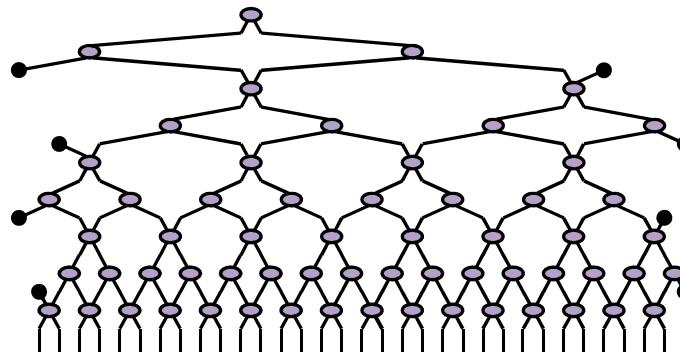


Recent progress in Entanglement Renormalization:



Guifre Vidal

In collaboration with:

Glen Evenbly (PhD student)

Robert Pfeifer (PhD student)

Miguel Aguado (MPQ Garching)

Robert Koenig (Caltech)

Ben Reichardt (IQC Waterloo)

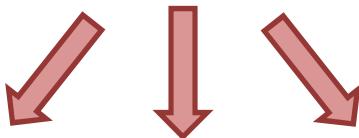
Outline

Foundations

Entanglement
Renormalization

MERA

(Multi-scale Entanglement
Renormalization Ansatz)



Applications

Quantum
Criticality

Topological
Order

Frustrated
Antiferromagnets

- MERA \leftrightarrow CFT

- Exact MERA
for quantum double
and string-net models

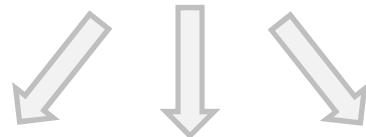
- Kagome lattice
- Triangular lattice

Foundations

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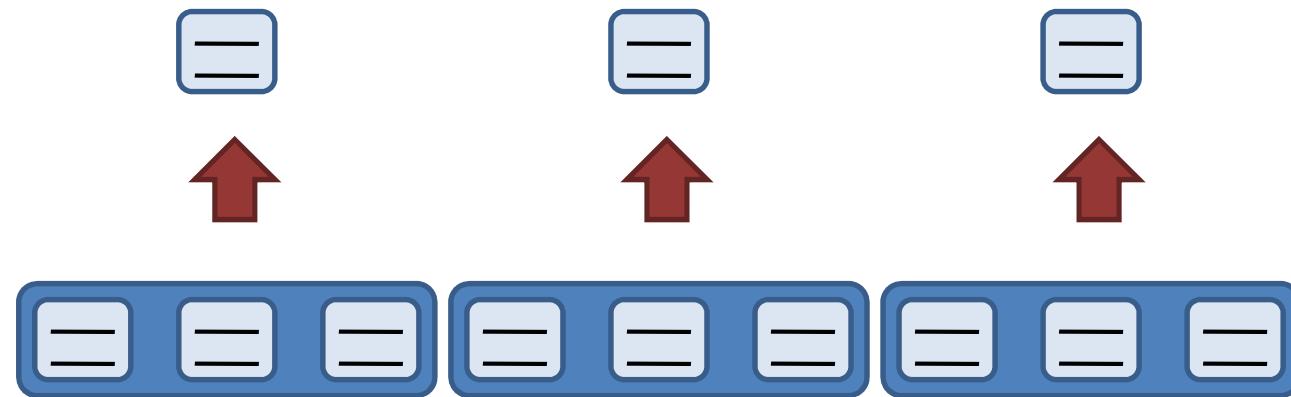
• MERA \leftrightarrow CFT

• Exact MERA
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• Kagome lattice
• Triangular lattice

Introduction

- real space RG transformation



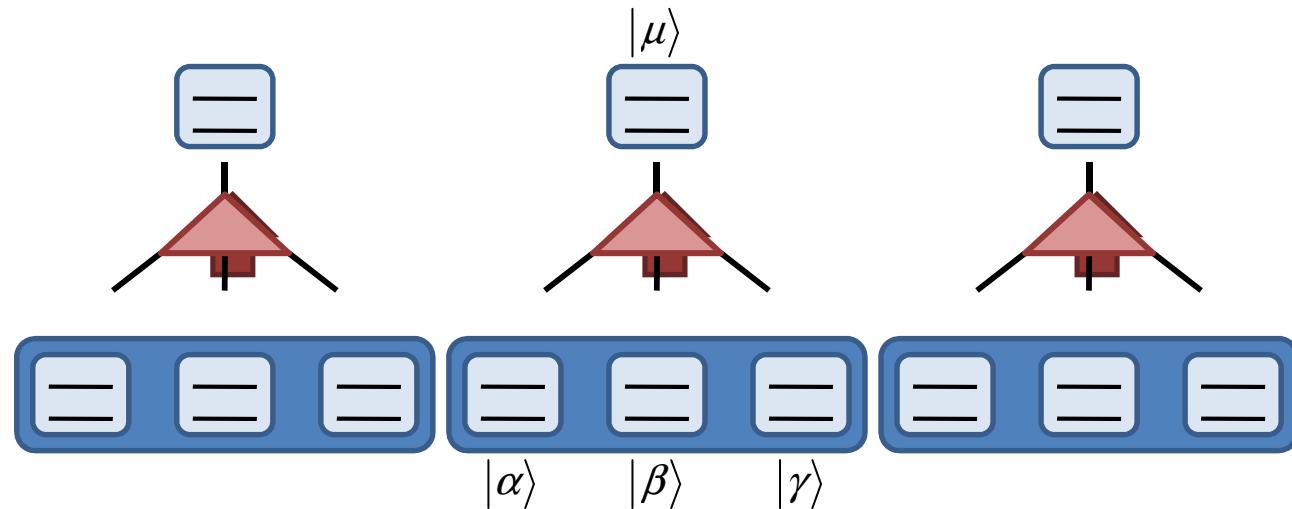
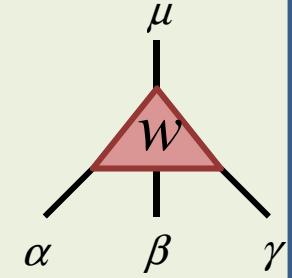
- Kadanoff (1960's): spin blocking

Introduction

- real space RG transformation

Change of basis

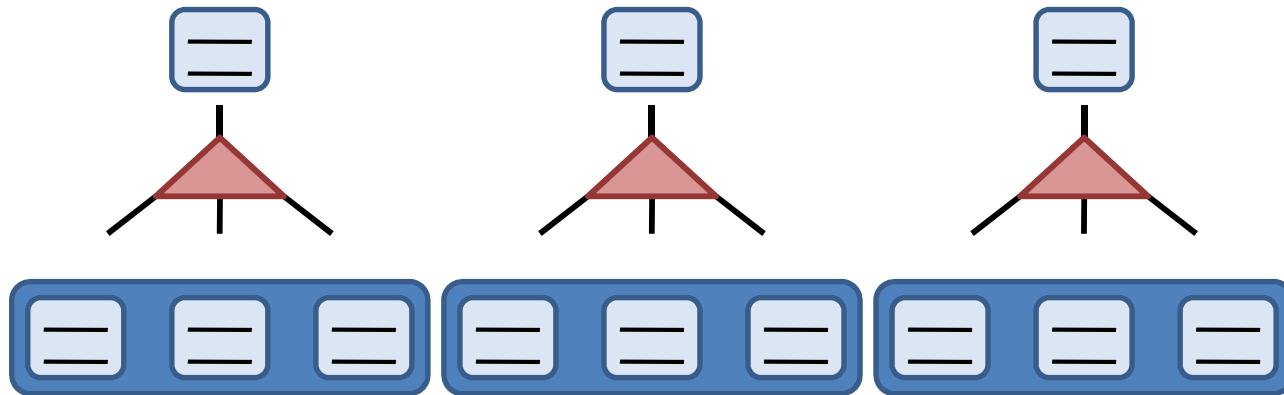
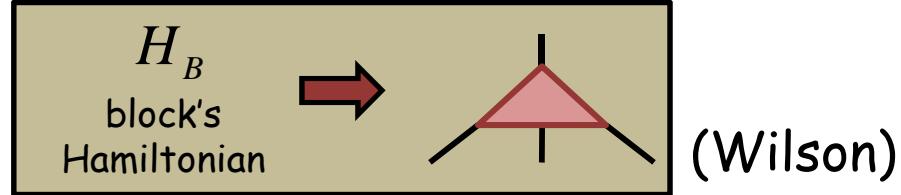
$$|\mu\rangle = \sum_{\alpha\beta\gamma} w_{\alpha\beta\gamma}^{\mu} |\alpha\rangle |\beta\rangle |\gamma\rangle$$



- Kadanoff (1960's): spin blocking
- Wilson (1970's): keep low energy subspace
Kondo impurity problem: Numerical Renormalization Group (NRG)

Introduction

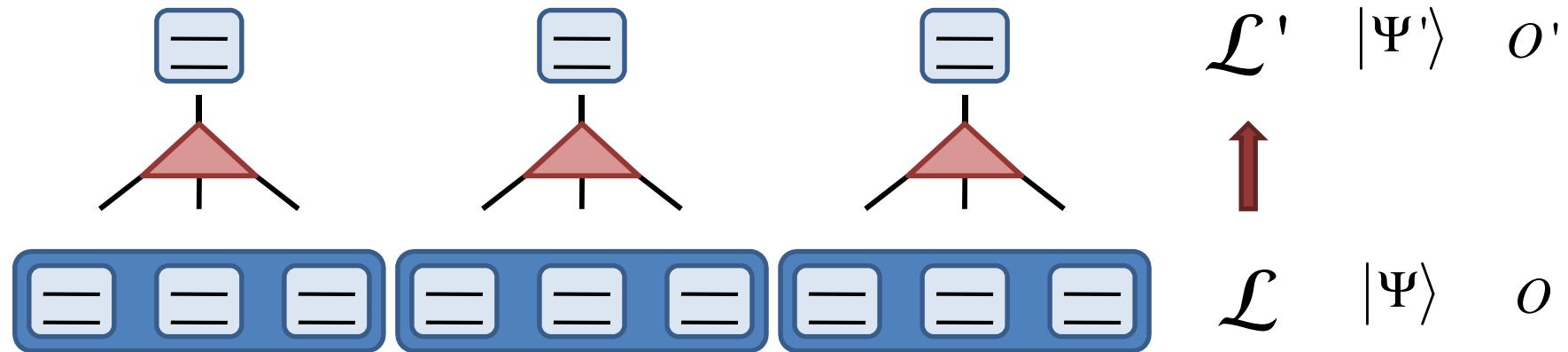
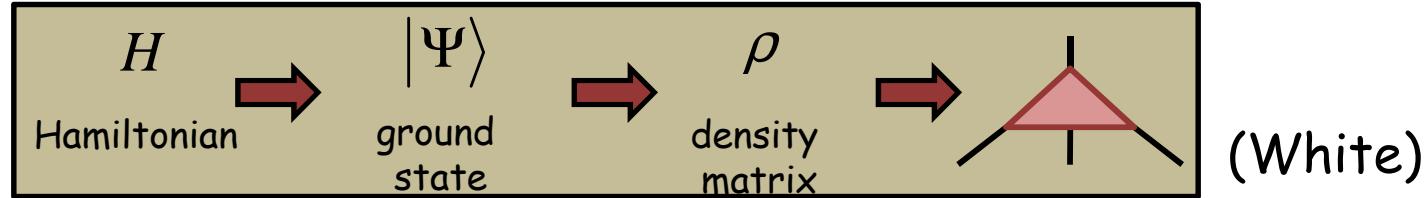
- real space RG transformation



- Kadanoff (1960's): spin blocking
- Wilson (1970's): keep low energy subspace
Kondo impurity problem: Numerical Renormalization Group (NRG)

