

Scaling laws and force balances in rotating convection: are they consistent?

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Keith Julien, Stefano Maffei



University of Colorado **Boulder**

PHYSICAL REVIEW FLUIDS 8, 093502 (2023)

Small scale quasigeostrophic convective turbulence at large Rayleigh number

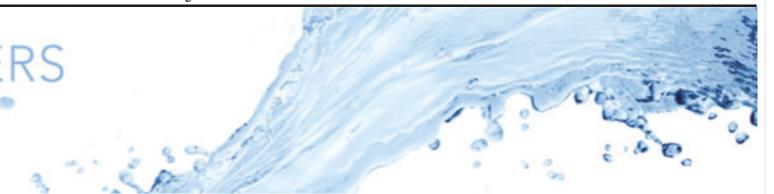
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J. Fluid Mech. (2024), vol. 981, A22, doi:10.1017/jfm.2024.78

JM PAPERS



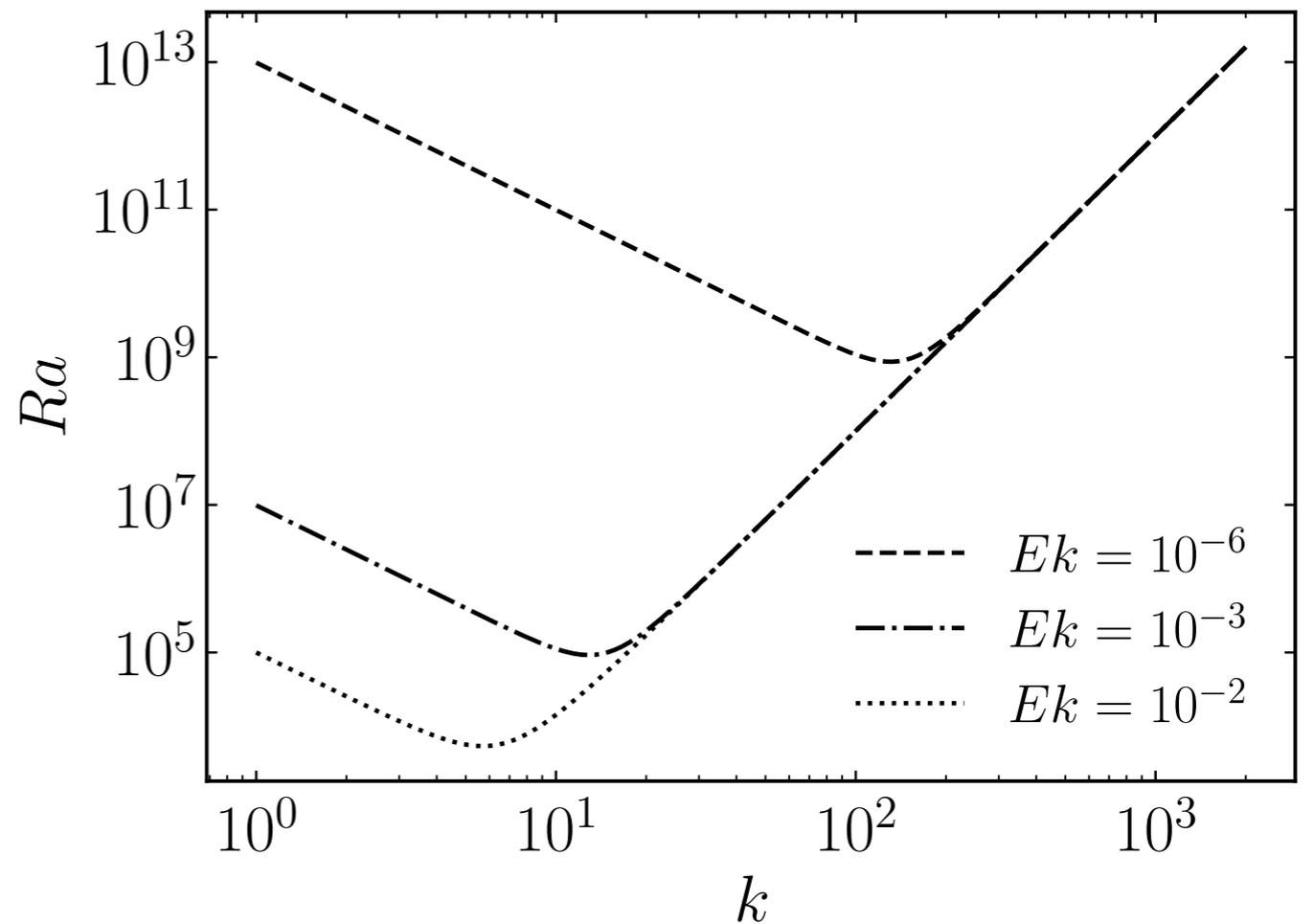
Asymptotic scaling relations for rotating spherical convection with strong zonal flows

Justin A. Nicoski¹, Anne R. O'Connor¹ and Michael A. Calkins^{1,†}

Linear Theory

- Mapped out in Chandrasekhar (1961):

$$Ra_m = \frac{(\pi^2 + k^2)^3 + \pi^2 Ek^{-2}}{k^2}$$

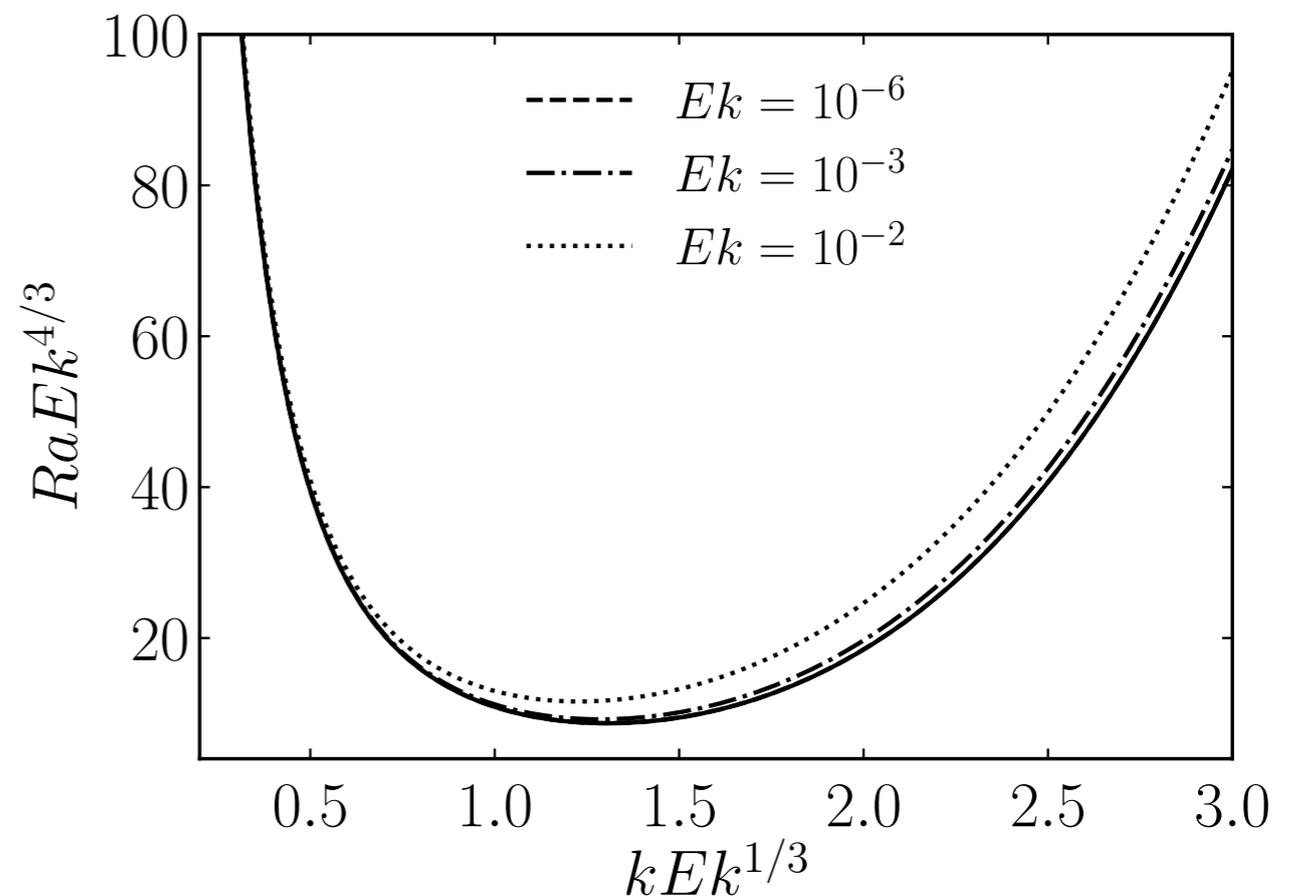


Linear Theory

- Most unstable mode is viscous length scale:

$$k = O(Ek^{-1/3})$$

$$\widetilde{Ra} = RaEk^{4/3}$$



$$\Rightarrow \widetilde{Ra}_m = \widetilde{k}^4 + \pi^2 \widetilde{k}^{-2}$$

All wavenumbers scale viscously

Nonlinear (QG) Theory

Julien, Knobloch, Werne & collaborators (1998, etc.)

Geostrophy at leading order:

$$\hat{\mathbf{z}} \times \mathbf{u} \approx -\nabla p$$

Dynamics at next order:

$$D_t \mathbf{u} + \dots \approx \dots \frac{\widetilde{Ra}}{Pr} \theta + \nabla_{\perp}^2 \mathbf{u}$$

$$\Rightarrow \mathbf{u} = O(Ek^{-1/3}) \quad \theta = O(Ek^{1/3}) \quad \ell = O(Ek^{1/3})$$

Scalings persist in turbulent regime, independent of geometry

Scaling with Rayleigh: CIA Balance

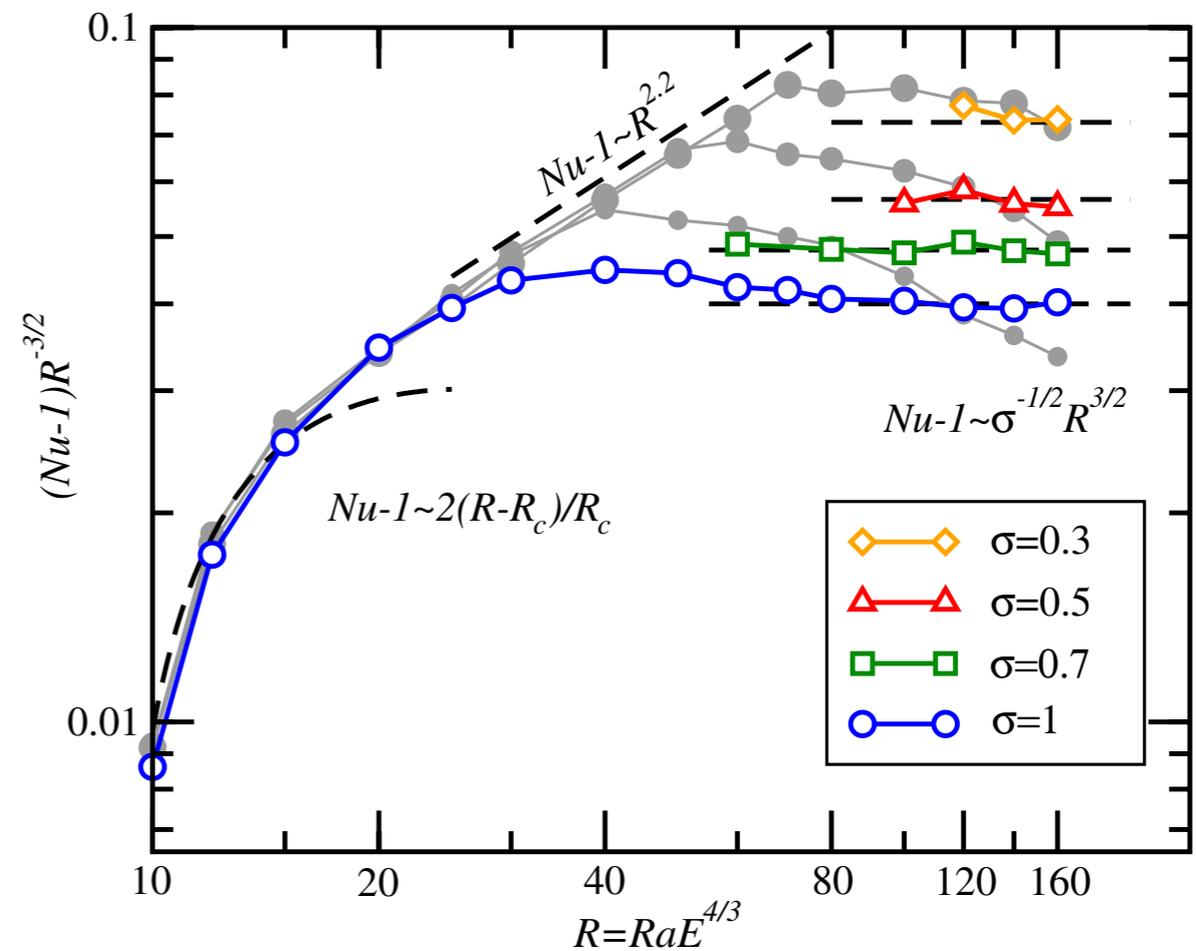
We really want a dependence on Rayleigh number:

$$\text{CIA: } u \sim \frac{\widetilde{Ra}}{Pr} Ek^{-1/3} \quad \ell \sim \sqrt{\frac{\widetilde{Ra}}{Pr} Ek^{1/3}}$$
$$\Rightarrow \widetilde{Re} \sim \frac{\widetilde{Ra}}{Pr}$$


Aurnou, Horn & Julien (PRR, 2020)
Oliver et al. (PRF, 2023)
Nicoski, O'Connor & Calkins (JFM, 2024)

QG Simulations

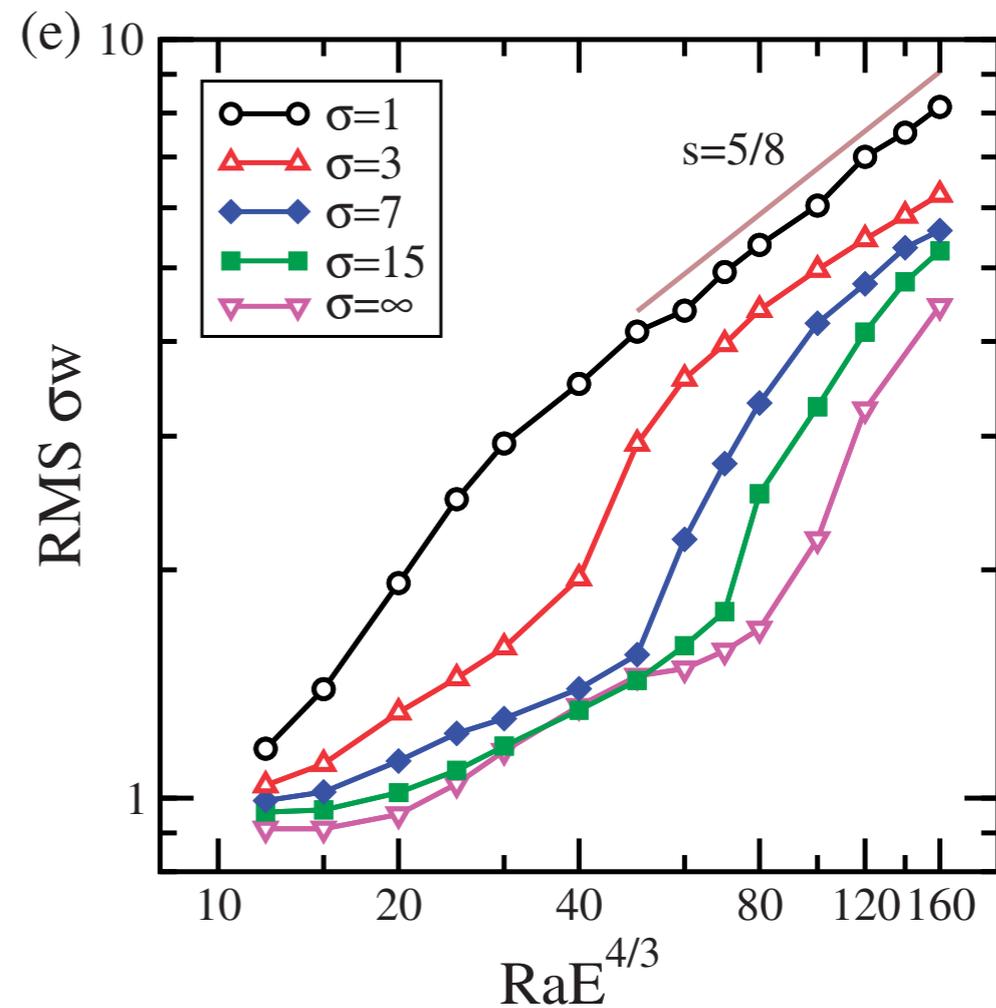
- Diffusion free scaling observed for heat transport



Julien et al. (PRL, 2012)

QG Simulations

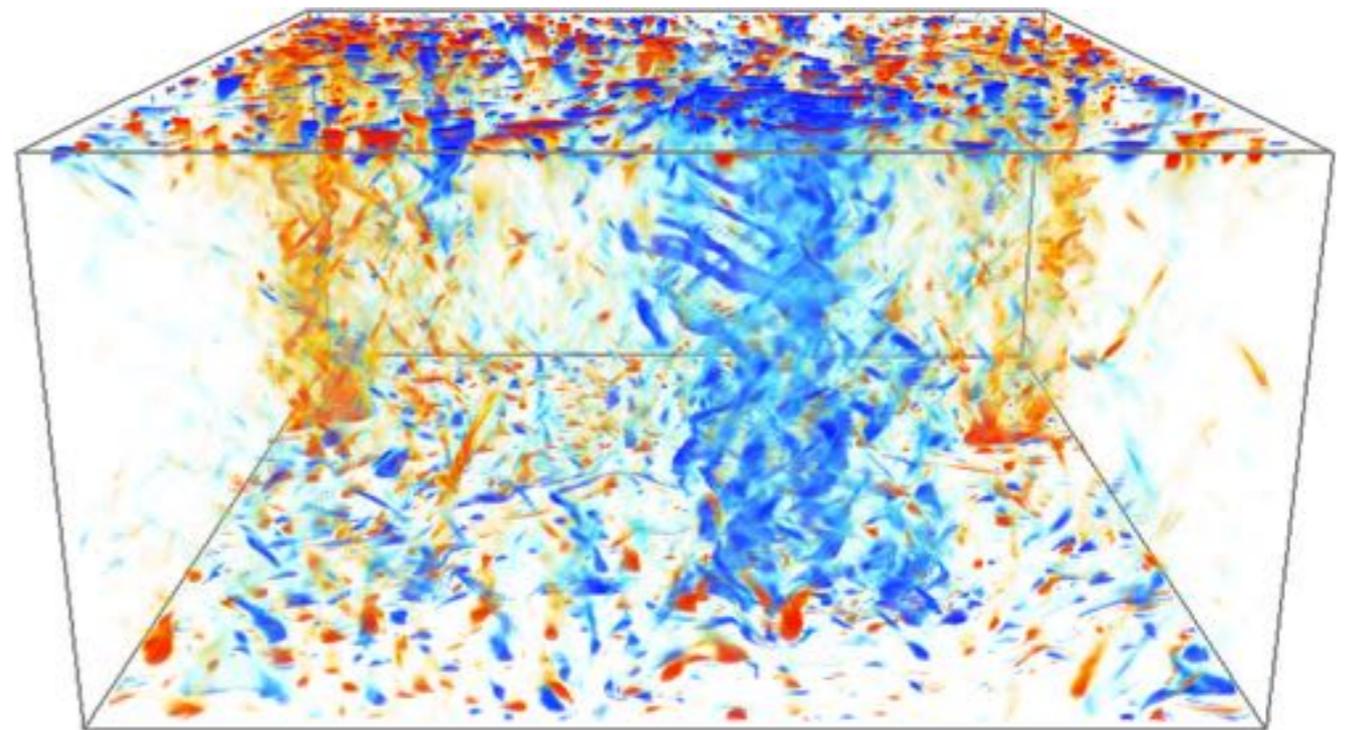
- Somewhat steeper scale observed for flow speeds.
- Prandtl number influences the observed scaling.



