

# Towards a unified description of ground and excited state properties: RPA vs GW

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Towards Reality in Nanoscale Materials V



Deutsche  
Forschungsgemeinschaft

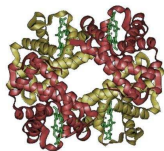
**DFG**



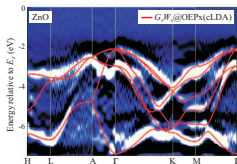
# Wish list for “optimum” electronic structure approach



$d/f$ -electrons  
(e.g. Cerium)



(bio)molecules



applicable across dimensionalities:

- molecules/clusters
- wires/tubes
- surfaces/films
- solids, extended systems

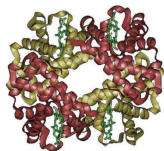
applicable across periodic table:

- from light to heavy elements
- including  $d/f$  electron physics/chemistry

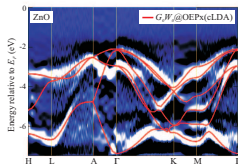
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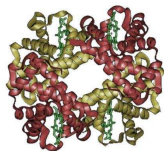
ground/excited state:

- consistent description
- parameter free
- free from pathologies e.g.:
  - ▶ self-interaction error
  - ▶ absence of van der Waals
  - ▶ band-gap problem
- computationally efficient
- gradients  $\Rightarrow$  structure relaxation
- ...

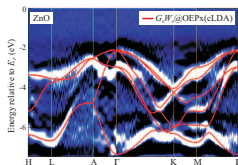
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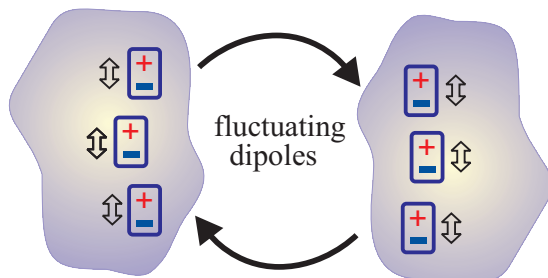
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# Treating van der Waals interactions



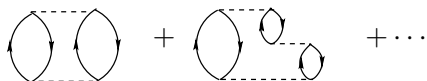
- difficult, because they require non-local treatment
- fluctuating dipoles described by dielectric function  $\epsilon(\mathbf{r}, \mathbf{r}', \omega)$
- correlation energy in random-phase approximation (RPA):

$$E_c^{\text{RPA}} = \frac{1}{2\pi} \int_0^\infty d\omega \text{Tr} [\ln(\epsilon(i\omega)) + (1 - \epsilon(i\omega))]$$

# Attractive features of the RPA

exact exchange + random phase approximation: (EX+cRPA)@reference

$$E_{\text{tot}}^{\text{EX+cRPA}} = T_s + E_{\text{ext}} + E_{\text{H}} + E_{\text{x}}^{\text{exact}} + E_{\text{c}}^{\text{RPA}}$$

$$E_{\text{c}}^{\text{RPA}} = \text{[diagram 1]} + \text{[diagram 2]} + \dots$$


- “Exact exchange” (with Kohn-Sham orbitals): OEPx
  - ▶ self-interaction error considerably reduced
- vdW interactions included automatically and seamlessly
- Screening taken into account
  - ▶ applicable to metals/small gap systems (in contrast to MP2)

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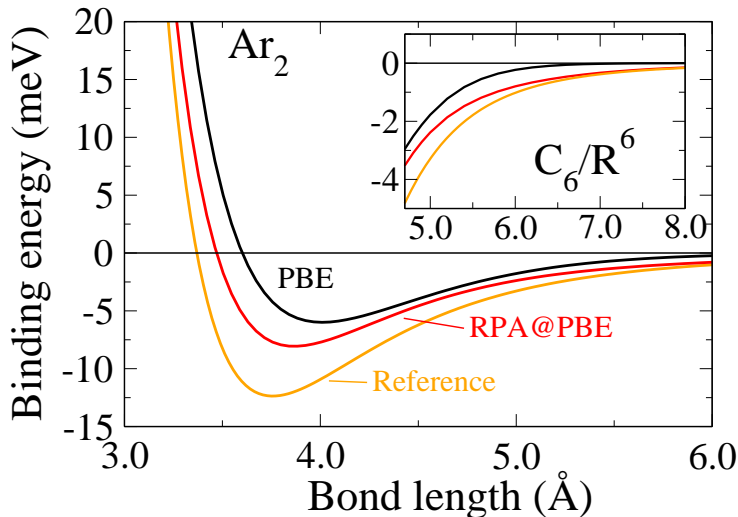
$$E_{\text{tot}}^{\text{EX+cRPA}} = T_s + E_{\text{ext}} + E_{\text{H}} + E_{\text{x}}^{\text{exact}} + E_{\text{c}}^{\text{RPA}}$$

$$E_{\text{c}}^{\text{RPA}} = \text{[diagram of diagrams]} + \dots$$

For an overview of RPA, see

- our recent review article:
  - “Exact Exchange + Random Phase Approximation” Ren, Joas, Rinke, Scheffler, arXiv:1203.5536v1
    - ▶ self-interaction error considerably reduced
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## van der Waals bonding



- ✓ Correct asymptotic behavior  $\Rightarrow$  crucial for large molecules
- ✗ Underbinding around the equilibrium distance



# The SOSEX correction to RPA

Digrammatic representation (motivated in coupled cluster context)

$$E_C^{\text{RPA+SOSEX}} =$$

The diagrammatic expansion is as follows:

- Row 1:  $E_C^{\text{RPA+SOSEX}} =$  (two bubbles) + (two bubbles with a vertex correction) + ...
- Row 2: + (2nd-order term) + (3rd-order term) + ...
- Row 3: = (two bubbles) + (two bubbles with a vertex correction)

- Arising from the anti-symmetric nature of the many-body wave function
- RPA+SOSEX is one-electron self-correlation free

D. L. Freeman, Phys. Rev. B **15**, 5512 (1977). A. Grüneis *et al.*, J. Chem. Phys. **131**, 154115 (2009). J. Paier *et al.*, J. Chem. Phys. **132**, 094103 (2010); Erratum: **133**, 179902 (2010).

# The concept of single excitation (SE) corrections

Rayleigh-Schrödinger perturbation theory:

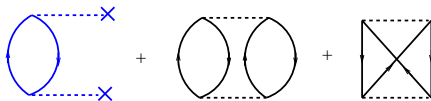
$$\hat{H} = \hat{H}^0 + \hat{H}'$$

$$\text{Ground-state energy: } E_0 = E_0^{(0)} + E_0^{(1)} + E_0^{(2)} + \dots$$

- zeroth-order:  $E_0^{(0)} = \langle \Phi_0 | \hat{H}^0 | \Phi_0 \rangle$
- 1st order:  $E_0^{(1)} = \langle \Phi_0 | \hat{H}' | \Phi_0 \rangle$
- 2nd order:

$$E_0^{(2)} = \sum_{n \neq 0} \frac{|\langle \Phi_0 | \hat{H}' | \Phi_n \rangle|^2}{E_0^{(0)} - E_n^{(0)}} = \underbrace{\sum_{i,a} \frac{|\langle \Phi_0 | \hat{H}' | \Phi_{i,a} \rangle|^2}{E_0^{(0)} - E_{i,a}^{(0)}}}_{\substack{\text{Single excitations} \\ = E_c^{\text{SE}}}} + \underbrace{\sum_{ij,ab} \frac{|\langle \Phi_0 | \hat{H}' | \Phi_{ij,ab} \rangle|^2}{E_0^{(0)} - E_{ij,ab}^{(0)}}}_{\substack{\text{Double excitations} \\ \text{MP2}}}$$

SE accounts for orbital relaxations



# Renormalized single excitations (rSE)

$$E_c^{\text{rSE}} = \text{2nd-order} + \text{3rd-order} + \dots$$

$f_{pq} = \langle p | \hat{f} | q \rangle$  : Hartree-Fock operator evaluated within Kohn-Sham orbitals

- $$E_c^{\text{SE}} = \sum_i^{\text{occ}} \sum_a^{\text{unocc}} \frac{|f_{ia}|^2}{\epsilon_i - \epsilon_a} \quad (\epsilon_i, \epsilon_a : \text{Kohn-Sham orbital energies})$$

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- “Diagonal rSE” : including only terms with  $i = j = \dots$  and  $a = b = \dots$

$$E_c^{rSE\text{-diag}} = \sum_i^{\text{occ}} \sum_a^{\text{unocc}} \frac{|f_{ia}|^2}{f_{ii} - f_{aa}} \quad \text{Problematic for weak interactions}$$

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$$E_c^{\text{rSE-diag}} = \sum_i^{\text{occ}} \sum_a^{\text{unocc}} \frac{|f_{ia}|^2}{f_{ii} - f_{aa}} \quad \text{Problematic for weak interactions}$$

- “full rSE” :  $E_c^{\text{rSE-full}} = \sum_i^{\text{occ}} \sum_a^{\text{unocc}} \frac{|\tilde{f}_{ia}|^2}{\tilde{f}_i - \tilde{f}_a}$

