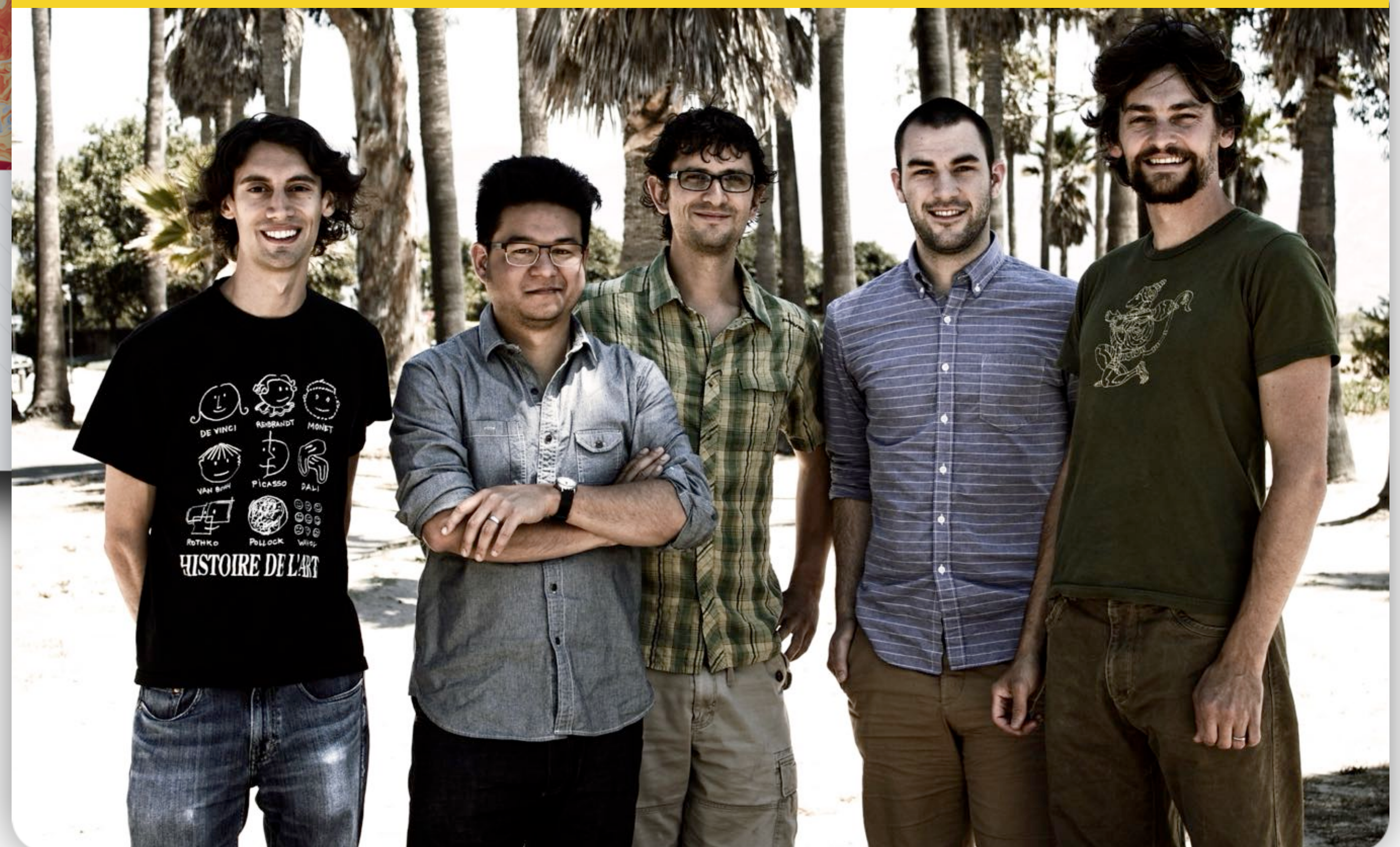
w/ Evan Anders

www.dedalus-project.org



Core developers

Daniel Lecoanet, Jeff Oishi, Geoff Vasil, Keaton Burns, Ben Brown



Log-Lattice Rayleigh Benard

Log-Fourier + Chebyshev

LFL resolution:

$$N_k \sim \log(L/\Delta x) \sim \log \text{Re}$$

Chebyshev resolution:

$$N_z = \frac{L_x}{\Delta x} \sim \text{Re}^{3/4}$$

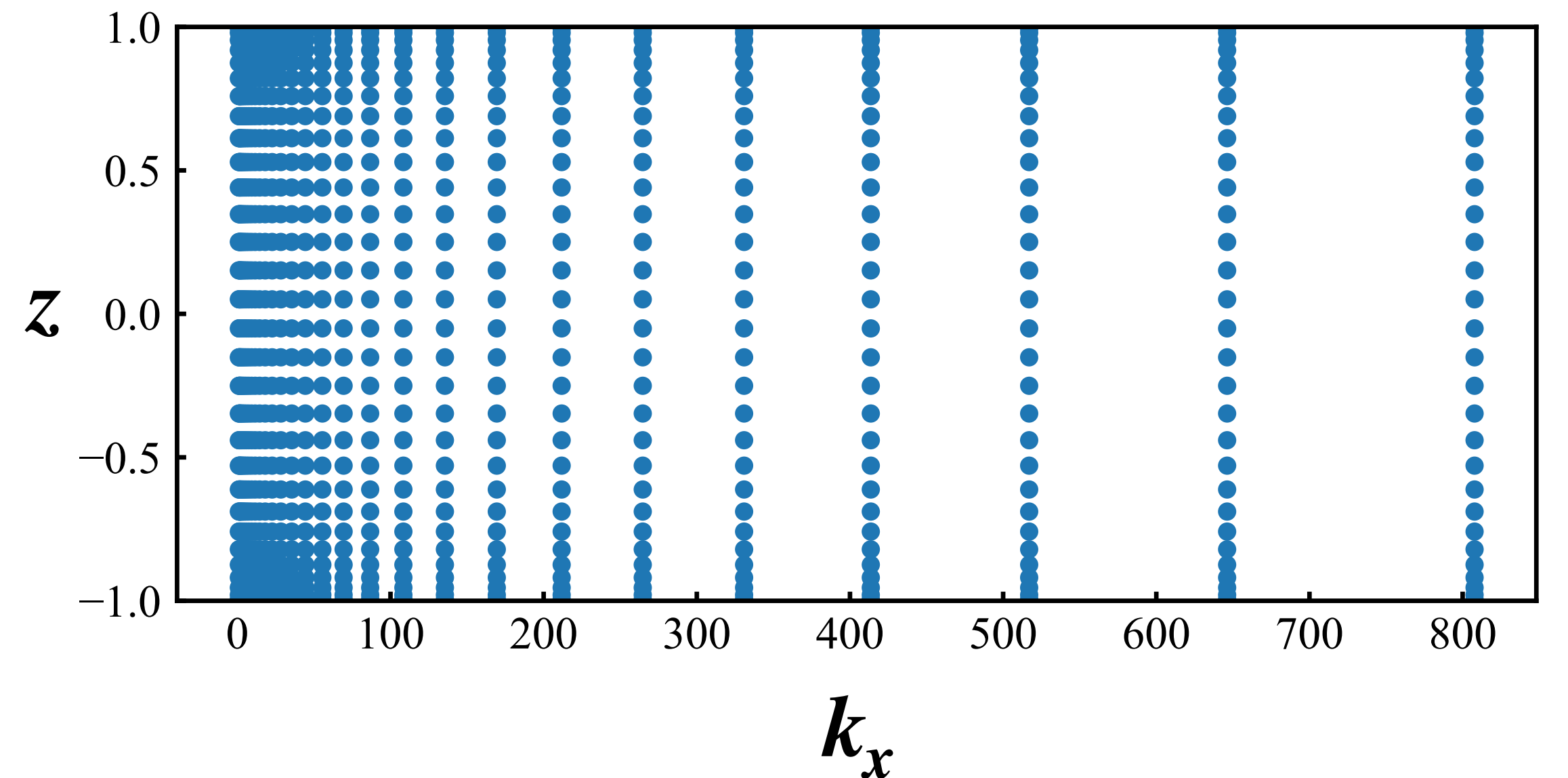
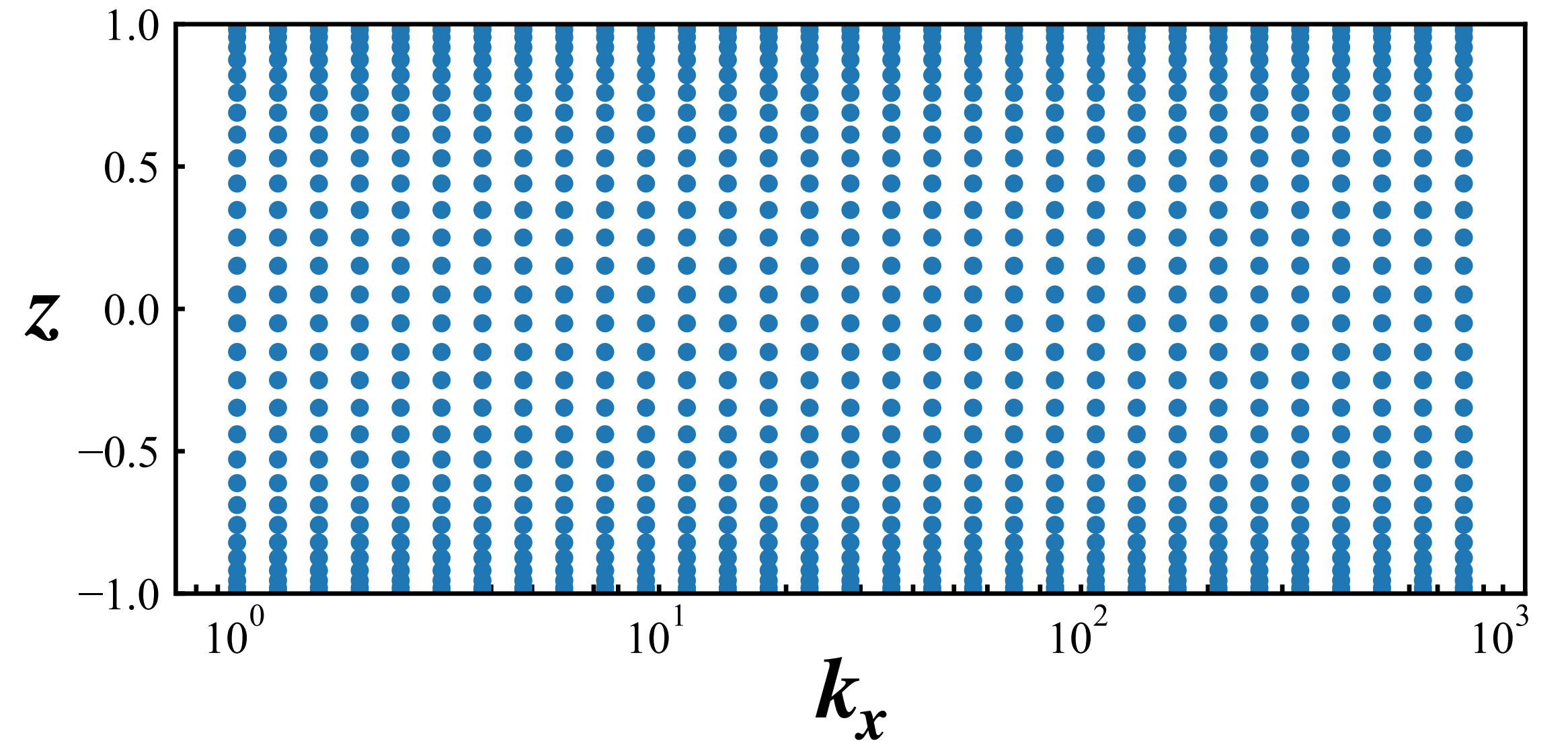
Temporal resolution:

$$N_t = \frac{T}{\Delta t} \sim \frac{L_x}{\Delta x} \sim \text{Re}^{3/4}$$

Total model DOF:

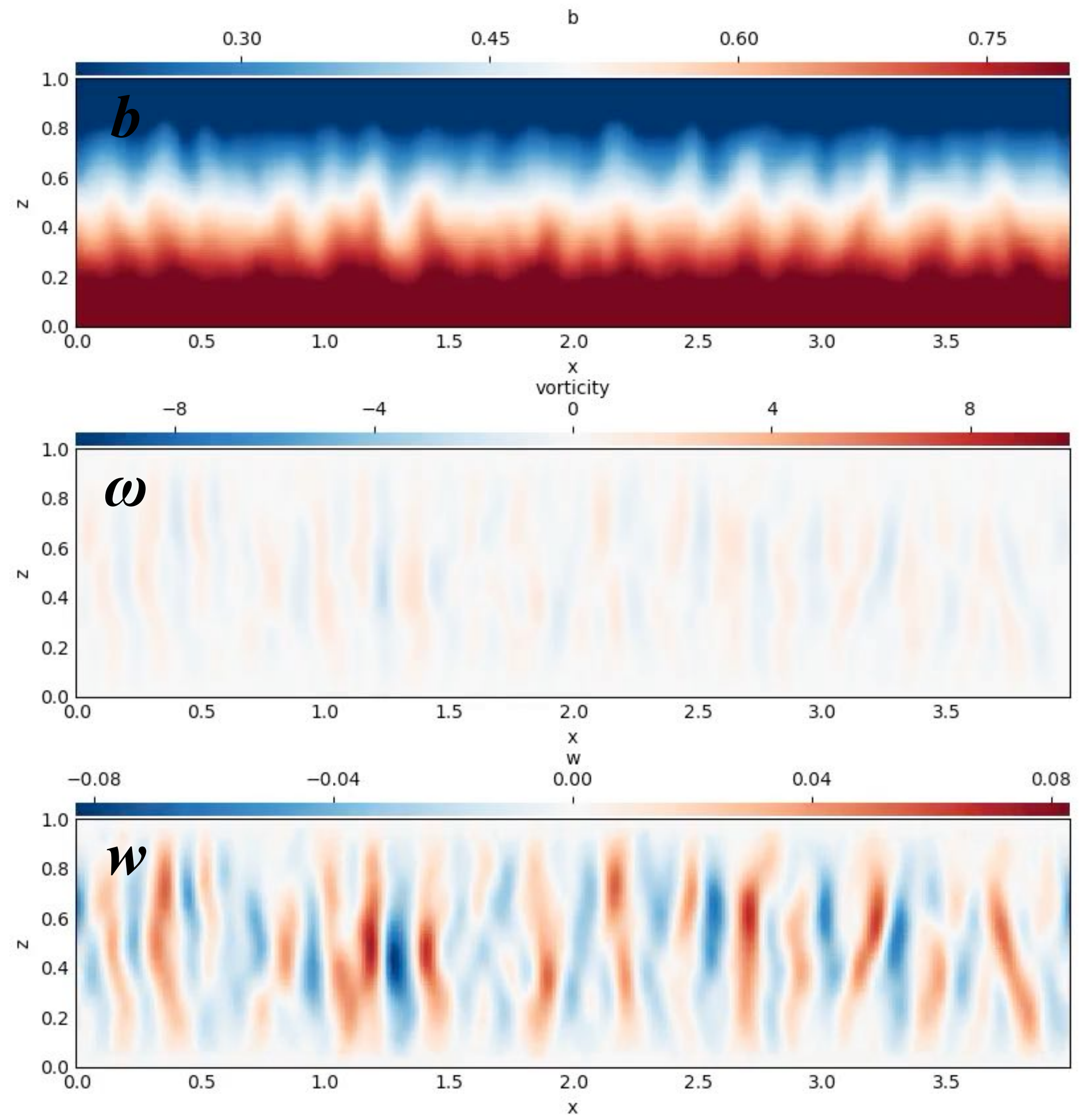
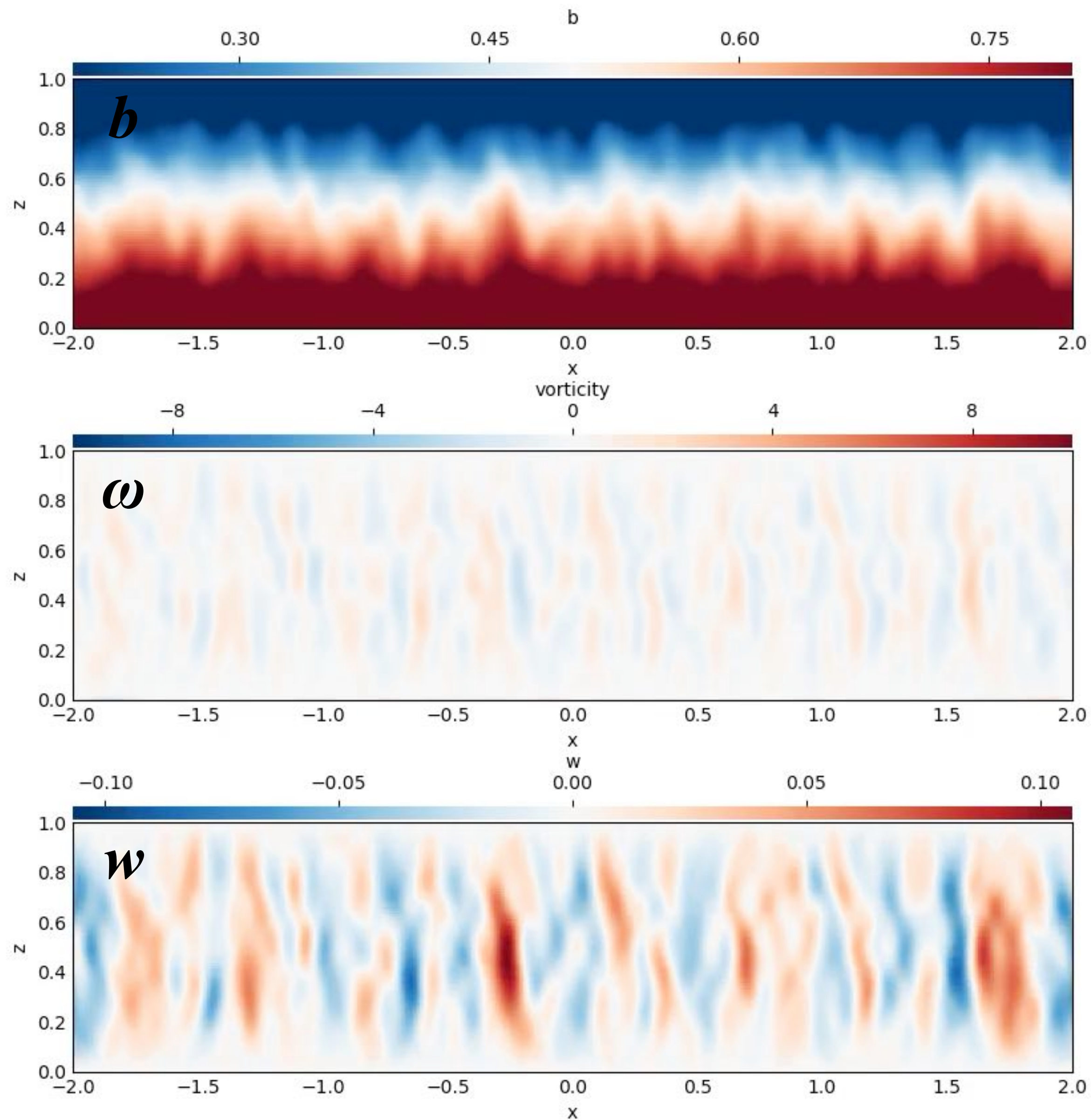
$$N_{\text{hom}} = N_k^d N_t \sim (\log \text{Re})^d \text{Re}^{3/4}$$

$$N_{\text{ch}} = N_k^{d-1} N_z N_t \sim (\log \text{Re})^{d-1} \text{Re}^{3/2}$$