Julien et al. (GAFD, 2012)

QG Simulations

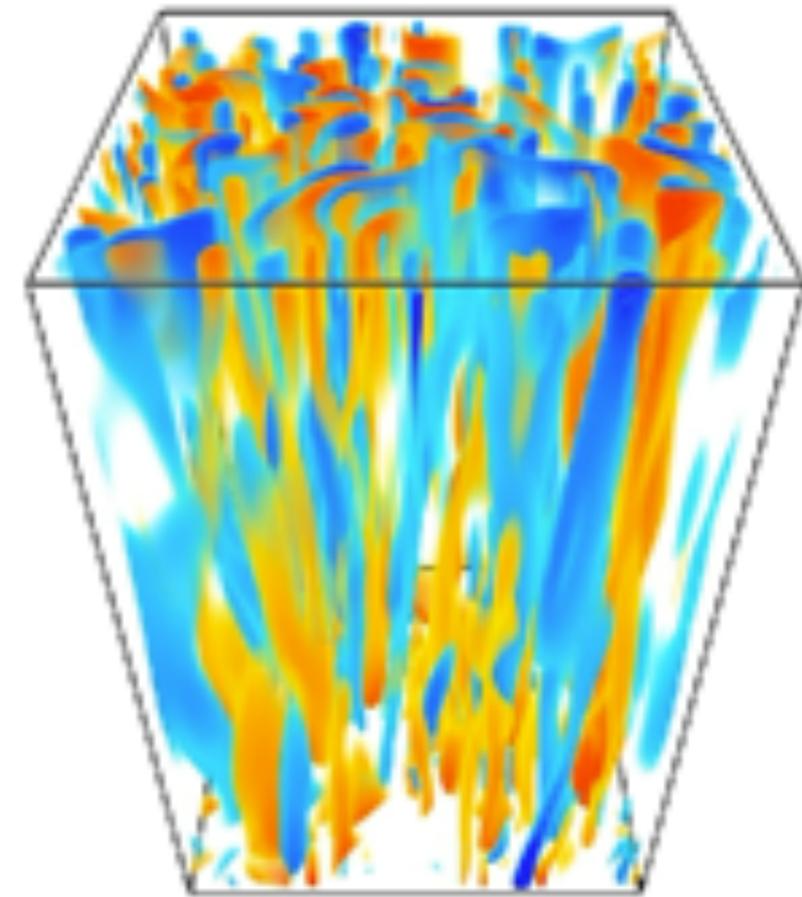
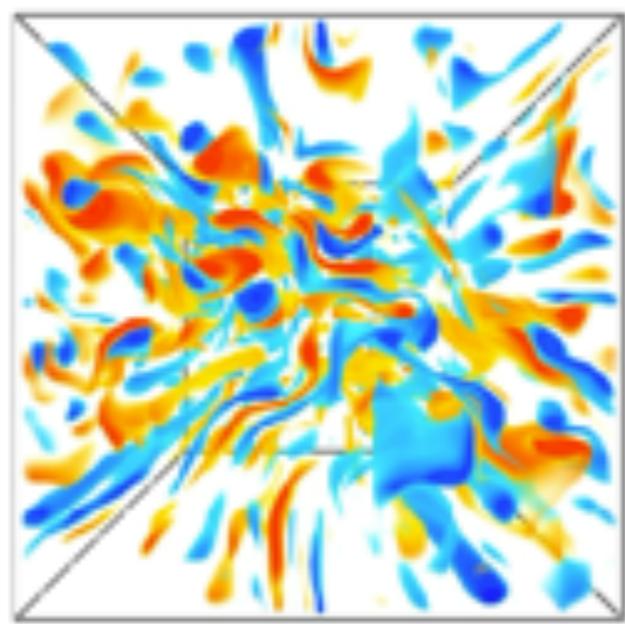
- Inverse cascade leads to large scale vortex (LSV)
- How does LSV influence scaling?



Rubio et al. (PRL, 2014)

QG Simulations

- Inverse cascade leads to large scale vortex (LSV)
- How does LSV influence scaling?

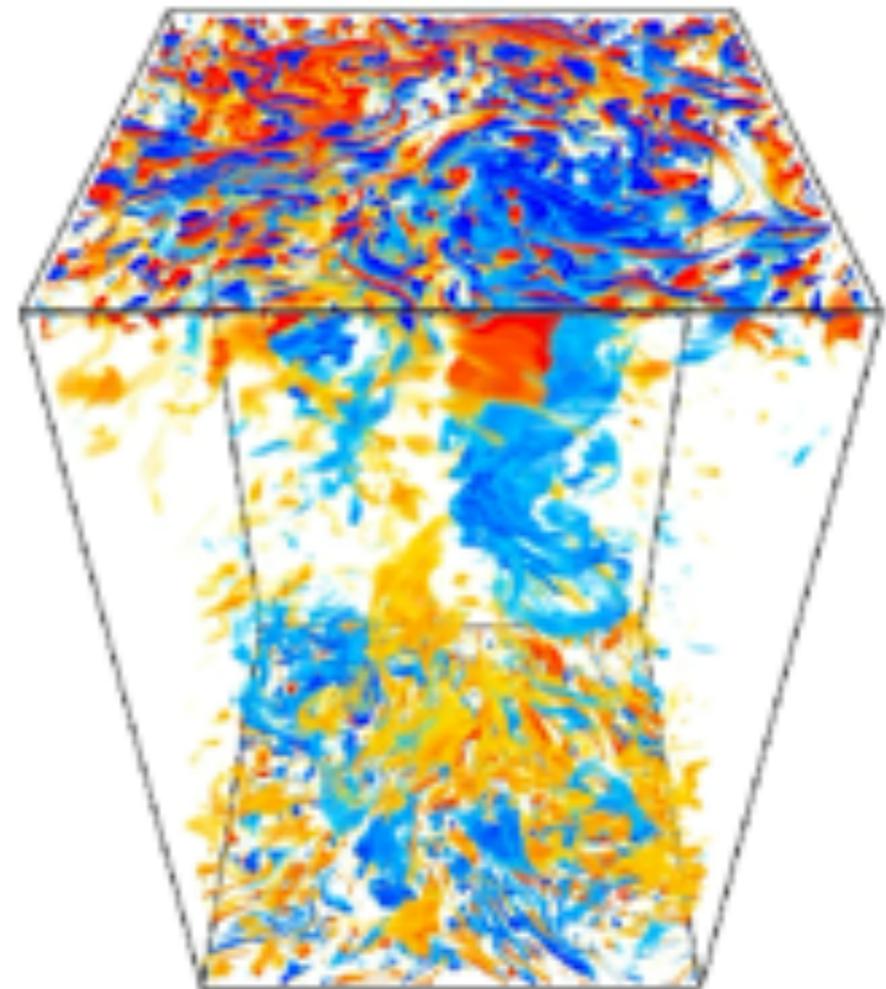


$$\widetilde{Ra} = 40$$

Maffei et al. (JFM, 2021)

QG Simulations

- Inverse cascade leads to large scale vortex (LSV)
- How does LSV influence scaling?

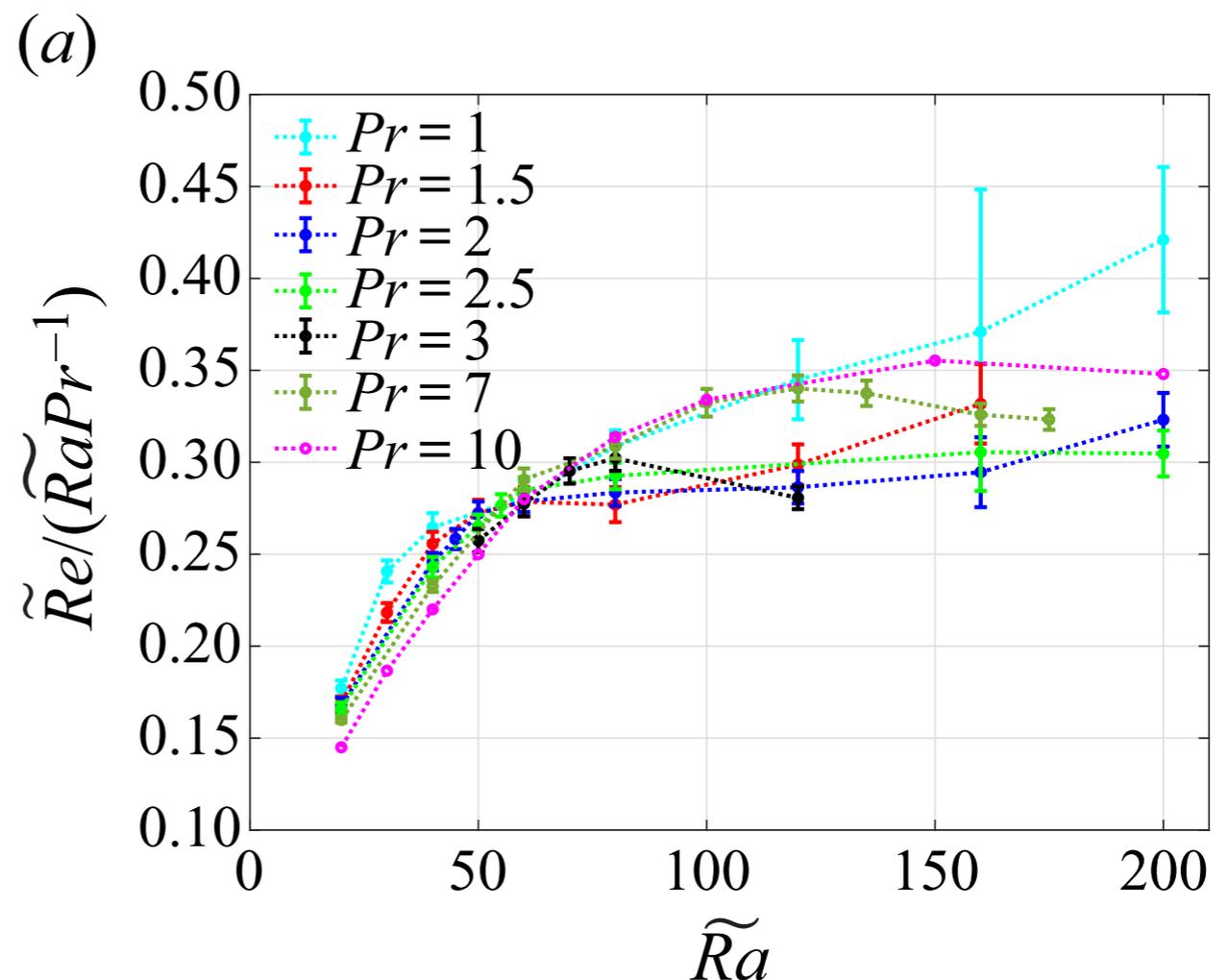


$$\widetilde{Ra} = 200$$

Maffei et al. (JFM, 2021)

QGG Simulations

- Inverse cascade leads to large scale vortex (LSV)
- How does LSV influence scaling?



Maffei et al. (JFM, 2021)

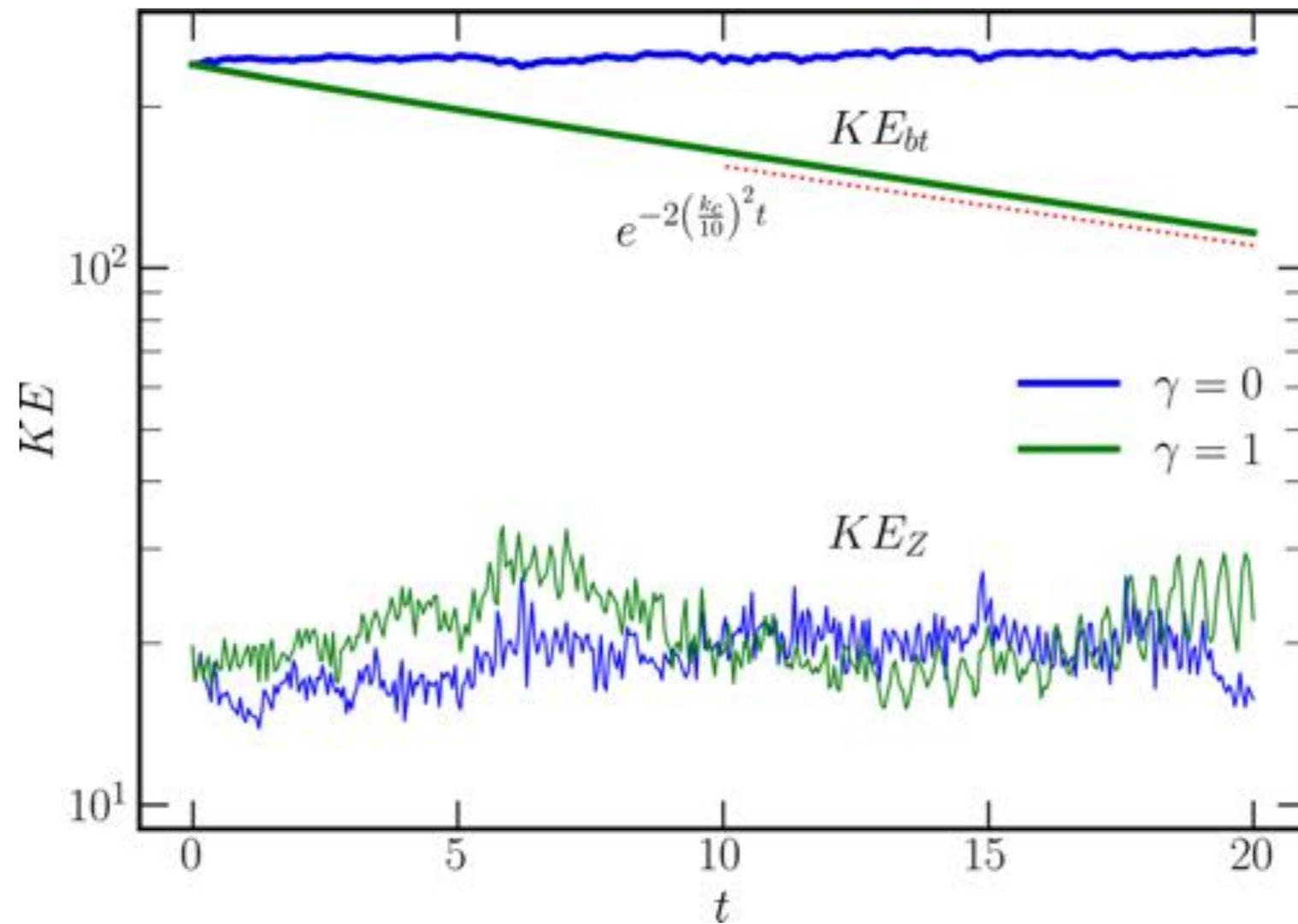
QG Simulations: No LSV

- We can “remove” the influence of the LSV:

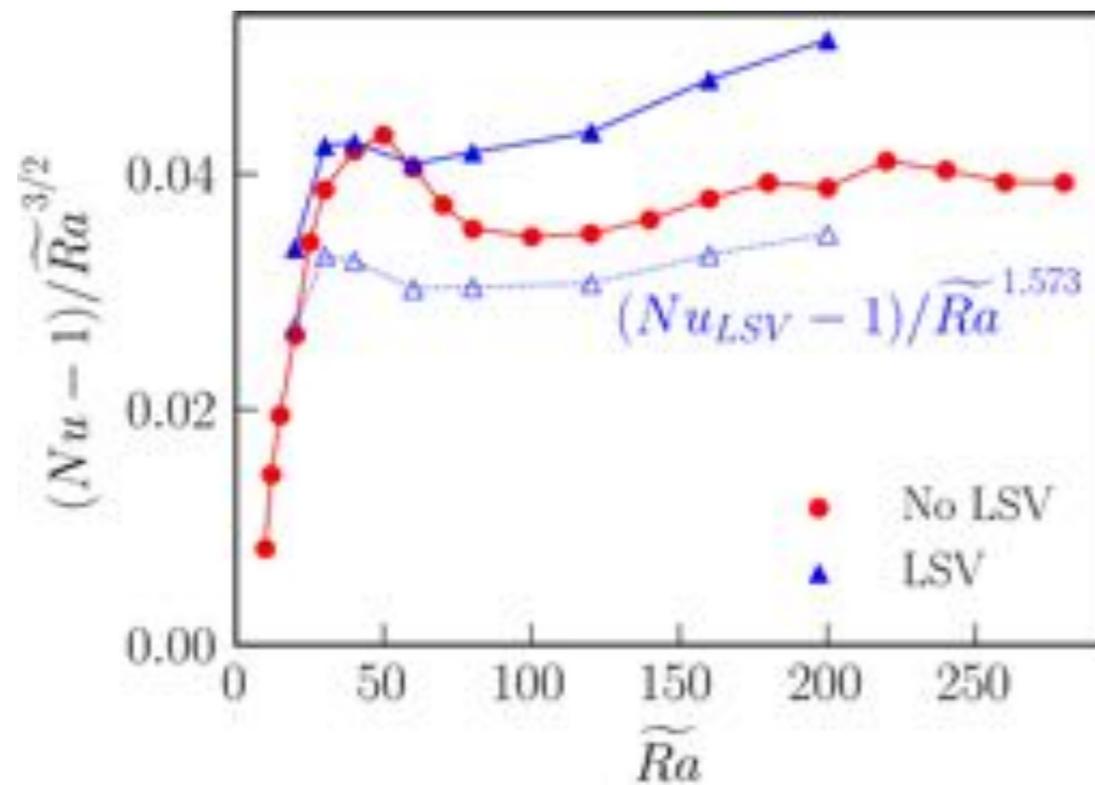
$$\partial_t \zeta + J[\psi, \zeta] - \gamma \langle J[\psi, \zeta] \rangle - \partial_z w = \nabla_{\perp}^2 \zeta,$$

$$\gamma = 1 \Rightarrow \partial_t \langle \zeta \rangle = \nabla_{\perp}^2 \langle \zeta \rangle$$

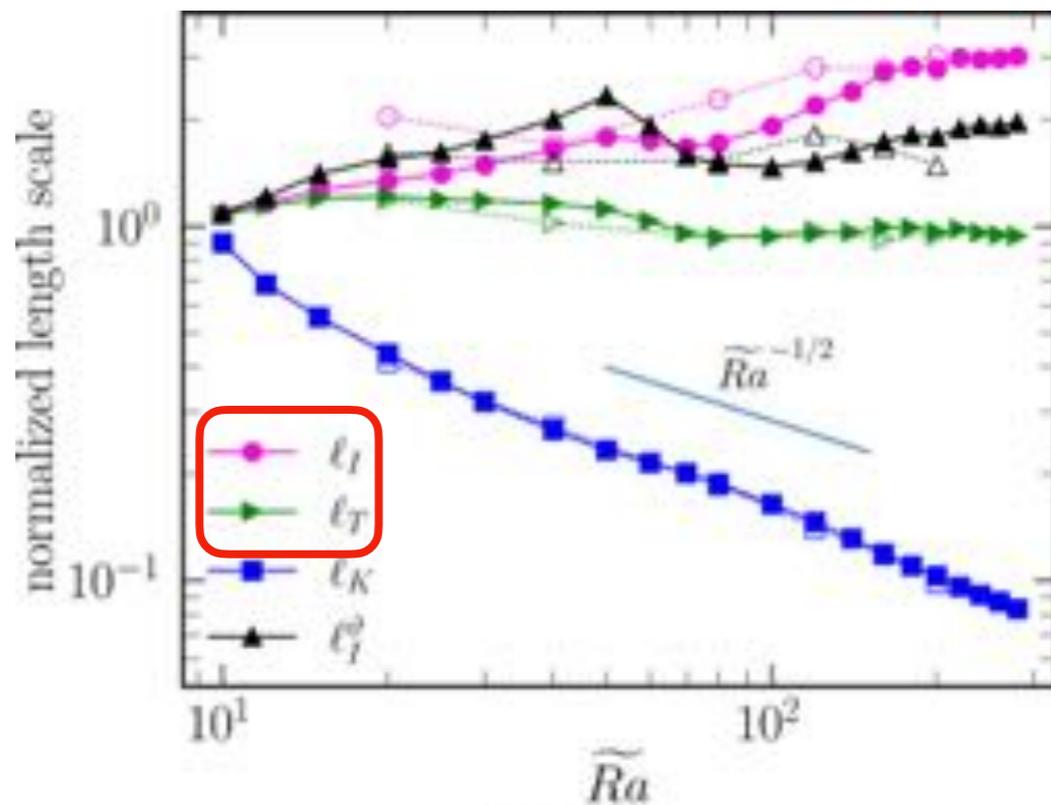
QG Simulations: No LSV



Global Transport



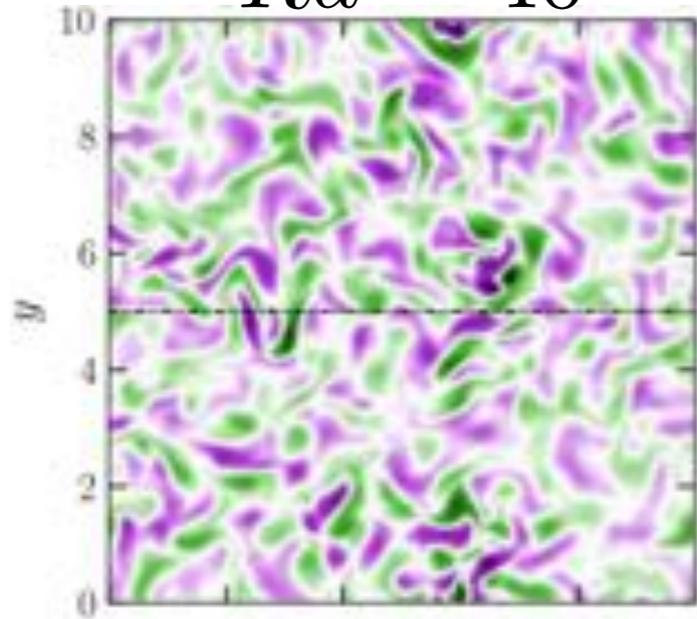
Length Scales



Length scales remain within the viscous range

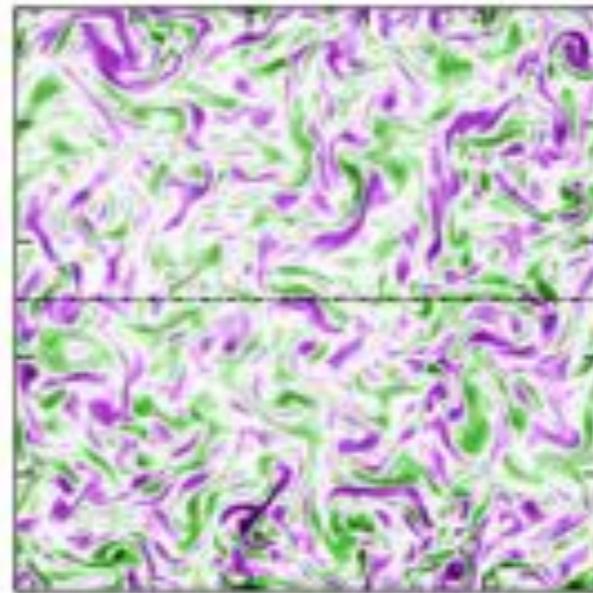
Flow Structure: vorticity

$\widetilde{Ra} = 40$



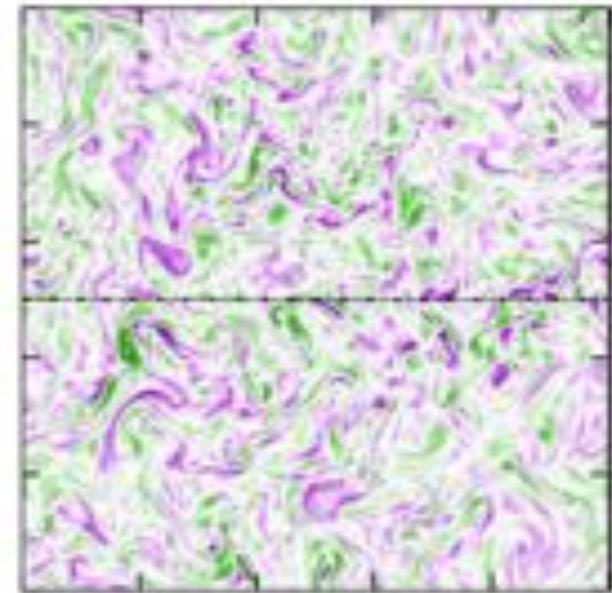
(a)

$\widetilde{Ra} = 160$



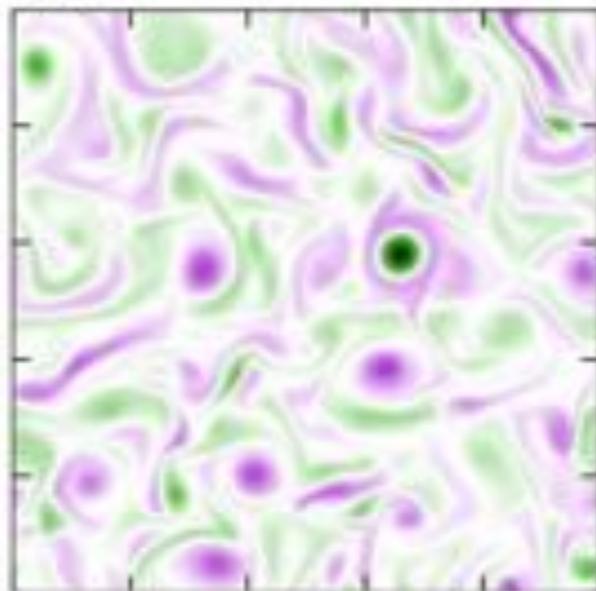
(b)

$\widetilde{Ra} = 280$

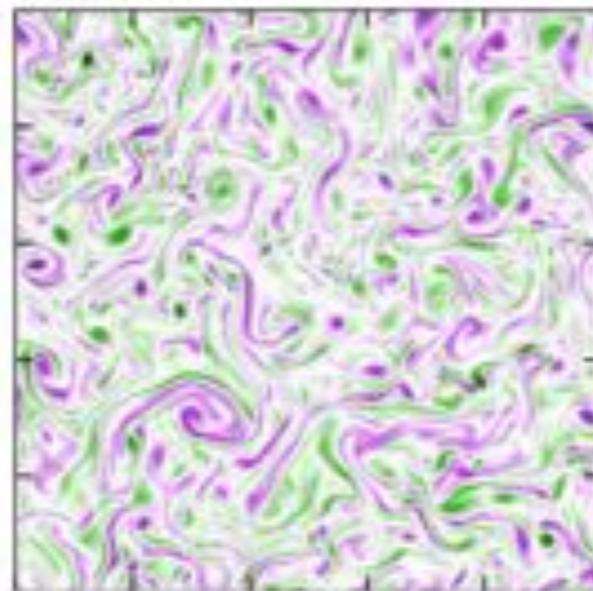


(c)

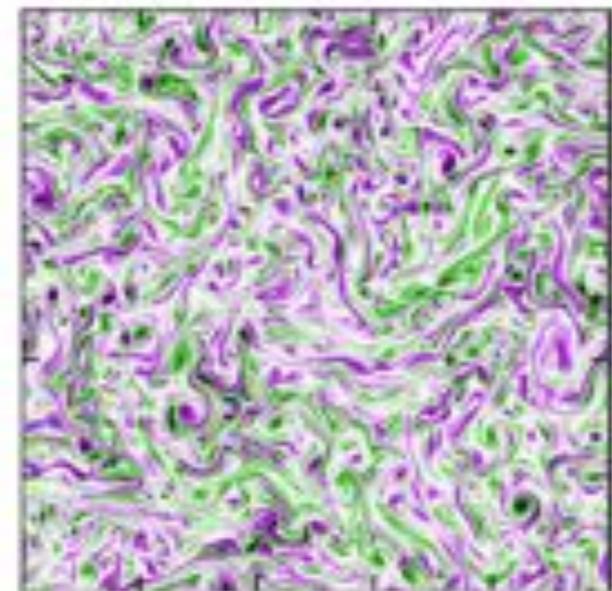
ℓ_I 
 ℓ_T 
 $10\ell_K$ 



ℓ_I 
 ℓ_T 
 $10\ell_K$ 

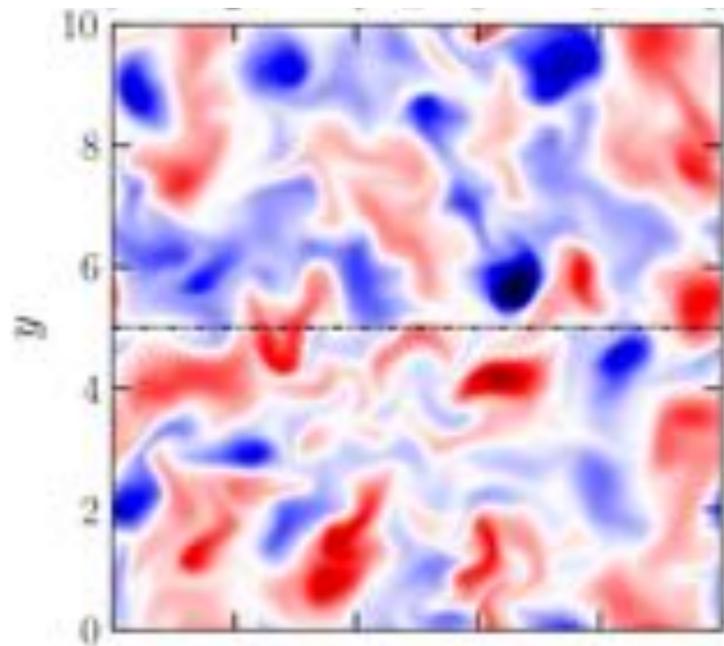


ℓ_I 
 ℓ_T 
 $10\ell_K$ 



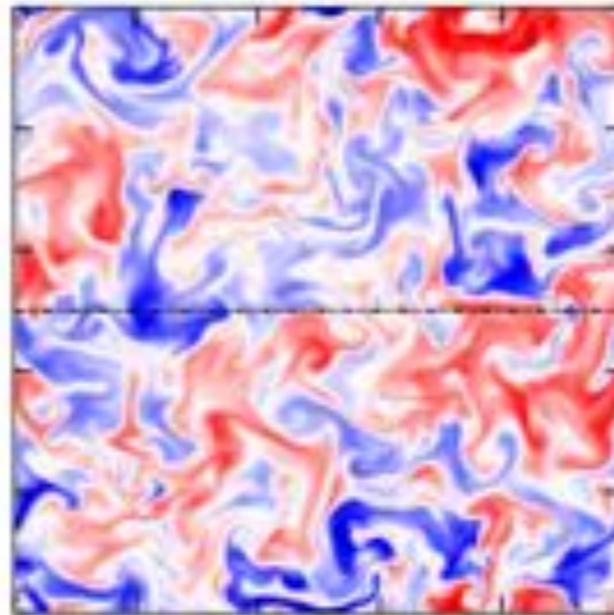
Flow Structure: temperature

$$\widetilde{Ra} = 40$$



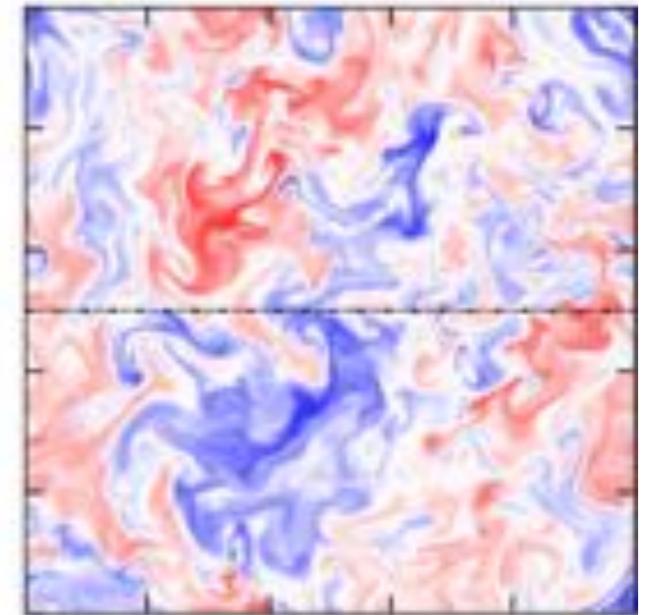
(a)

$$\widetilde{Ra} = 160$$



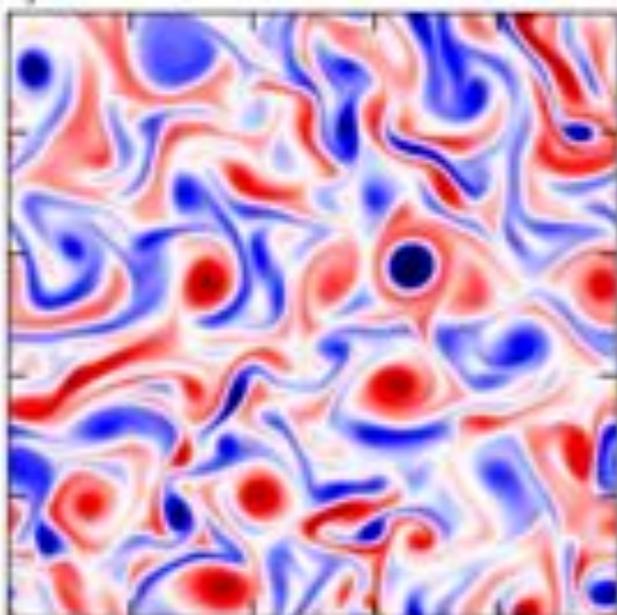
(b)

$$\widetilde{Ra} = 280$$

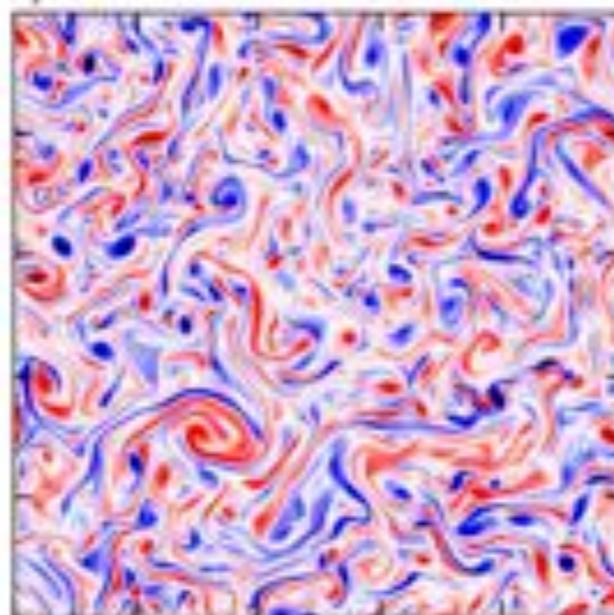


(c)