Introduction

- real space RG transformation



- Kadanoff (1960's): spin blocking

$$\langle \Psi | O | \Psi \rangle = \langle \Psi' | O' | \Psi' \rangle$$

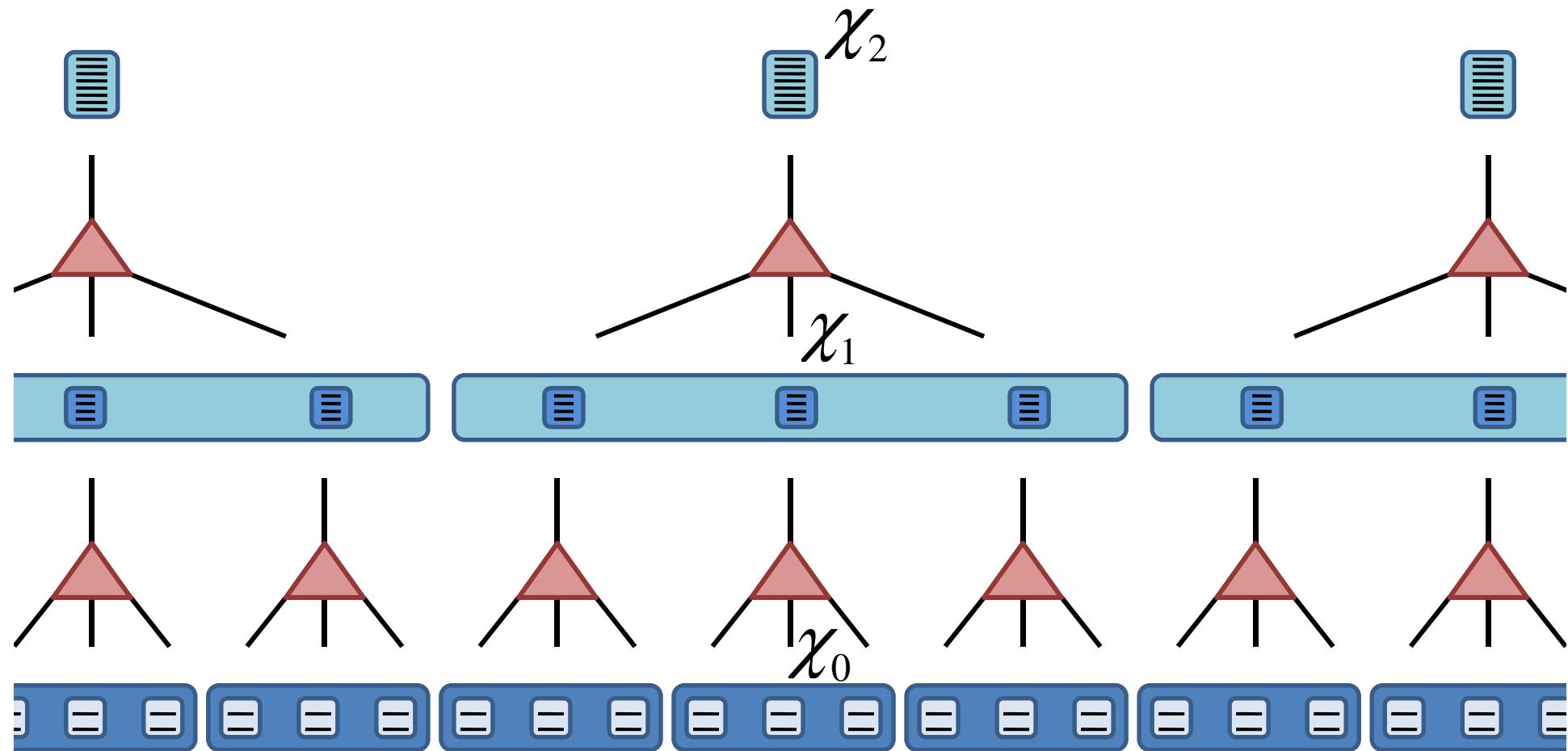
- Wilson (1970's): keep low energy subspace
Kondo impurity problem: Numerical Renormalization Group (NRG)

- White (1990's): keep local support of ground state reduce density matrix
Density Matrix Renormalization Group (DMRG)

Introduction

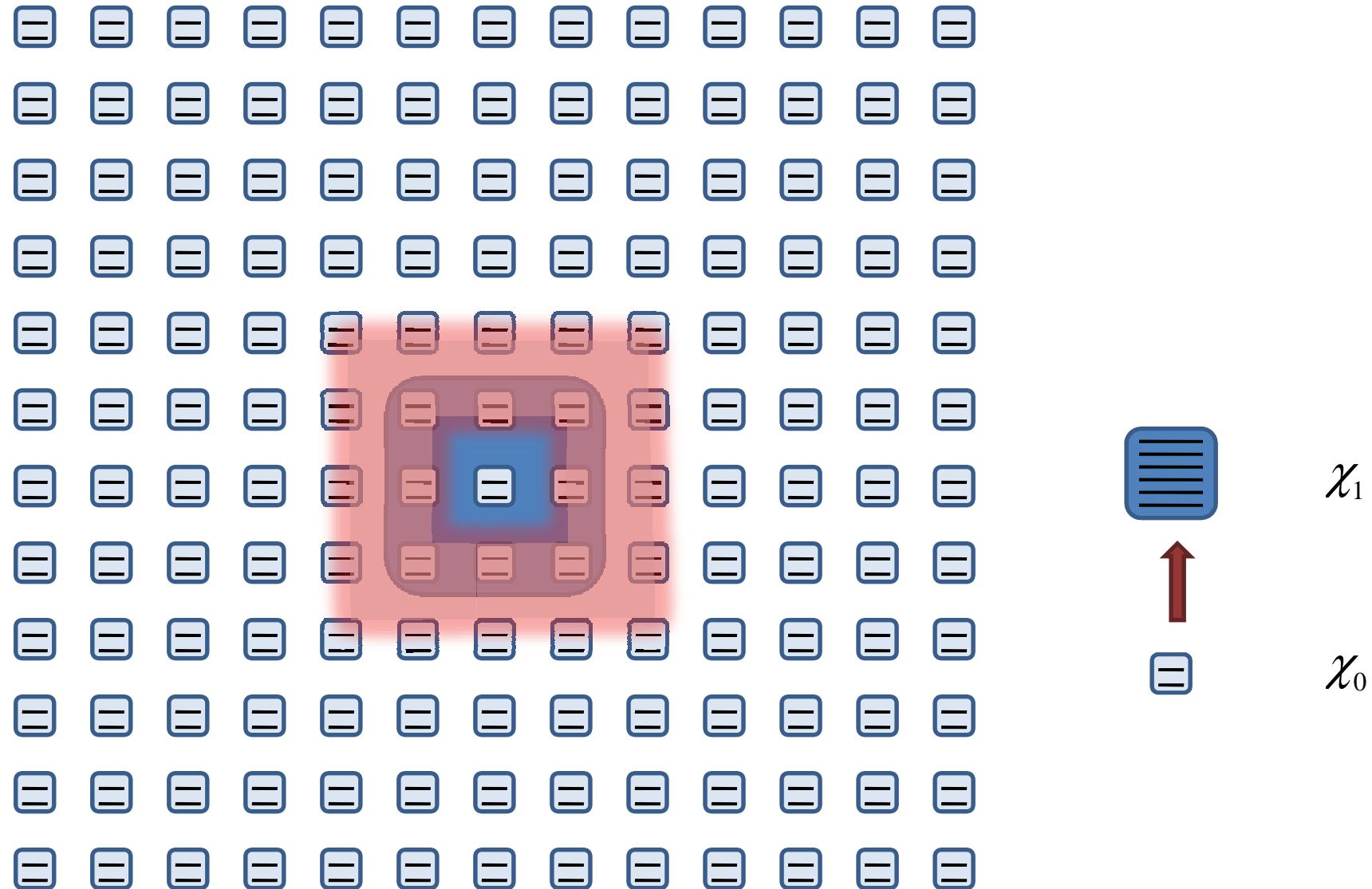
- entanglement in 1D critical

$$S(\rho_L) \sim k \log L \Rightarrow \chi_\tau \sim e^\tau$$



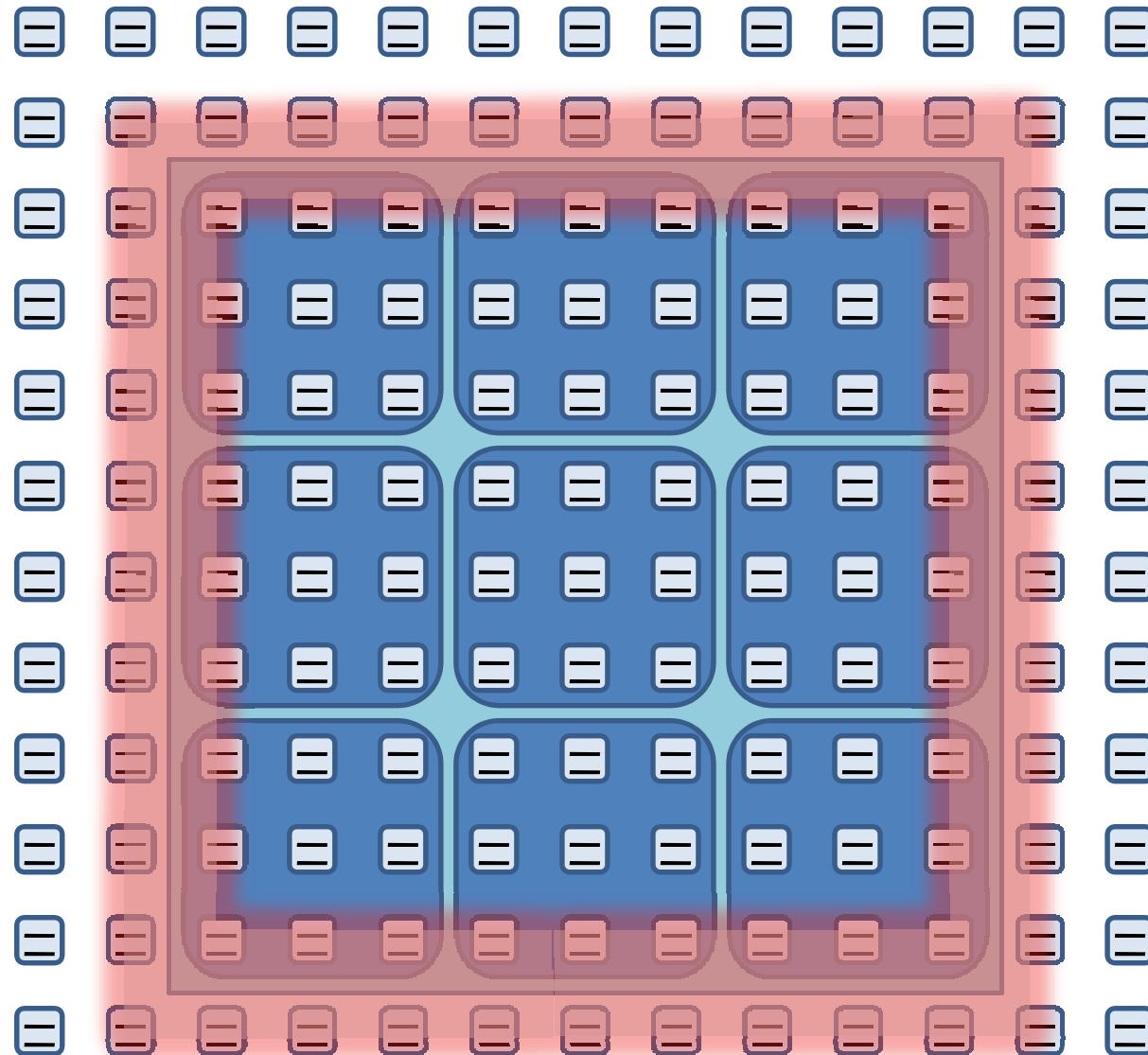
Introduction

- entanglement in 2D: boundary law $S(\rho_{L \times L}) \sim L$

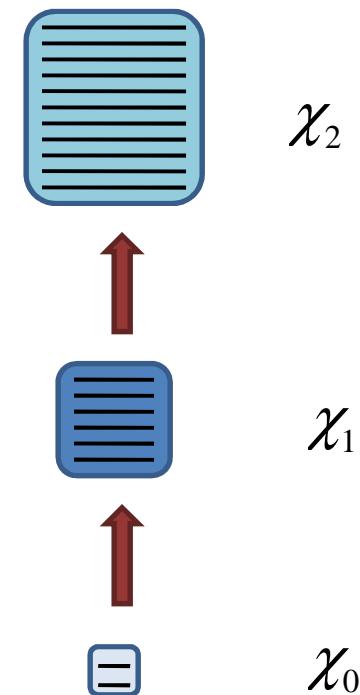


Introduction

- entanglement in 2D: boundary law $S(\rho_{L \times L}) \sim L$

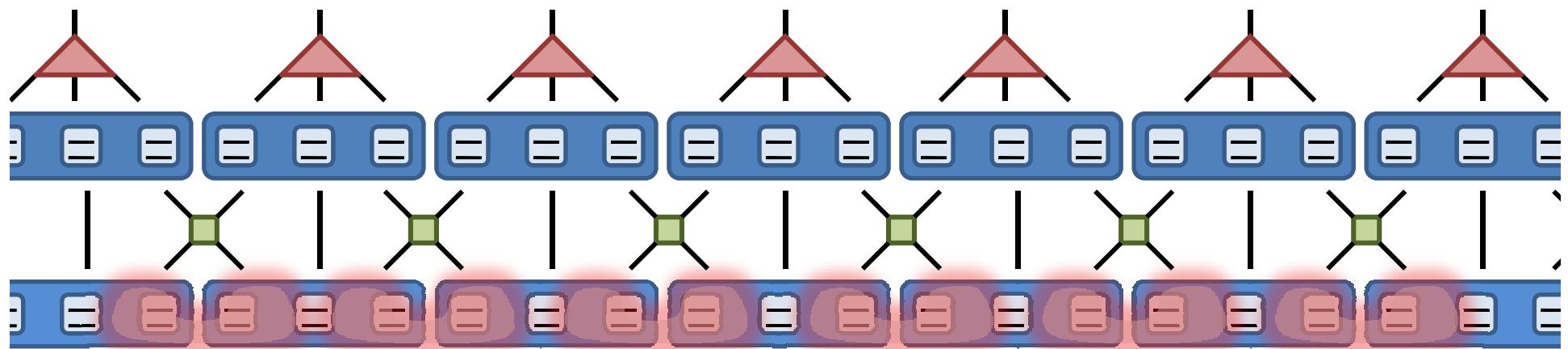
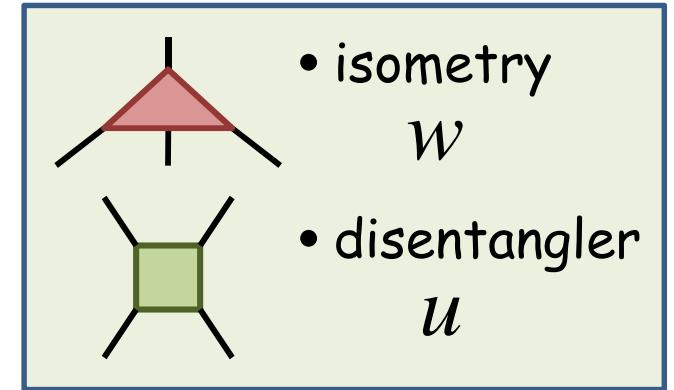