Log-Fourier Rayleigh-Benard

$$\text{Ra} = 10^8$$

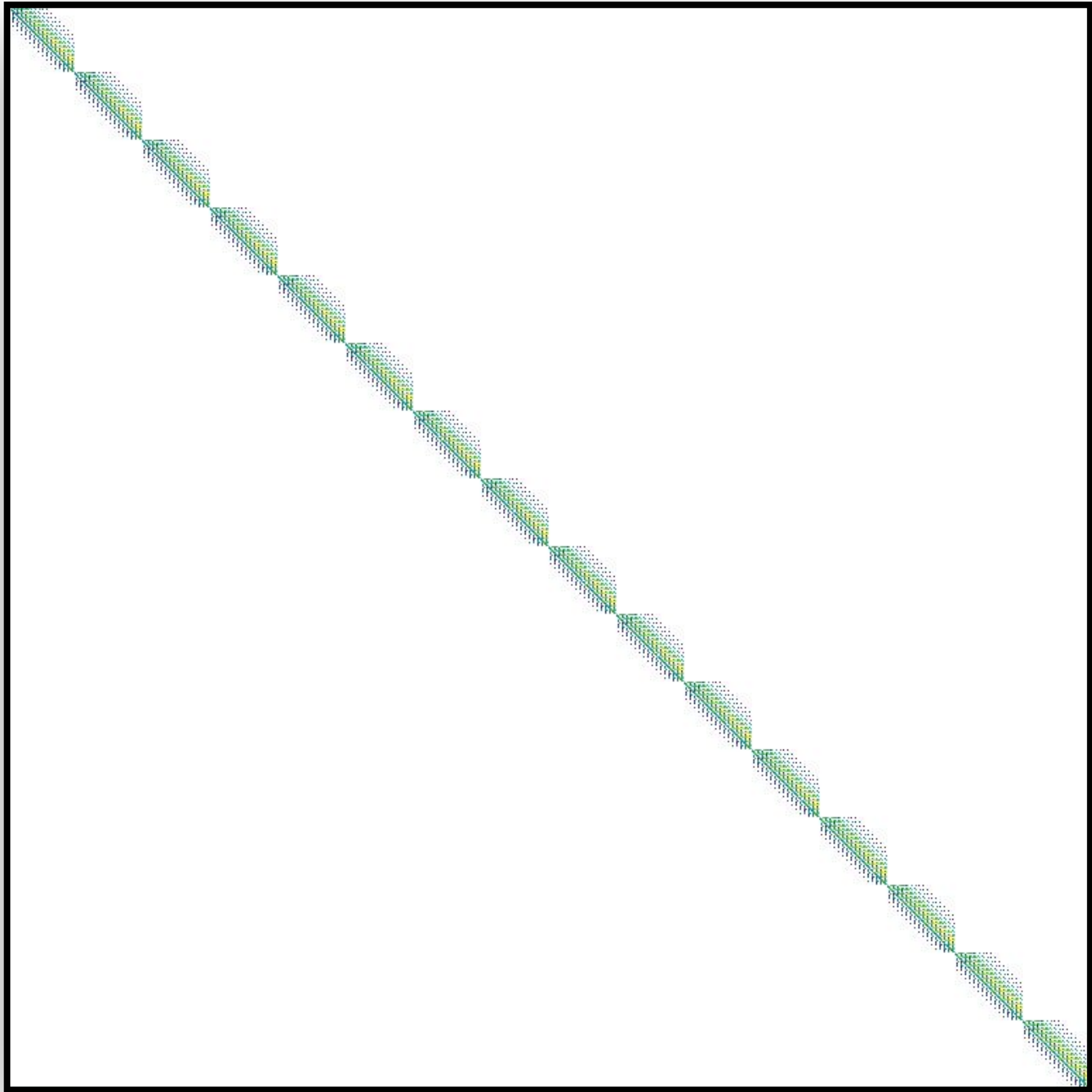


Log-Fourier Rayleigh-Benard

$\lambda = 1.16$

Ra	DNS resolution (Zhu et al. 2018)	Log lattice resolution
10^{10}	2,048 x 1,024	112 x 1,024
10^{11}	4,096 x 2,048	128 x 2,048
10^{12}	8,192 x 4,096	128 x 4,096
10^{13}	12,288 x 6,144	136 x 4,096
10^{14}	20,480 x 10,240	144 x 8,192
10^{15}		152 x 16,384
10^{16}		156 x 24,576
10^{17}		160 x 32,768

Sparse Chebyshev solver
(Dedalus)