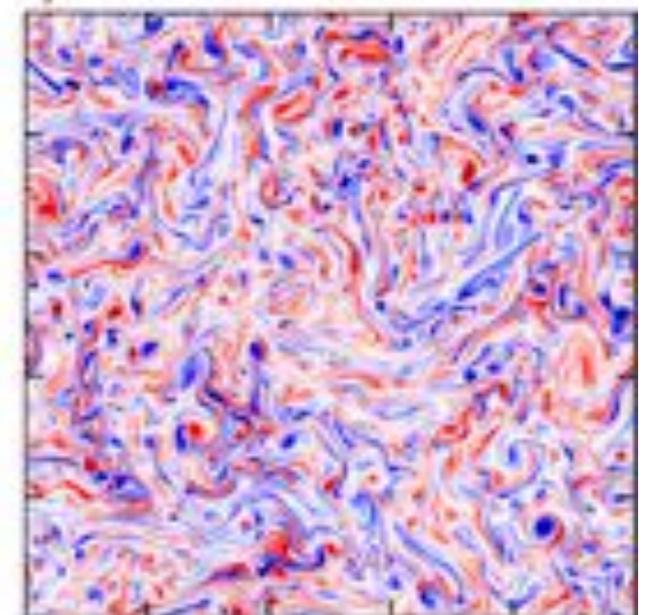
ϵ_T^* ———



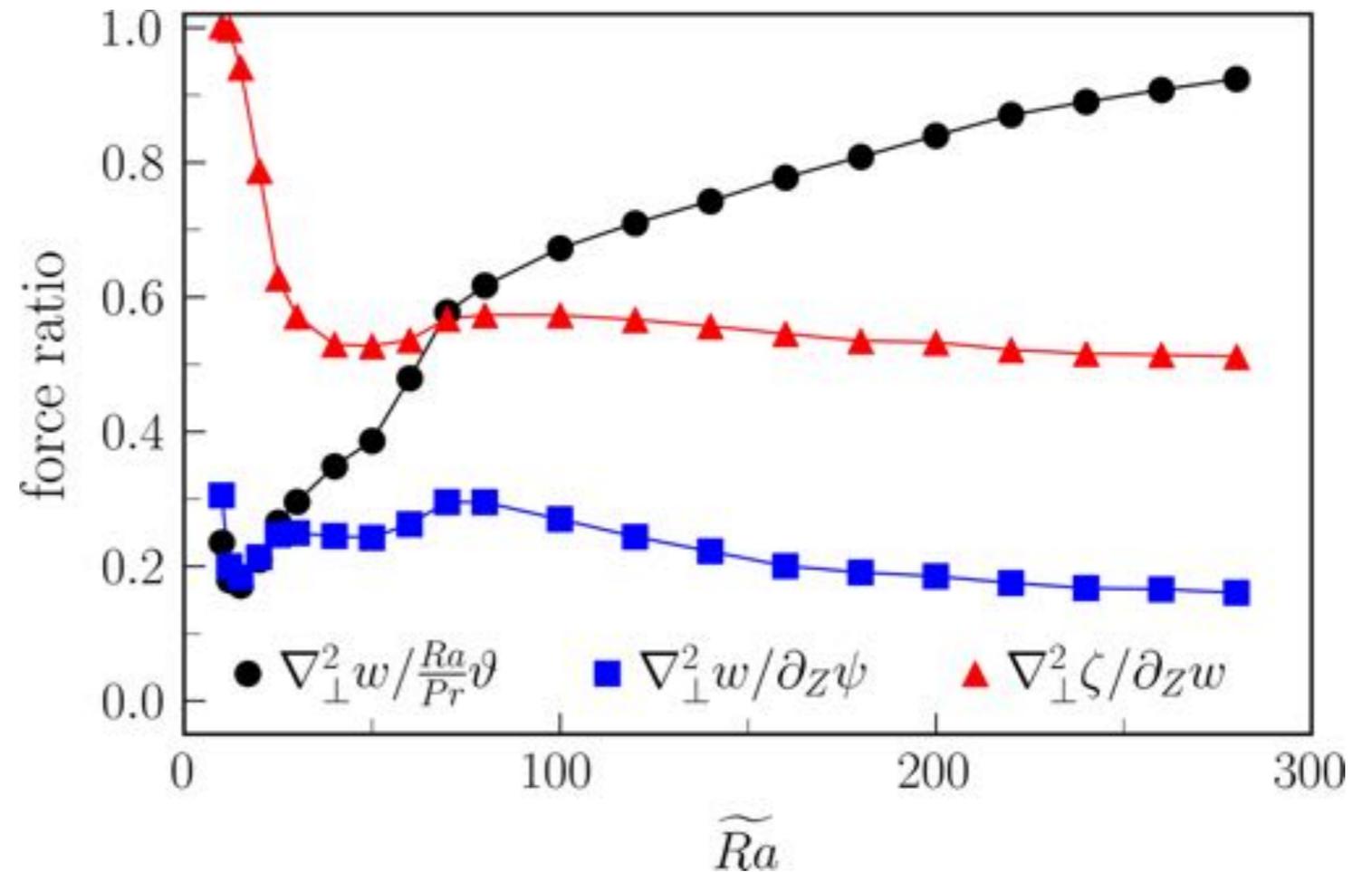
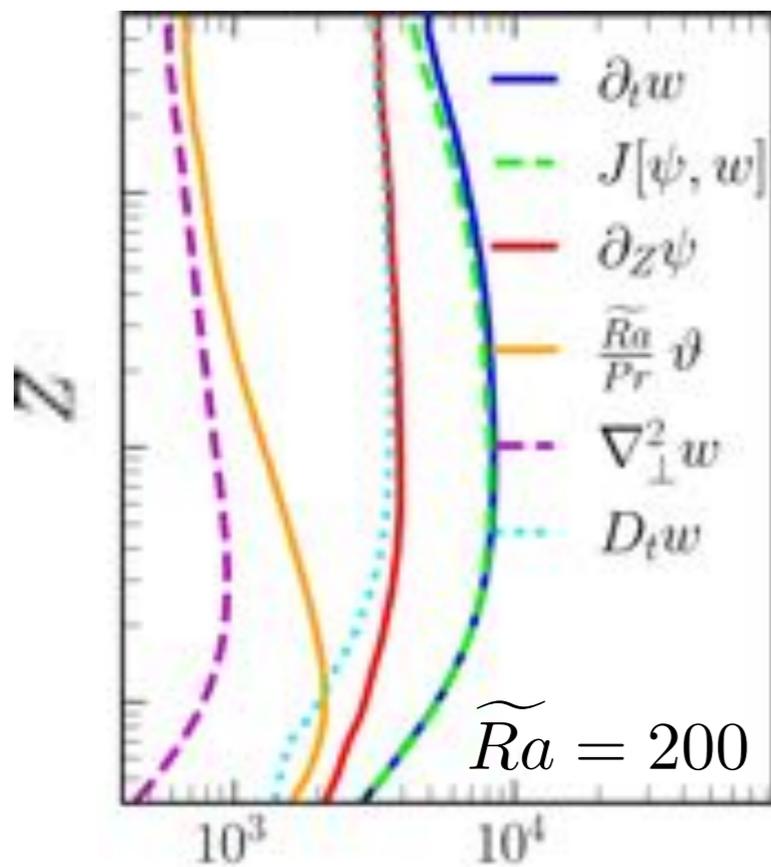
ϵ_T^* ———



ϵ_T^* ———



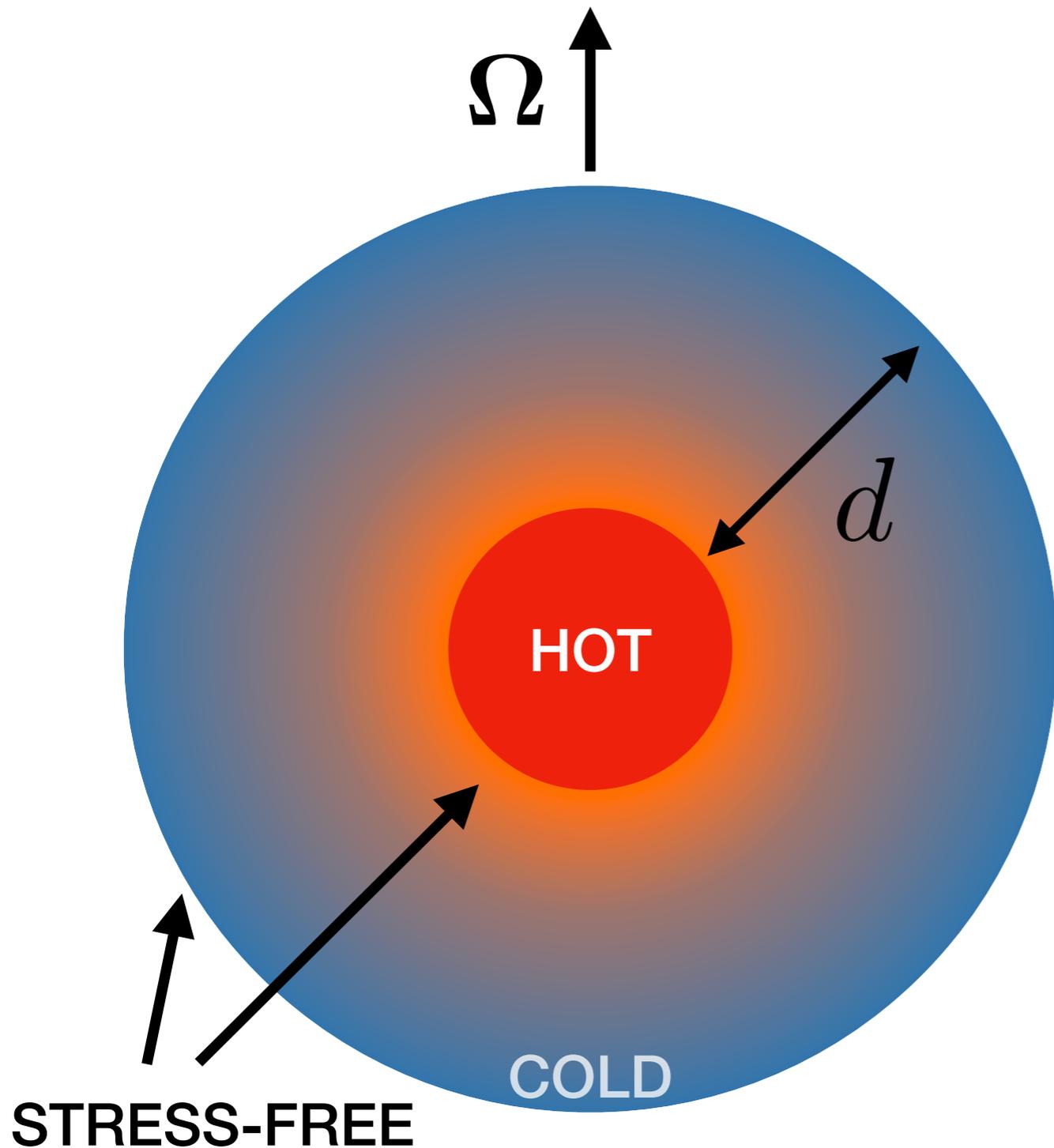
QGG Balances



Balances & length scales are not consistent with CIA

Similar in spherical geometries?

Spherical Convection



Ekman number: Ek

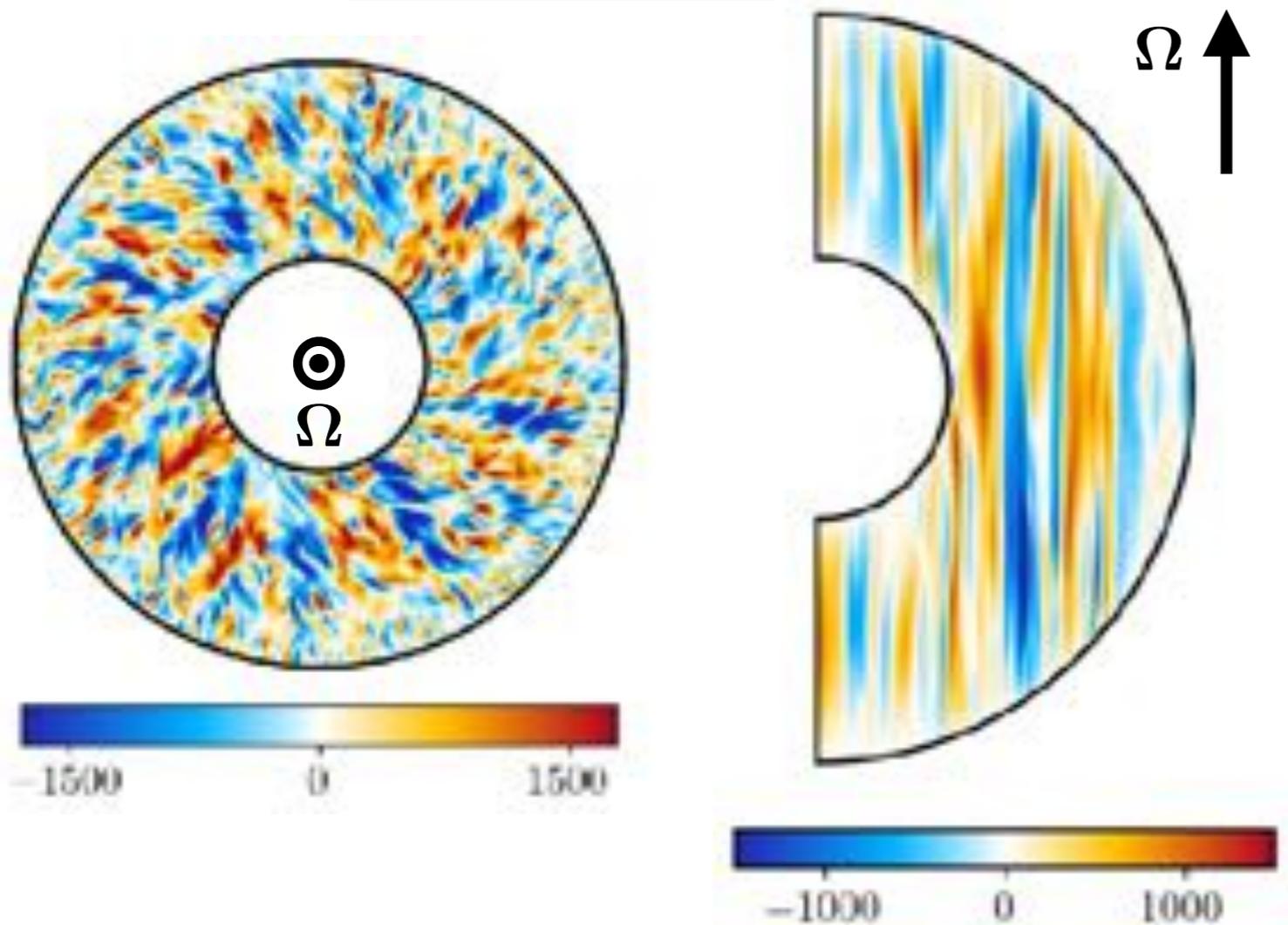
Reduced Rayleigh number:

$$\widetilde{Ra} = RaEk^{4/3}$$

DNS: $10^{-6} \ll Ek \ll 10^{-3}$

Flow Decomposition

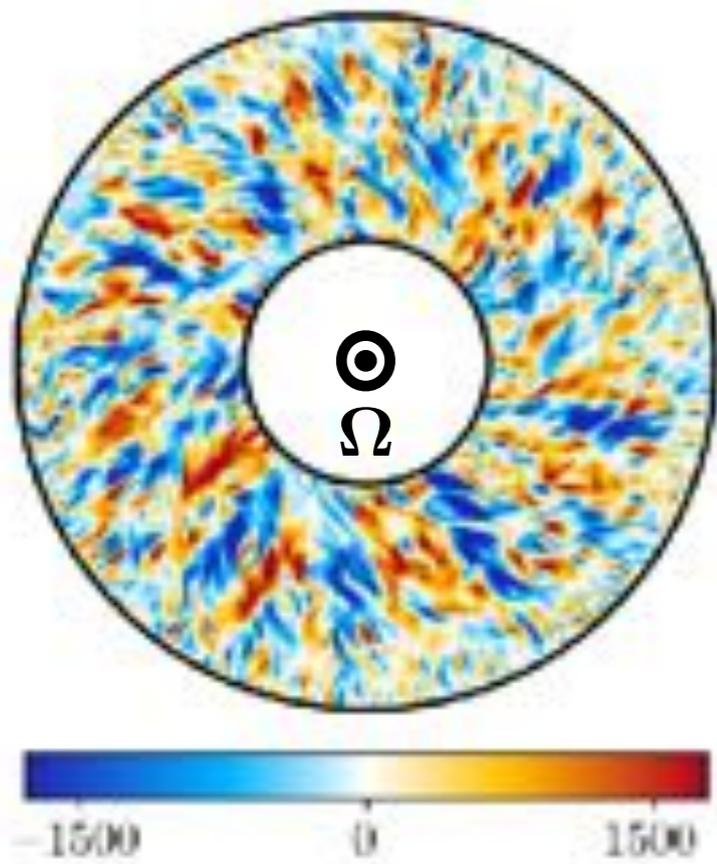
Convection



$$\mathbf{u}(r, \theta, \phi, t) = \bar{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Flow Decomposition