$$\Rightarrow \chi_\tau = e^{e^\tau}$$



Entanglement Renormalization

Vidal, Phys. Rev. Lett. 99, 220405 (2007)

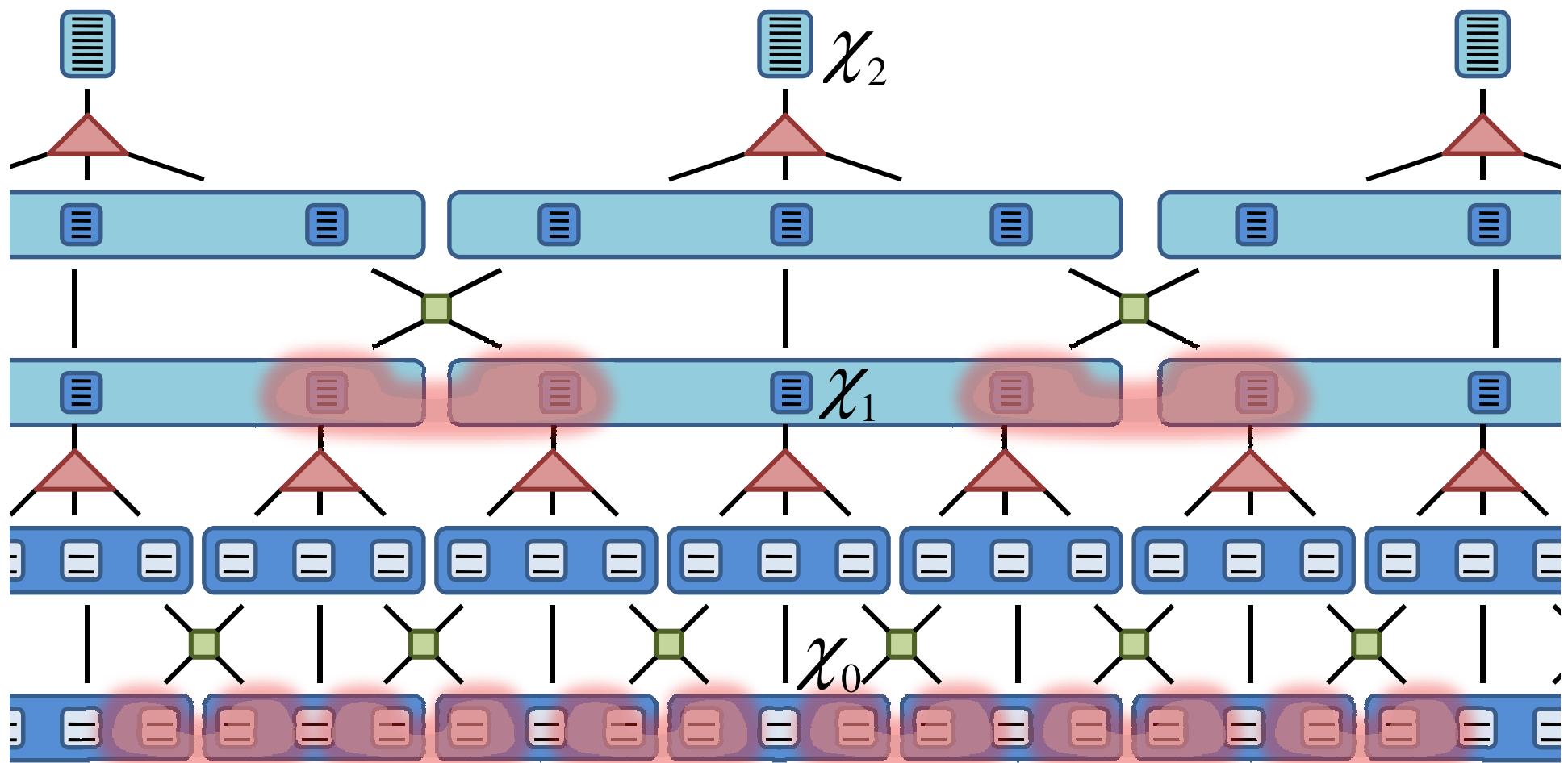
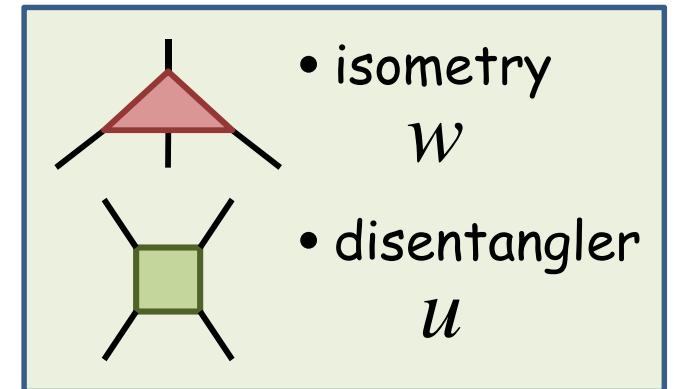


Entanglement Renormalization

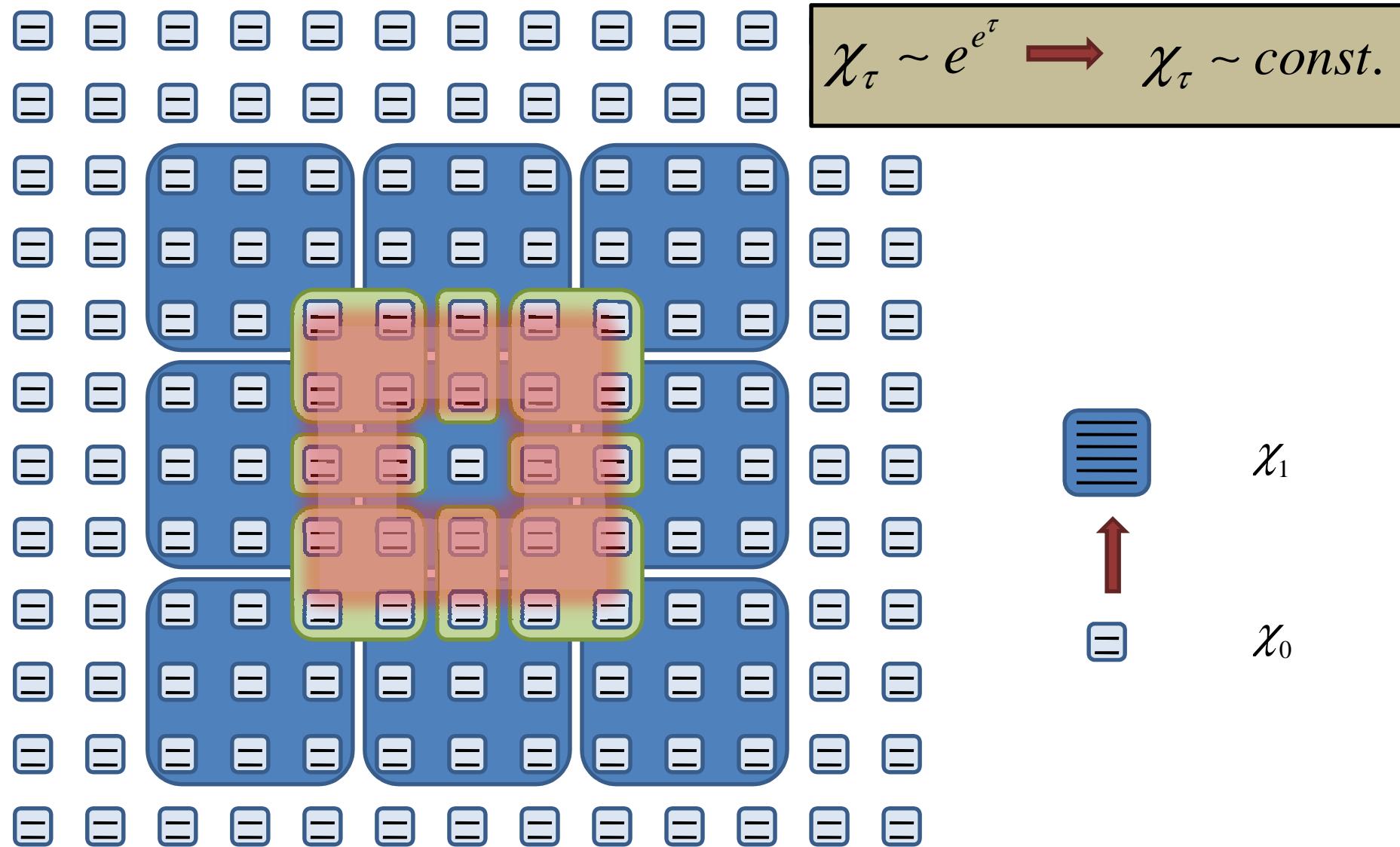
Vidal, Phys. Rev. Lett. 99, 220405 (2007)

$$\chi_\tau \sim e^\tau \rightarrow \chi_\tau \sim \text{const.}$$

(1D critical)

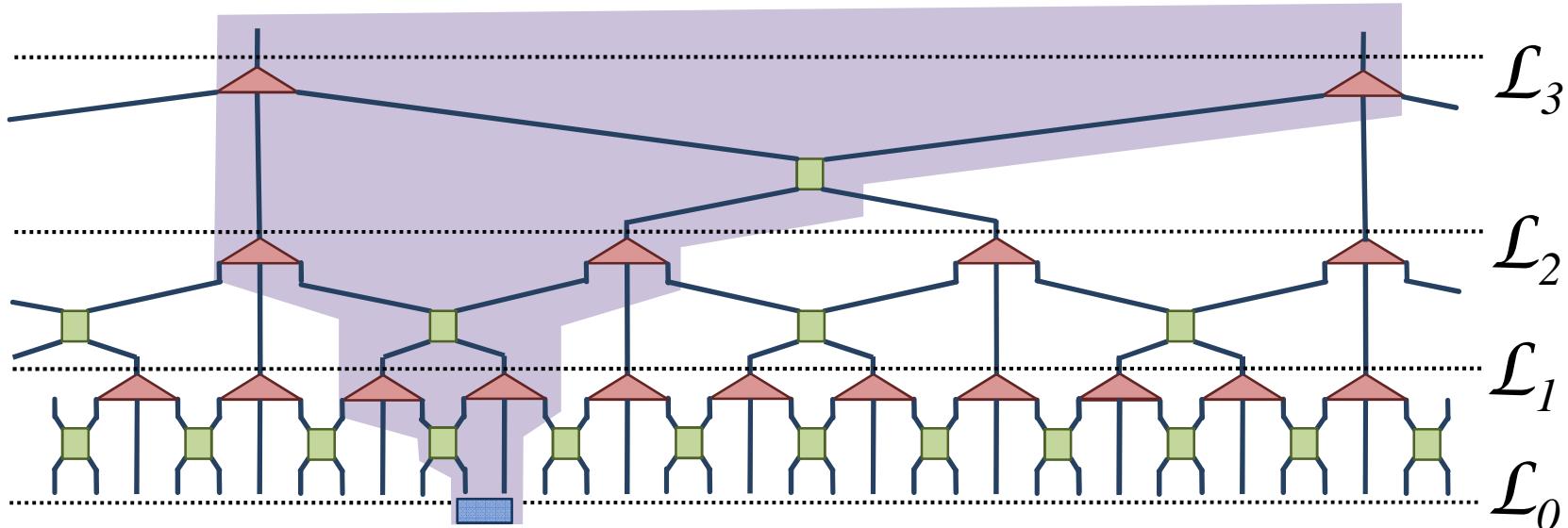


- entanglement in 2D: boundary law $S(\rho_{L \times L}) \sim L$



Entanglement Renormalization

Vidal, Phys. Rev. Lett. 99, 220405 (2007)



disentangler

$$u \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} = \begin{array}{c} | \\ | \end{array} I$$