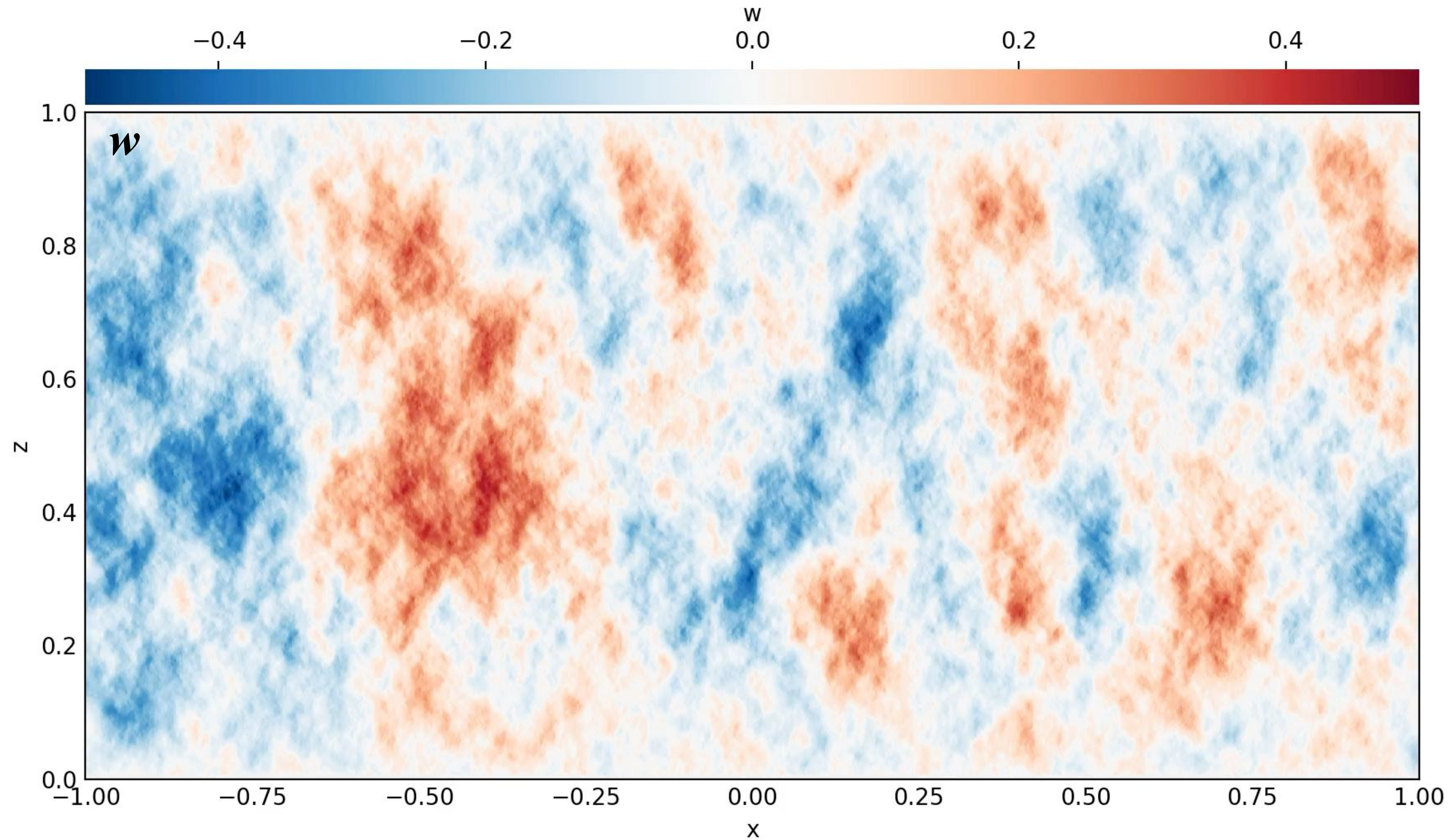


Burns et al., PRR (2020)