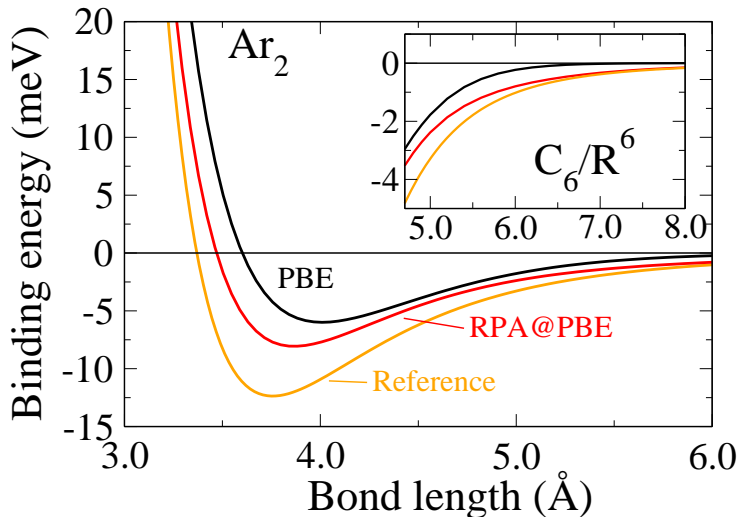
Diagonalize  $f_{ij}$  and  $f_{ab}$  blocks separately  $\rightarrow \tilde{f}_i, \tilde{f}_a$ , and transform  $f_{ia} \rightarrow \tilde{f}_{ia}$

# Renormalized 2nd-order Perturbation Theory (r2PT)

$$\begin{aligned}
 E_c^{\text{RPA+SOSEX+rSE}} = & \quad \text{[RPA diagrams]} + \dots \quad (= \text{RPA}) \\
 & + \quad \text{[SOSEX diagrams]} + \dots \quad (= \text{SOSEX}) \\
 & + \quad \text{[rSE diagrams]} + \dots \quad (= \text{rSE}) \\
 & \quad \quad \quad \text{2nd-order} \qquad \qquad \quad \text{3rd-order} \\
 & \quad \quad \quad (= \text{2PT})
 \end{aligned}$$

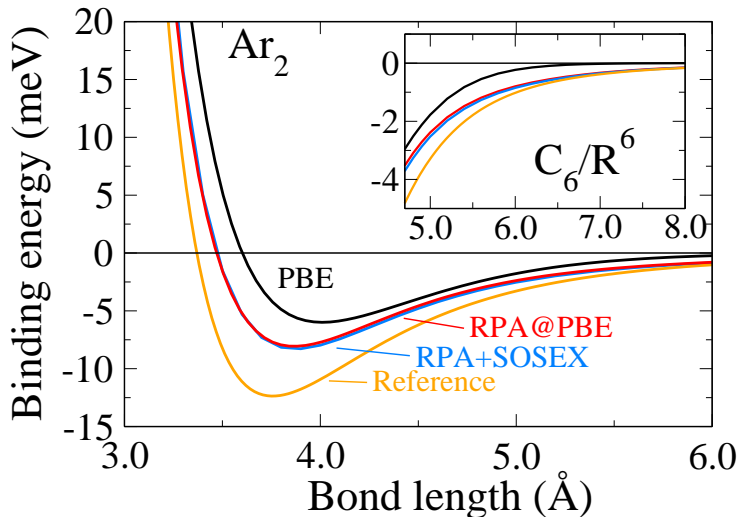
“r2PT” = “RPA+SOSEX+rSE”

## van der Waals bonding



- ✓ Correct asymptotic behavior  $\Rightarrow$  crucial for large molecules
- ✗ Underbinding around the equilibrium distance

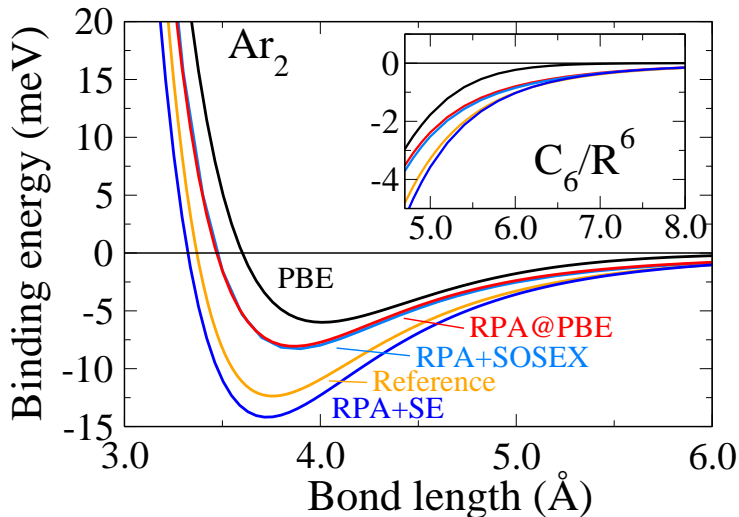
## van der Waals bonding



- SOSEX has almost no effect



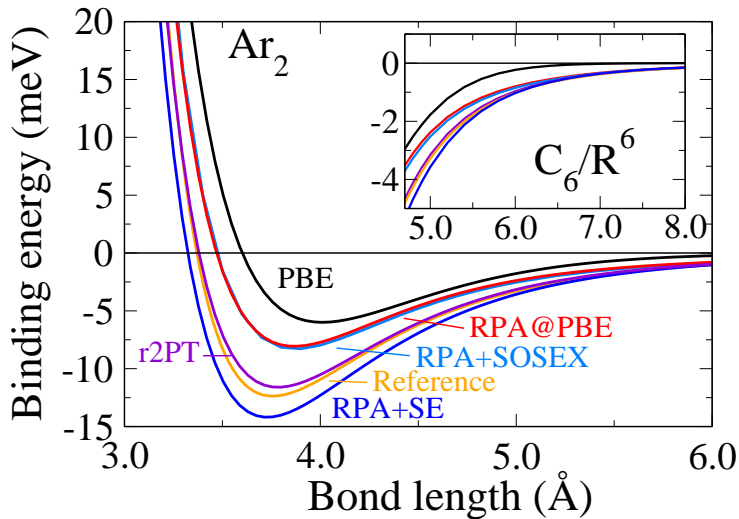
## van der Waals bonding



- much improved binding and asymptotics

Ren, Tkatchenko, Rinke, and Scheffler, Phys. Rev. Lett. 106, 153003 (2011)

## van der Waals bonding

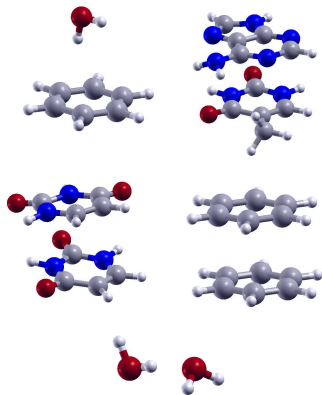


- best overall performance

# Van der Waals interactions - the S22 benchmark set

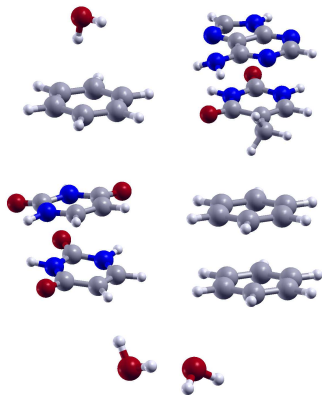
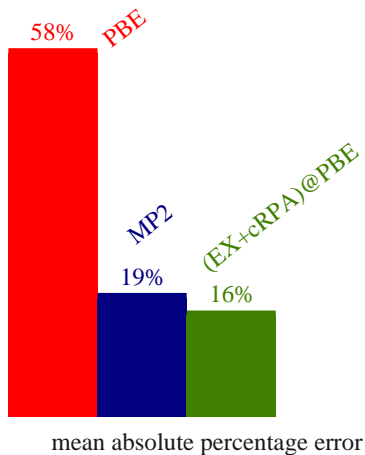
## S22 set of binding energies

- 22 representative dimers:
  - ▶ 7 hydrogen bonded
  - ▶ 8 dispersion bonded
  - ▶ 7 mixed character
- CCSD(T) (“gold standard”) reference values

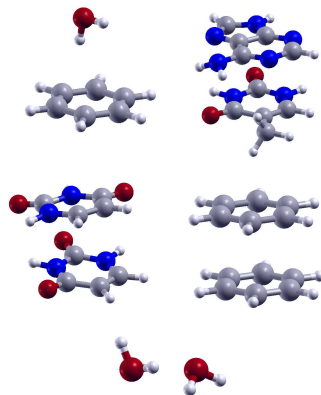
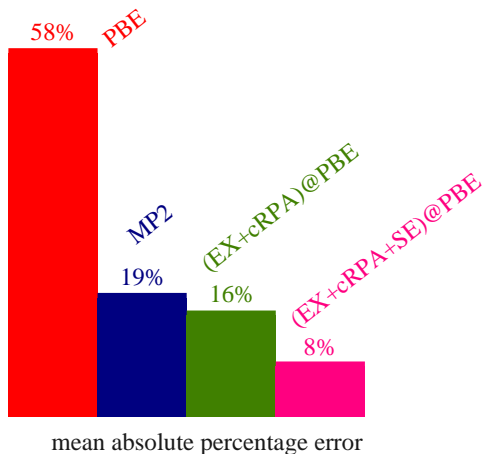


Jurečka *et al.*, PCCP **8**,1985 (2006)

# Van der Waals interactions - the S22 benchmark set

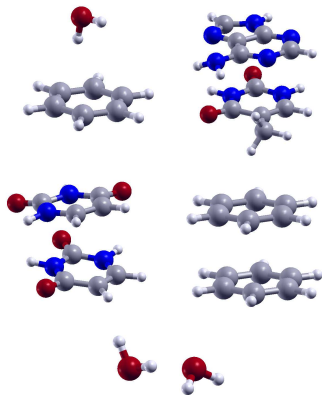
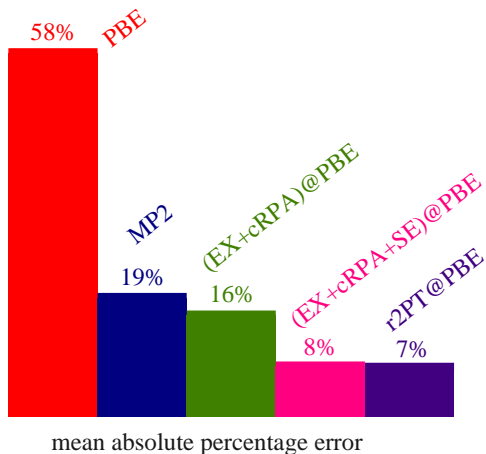