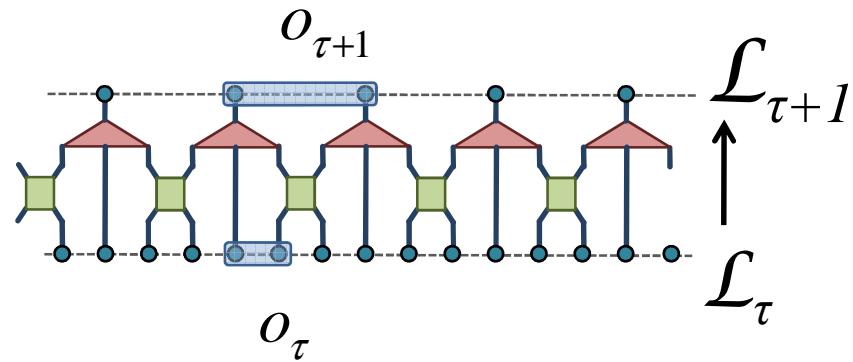
$$u^\dagger \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} = \begin{array}{c} | \\ | \end{array}$$

isometry

$$w \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} = \begin{array}{c} | \\ | \end{array} I$$

$$w^\dagger \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} = \begin{array}{c} | \\ | \end{array}$$

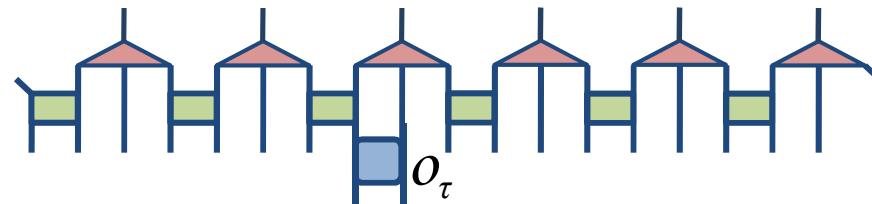
- transformation of local operators ?



Local operators

$$\mathcal{L}_\tau \longrightarrow \mathcal{L}_{\tau+1}$$

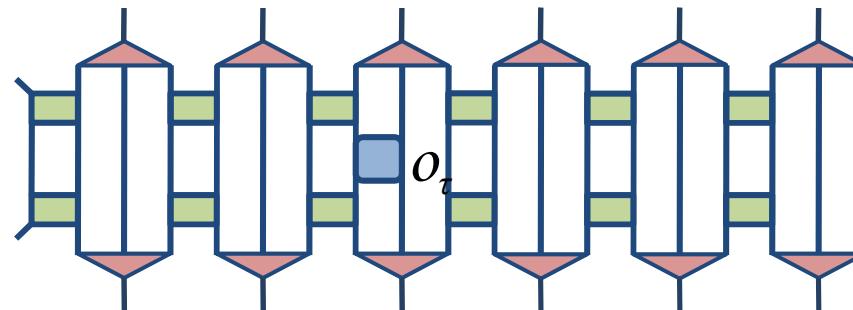
$$O_\tau \rightarrow O_{\tau+1}$$



Local operators

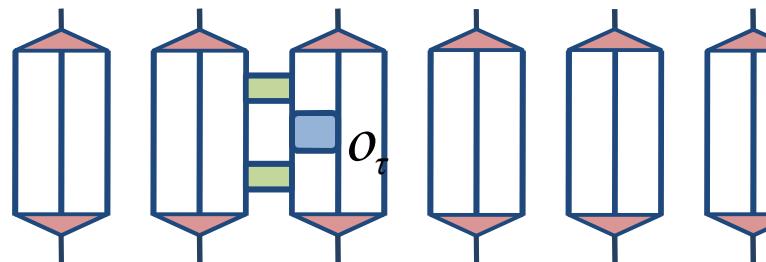
$$\mathcal{L}_\tau \rightarrow \mathcal{L}_{\tau+1}$$

$$O_\tau \rightarrow O_{\tau+1}$$



disentangler

$$u \quad u^\dagger = I$$

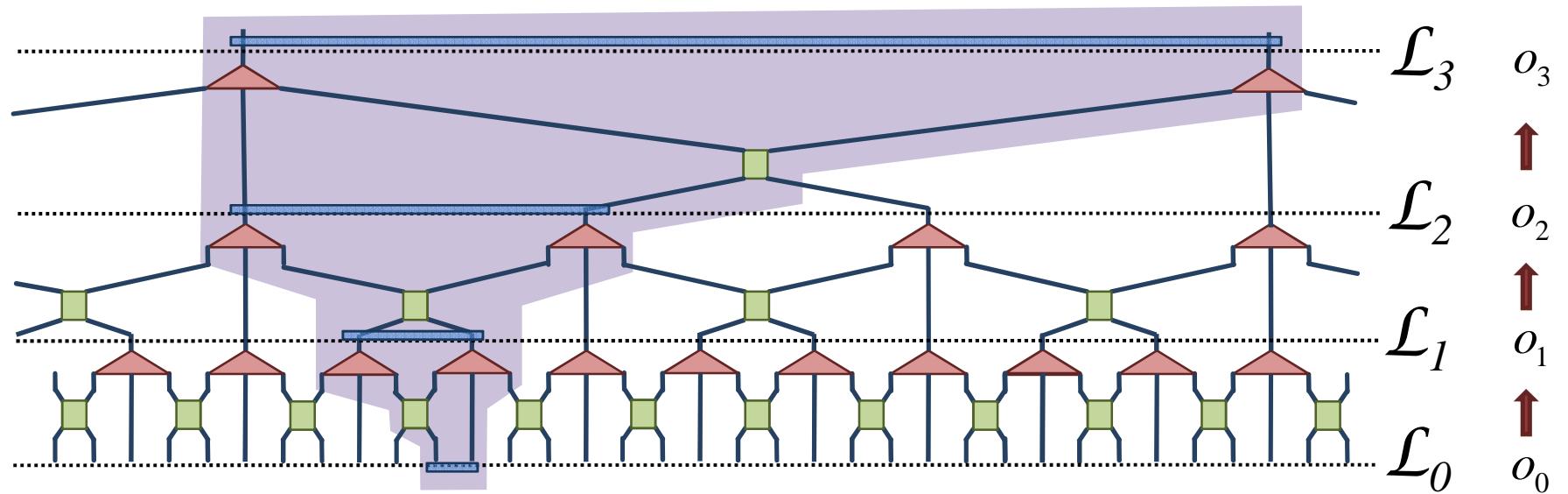


isometry

$$w \quad w^\dagger = I$$

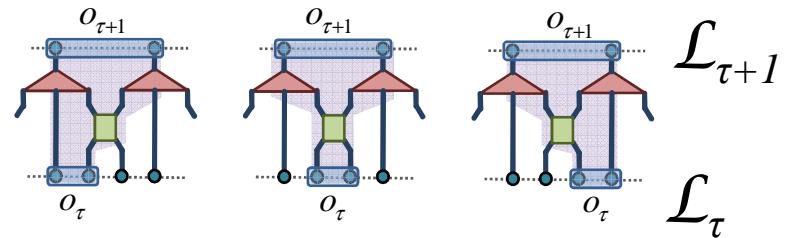
Ascending superoperator

- RG transformation



- ascending superoperator $o_{\tau+1} = \mathcal{A}(o_\tau)$

$$o_0 \rightarrow o_1 \rightarrow o_2 \rightarrow \dots$$



$o_{\tau+1} = \mathcal{A}_L(o_\tau)$

$o_{\tau+1} = \mathcal{A}_C(o_\tau)$

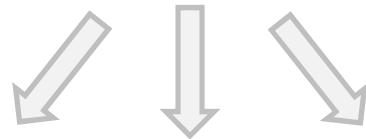
$o_{\tau+1} = \mathcal{A}_R(o_\tau)$

Foundations

Entanglement
Renormalization

MERA

(Multi-scale Entanglement
Renormalization Ansatz)



Applications

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- MERA \leftrightarrow CFT

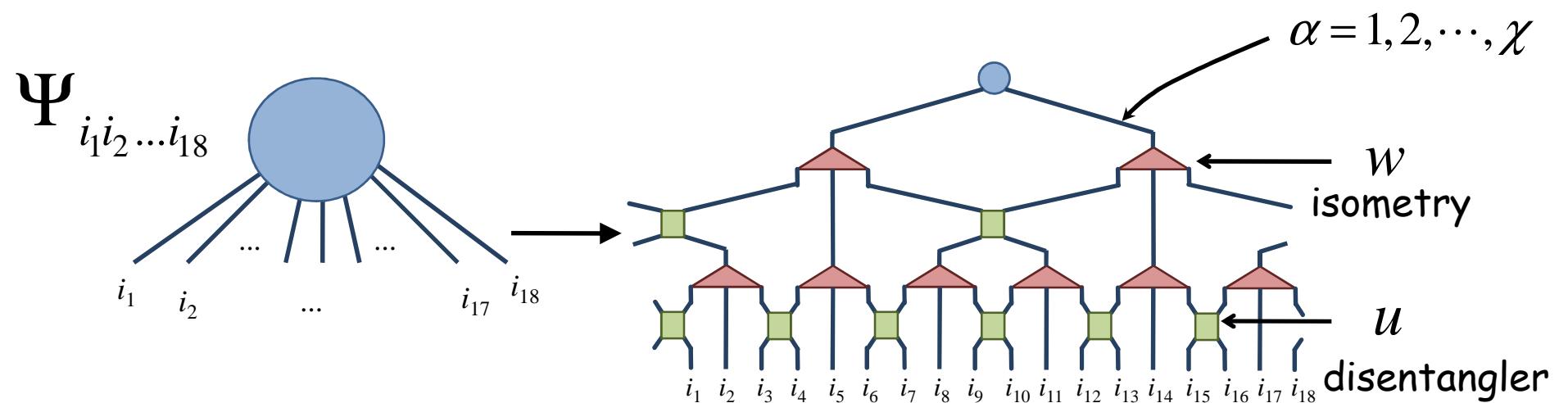
- Exact MERA
for quantum double
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- Kagome lattice
- Triangular lattice

MERA (Multi-scale entanglement renormalization ansatz)

Vidal, Phys. Rev. Lett. 101, 110501 (2008)

$$|\Psi\rangle = \sum_{i_1 i_2 \dots i_{18}} \Psi_{i_1 i_2 \dots i_{18}} |i_1 \otimes i_2 \otimes \dots \otimes i_{18}\rangle$$



MERA (Multi-scale entanglement renormalization ansatz)

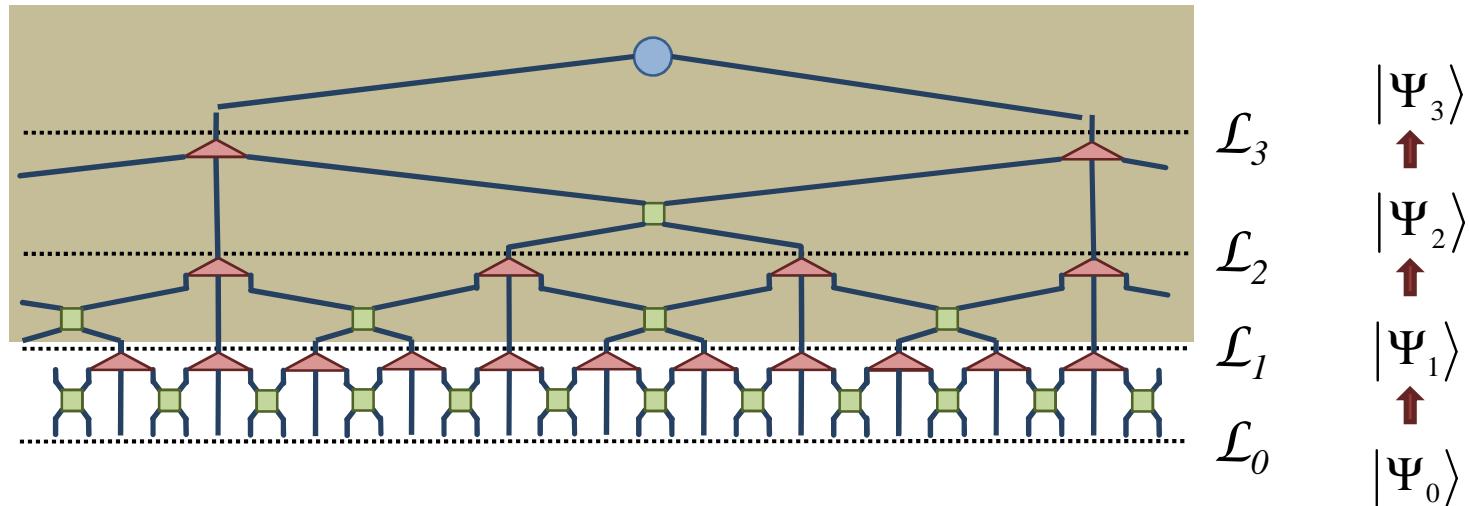
- Optimization algorithms

(Evenbly, Vidal, arXiv:0707.1454v3)