Log-Fourier Rayleigh-Benard

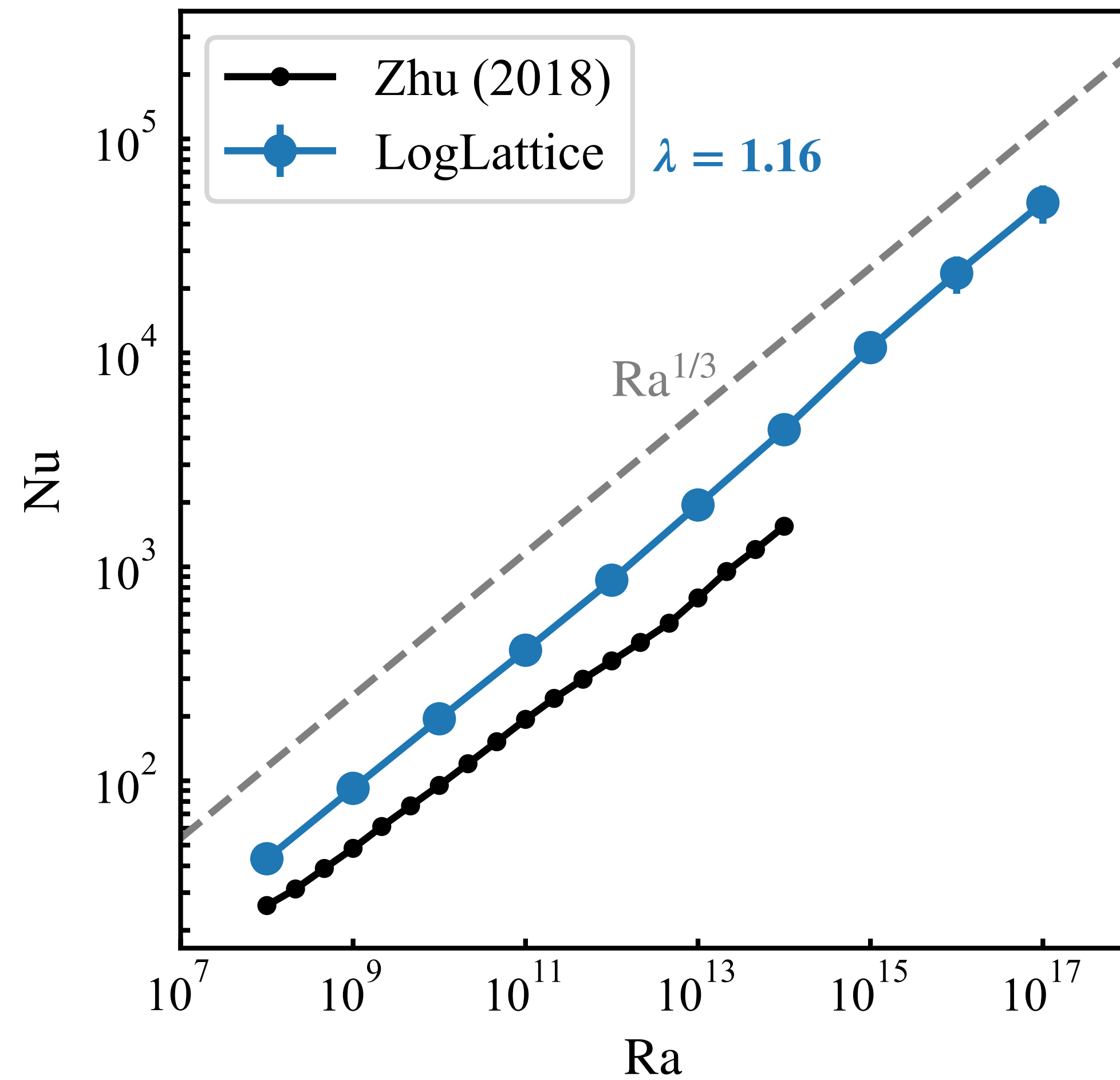
$$\text{Ra} = 10^{14}$$

$$\lambda = 1.16$$

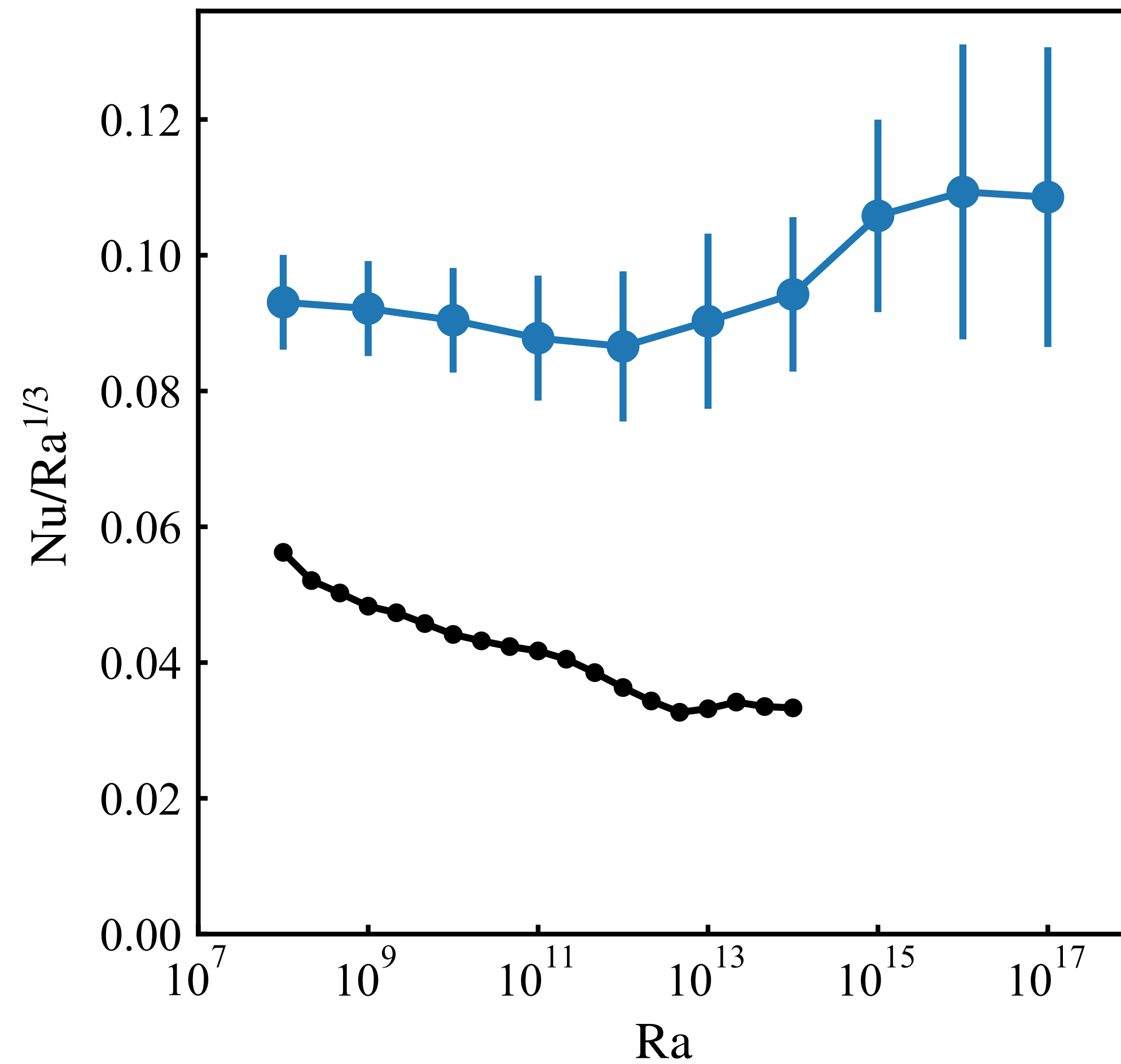


Log-Fourier Rayleigh-Benard

Nusselt scaling



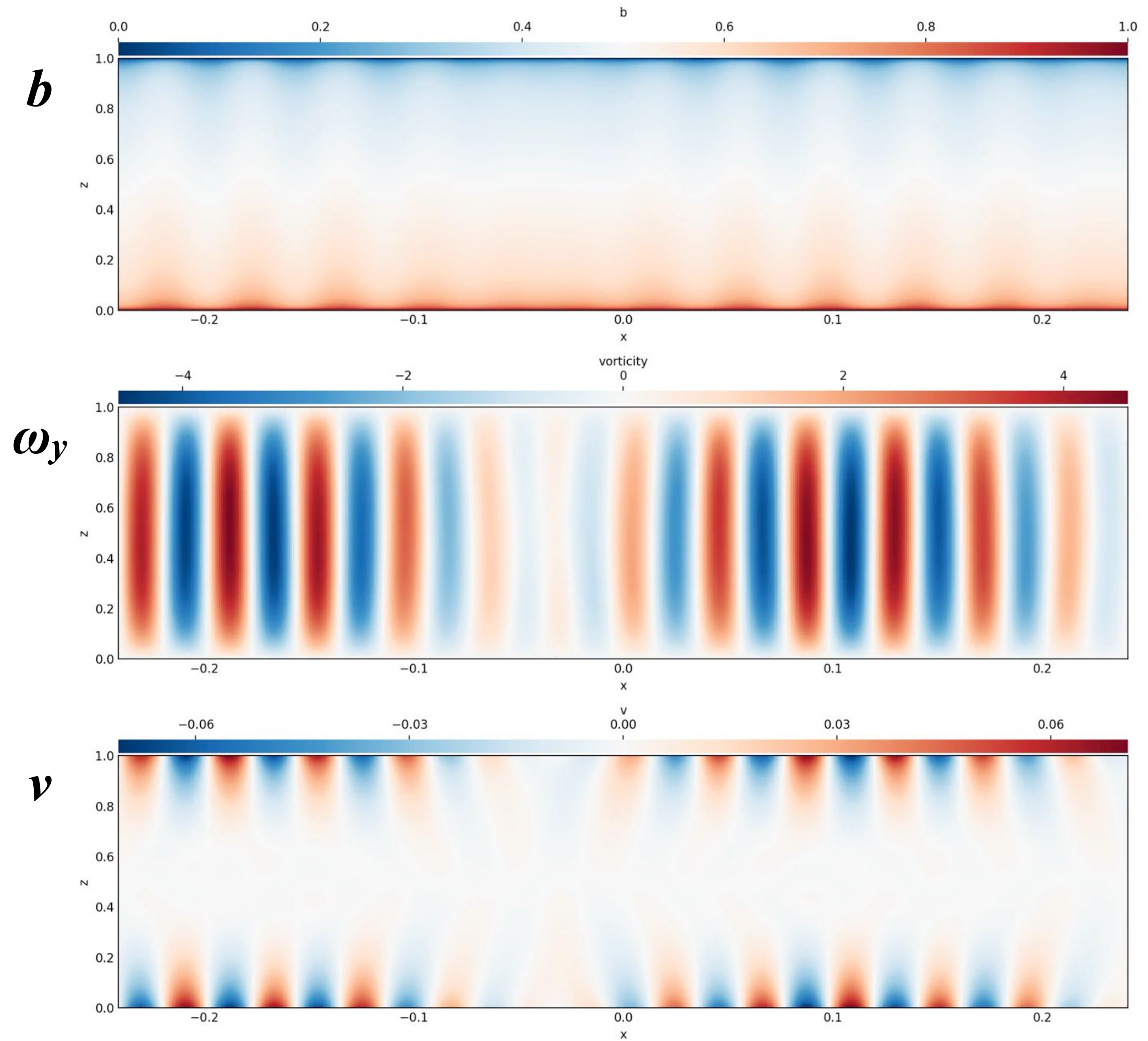
Nusselt scaling (compensated)