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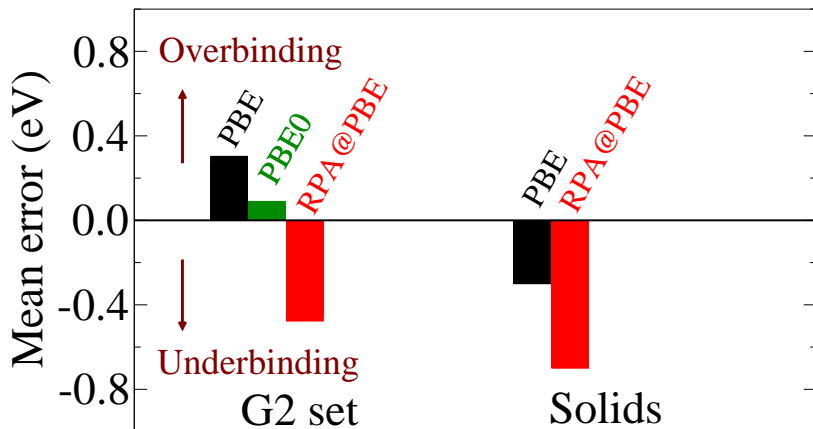
X. Ren, P. Rinke, A. Tkatchenko, M. Scheffler,  
Phys. Rev. Lett. **106**, 153003 (2011)

# Van der Waals interactions - the S22 benchmark set



X. Ren, P. Rinke, A. Tkatchenko, M. Scheffler,  
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## Atomization energies

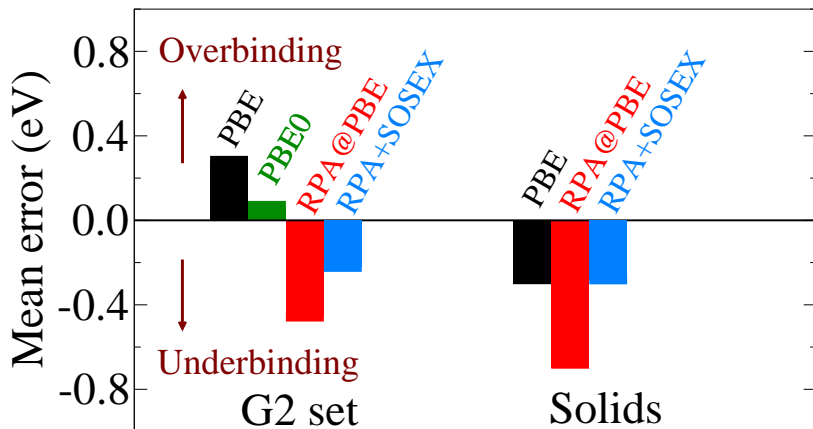


- RPA underestimates bond strengths

Paier *et al.*, J. Chem. Phys. **132**, 094103 (2010); **133**, 179902 (2010)

Harl, Schimka, and Kresse, Phys. Rev. B **81**, 115126 (2010)

## Atomization energies



- SOSEX alleviates the underbinding problem

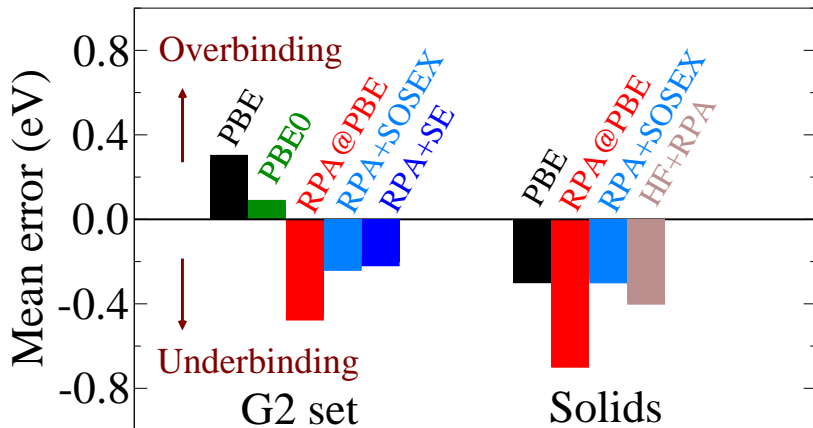
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Paier *et al.*, arXiv:1111.0173



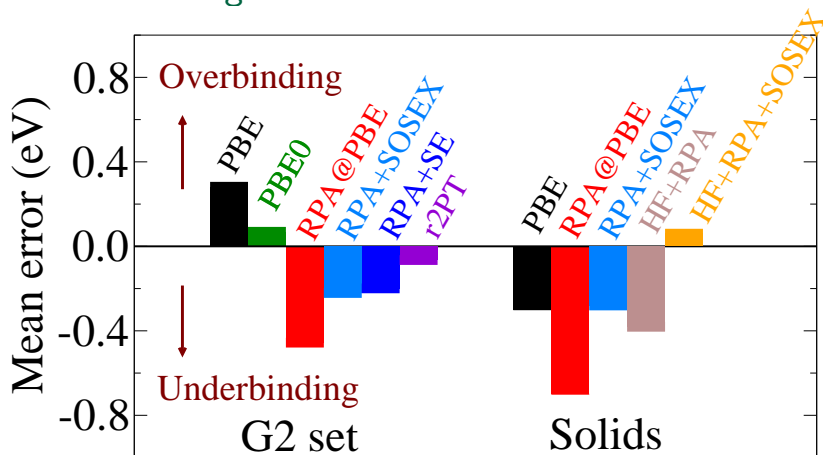
# Atomization energies



- “HF+RPA”: EX@HF + RPA@PBE similar to  
“RPA+SE”: (EX+RPA+SE)@PBE

Paier, Ren, Rinke, Scuseria, Grüneis, Kresse, and Scheffler, NJP **14**, 043002 (2012)  
Ren, Tkatchenko, Rinke, and Scheffler, Phys. Rev. Lett. **106**, 153003 (2011)

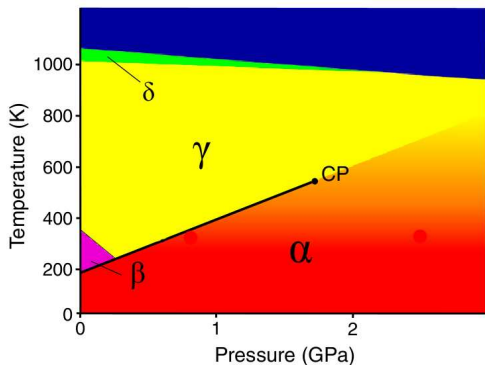
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# The $f$ -electron system Cerium

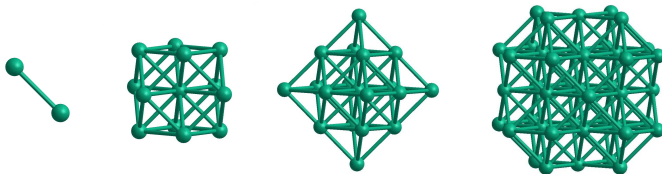


- iso-structural (fcc-fcc)  $\alpha$ - $\gamma$  phase transition
- accompanied by large (15%-17%) volume collapse

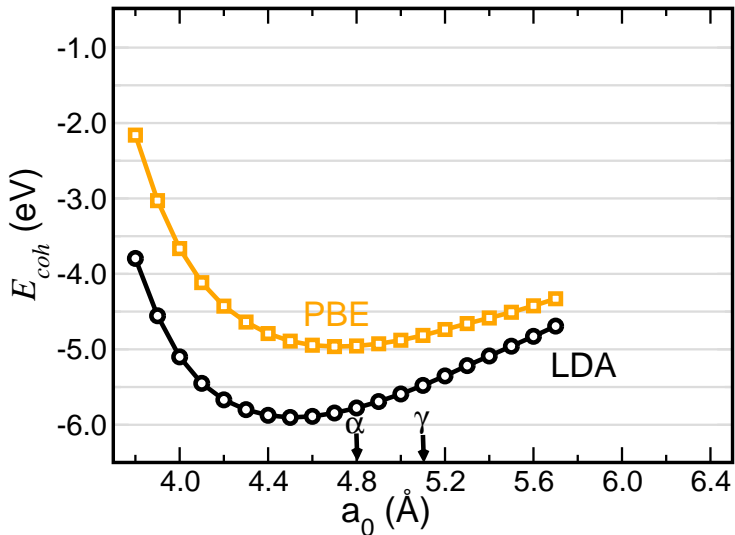
- LDA/GGA capture only low-volume ( $\alpha$ ) phase
- LDA+ $U$ /SIC-LDA capture only high-volume ( $\gamma$ ) phase
- LDA+DMFT so far restricted to high temperature