Convection

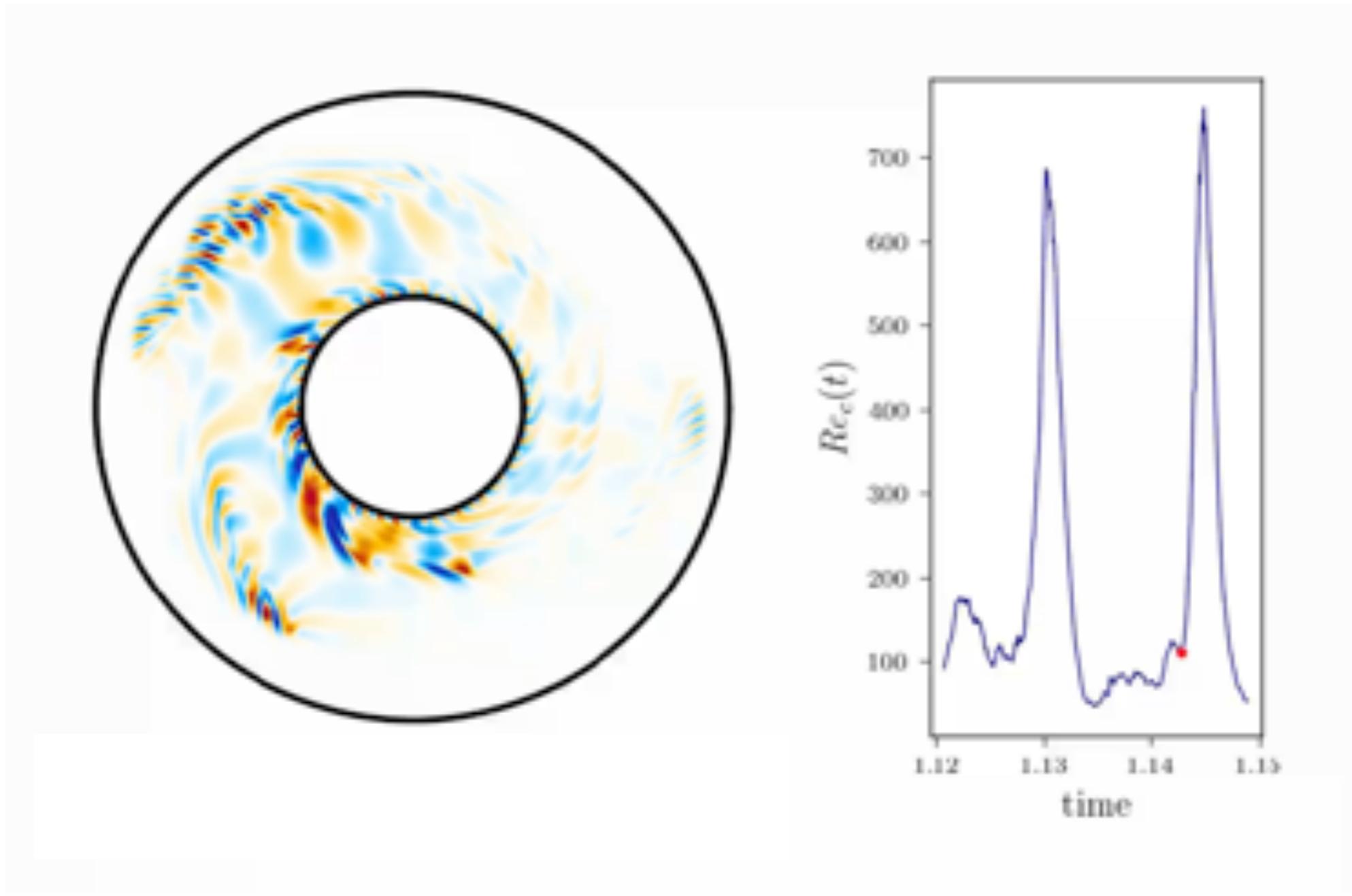


Zonal Flow



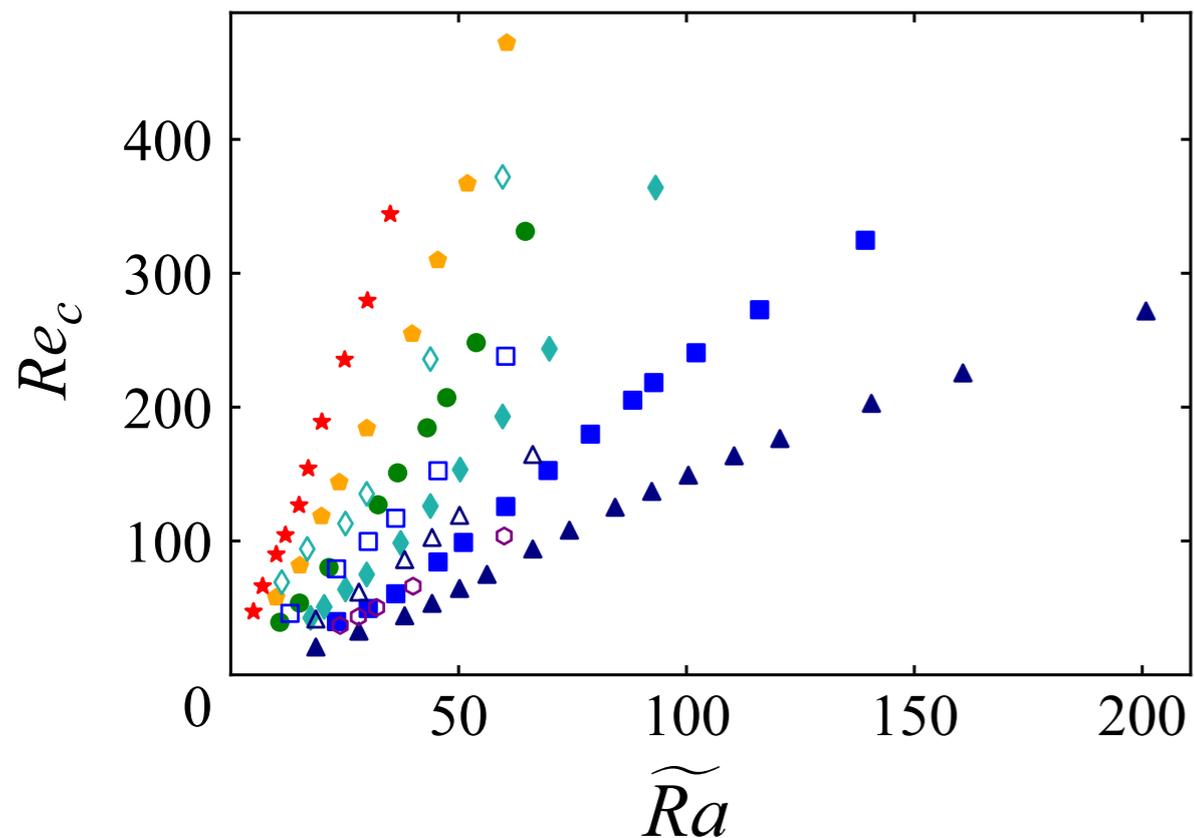
$$\mathbf{u}(r, \theta, \phi, t) = \bar{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Typical Dynamics



$$\mathbf{u}(r, \theta, \phi, t) = \bar{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Convective Speeds

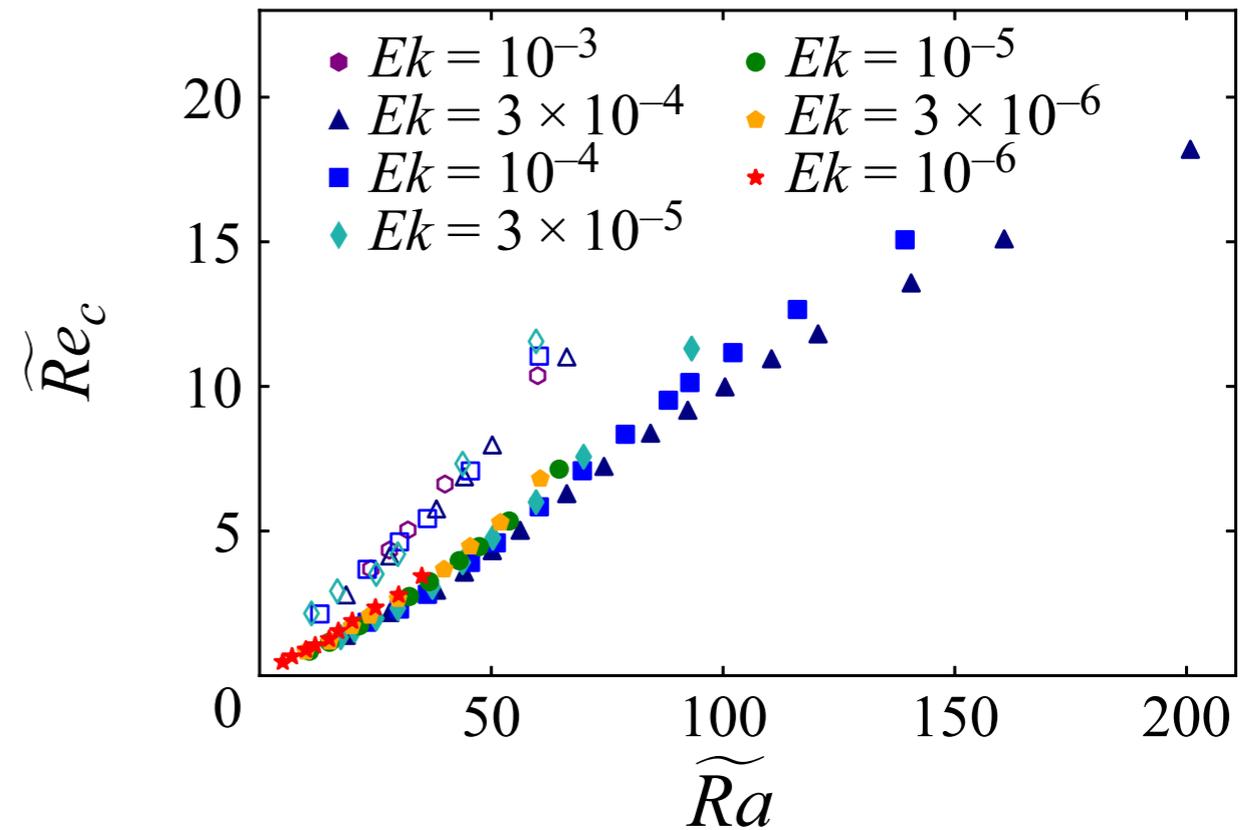
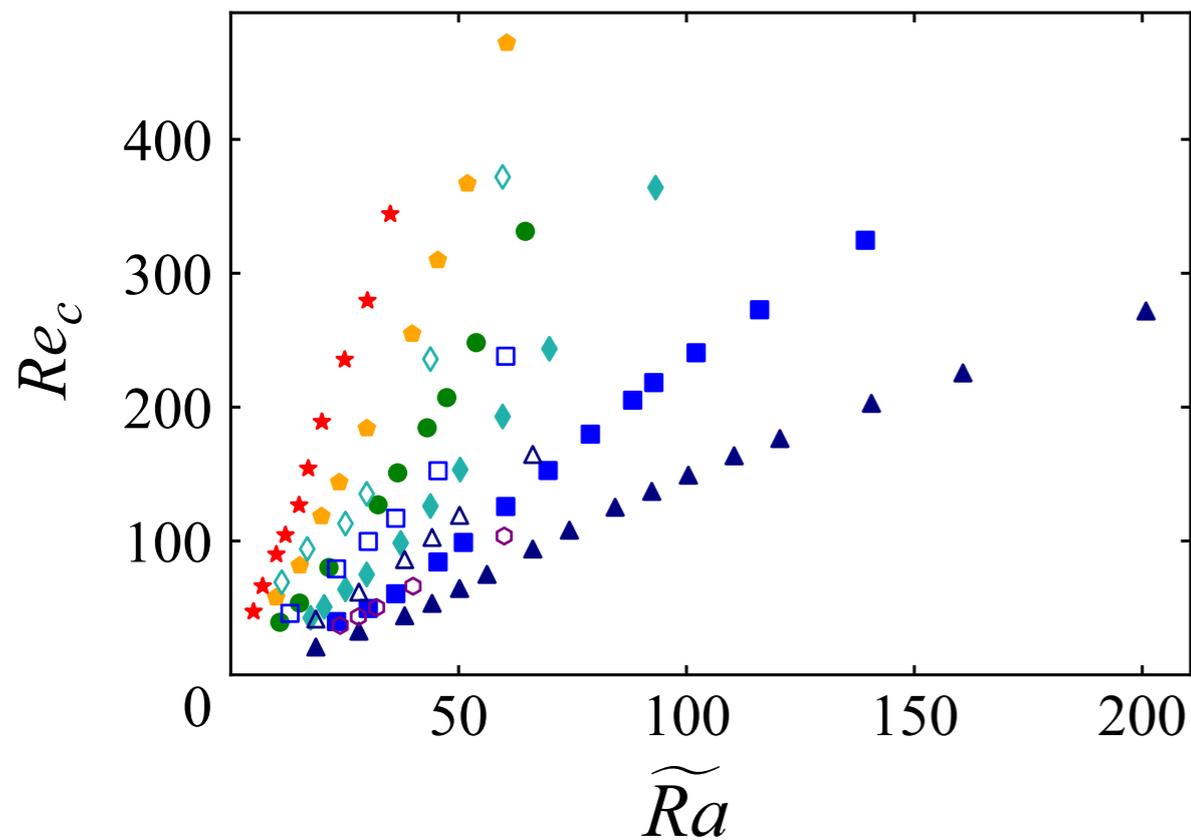


Geostrophy on
viscous length scale:

$$Re_c = O\left(Ek^{-1/3}\right)$$

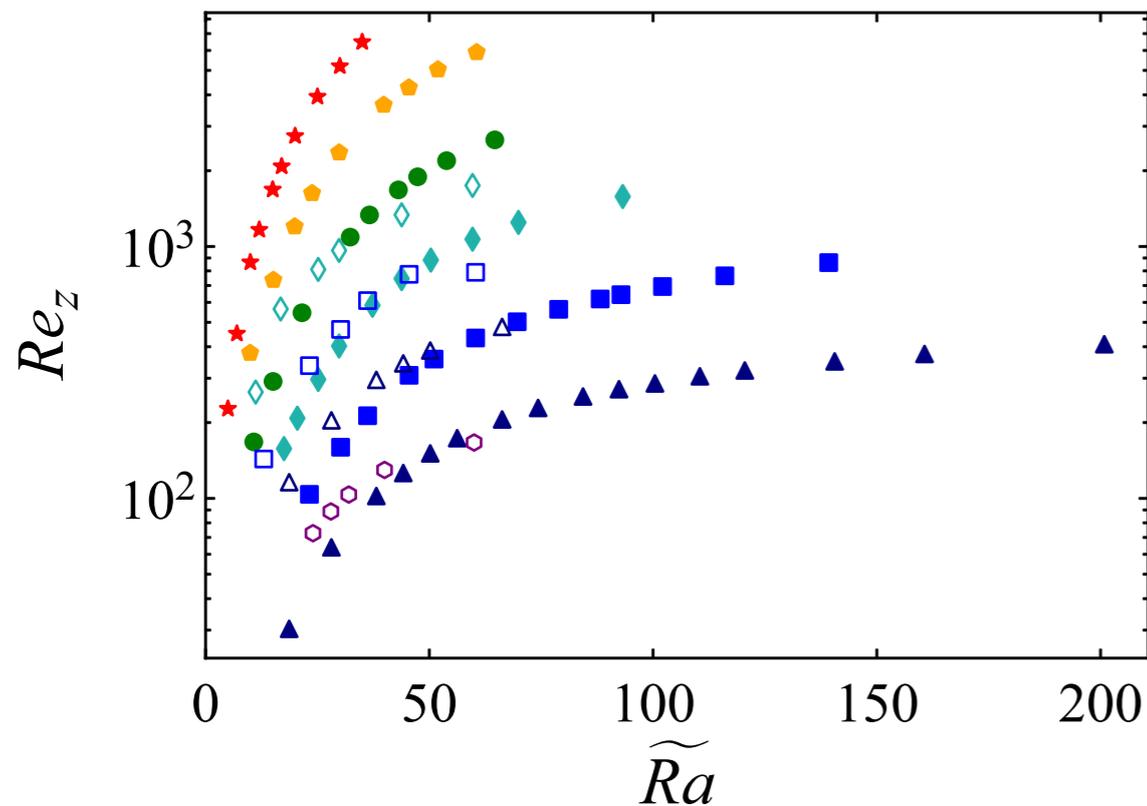
$$\mathbf{u}(r, \theta, \phi, t) = \bar{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Convective Speeds



$$\mathbf{u}(r, \theta, \phi, t) = \overline{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Zonal Flow Speeds