Log-Lattice Rotating Convection

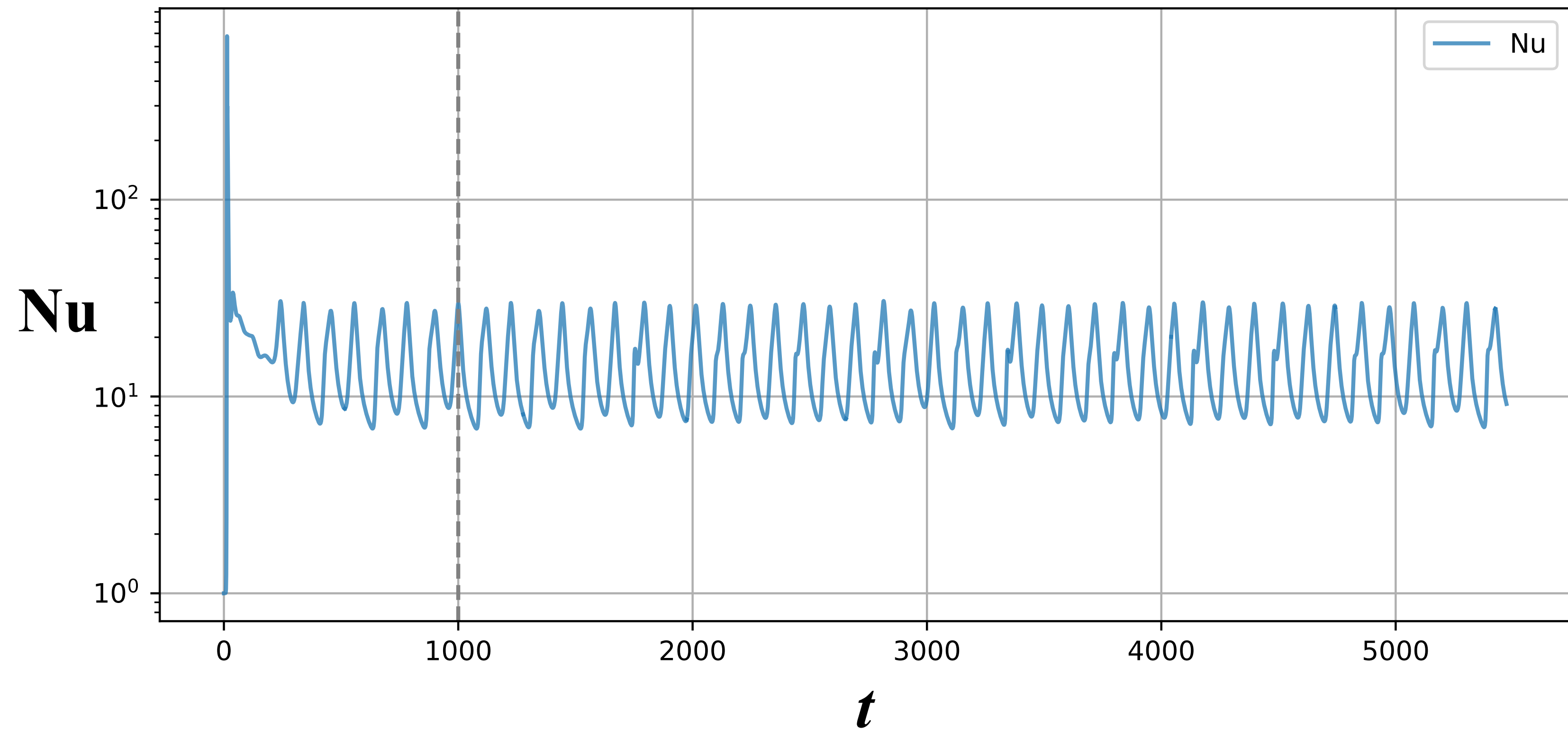
Log-Fourier Rotating Rayleigh-Benard

$$\begin{aligned}\widetilde{\text{Ra}} &= 40 \\ \text{Ek} &= 10^{-6} \\ L_x &= 10\ell_c \\ \lambda &= 1.16\end{aligned}$$



Nusselt number in Log-Lattice RRBC

$$\begin{aligned}\widetilde{\text{Ra}} &= 40 \\ \text{Ek} &= 10^{-6} \\ L_x &= 10\ell_c \\ \lambda &= 1.16\end{aligned}$$