# Cerium from first principles

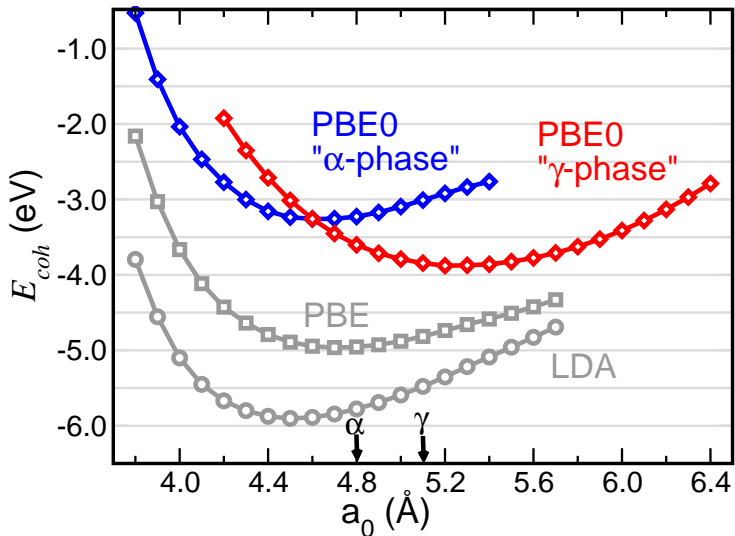
- All electrons are treated on the same quantum mechanical level
- hybrid functional PBE0
  - ▶ 25% exact exchange
  - ▶ reduces self-interaction error
- (EX+cRPA)@PBE0
  - ▶ physically meaningful screening
  - ▶ compatible with  $GW$  approach
  - ▶ cluster extrapolation approach



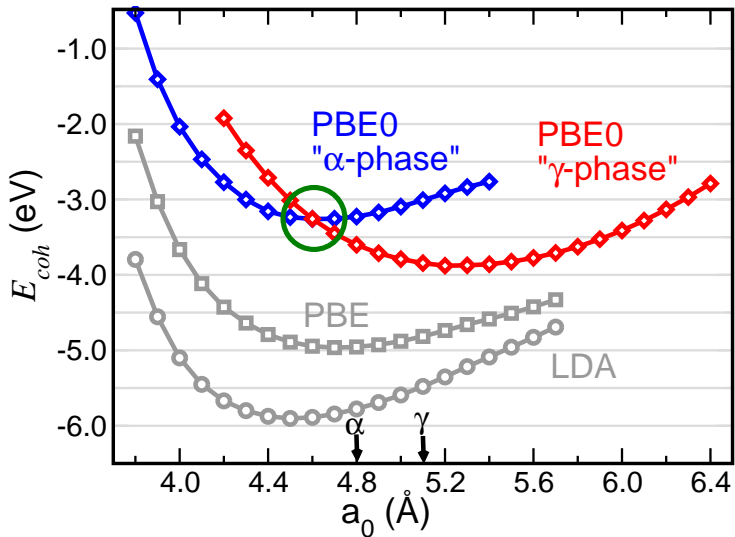
# Cerium - cohesive energy



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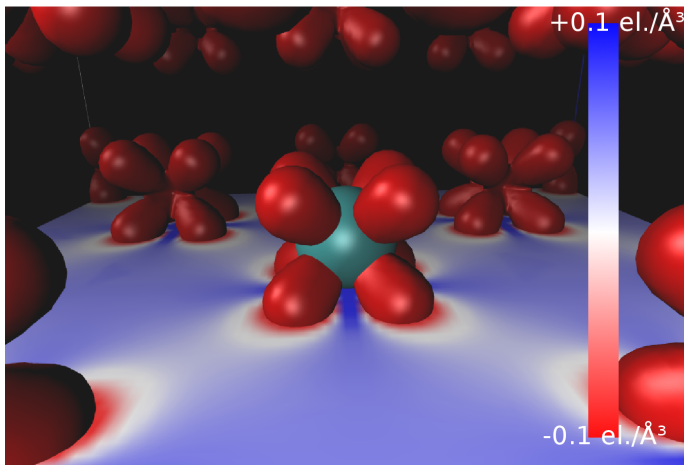


# Cerium - cohesive energy



# Cerium - electronic structure

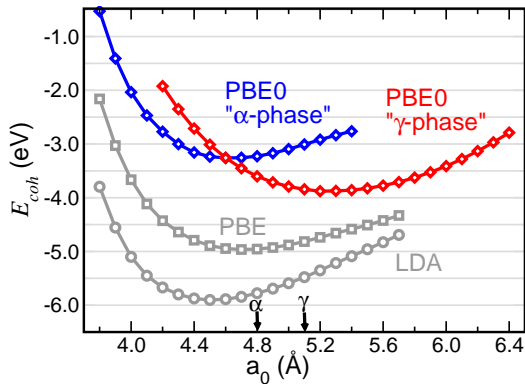
$n(\alpha) - n(\gamma)$  at  $a=4.6\text{\AA}$



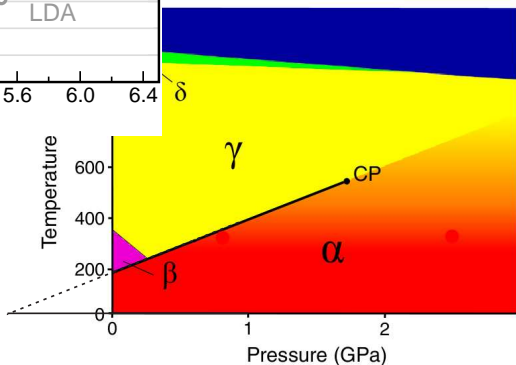
- $f$ -electrons more localized in  $\gamma$ -phase



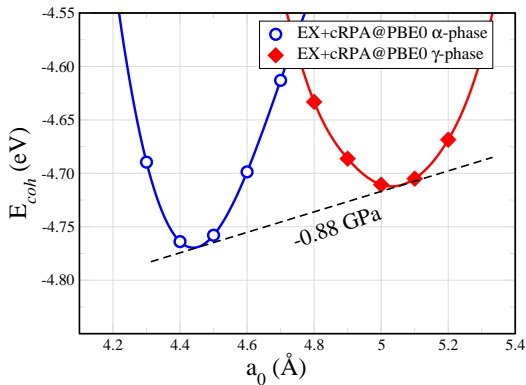
# Cerium - cohesive energy



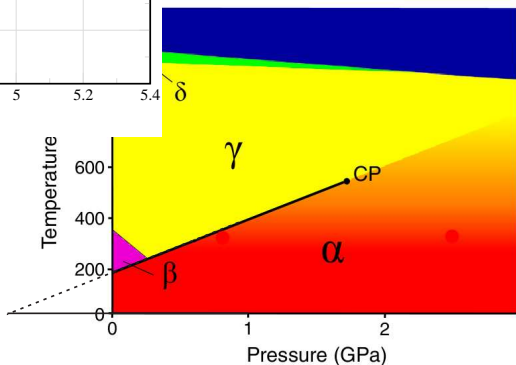
● wrong phase ordering



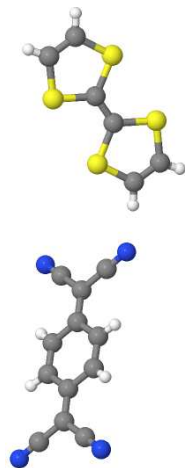
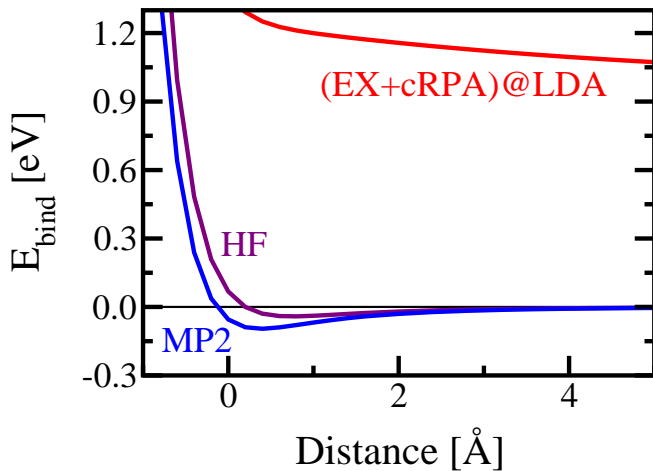
# Cerium - cohesive energy



- correct phase ordering
- results for a  $Ce_{19}$  cluster



# TTF/TCNQ dimer at infinite separation



# Electron gas at TTF/TCNQ interface

- TTF and TCNQ crystals have large band gap
- **but** 2 dimensional electron gas observed at interface

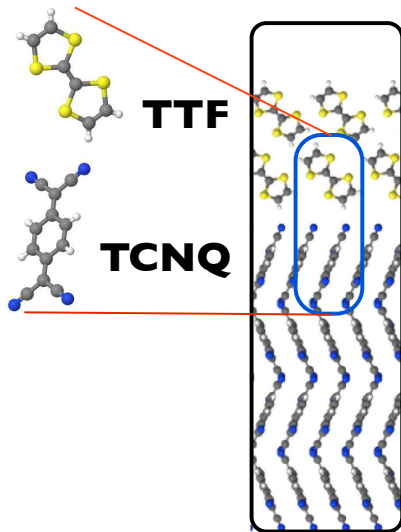
Metallic conduction at organic charge-transfer interfaces

HELENA ALVES, ANNA S. MOLINARI, HANGXING XIE AND ALBERTO F. MORPURGO

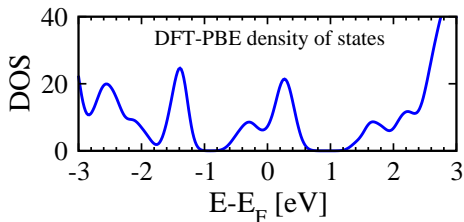
Kavli Institute of Nanoscience, Delft University of Technology, Lorentzweg 1, 2628 CJ Delft, The Netherlands

Nat. Mat. 7, 574 (2008)

What is the origin?