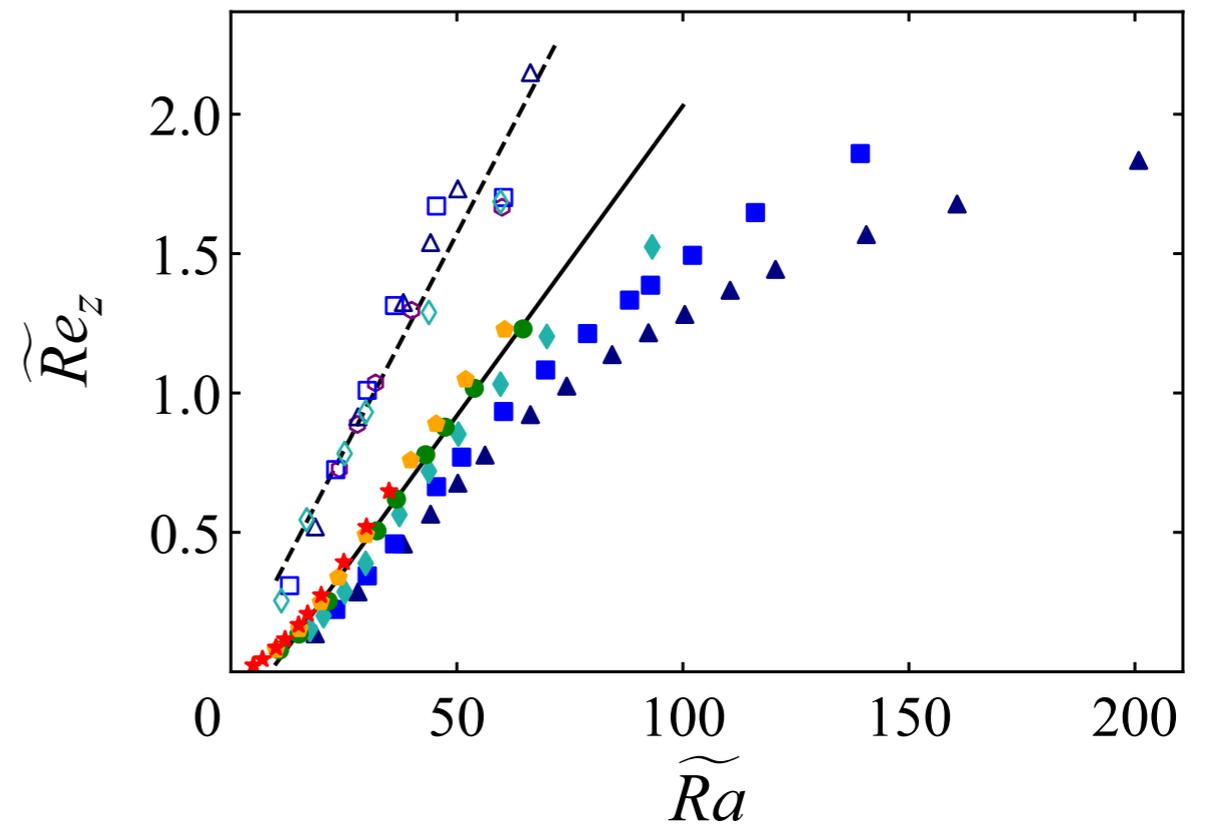
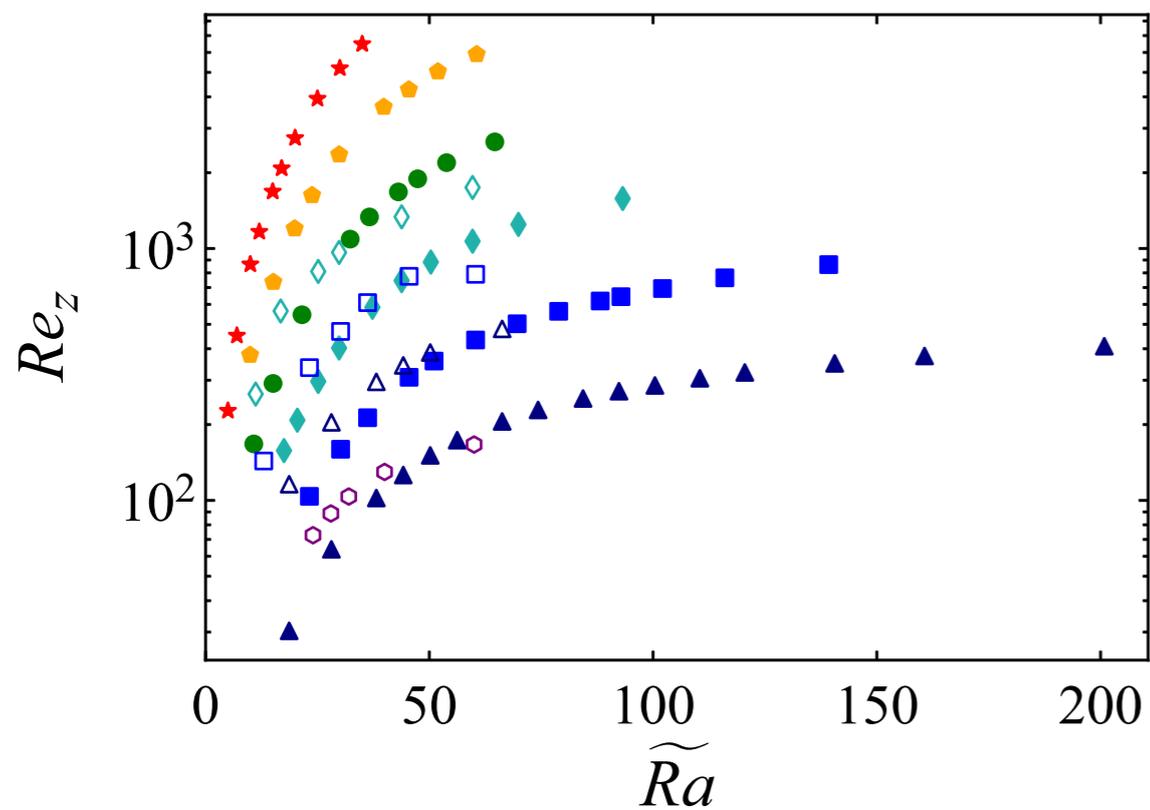


Viscosity/Reynolds stress
balance:

$$Re_z = O\left(Ek^{-2/3}\right)$$

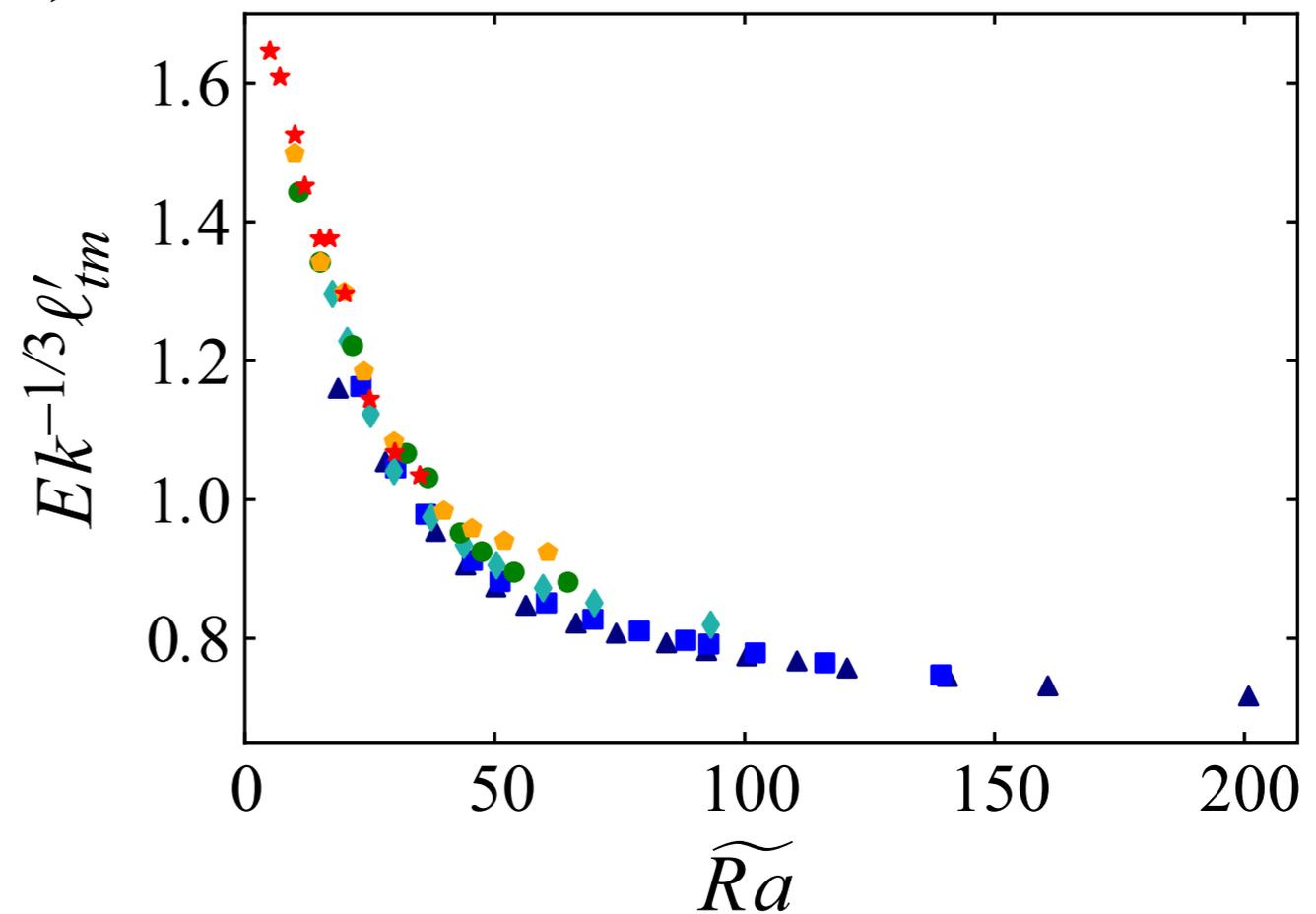
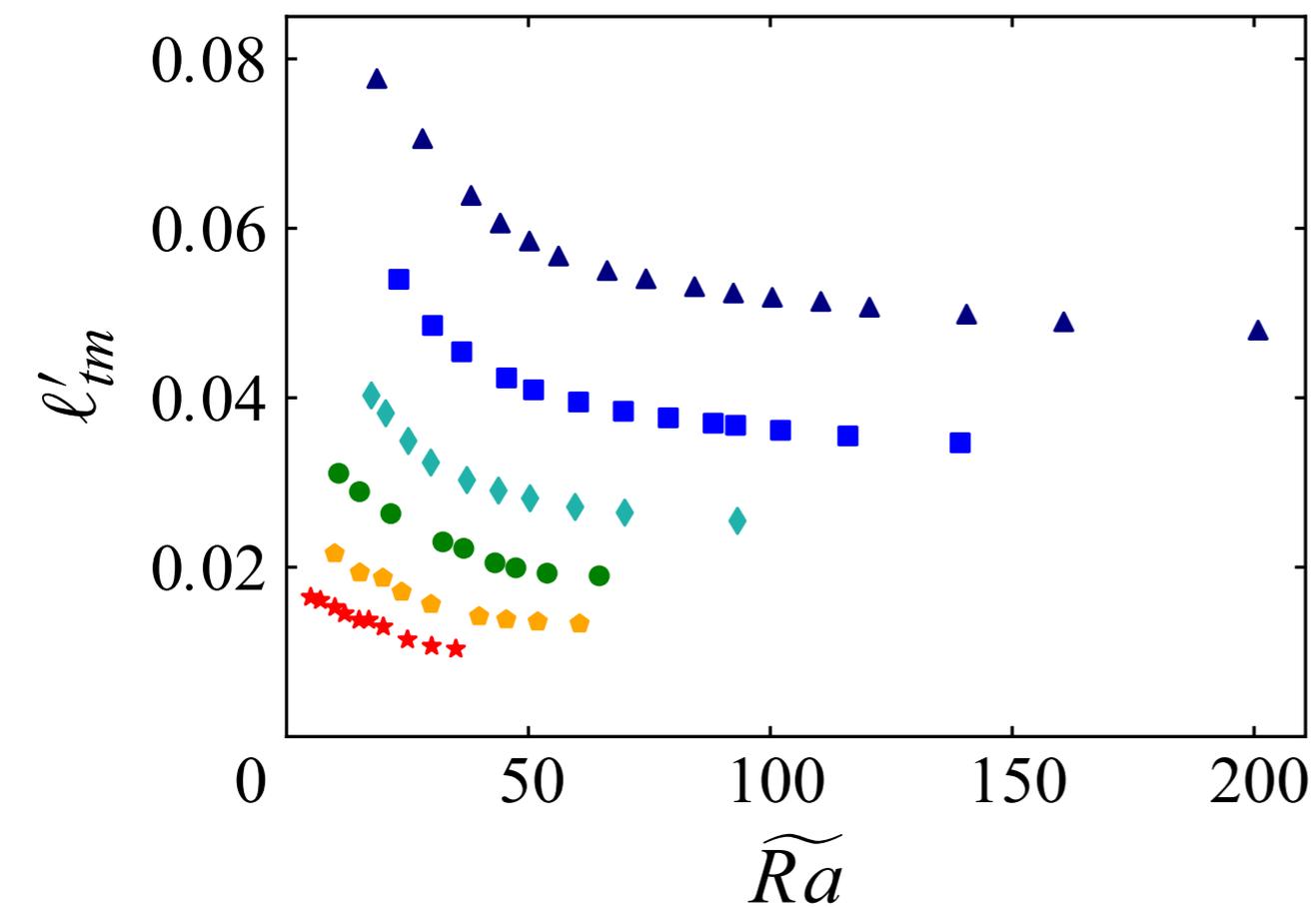
$$\mathbf{u}(r, \theta, \phi, t) = \overline{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Zonal Flow Speeds

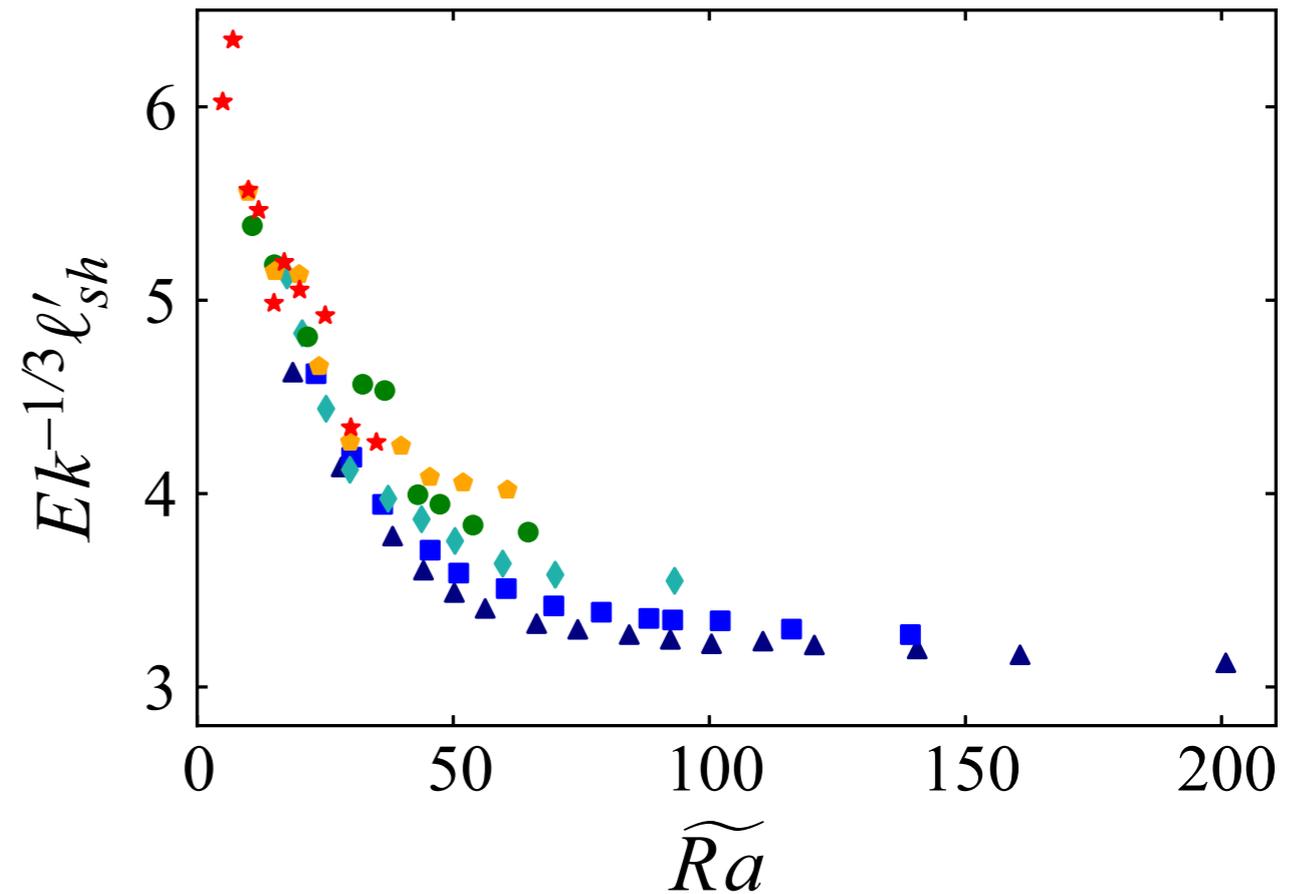
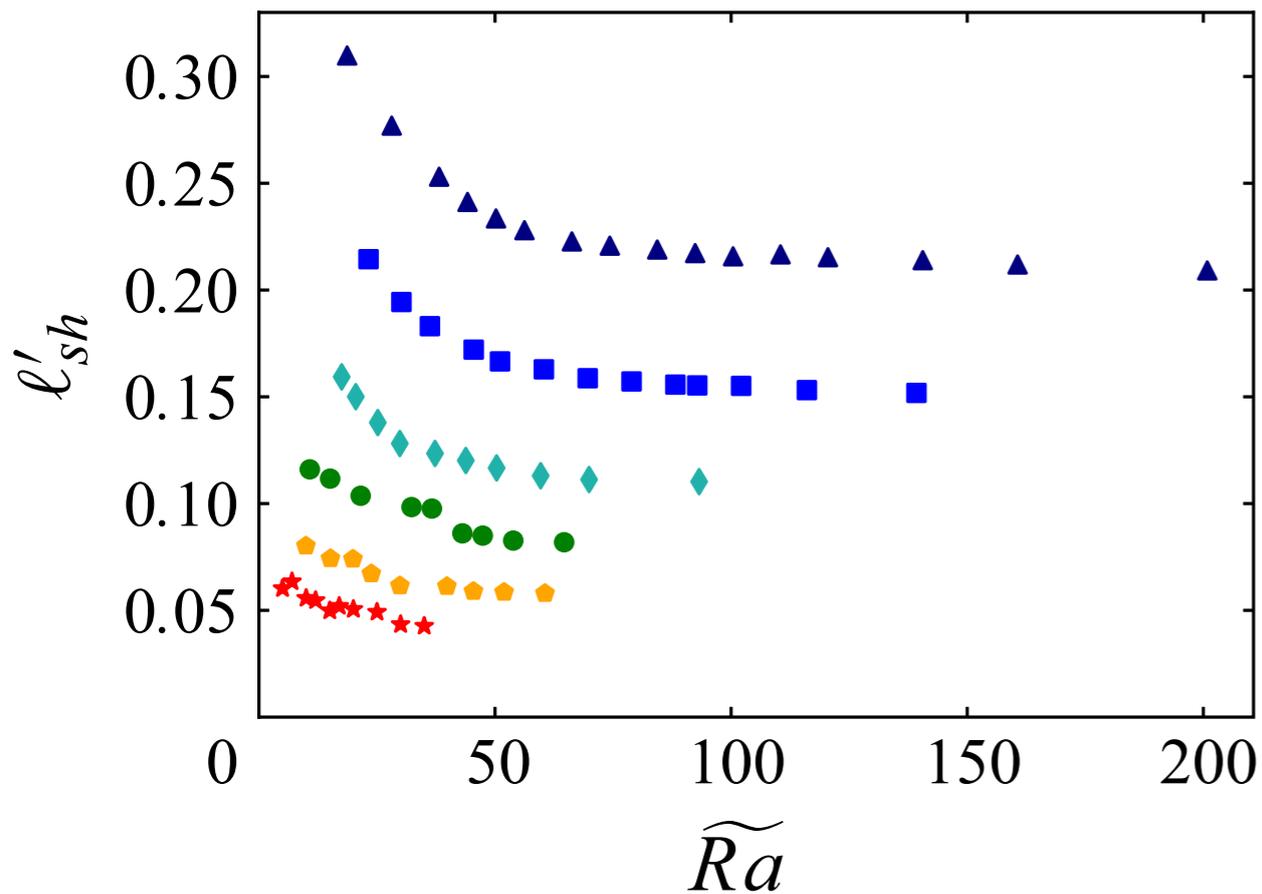


$$\mathbf{u}(r, \theta, \phi, t) = \overline{\mathbf{u}}(r, \theta, t) + \mathbf{u}'(r, \theta, \phi, t)$$