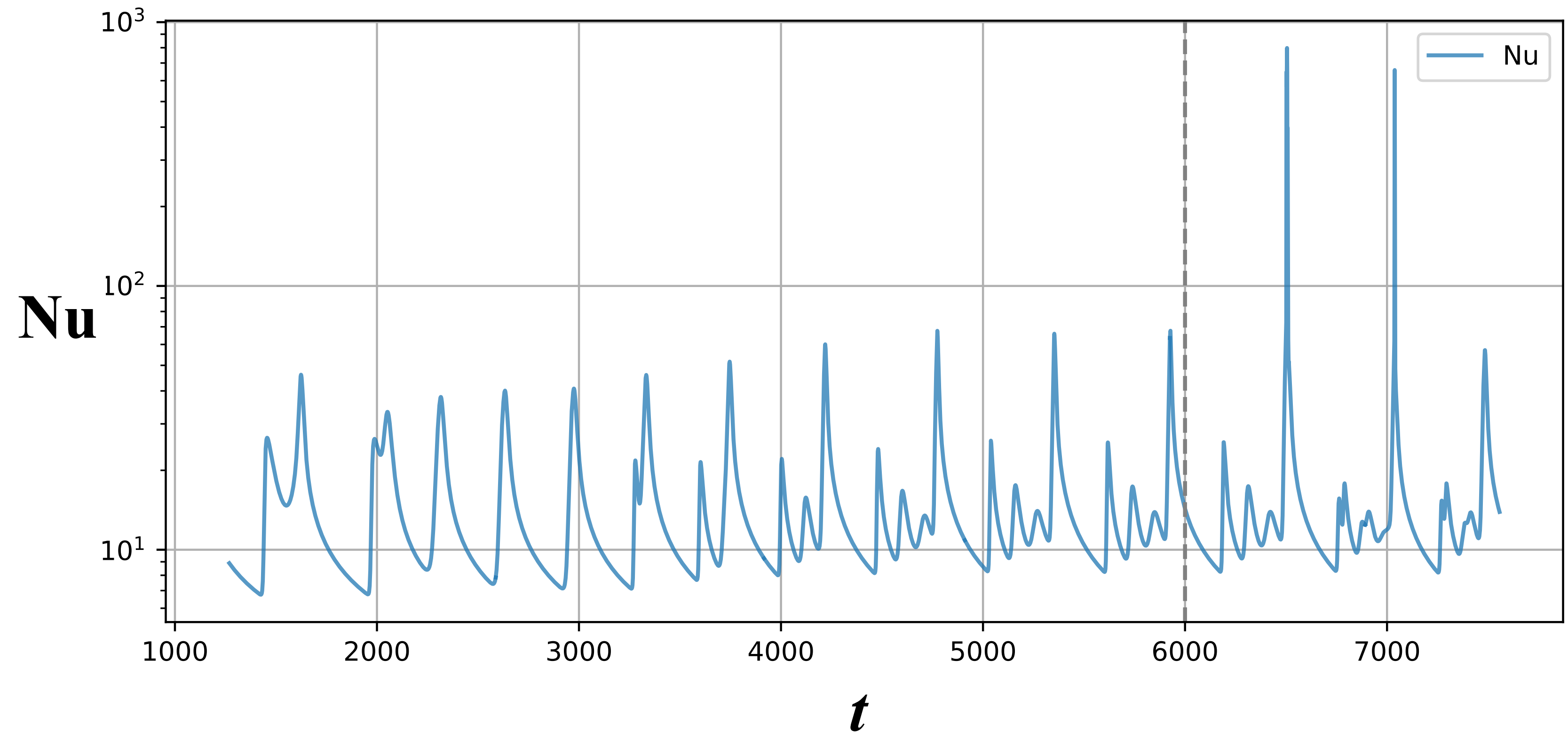


Measured: $\langle \text{Nu} \rangle_t = 16.0 \pm 6.2$

Reference: $\text{Nu}_{\text{AvK}} = 15.8 \pm 1.0$

Strong intermittency at lower Ekman

$$\begin{aligned}\widetilde{\text{Ra}} &= 40 \\ \text{Ek} &= 10^{-7} \\ L_x &= 10\ell_c \\ \lambda &= 1.16\end{aligned}$$



Measured: $\langle \text{Nu} \rangle_t = 13.1 \pm 8.3$

Reference: $\text{Nu}_{\text{AvK}} = 13.2 \pm 0.7$

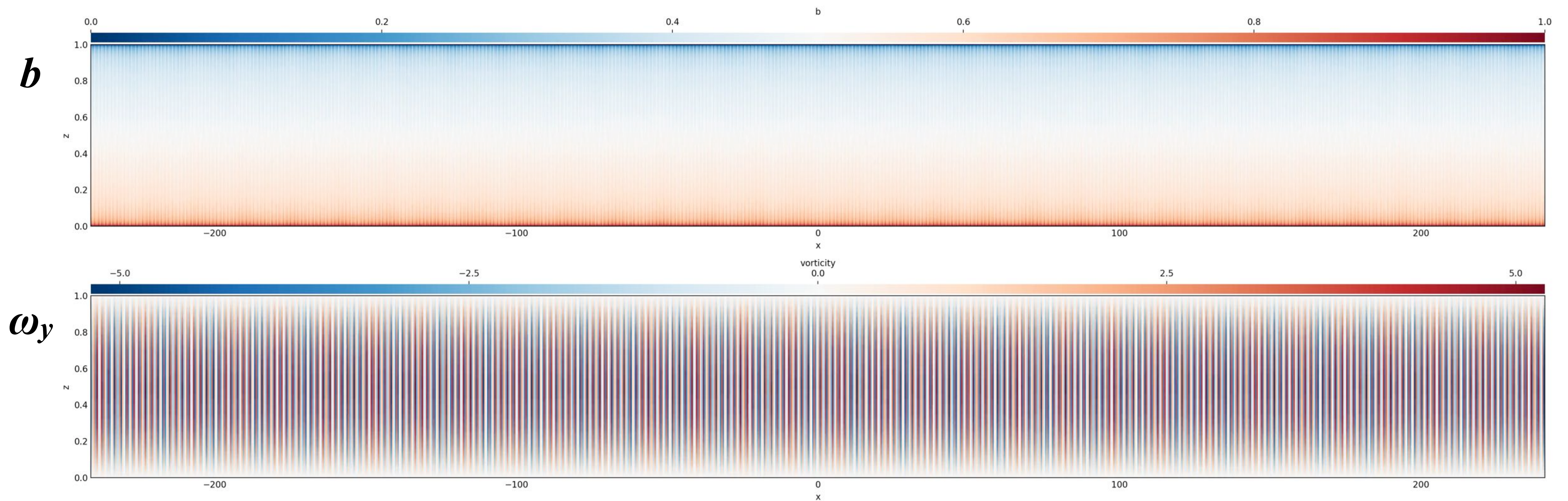
Extreme domain sizes

$$\widetilde{\text{Ra}} = 40$$

$$L_x = 10,000 \ell_c$$

$$\text{Ek} = 10^{-6}$$

$$\lambda = 1.16$$



$$\text{Measured: } \langle \text{Nu} \rangle_t = 14.9 \pm 6.7$$

Next steps

Lower Ekman

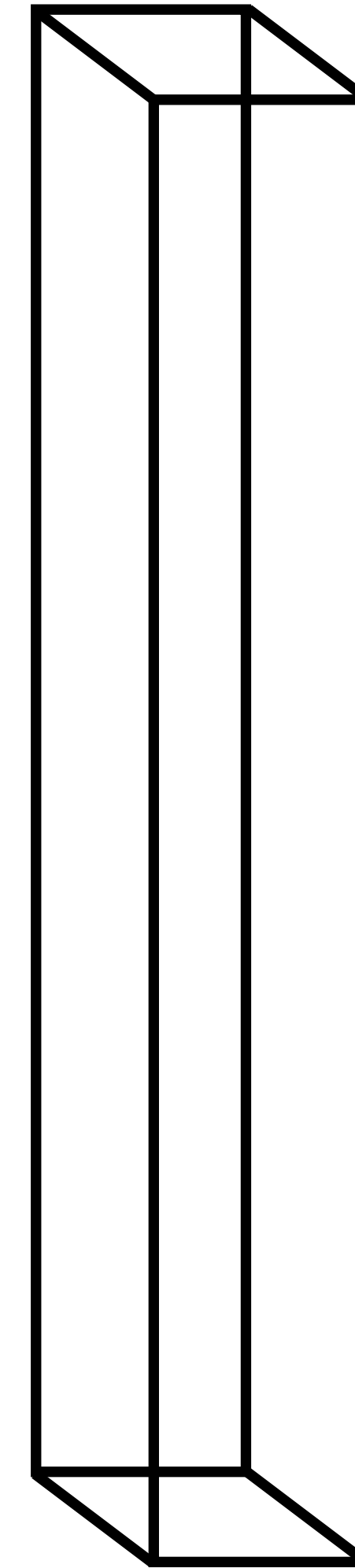
- Artifacts from 2.5D?
- Adaptive timestepping
- Examine large-scales in GT

Extend to 3D

- Even more relative speedup vs DNS
- Difficult parallelization

Optimize lattice

- Vary λ , a , b
- Different lattices
- Flow dependent



$\text{Re}^{3/2}$ vs Re^3

Summary

- **Logarithmic Fourier lattices**
 - Alternative to DNS for highly multiscale problems
 - Promising results in homogeneous flows
 - Can mix with standard bases for bounded domains
- **Results for RBC**
 - Can scale to $Ra = 10^{17}$ w/ sparse Chebyshev
 - General agreement with Malkus scaling
- **Results for RRBC**
 - Highly intermittent flows in 2.5D
 - Average statistics match DNS at modest Ek
 - Scales to large domains $L_x \gg \ell_c$
- **Ongoing work:**
 - Optimize λ , triad couplings, etc.
 - Extend to 3D, modifications for coherent structures