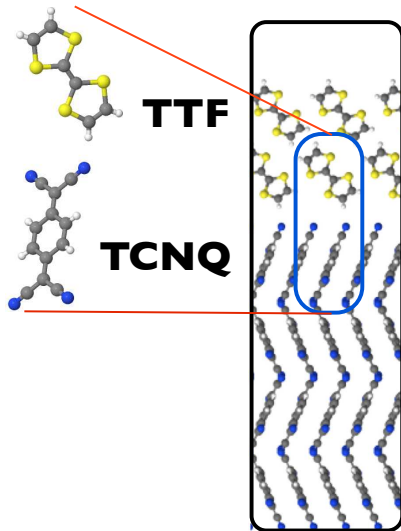


# Electron gas at TTF/TCNQ interface

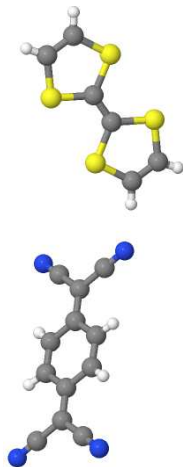
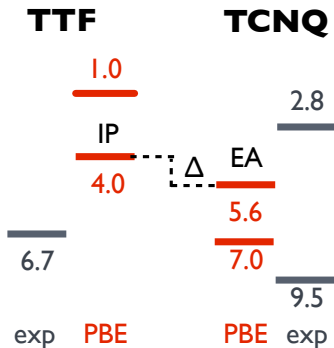


DFT-PBE:

- predicts metallic interface
- 0.12 e/molecule charge transfer to TCNQ
- in seeming agreement with experiment



# TTF/TCNQ dimer at infinite separation



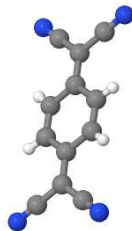
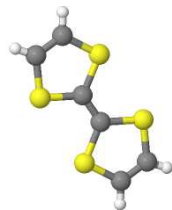
- $IP(TTF) > EA(TCNQ)$
- erroneous charge transfer

# TTF/TCNQ dimer at infinite separation

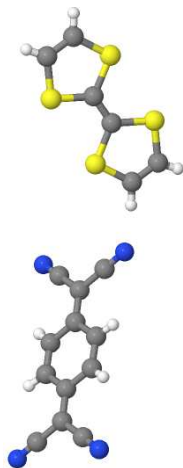
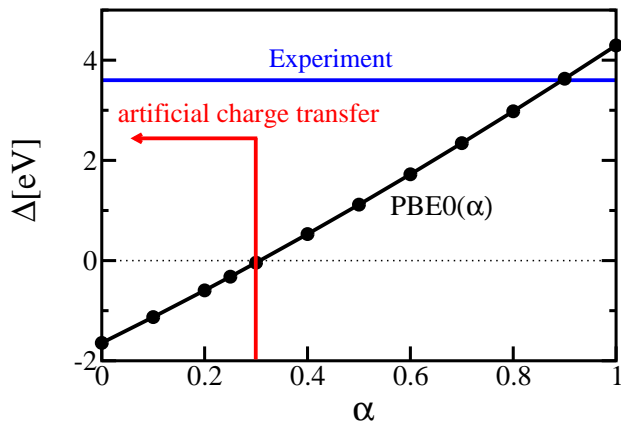
- culprit: self-interaction error in PBE
- add Hartree-Fock (HF) exchange to remove self-interaction

⇒ PBE0-like hybrid functional:

$$E_{xc}(\alpha) = \alpha(E_x^{\text{HF}} - E_x^{\text{PBE}}) + E_c^{\text{PBE}}$$



# TTF/TCNQ dimer at infinite separation

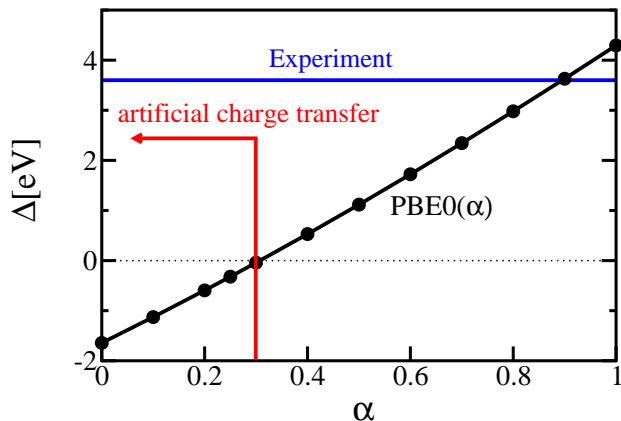


- PBE0-like hybrid functional:

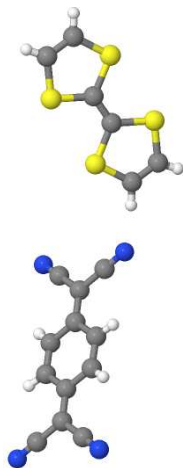
$$E_{xc}(\alpha) = \alpha(E_x^{\text{HF}} - E_x^{\text{PBE}}) + E_c^{\text{PBE}}$$



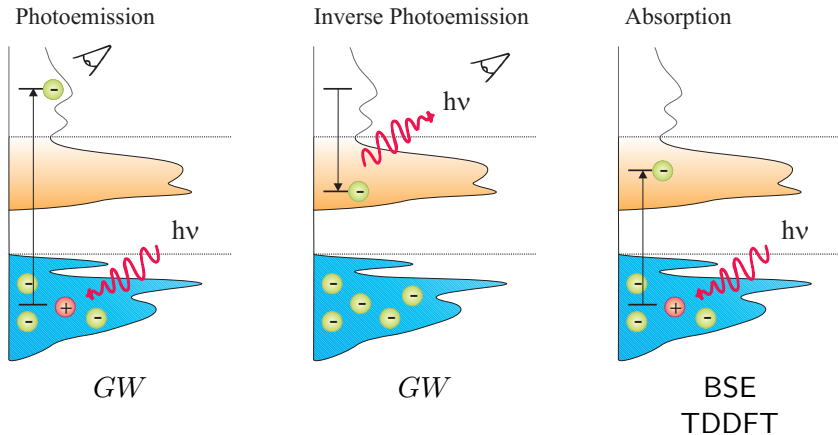
# TTF/TCNQ dimer at infinite separation



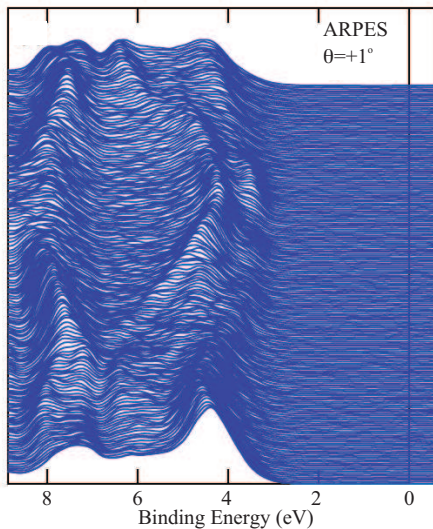
How to choose  $\alpha$ ?



# Band structures: photo-electron spectroscopy

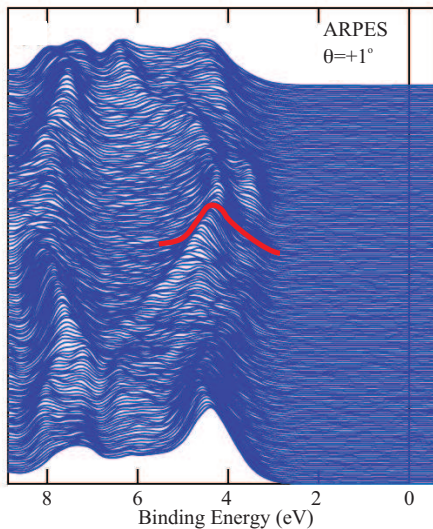


# Experimental: angle-resolved photoemission spectroscopy



Masaki Kobayashi, PhD dissertation

# Experimental: angle-resolved photoemission spectroscopy

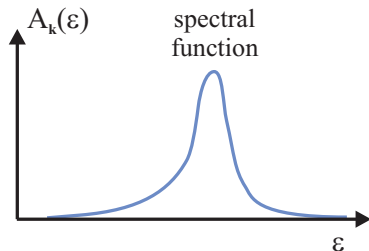


Masaki Kobayashi, PhD dissertation

# ARPES vs $GW$ : the quasiparticle concept

## Quasiparticle:

- single-particle like excitation



# ARPES vs $GW$ : the quasiparticle concept

## Quasiparticle:

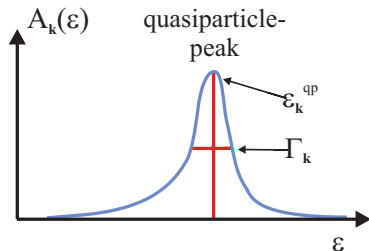
- single-particle like excitation

- $$A_{\mathbf{k}}(\epsilon) = \text{Im}G_{\mathbf{k}}(\epsilon) \approx \frac{Z_{\mathbf{k}}}{\epsilon - (\epsilon_{\mathbf{k}} + i\Gamma_{\mathbf{k}})}$$

$\epsilon_{\mathbf{k}}$  : excitation energy

$\Gamma_{\mathbf{k}}$  : lifetime

$Z_{\mathbf{k}}$  : renormalisation



# ARPES vs $GW$ : the quasiparticle concept

## Quasiparticle:

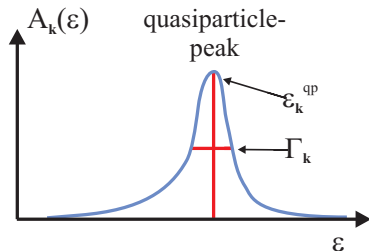
- single-particle like excitation

- $A_{\mathbf{k}}(\epsilon) = \text{Im}G_{\mathbf{k}}(\epsilon) \approx \frac{Z_{\mathbf{k}}}{\epsilon - (\epsilon_{\mathbf{k}} + i\Gamma_{\mathbf{k}})}$

$\epsilon_{\mathbf{k}}$  : excitation energy

$\Gamma_{\mathbf{k}}$  : lifetime

$Z_{\mathbf{k}}$  : renormalisation



- Density-Functional Theory:  $G_0$

- $GW$ :  $\chi_0 = iG_0G_0 \rightarrow W_0 = v/(1 - \chi_0v) \rightarrow \Sigma_0 = iG_0W_0$

$$G^{-1} = G_0^{-1} - \Sigma_0$$

# The screened Coulomb interaction $W$

Work with screened coulomb interaction:

$$W(\mathbf{r}, \mathbf{r}', t) = \int d\mathbf{r}'' \frac{\varepsilon^{-1}(\mathbf{r}, \mathbf{r}'', t)}{|\mathbf{r}'' - \mathbf{r}'|} \quad \varepsilon(\mathbf{r}, \mathbf{r}'', t): \text{ dielectric function}$$

The challenge is to calculate  $W(\mathbf{r}, \mathbf{r}'', t)$  from first principles!