Taylor Microscale

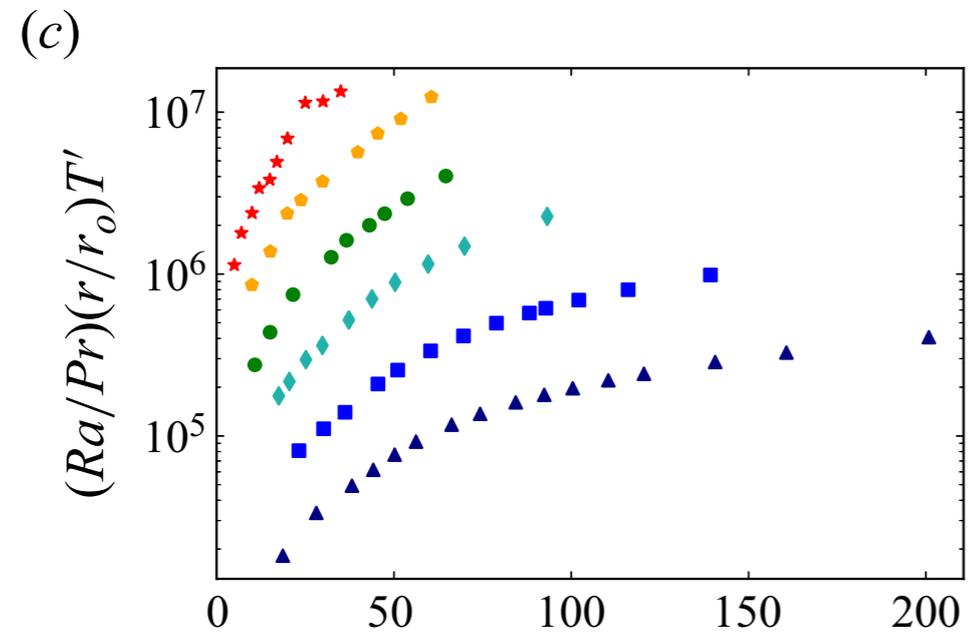


Integral Length Scale



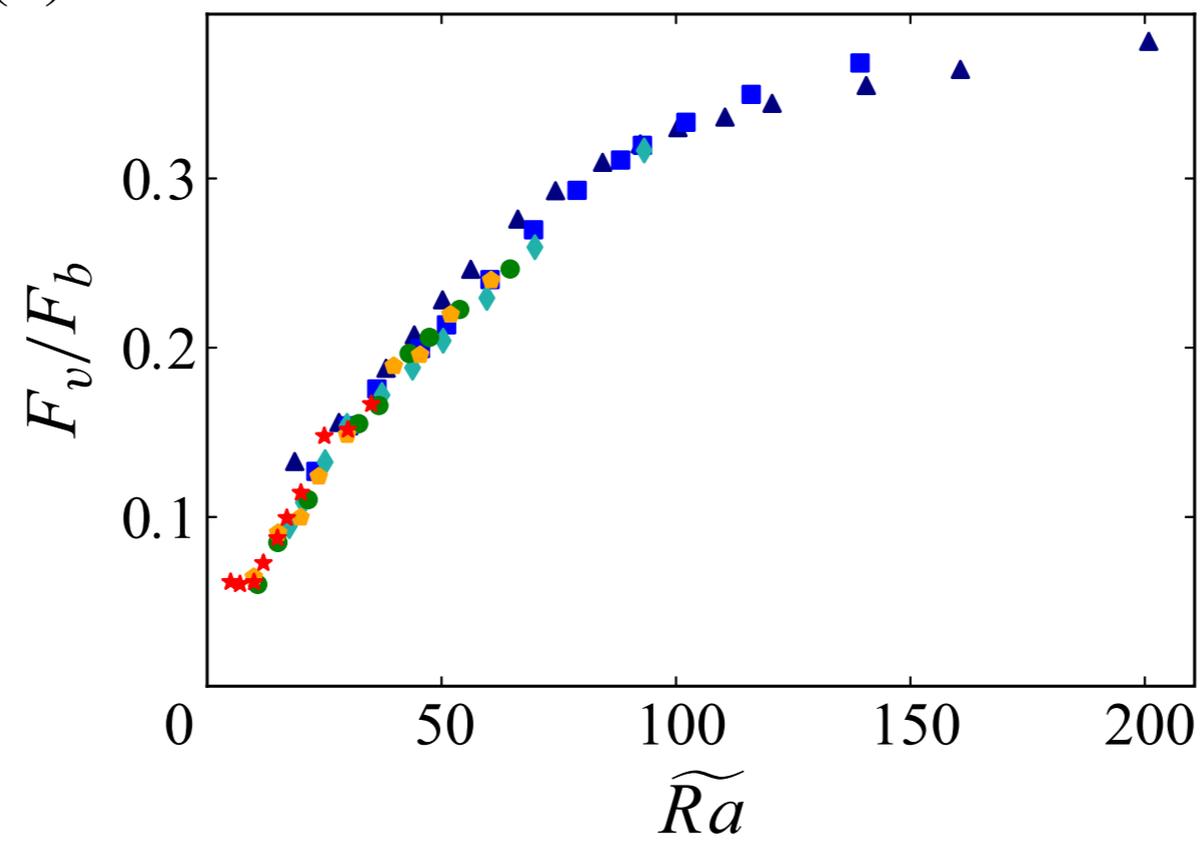
Energetically-dominant length scales: $\ell = O\left(Ek^{1/3}\right)$

Force Scalings



Force Scalings

(a)



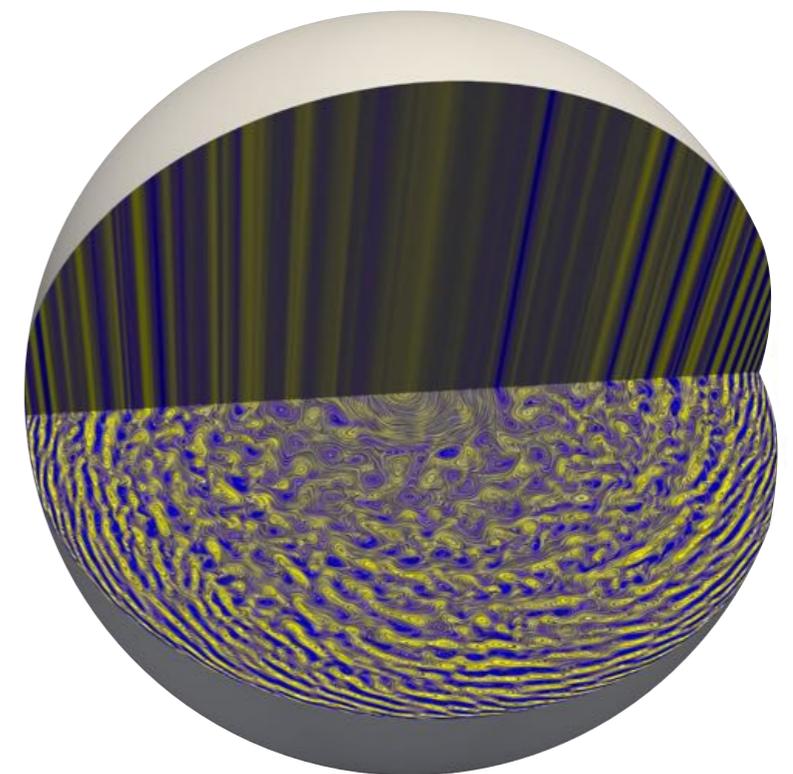
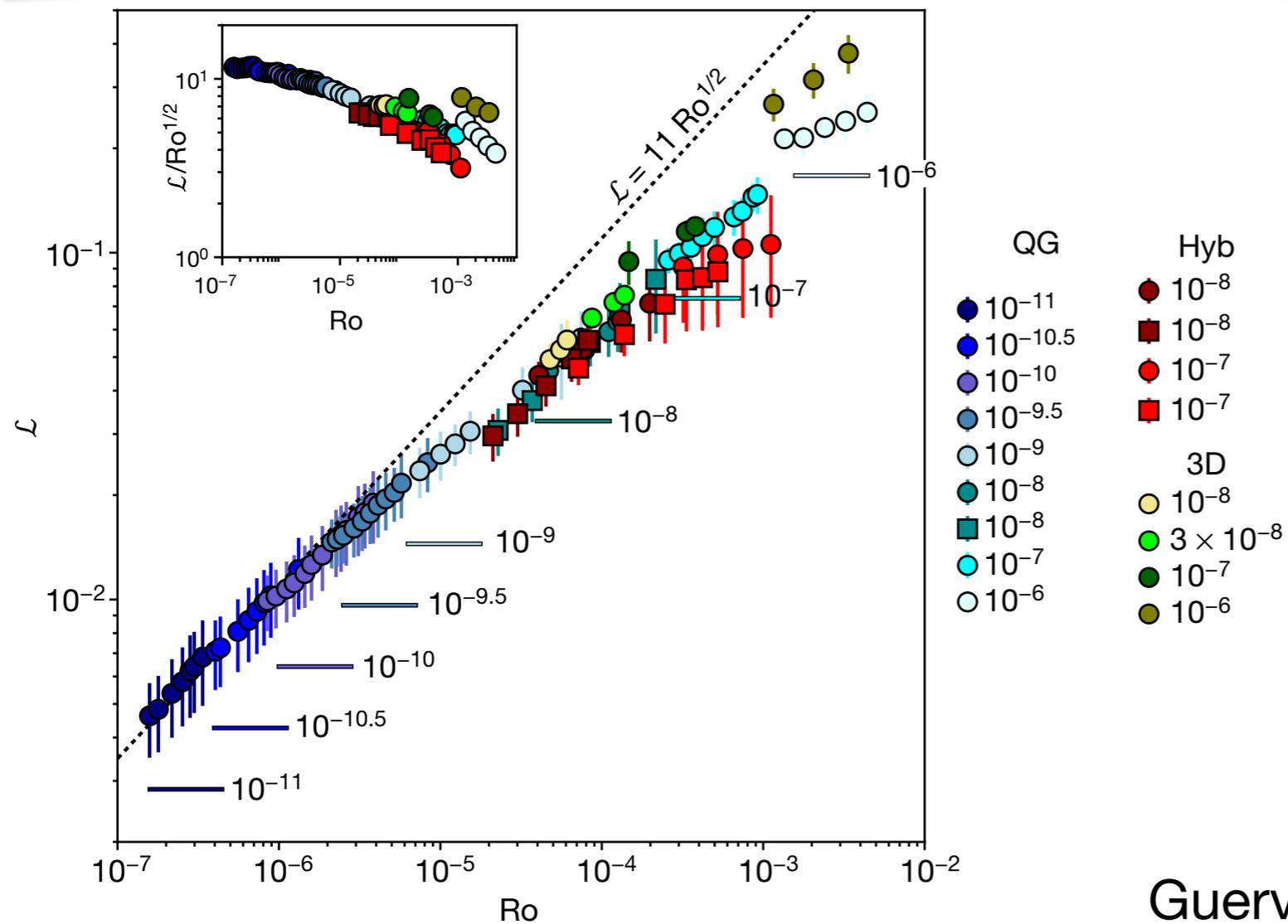
Related Work

LETTER

<https://doi.org/10.1038/s41586-019-1301-5>

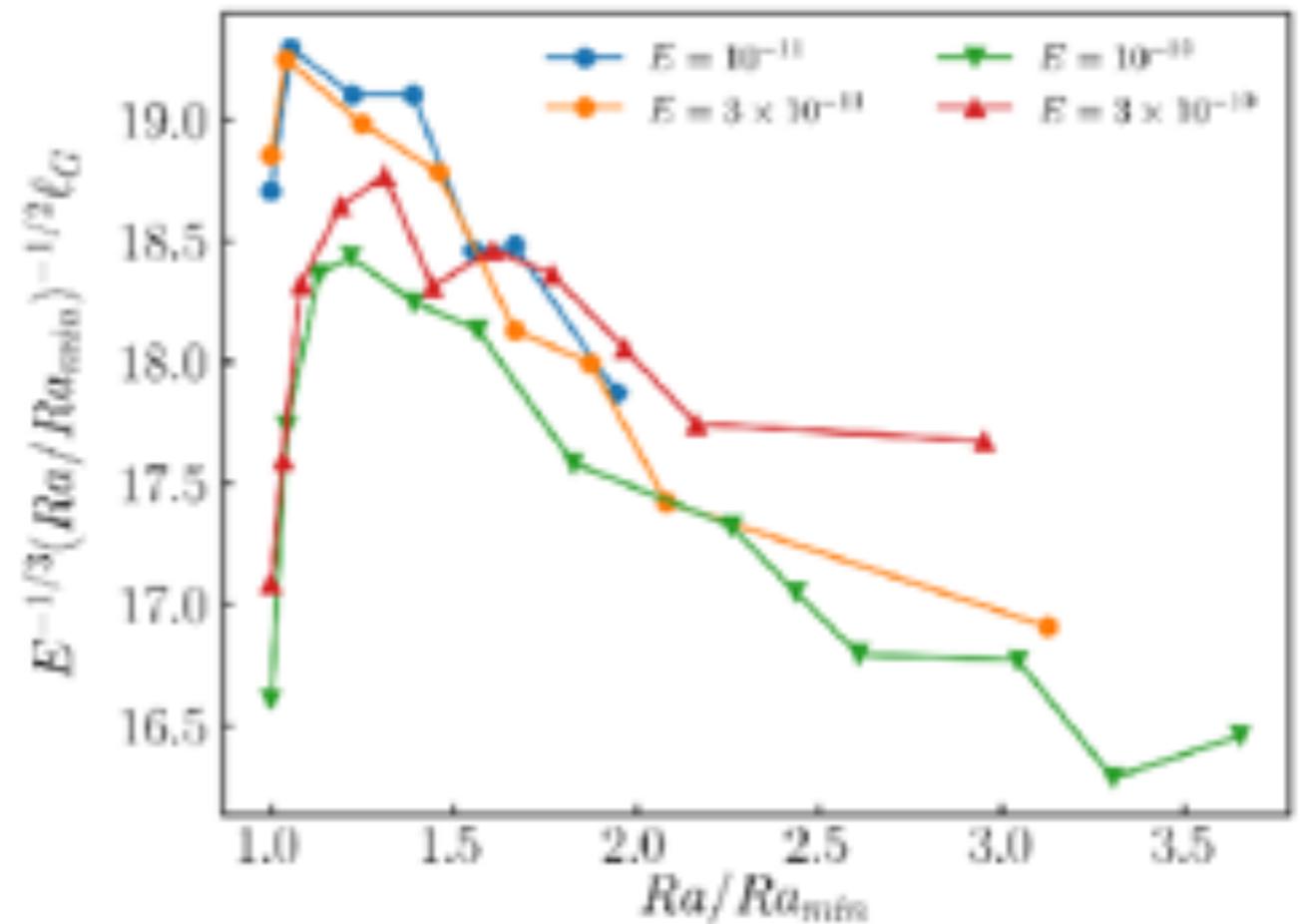
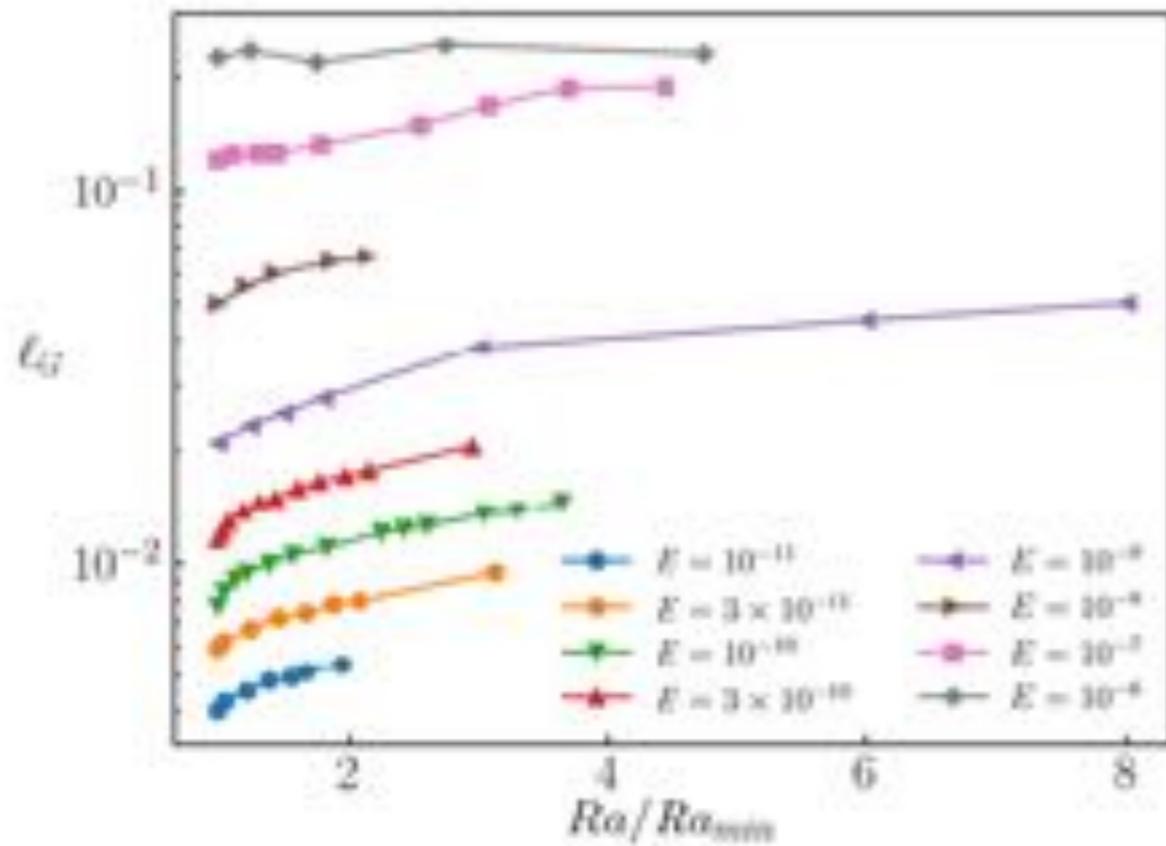
Turbulent convective length scale in planetary cores