H local Hamiltonian

$$\min_{u,w} \langle \Psi | H | \Psi \rangle \Rightarrow (u, w)$$

- RG transformation for ground states

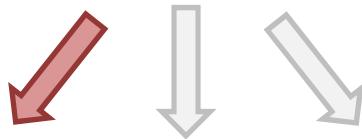


Foundations

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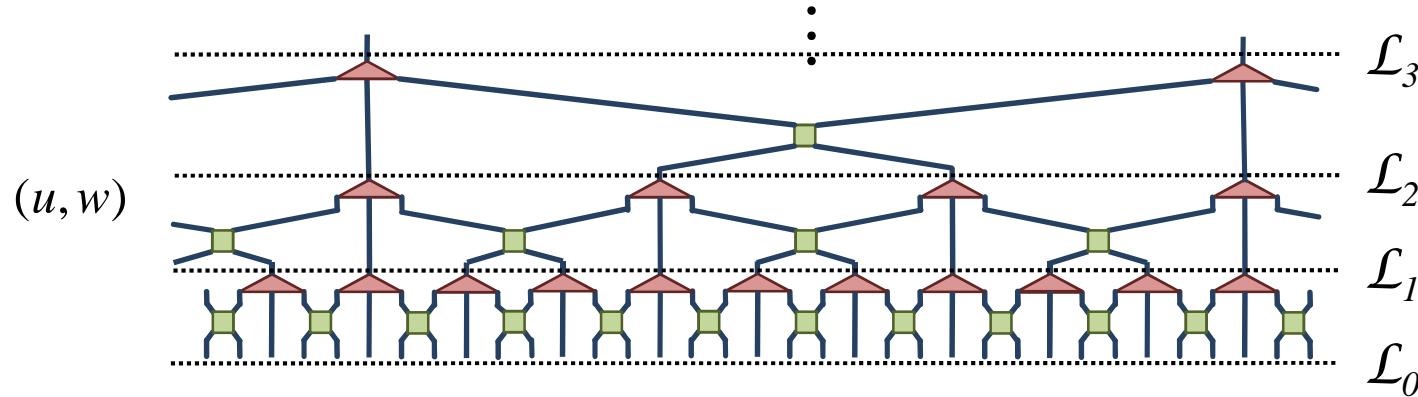
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Quantum Criticality

Pfeifer, Evenbly, Vidal, arXiv:0810.0580

- scale invariant MERA:



- scaling superoperator:

$$o_{\tau+1} = \mathcal{S}(o_\tau)$$
$$\text{---} = \frac{1}{3} \left\{ \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right\}$$

The diagram shows a scaling superoperator \mathcal{S} mapping an operator o_τ to $o_{\tau+1}$. It is represented by a sum of three terms, each involving a central tensor (blue square) and two vertical legs (yellow rectangles). The central tensor has three ports: top-left, top-right, and bottom. The first term has the top-left port connected to the bottom of the left leg and the top-right port connected to the bottom of the right leg. The second term has the top-left port connected to the top of the left leg and the top-right port connected to the top of the right leg. The third term has the top-left port connected to the top of the left leg and the top-right port connected to the bottom of the right leg.

Quantum Criticality

- scaling superoperator S is a fixed point RG map

What are the scaling operators/dimensions of the theory?

- scaling operators ϕ_α
- scaling dimensions Δ_α $\Delta_\alpha \equiv -\log_3 \lambda_\alpha$

$$\phi_\alpha \rightarrow 3^{-\Delta_\alpha} \phi_\alpha \rightarrow 3^{-2\Delta_\alpha} \phi_\alpha \rightarrow \dots$$

- From spectral decomposition of the scaling superoperator:

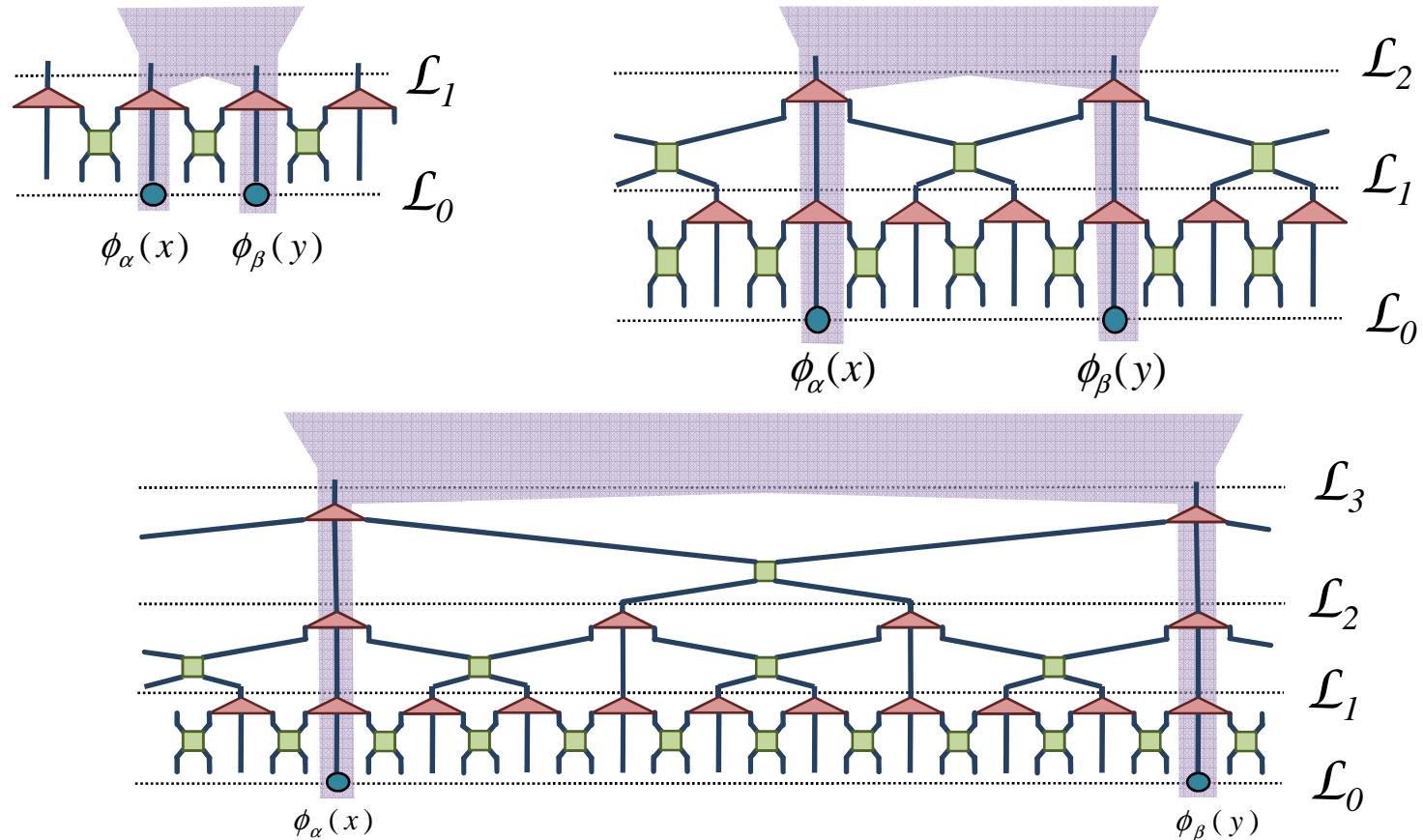
Giovannetti, Montangero, Fazio, Phys. Rev. Lett. 101, 180503 (2008)
Pfeifer, Evenbly, Vidal, arXiv:0810.0580

$$S(\phi_\alpha) = \lambda_\alpha \phi_\alpha$$

Quantum Criticality

Pfeifer, Evenbly, Vidal, arXiv:0810.0580

- local observables: $\langle \phi_\alpha(x) \rangle = \delta_{\alpha\mathbb{I}}$
- two-point correlator: $\langle \phi_\alpha(x) \phi_\alpha(y) \rangle = (\lambda_\alpha)^{2\log_3|x-y|} \langle \phi_\alpha(1) \phi_\alpha(0) \rangle = \frac{\langle \phi_\alpha(1) \phi_\alpha(0) \rangle}{|x-y|^{2\Delta_\alpha}}$



Quantum Criticality

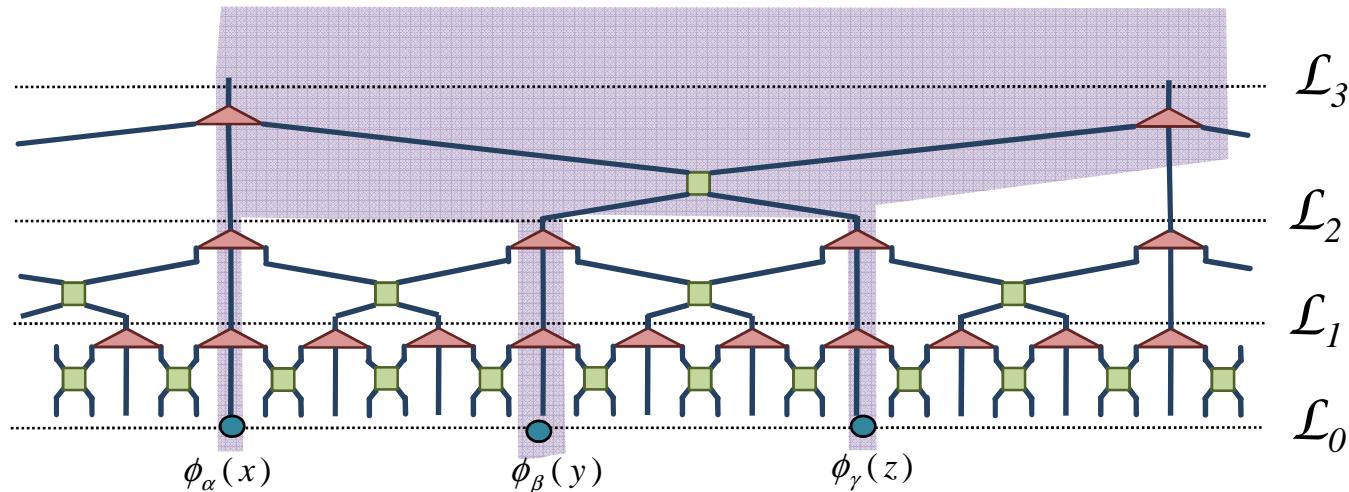
Pfeifer, Evenbly, Vidal, arXiv:0810.0580

- two-point correlator

$$\langle \phi_\alpha(x) \phi_\beta(y) \rangle = \frac{C_{\alpha\beta}}{|x - y|^{\Delta_\alpha + \Delta_\beta}}$$

$$C_{\alpha\beta} = \langle \phi_\alpha(0) \phi_\beta(1) \rangle \\ = \text{tr}(\rho^{(2)}(\phi_\alpha \otimes \phi_\beta))$$

- three-point correlator



$$\langle \phi_\alpha(x) \phi_\beta(y) \phi_\gamma(z) \rangle = \frac{C_{\alpha\beta\gamma}}{|x - y|^{\Delta_\alpha + \Delta_\beta - \Delta_\gamma} |y - z|^{\Delta_\beta + \Delta_\gamma - \Delta_\alpha} |z - x|^{\Delta_\gamma + \Delta_\alpha - \Delta_\beta}}$$