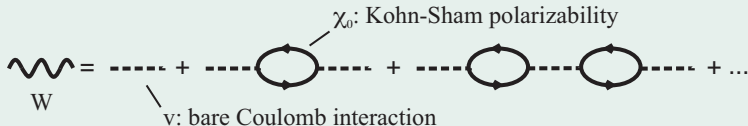
# The screened Coulomb interaction $W$

Work with screened coulomb interaction:

$$W(\mathbf{r}, \mathbf{r}', t) = \int d\mathbf{r}'' \frac{\varepsilon^{-1}(\mathbf{r}, \mathbf{r}'', t)}{|\mathbf{r}'' - \mathbf{r}'|} \quad \varepsilon(\mathbf{r}, \mathbf{r}'', t): \text{ dielectric function}$$

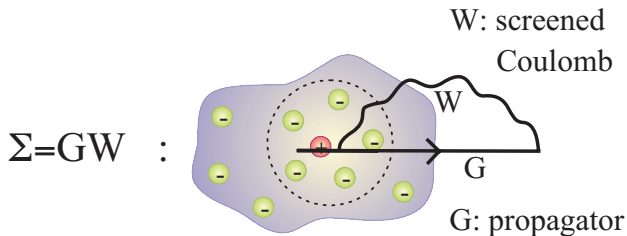
The challenge is to calculate  $W(\mathbf{r}, \mathbf{r}'', t)$  from first principles!

- we use **random-phase approximation (RPA)** for  $W$ :



- formally scales with (system size)<sup>4</sup>

# From screening to quasiparticle energies



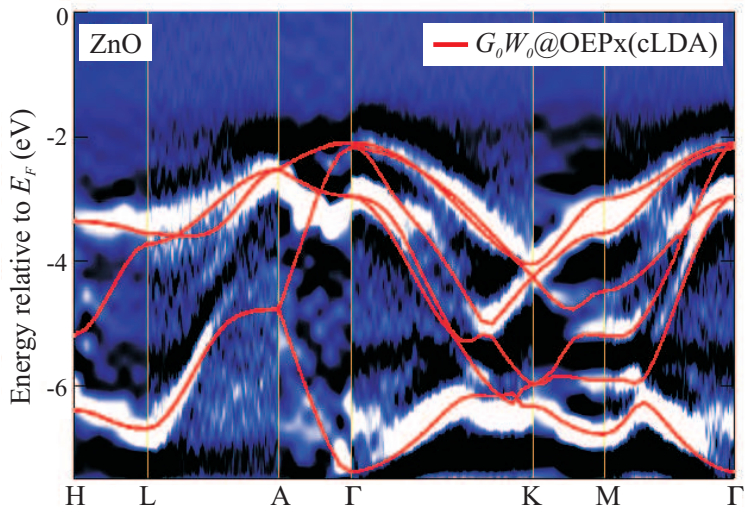
## Quasiparticle energies: $G_0W_0$ scheme

- electron addition and removal energies
- correction to DFT eigenvalues  $\epsilon_{n\mathbf{k}}^{\text{DFT}}$ :

$$\epsilon_{n\mathbf{k}}^{qp} = \epsilon_{n\mathbf{k}}^{\text{DFT}} + \sum_{n\mathbf{k}} G_0 W_0 (\epsilon_{n\mathbf{k}}^{qp}) - v_{n\mathbf{k}}^{xc}$$

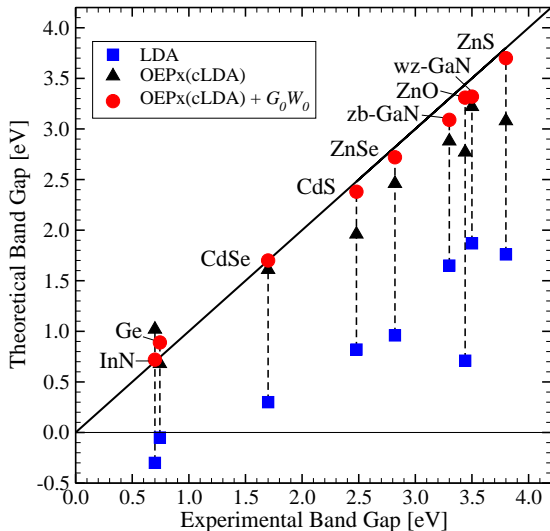
- includes image effect

# ARPES – GW: wurtzite ZnO



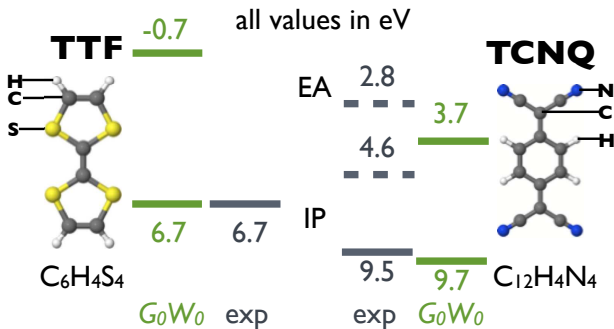
Yan, Rinke, Winkelkemper, Qteish, Bimberg, Scheffler, Van de Walle,  
Semicond. Sci. Technol. **26**, 014037 (2011)

# $G_0W_0$ Band Gaps



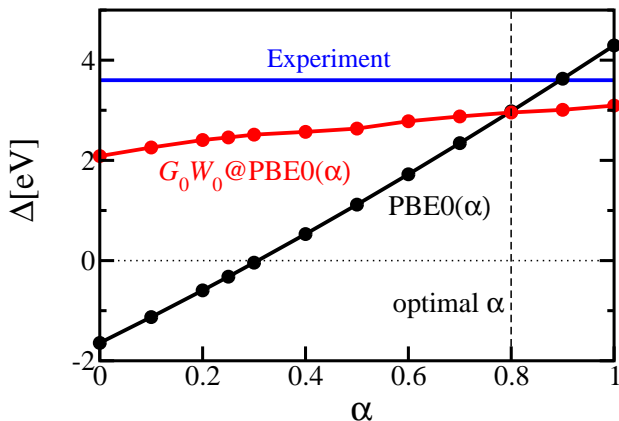
Rinke *et al.* *New J. Phys.* **7**, 126 (2005), *phys. stat. sol. (b)* **245**, 929 (2008)

# Performance of $G_0W_0$

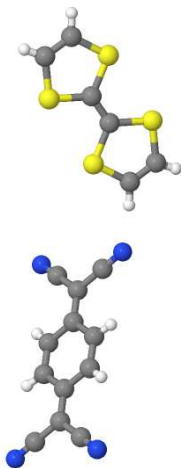


Exp.: JACS **112**, 3302 (1990), J. Chem. Phys. **60**, 1177 (1974),  
 Adv. Mat. **21**,1450 (2009),  
 Mol. Cryst. Liquid Cryst. Inc. Nonlin. Opt. **171**, 255 (1989)

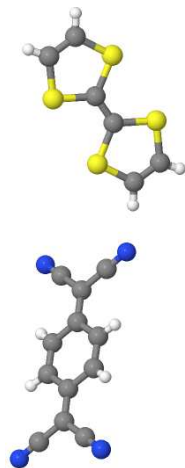
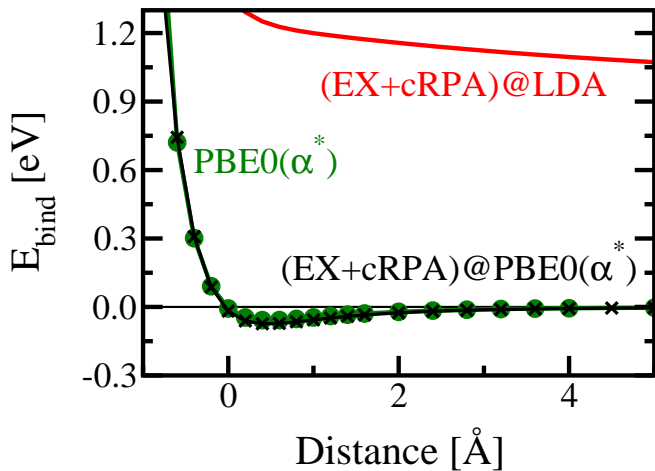
# TTF/TCNQ dimer at infinite separation