Céline Guervilly^{1*}, Philippe Cardin² & Nathanaël Schaeffer²



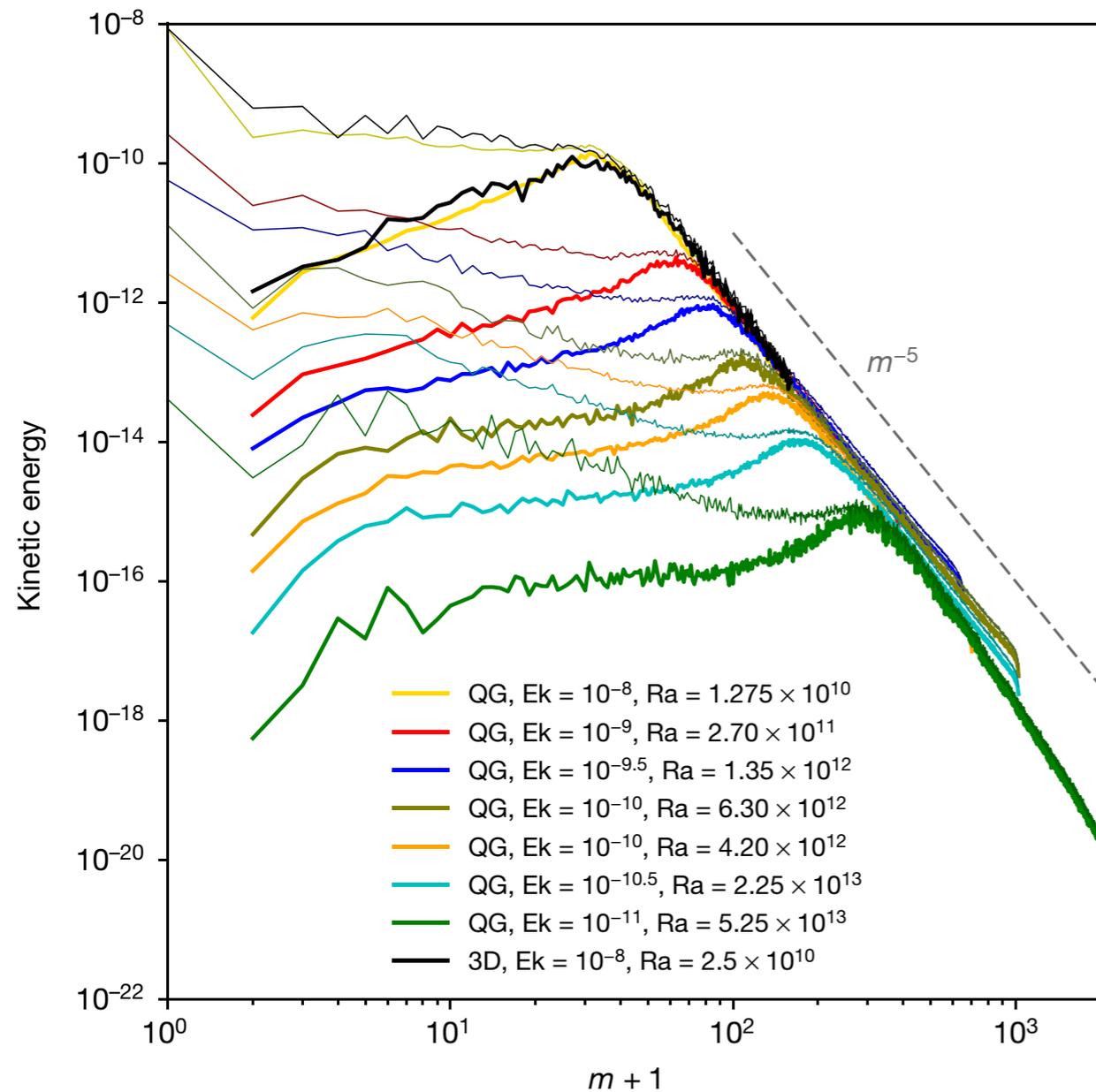
Guervilly, Cardin & Schaeffer (Nat, 2019)

Related Work

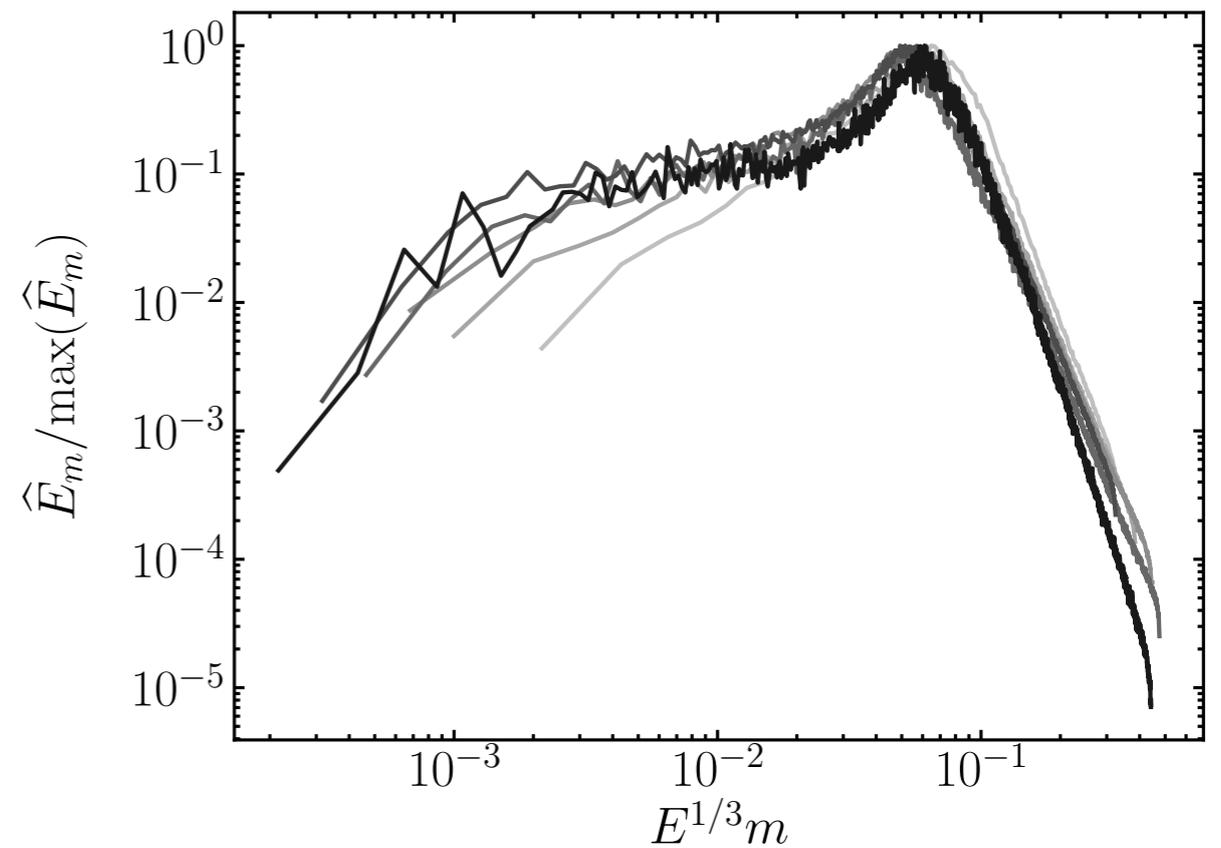
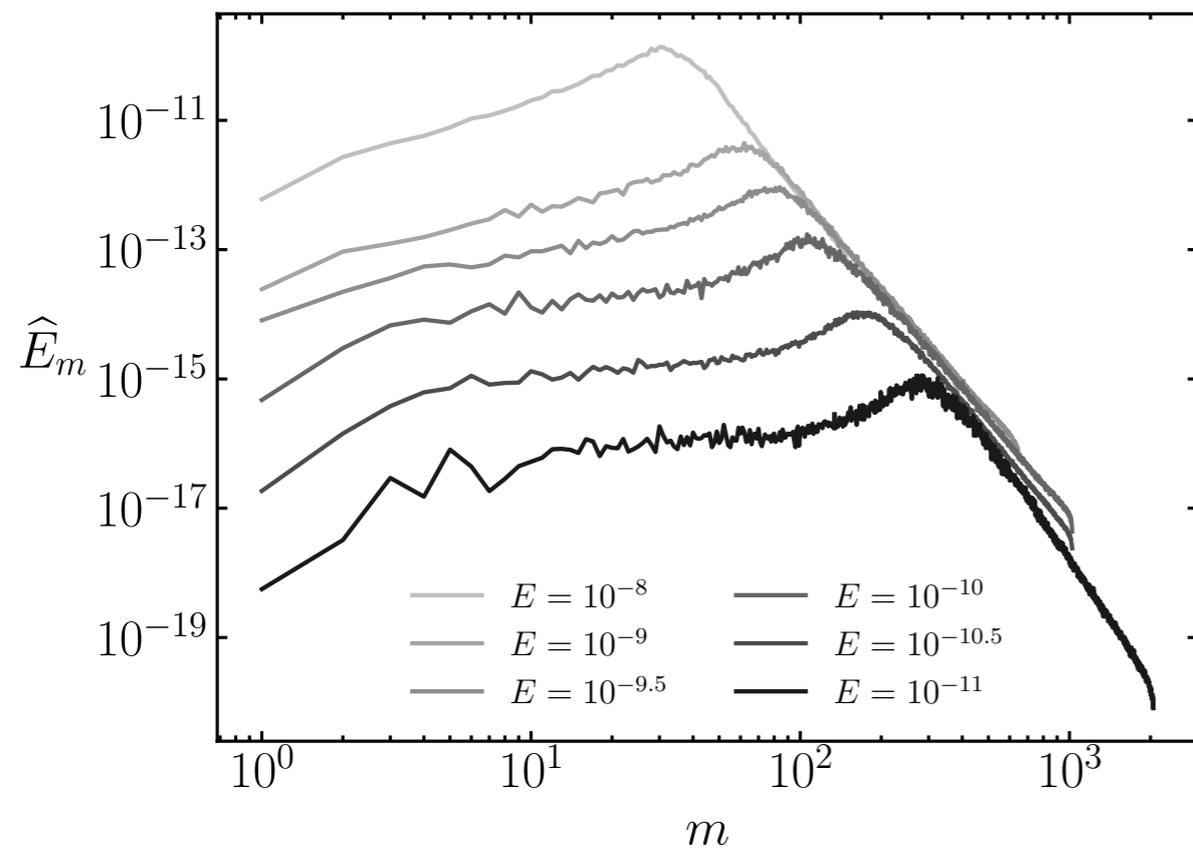


Data from: Guervilly, Cardin & Schaeffer (Nat, 2019)

Related Work



Related Work



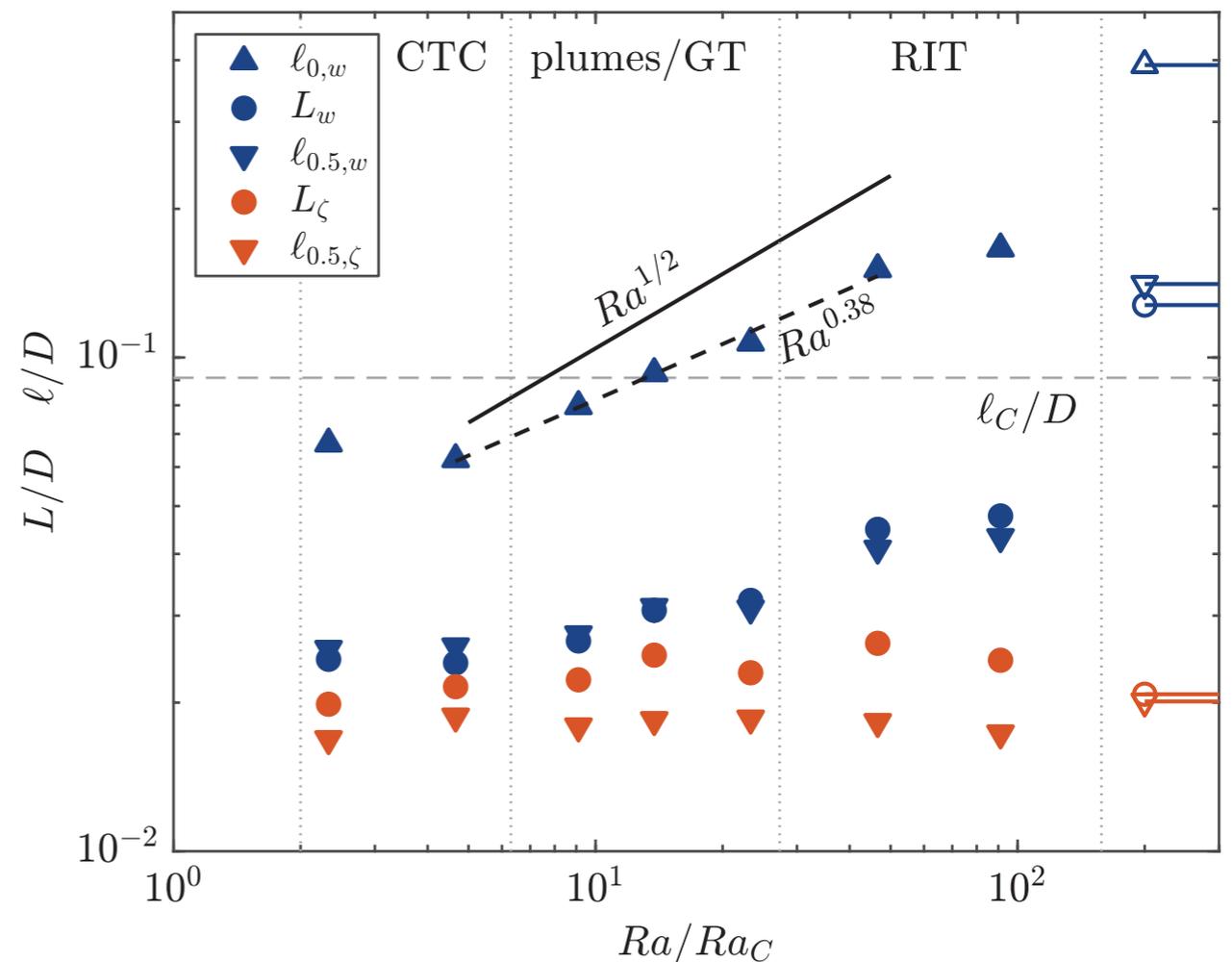
Data from: Guervilly, Cardin & Schaeffer (Nat, 2019)

Related Work

Velocimetry in rapidly rotating convection: Spatial correlations, flow structures and length scales^(a)

MATTEO MADONIA, ANDRÉS J. AGUIRRE GUZMÁN, HERMAN J. H. CLERCX and RUDIE P. J. KUNNEN^(b) 

- Length scale grows less quickly than CIA.
- Taylor microscale is approximately constant.



Madonia et al. (EPL, 2021)

Summary

- The viscous length scale persists in the turbulent regime of rotating convection, independent of geometry.
- While diffusion-free scalings may approximately describe data, the reason for this is unclear.
 - The importance of the viscous force is likely tied to saturation of interior temperature gradient.
- Spherical data: scalings are all consistent with plane layer asymptotic behavior, zonal flow flows exhibit asymptotic dependence on Ekman number.