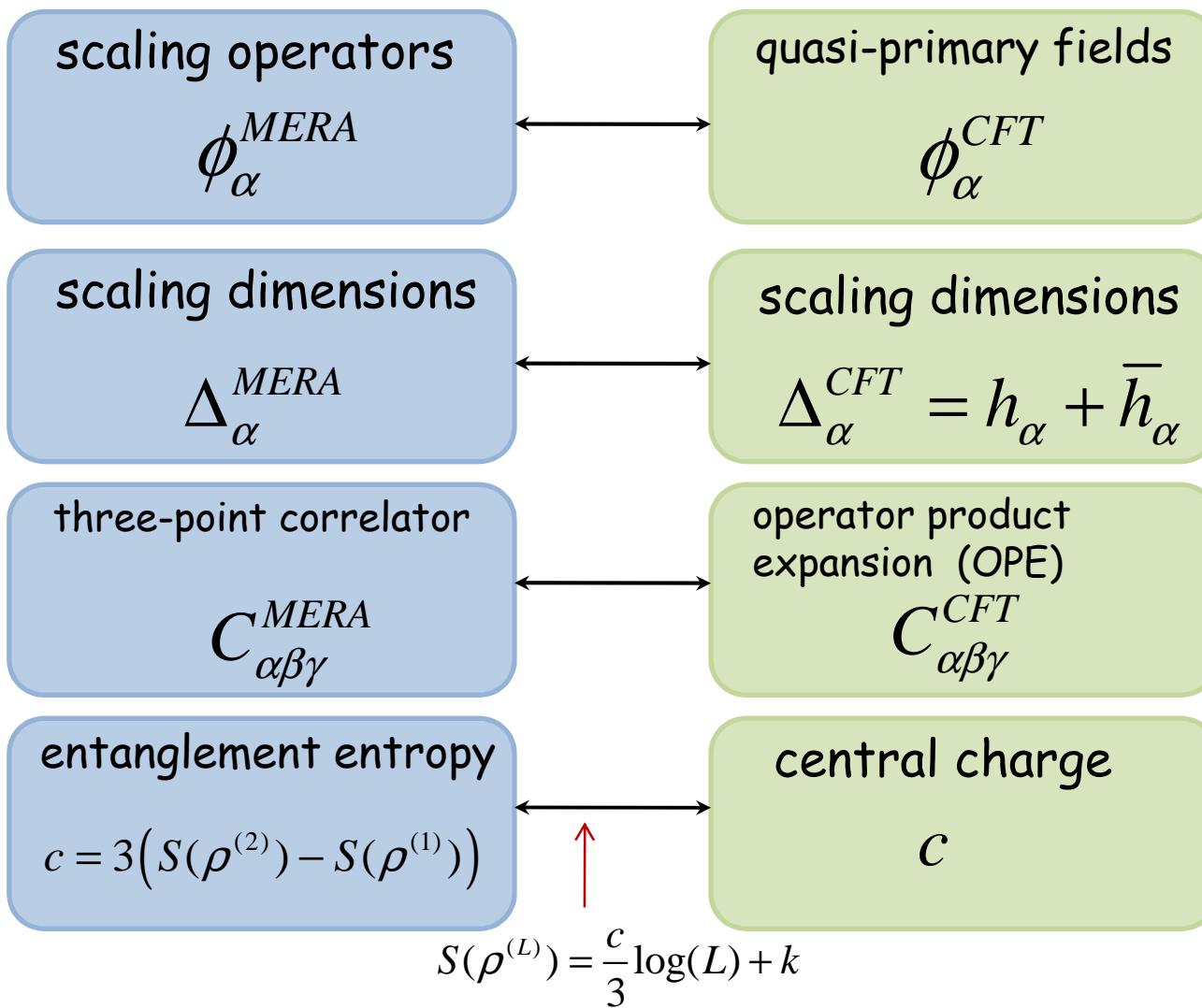
$$C_{\alpha\beta\gamma} = 2^{\Delta_\alpha - \Delta_\beta + \Delta_\gamma} \text{tr}(\rho^{(3)}(\phi_\alpha \otimes \phi_\beta \otimes \phi_\gamma))$$

Quantum Criticality

Pfeifer, Evenbly, Vidal, arXiv:0810.0580

scale invariant MERA
(lattice)

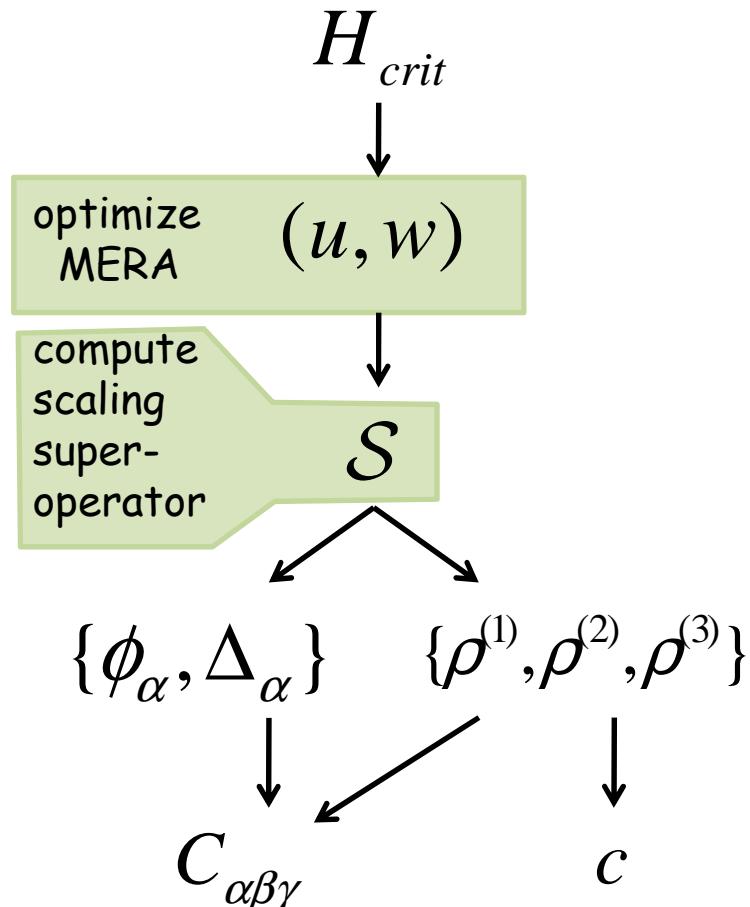
CFT
(continuum)



Quantum Criticality

Pfeifer, Evenbly, Vidal, arXiv:0810.0580

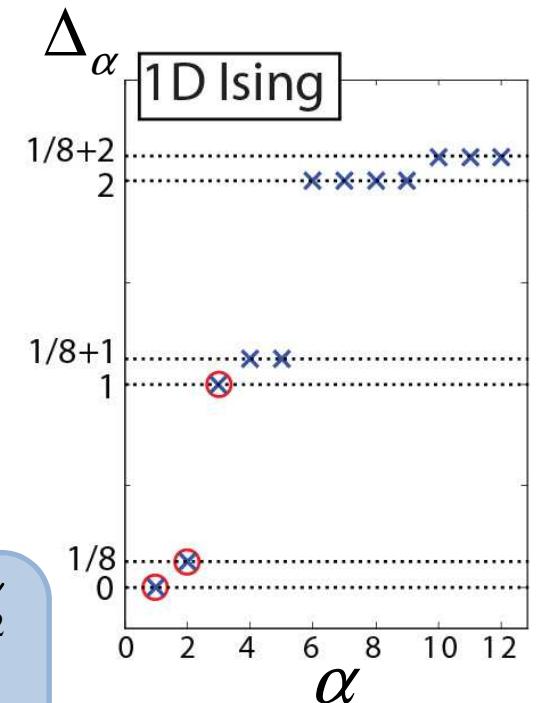
Given a critical hamiltonian on the lattice



example 1D quantum Ising:

$$H_{\text{Ising}} = \sum_i [\sigma_x^{[i]} \sigma_x^{[i+1]} + \sigma_z^{[i]}]$$

$$\begin{aligned} \Delta_{\mathbb{I}} &= 0 \\ \Delta_\sigma &= 0.124997 \quad (0.002\%) \\ \Delta_\varepsilon &= 1.0001 \quad (0.01\%) \end{aligned}$$



$$\begin{aligned} C_{\alpha\beta\mathbb{I}} &= \delta_{\alpha\beta} & C_{\sigma\sigma\varepsilon} &= \tfrac{1}{2} \\ C_{\sigma\varepsilon\varepsilon} &= C_{\sigma\sigma\sigma} = C_{\varepsilon\varepsilon\varepsilon} & &= 0 \\ &&&\quad (\pm 5 \times 10^{-4}) \end{aligned}$$

$$C_{\alpha\beta\gamma} = 2^{\Delta_\alpha - \Delta_\beta + \Delta_\gamma} \text{tr}(\rho^{(3)}(\phi_\alpha^{(1)} \otimes \phi_\beta^{(1)} \otimes \phi_\gamma^{(1)}))$$

$$c = 0.5008 \quad (0.15\%)$$

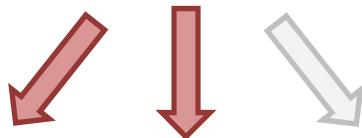
$$c = 3(S(\rho^{(2)}) - S(\rho^{(1)}))$$

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Topological Order

Exact MERA representation for

- quantum double models

Aguado, Vidal, Phys. Rev. Lett. 100, 070404 (2008)

- Levin-Wen's string net models

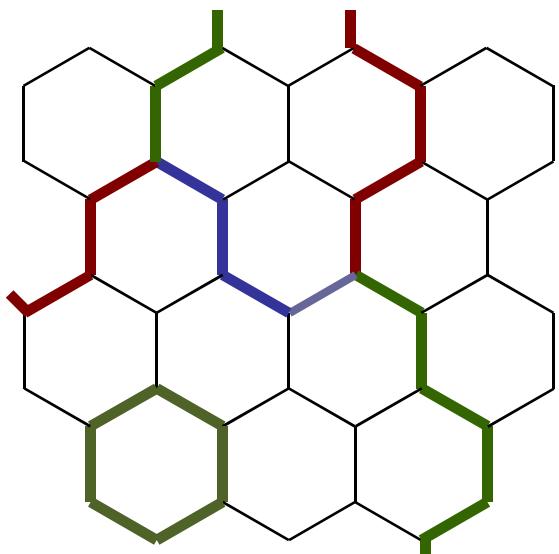
Koenig, Reichardt, Vidal, arXiv:0806.4583

Topological Order

Koenig, Reichardt, Vidal, arXiv:0806.4583

- Exact MERA representation for Levin-Wen's string net models

Levin, Wen, Phys. Rev. B71 (2005) 045110



- string types $i = \{1, 2, \dots, N\}$

- fusion rules δ_{ijk}

- 6-j symbols F_{klm}^{ijm}

- fixed-point Hamiltonian

$$H = - \sum_{\prec} A_{\prec} - \sum_{\circlearrowleft} B_{\circlearrowleft}$$

δ_{ijk} F_{klm}^{ijm}

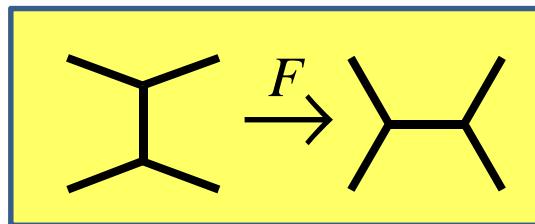
'physical'
subspace

Topological Order

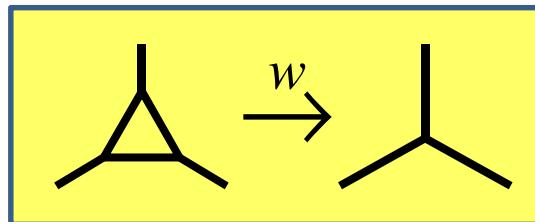
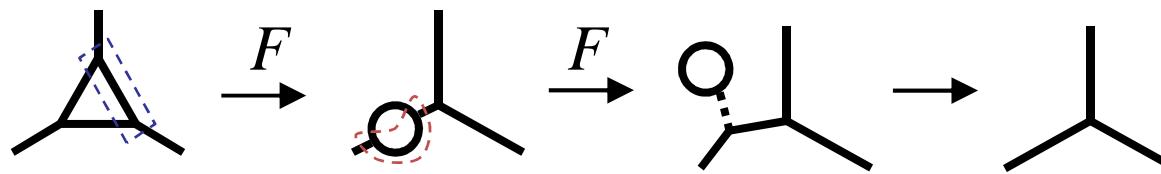
Koenig, Reichardt, Vidal, arXiv:0806.4583

- Exact MERA representation for Levin-Wen's string net models

The tensor $F_{k l n}^{i j m}$ can be used to define an "F-move"



Disentangler u ($= F$)