- choose  $\alpha$  to match  $G_0W_0$

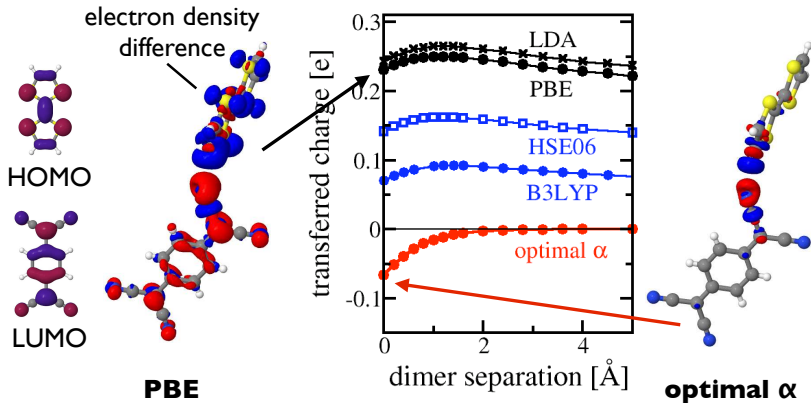


# TTF/TCNQ dimer – RPA binding curve



# TTF/TCNQ dimer – implications for interface

- no charge transfer, but small charge rearrangement
- 2D electron gas not due to charge transfer

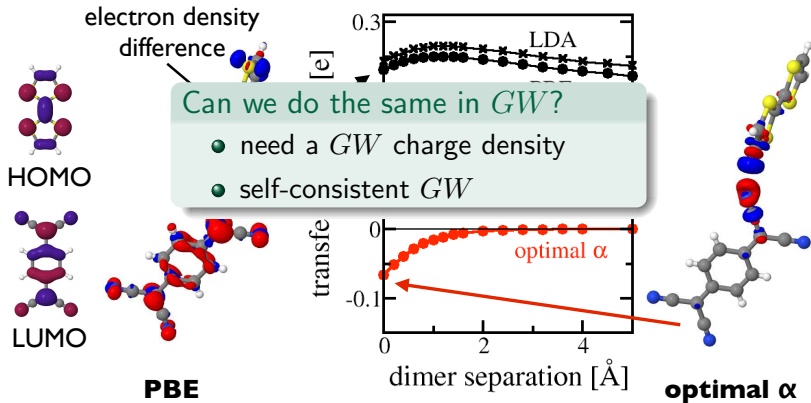


- *optimal- $\alpha$*  calculations for interface in progress



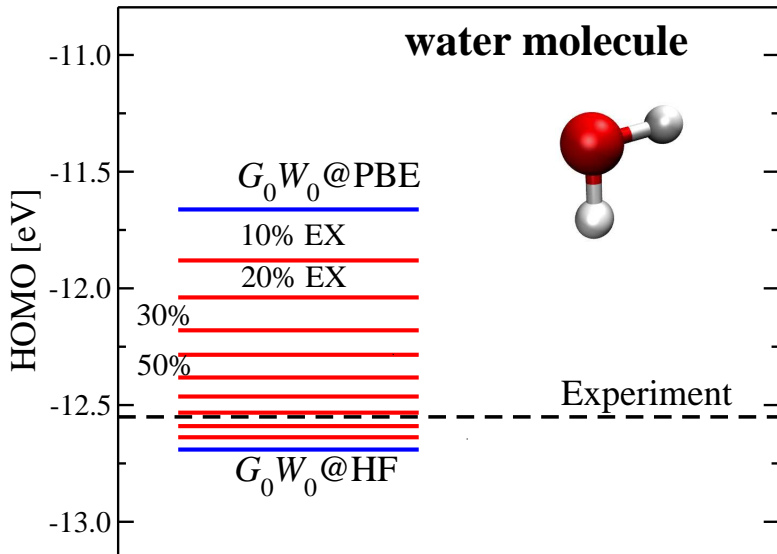
# TTF/TCNQ dimer – implications for interface

- no charge transfer, but small charge rearrangement
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- *optimal- $\alpha$*  calculations for interface in progress

# $G_0W_0$ for water



# GW – The Issue of Self-Consistency

Hedin's GW equations:

$$G(1,2) = G_0(1,2)$$

$$\Gamma(1,2,3) = \delta(1,2)\delta(1,3)$$

$$P(1,2) = -iG(1,2)G(2,1^+)$$

$$W(1,2) = v(1,2) + \int v(1,3)P(3,4)W(4,2)d(3,4)$$

$$\Sigma(1,2) = iG(1,2)W(2,1)$$

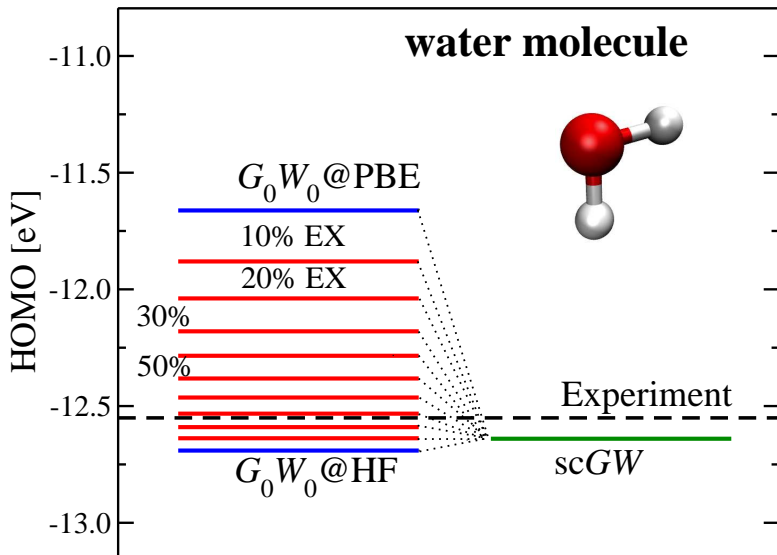
Dyson's equation:

$$G^{-1}(1,2) = G_0^{-1}(1,2) - \Sigma(1,2)$$

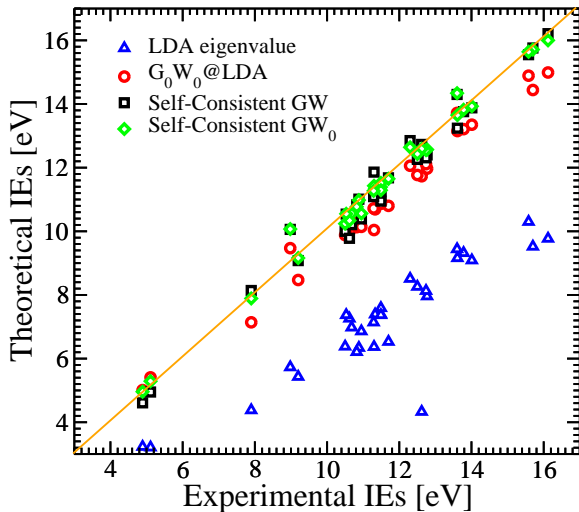
self-  
consistency

self-  
consistency

# scGW for water



# scGW and ionization potentials



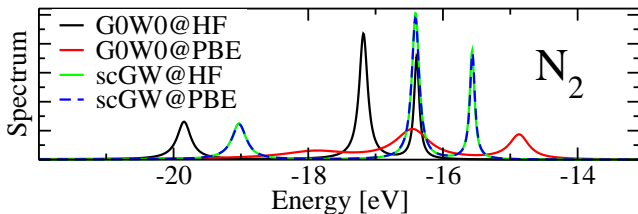
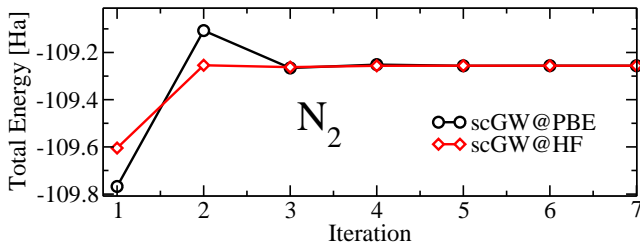
deviation from exp.

LDA	40.0%
$G_0W_0$ @LDA	5.6%
scGW	2.9%
$GW_0$ @LDA	1.5%

set taken from Rostgaard, Jacobsen, and Thygesen, PRB **81**, 085103, (2010)

## Dependence on the starting point: $N_2$

Galitskii-Migdal formula:  $E_{\text{tot}} = -\frac{i}{2} \int \text{Tr} [(\omega + h_0)G(\omega)] \frac{d\omega}{2\pi}$

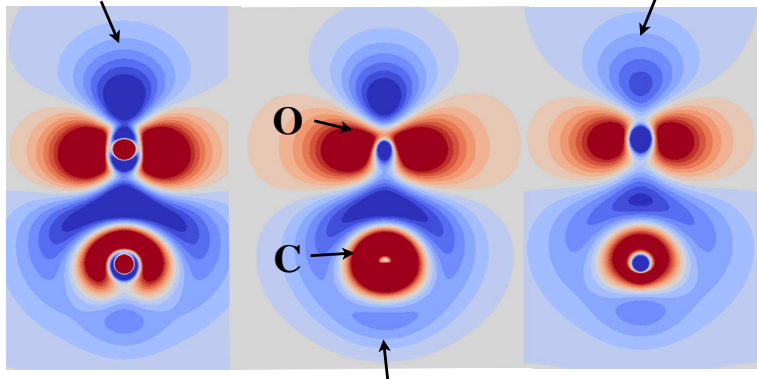


At self-consistency **no** dependence on input Green's function!