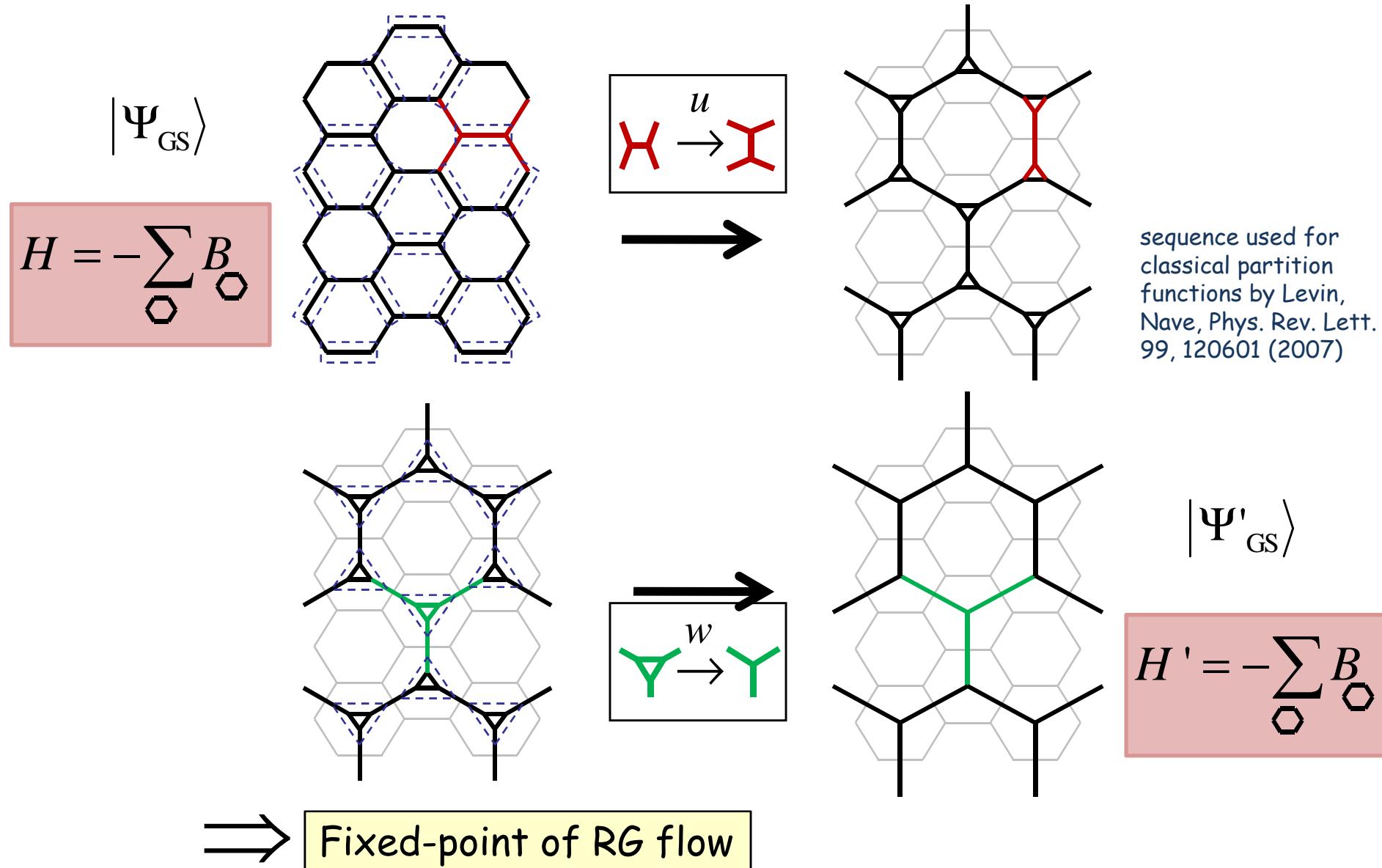


Isometry w

Topological Order

Koenig, Reichardt, Vidal, arXiv:0806.4583

- Exact MERA representation for Levin-Wen's string net models



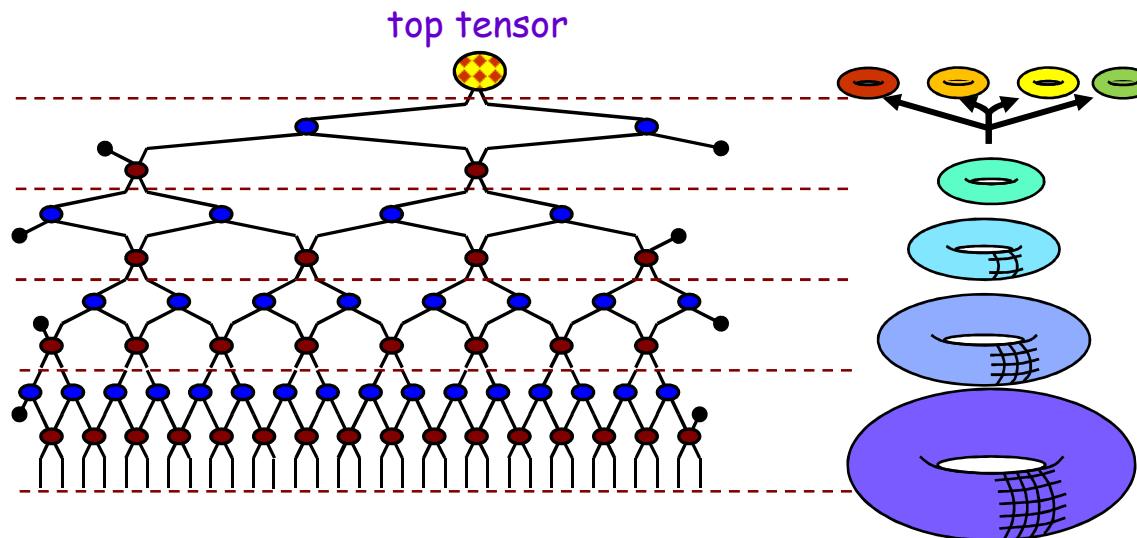
Topological Order

- MERA for ground states with topological order

degenerate ground states

$$|\Psi_{gr}^{(1)}\rangle \quad |\Psi_{gr}^{(2)}\rangle \quad |\Psi_{gr}^{(3)}\rangle \quad |\Psi_{gr}^{(4)}\rangle$$

Topological information
is stored in the top
tensor of the MERA



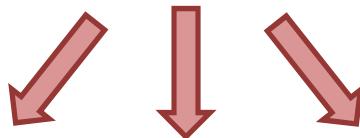
$$|\Psi_{gr}^{(1)}\rangle \quad |\Psi_{gr}^{(2)}\rangle \quad |\Psi_{gr}^{(3)}\rangle \quad |\Psi_{gr}^{(4)}\rangle$$

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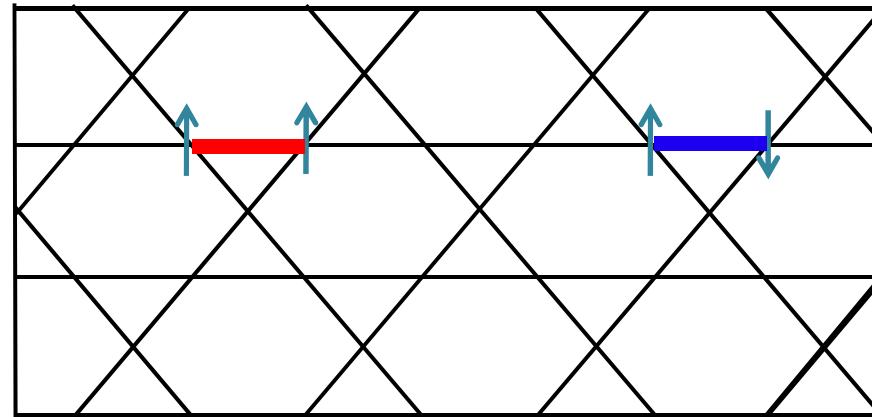
- Kagome lattice
- Triangular lattice

Frustrated Antiferromagnets

Evenbly, Vidal, in preparation



Glen
Evenbly

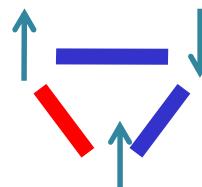


Hamiltonian: $H = \sum_{\langle i, j \rangle} S_i \cdot S_j$

- Favours antiferromagnetic alignment



- Geometrically frustrated



Frustrated Antiferromagnets

- AF Heisenberg Model on Kagome Lattice

What is the ground state? Open question...

- Valence Bond Crystal

Marston, Zeng (1991), Syromyatnikov, Maleyev (2002) Nikolic, Senthil (2003), Budnik, Auerbach (2004)

Rajiv R. P. Singh, David A. Huse, Phys. Rev. B **76**, 180407 (2007):
series expansion

- Spin Liquid

Hermele, Senthil, Fisher (2005)

Y. Ran, M. Hermele, P. A. Lee, X. Wen, Phys. Rev. Lett. **98**, 117205 (2007)
Gutzwiller ansatz

H. C. Jiang, Z. Y. Weng, D. N. Sheng, Phys. Rev. Lett. **101**, 117203 (2008)
DMRG

Frustrated Antiferromagnets

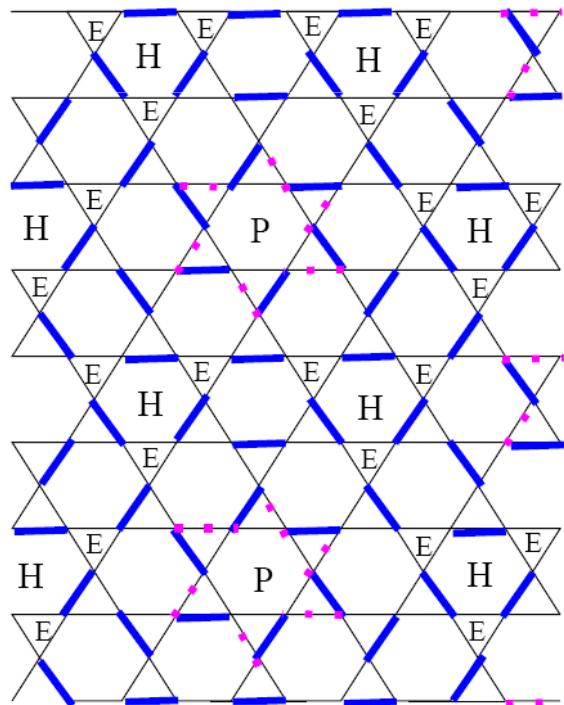
- AF Heisenberg Model on Kagome Lattice

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Rajiv R. P. Singh, David A. Huse, Phys. Rev. B **76**, 180407 (2007):
series expansion



Order	VBC (∞ system)	36-site VBC
0	-0.375	-0.375
1	-0.375	-0.375
2	-0.42187	-0.42187
3	-0.42578	-0.42578
4	-0.43155	-0.43400
5	-0.43208	-0.43624
∞	-0.433 ± 0.001	

Frustrated Antiferromagnets

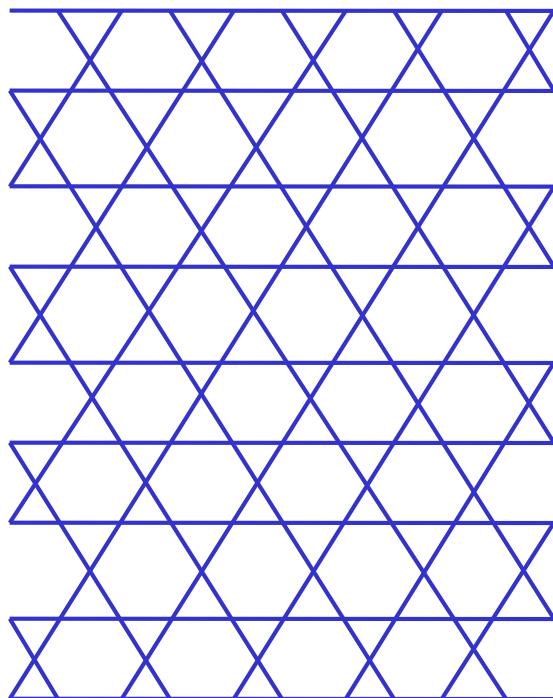
- AF Heisenberg Model on Kagome Lattice

H. C. Jiang, Z. Y. Weng, D. N. Sheng, Phys. Rev. Lett. **101**, 117203 (2008)

- Spin Liquid (DMRG)

Y. Ran, M. Hermele, P. A. Lee, X. Wen, Phys. Rev. Lett. **98**, 117205 (2007)

- Spin Liquid (Gutzwiller ansatz)



DMRG (Jiang)

N	E_0/N
48	-0.43591
108	-0.43111

Gutzwiller (Ran)

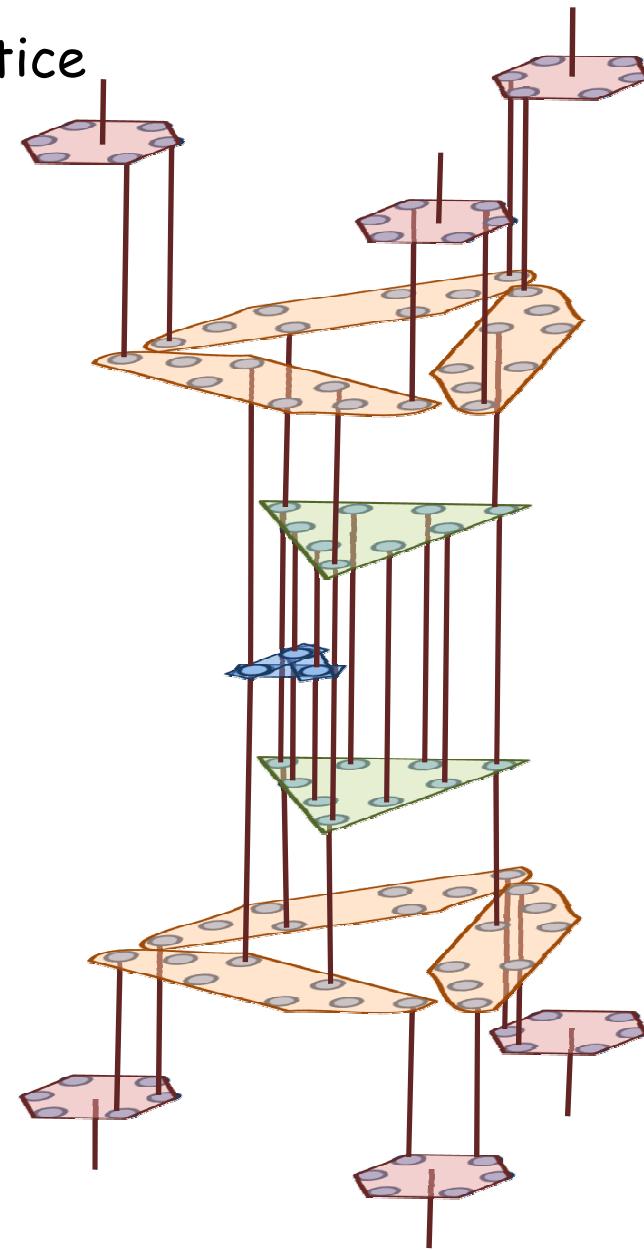
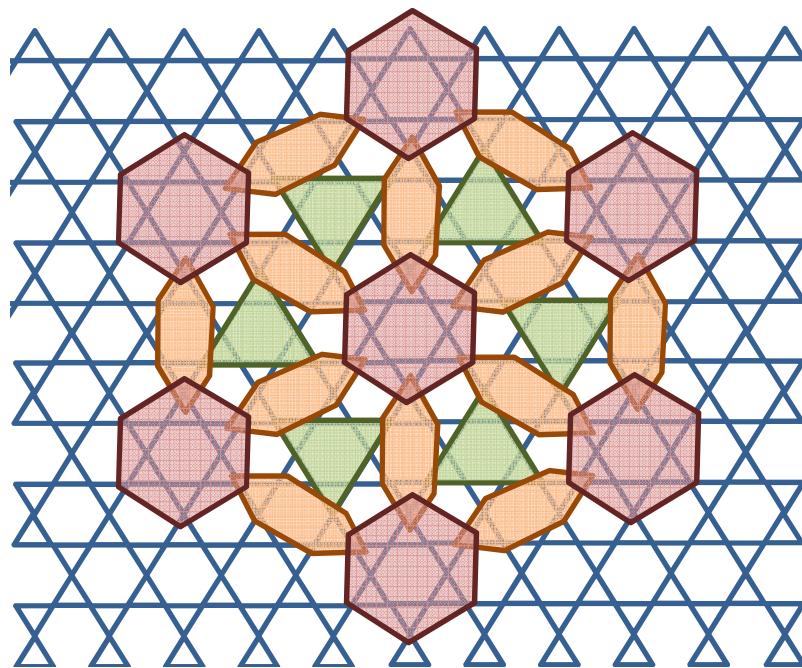
N	E_0/N
192	-0.42866(2)
432	-0.42863(2)

Frustrated Antiferromagnets

- AF Heisenberg Model on Kagome Lattice

MERA ?

Movie!!!



Frustrated Antiferromagnets

- AF Heisenberg Model on Kagome Lattice

Series Expansion (Singh)

Order	VBC (∞ system)	36-site VBC
0	-0.375	-0.375
1	-0.375	-0.375
2	-0.42187	-0.42187
3	-0.42578	-0.42578
4	-0.43155	-0.43400
5	-0.43208	-0.43624
∞	-0.433±0.001	

DMRG (Jiang)

N	E_0/N
48	-0.43591
108	-0.43111

Gutzwiller (Ran)

N	E_0/N
192	-0.42866(2)
432	-0.42863(2)

Entanglement Renormalization

N	E_0/N
∞	-0.4315

36-site Exact Diag (Leung)

N	E_0/N
36	-0.43837

Frustrated Antiferromagnets

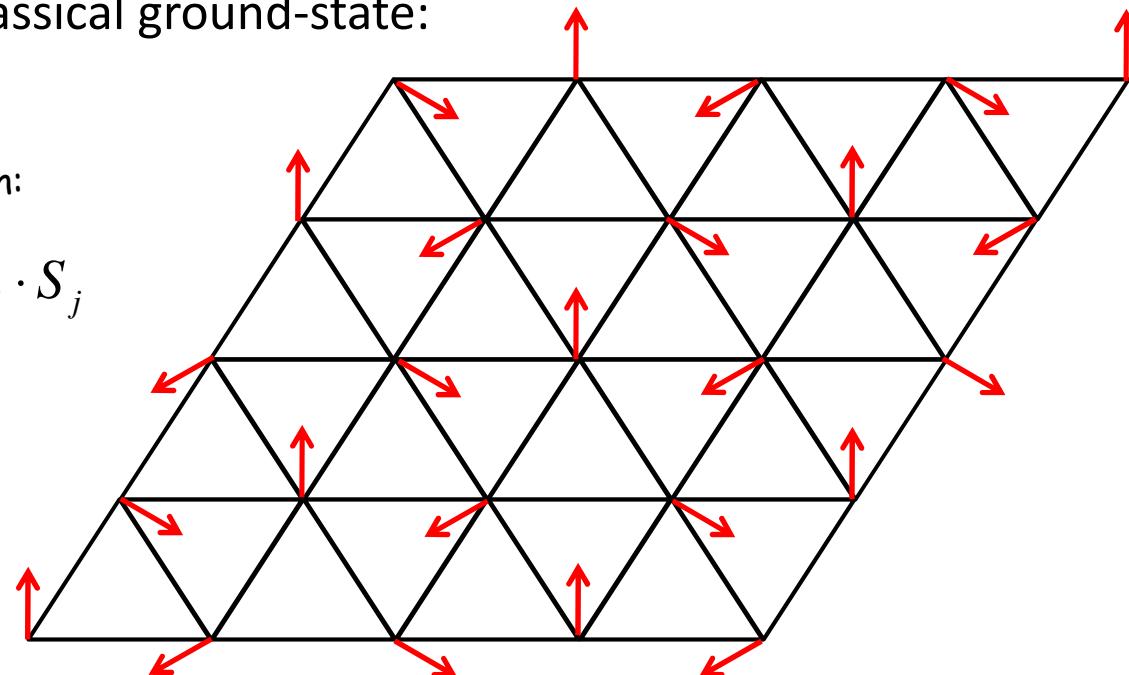
- AF Heisenberg Model on Triangular Lattice

Classical ground-state:

Hamiltonian:

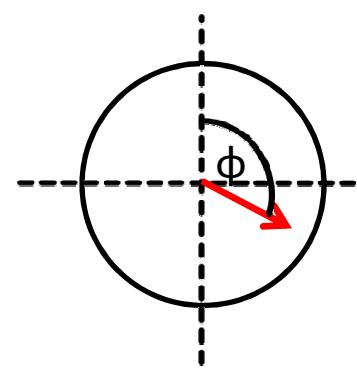
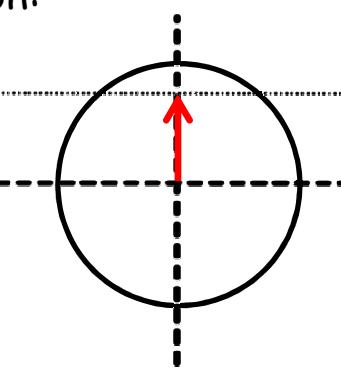
$$H = \sum_{\langle i,j \rangle} S_i \cdot S_j$$

MERA ?



Magnetization:

60%



$$\phi = 120^\circ \pm 0.1$$

Frustrated Antiferromagnets

- AF Heisenberg Model on Triangular Lattice

(from W. Zheng, J. Fjaerestad, R. Singh, R. McKenzie, R. Coldea, Phys. Rev. B **74**, 224420 (2006))

Method	E_0/N	Mag
DMRG	-0.5442	
GFQMC	-0.5458(1)	41%
VQMC, RVB	-0.5357	0%
Series Expansion	-0.5502(4)	39%
VQMC, BCS+Neel	-0.532(1)	72%
SWT + 1/S	-0.5466	50%
Entanglement Renormalization	-0.5451	60%

- ER gives ***exact upper bound*** for ground energy of the infinite system
(not based upon finite size scaling)

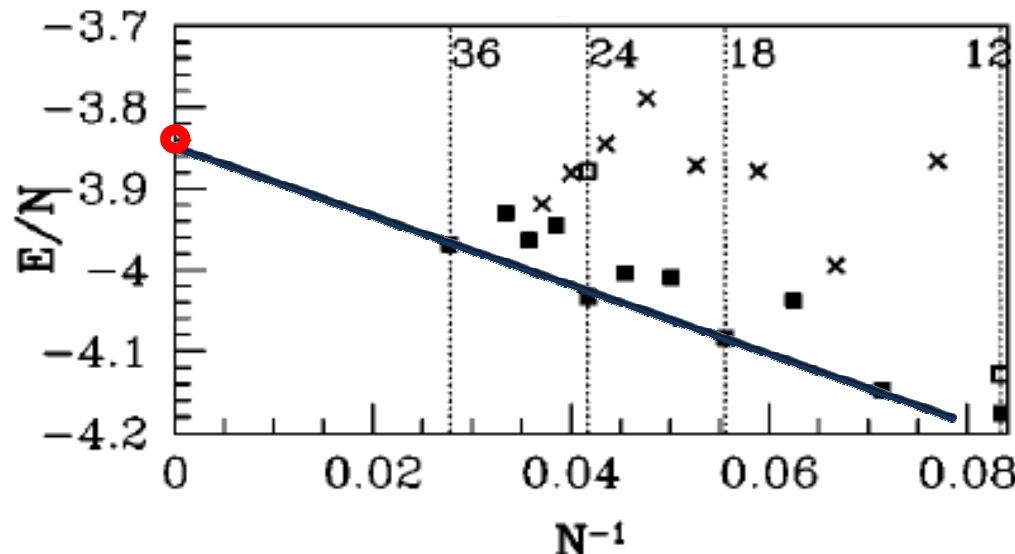
Frustrated Antiferromagnets

- Multiple spin exchange (MSE) on Triangular Lattice

$$H = J_2 \sum_{\text{---}} P_{ij} + J_4 \sum_{\square} (P_{ijkl} + P_{lkji})$$

- G. Misguich, C. Lhuillier, B. Bernu, and C. Waldtmann, Phys. Rev. B **60**, 1064 (1999)
- exact diagonalization (36 sites)

Spin Liquid phase ($J_2 = -2, J_4 = 1$)



Entanglement
Renormalization

N	E_0/N
∞	-3.84

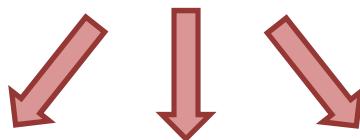
Summary

Foundations

Entanglement
Renormalization

MERA

(Multi-scale Entanglement
Renormalization Ansatz)



Applications

Quantum
Criticality

Topological
Order

Frustrated
Antiferromagnets

- MERA \leftrightarrow CFT

- Exact MERA
for quantum double
and string-net models

- Kagome lattice
- Triangular lattice

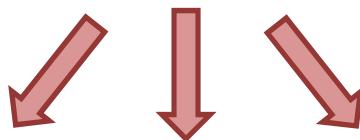
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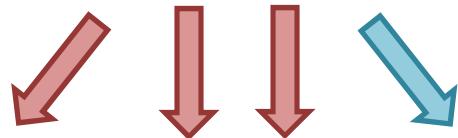
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Fermions ?

• MERA \leftrightarrow CFT

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• Kagome lattice
• Triangular lattice