## The scGW density

$$\rho(PBE) - \rho(HF)$$

$$\rho(scGW) - \rho(HF)$$



$$\rho(CCSD) - \rho(HF)$$

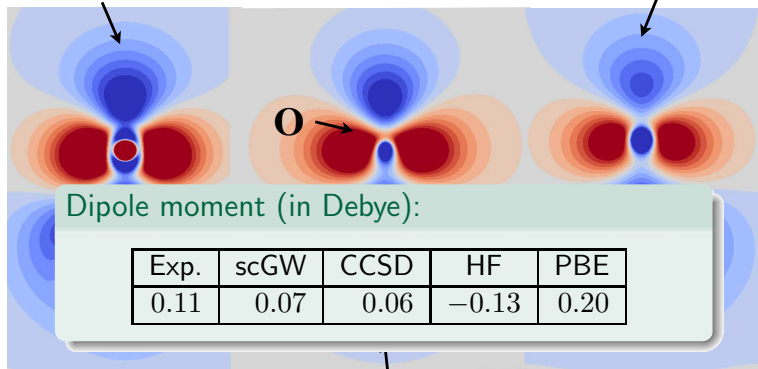
- density from Green's function:  $\rho(\mathbf{r}) = -i \sum_{\sigma} G_{\sigma\sigma}(\mathbf{r}, \mathbf{r}, \tau = 0^+)$

Caruso, Ren, Rinke, Rubio, Scheffler: arXiv:1202.3547v1

## The scGW density

$$\rho(PBE) - \rho(HF)$$

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Caruso, Ren, Rinke, Rubio, Scheffler: arXiv:1202.3547v1



## CO – ground state properties

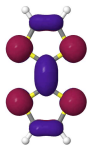
CO	$d$	$\nu_{\text{vib}}$	$\mu$	$E_{\text{b}}$
Exp. (NIST database)	1.128	2169	0.11	11.11
sc-GW	1.118	2322	0.07	10.19
$G_0W_0$ @HF	1.119	2647	–	11.88
$G_0W_0$ @PBE	1.143	2322	–	12.16
(EX+cRPA)@HF	1.116	2321	–	10.19
(EX+cRPA)@PBE	1.135	2115	–	10.45
PBE	1.137	2128	0.20	11.67
HF	1.102	2448	-0.13	7.63

- Galitskii-Migdal formula:

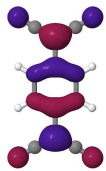
$$E_{\text{GM}} = -i \int \frac{d\omega}{2\pi} \text{Tr} \{ [\omega + h_0] G(\omega) \} + E_{\text{ion}}$$

Caruso, Ren, Rinke, Rubio, Scheffler: arXiv:1202.3547v1

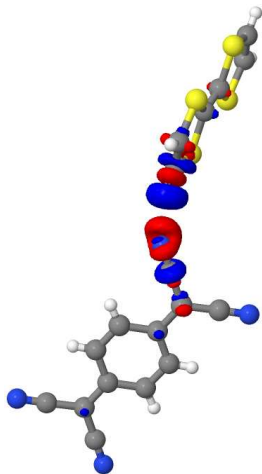
# TTF/TCNQ dimer revisited



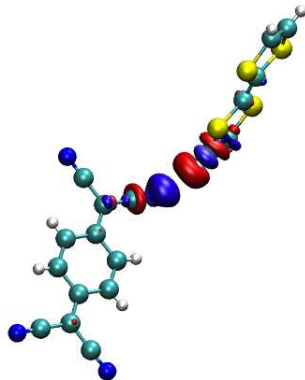
HOMO



LUMO



**opt  $\alpha$**



**scGW**

# RPA and $GW$ – a consistent pair

$$\Sigma = \text{[Diagram 1]} + \text{[Diagram 2]} + \dots$$

**GW**

many-body perturbation theory

**RPA**

density-functional theory

**Dyson's equation:**

$$G^{-1} = G_0^{-1} - \Sigma[G]$$

**total energy:**

$$E = E[G]$$

e.g. Galitskii-Migdal

**optimized effective potential:**

$$G_0 [\Sigma - v_{xc}^{\text{OEP}}] G_0 = 0$$

**Adiabatic connection fluctuation  
dissipation theorem (ACFDT):**

$$E_{xc} = E_{xc}[\chi_0] = E_{xc}[-iG_0G_0]$$

## RPA and $GW$ – a consistent pair

$$\Sigma = \text{GW} + \text{RPA} + \dots$$

man scGW versus scRPA:

- RPA also iterated to self-consistency (scRPA)

**Do MBPT and DFT differ for the same diagrams?**

**total energy:**

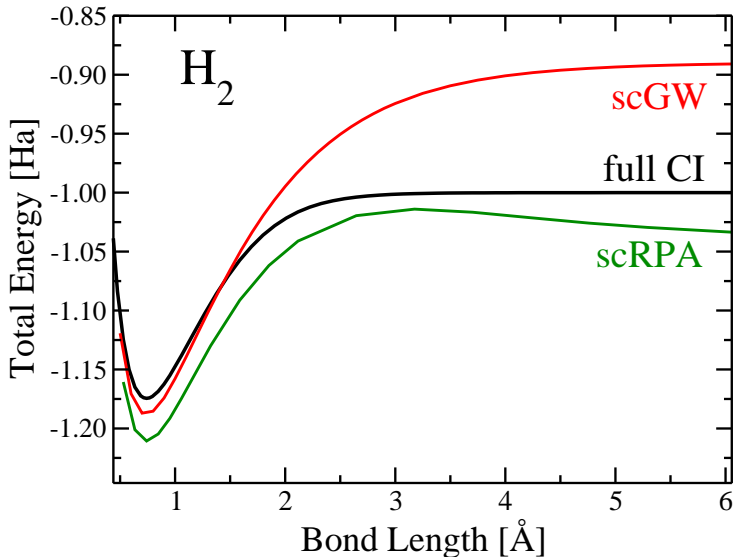
$$E = E[G]$$

e.g. Galitskii-Migdal

**Adiabatic connection fluctuation  
dissipation theorem (ACFDT):**

$$E_{xc} = E_{xc}[\chi_0] = E_{xc}[-iG_0G_0]$$

# The most strongly correlated system



full CI: L. Wolniewicz, J. Chem. Phys. **99**, 1851 (1993)

# Kinetic energy

- correlation energy in  $G_0W_0$ :

$$E_c^{G_0W_0} = - \int_0^\infty \frac{d\omega}{2\pi} \text{Tr} \left[ \sum_{n=2}^{\infty} (\chi_0(i\omega)v)^n \right]$$

- correlation energy in RPA:

$$E_c^{\text{RPA}} = - \int_0^\infty \frac{d\omega}{2\pi} \text{Tr} \left[ \sum_{n=2}^{\infty} \boxed{\frac{1}{n}} (\chi_0(i\omega)v)^n \right]$$

- adiabatic connection recovers kinetic energy

# Kinetic and correlation energy

- RPA kinetic energy contribution:

$$T_c^{\text{RPA}} = E_c^{\text{RPA}} - E_c^{G_0W_0}$$

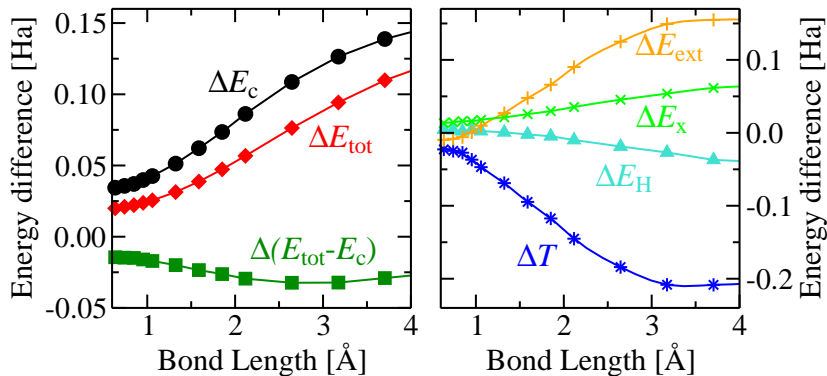
- RPA kinetic energy:

$$T^{\text{RPA}} = T_0 + T_c^{\text{RPA}}$$

- RPA Coulomb correlation:

$$U_c^{\text{RPA}} = E_c^{\text{RPA}} - T_c^{\text{RPA}}$$

## H<sub>2</sub> – energy contributions



- Coulomb correlation energy makes the difference



# Acknowledgements

## RPA+SOSEX/SE/R2PT

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**Xinguo Ren**

Alex Tkatchenko

Joachim Paier

Andreas Grüneis

Georg Kresse

Gus Scuseria

## TTF-TCNQ

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Viktor Atalla

Mina Yoon

## scRPA

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Daniel Rohr

Maria Hellgren

## scGW/Cerium

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Fabio Caruso/Marco Casadei

Xinguo Ren

Angel Rubio

## Support and Ideas

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Matthias Scheffler

## FHI-aims code

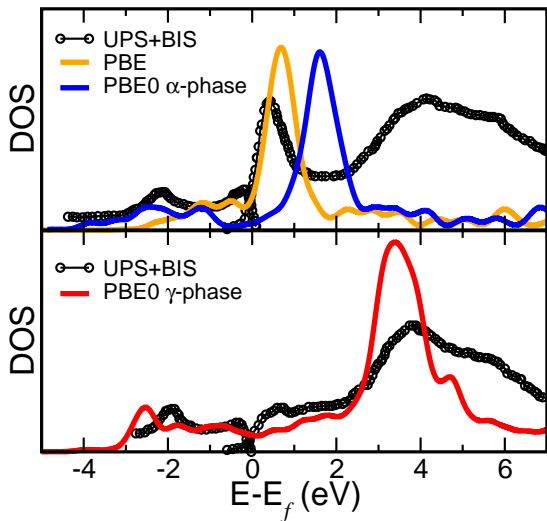
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Volker Blum

the whole FHI-aims developer team



# Cerium - spectroscopy



- description of  $\gamma$ -phase ok –  $\alpha$ -phase problematic

UPS: Wieliczka *et al.* PRB 29, 3028 (1984) – BIS: Wuilloud *et al.* PRB