

# Nonlinear Integrate & Fire Neuron Models

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- Joint work with: J. Carrillo (UAB), M. Gualdani (UT Austin), M. Schonbek (UCSC)

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- Previous work: M. Cáceres, J. Carrillo, B. Perthame (2011)

# The model

Unknown:  $V$  = membrane potential

## Basic LIF model

$$C_m \frac{dV}{dt} = -g_L(V - V_L) + I(t)$$

- $\tau_m = C_m/g_L \approx 10ms$   
 $V_L \approx -70mV$ .
- $I(t)$ : the external input current.

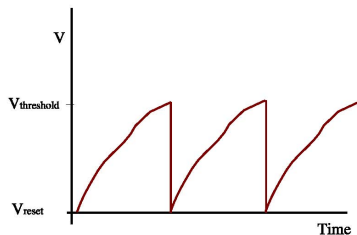
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 $V_L \approx -70mV$ .
- $I(t)$ : the external input current.
- Firing voltage: threshold value  
 $V_F \approx -50mV$ .
- Reset voltage: discharged value  
 $V_R \approx -60mV$ .

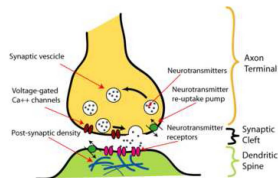


## Stochastic Synapse Model

$$I(t) = J_E \sum_{i=1}^{C_E} \sum_j \delta(t - t_{Ej}^i) - J_I \sum_{i=1}^{C_I} \sum_j \delta(t - t_{Ij}^i)$$

### Parameters

- Two different Neuron-types: **Inhibitory** (I) or **Excitatory** (E).  
Strength of the Synapse:  $J$ .  
Number of presynaptic neurons:  $C$ .
- Spiking times:  $t_j^i$  = time of the  $j^{\text{th}}$ -spike coming from the  $i^{\text{th}}$ -presynaptic neuron.
- Stochastic Assumption: Neurons spike according to a Poisson process with constant probability of emitting a spike per unit time  $\nu$ .



# Diffusion Approximation

Mean and Variance of  $I(t)$ :

$$\mu_C = b\nu \quad \text{with} \quad b = C_E J_E - C_I J_I \quad \text{and} \quad \sigma_C^2 = (C_E J_E^2 + C_I J_I^2)\nu.$$

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- 2 The value of  $N(t)$  is then computed as the flux of probability that neurons cross the threshold voltage  $V_F$  per unit time.
- 3 Average-excitatory (-inhibitory resp.) if  $b > 0$  ( $b < 0$  resp.).

# PDE model (special case)

- Ito's rule: PDE for the evolution of the **probability density**  $p(v, t) \geq 0$  of finding neurons at a voltage  $v \in (-\infty, V_F]$  at a time  $t \geq 0$ .

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- **Boundary conditions:**  $p(V_F, t) = 0$ ,  $p(-\infty, t) = 0$
- **Initial data:**  $p(v, 0) = p_I(v)$ .



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  - There exists initial data such that develop blow up at finite time
  - There are cases of non-existence or non-uniqueness for the steady state
- Concentrate, for now, on the inhibitory case ( $b < 0$ )

- Change of variables: Fokker-Plank to heat equation

$$y = e^t v, \quad \tau = \frac{1}{2}(e^{2t} - 1), \quad p(v, t) = e^t w\left(e^t v, \frac{1}{2}(e^{2t} - 1)\right)$$

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$$u(x, \tau) = w(y, \tau), \quad x = y - b \int_0^\tau M(s)\alpha^{-1}(s) ds.$$



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- Change of variables: Get rid of  $w_y$  term:  
 $u(x, \tau) = w(y, \tau), \quad x = y - b \int_0^\tau M(s)\alpha^{-1}(s) ds.$
- Get a free boundary Stefan problem with source
- Related to a price formation equation  
(Lasry-Lions, G.-Gualdani, Chayes-G.-Gualdani-Kim,  
Caffarelli-Markowich-Pietschmann)

# Relation to Stefan problem

$$(P) \left\{ \begin{array}{l} u_t = u_{xx} + M(t)\delta_{x=s_1(t)}, \quad x < s(t), t > 0, \\ s(t) = -b \int_0^t M(s)\alpha^{-1}(s) ds, \quad t > 0, \\ M(t) = - \left. \frac{\partial u}{\partial x} \right|_{x=s(t)}, \quad t > 0, \end{array} \right.$$

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- 3 If  $b < 0$ , the free boundary  $s(t)$  is an increasing function of time, and  $s(t)$ ,  $s_1(t)$  never cross.

- Green's function:  $G(x, t, \xi, \tau) = \frac{1}{[4\pi(t-\tau)]^{1/2}} \exp\left\{-\frac{|x-\xi|^2}{4(t-\tau)}\right\}$



# Integral formulation *a la* Friedman

- Green's function:  $G(x, t, \xi, \tau) = \frac{1}{[4\pi(t-\tau)]^{1/2}} \exp\left\{-\frac{|x-\xi|^2}{4(t-\tau)}\right\}$
- Duhamel's formula:

$$u(x, t) = \underbrace{\int_{-\infty}^0 G(x, t, \xi, 0) u_I(\xi) d\xi}_{\text{homogeneous heat equation}} - \underbrace{\int_0^t M(\tau) G(x, t, s(\tau), \tau) d\tau}_{\text{Takes care of free boundary}} + \underbrace{\int_0^t M(\tau) G(x, t, s_1(\tau), \tau) d\tau}_{\text{Delta function at } s_1(t)}$$

$$\begin{aligned}M(t) &= \partial_x u(s(t), t) \\&= -2 \int_{-\infty}^0 G(s(t), t, \xi, 0) u'_l(\xi) d\xi \\&\quad + 2 \int_0^t M(\tau) G_x(s(t), t, s(\tau), \tau) d\tau \\&\quad - 2 \int_0^t M(\tau) G_x(s(t), t, s_1(\tau), \tau) d\tau \\&=: \Theta(M)(t)\end{aligned}$$

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**Fixed point argument:**  $\Theta : C_{\sigma, m} \rightarrow C_{\sigma, m}$ , where

$$C_{\sigma, m} := \{M \in \mathcal{C}([0, \sigma]) : \|M\| \leq m\}$$

$$\|M\| := \sup_{0 \leq t \leq \sigma} |M(t)|$$

# Integral formula for $M$

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Contraction for  $\sigma$  small, where  $\sigma = \sigma(\|u'_l\|, v_R)$

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Corollary - short time existence for  $u$

$\exists!$  “regular” solution  $u$  for  $t \in (0, T)$

# Global existence ( $b < 0$ )

## Lemma - no break up time

$\exists \mu > 0$  small enough such that, for any  $t_0 > 0$ , if

$$\sup_{x \in (-\infty, s(t_0 - \mu)]} |u_x(x, t_0 - \mu)| < \infty,$$

then also

$$\sup_{t_0 - \mu < t < t_0} M(t) < \infty.$$

Here  $\mu$  is independent of  $t_0$ .

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## Theorem ( $b < 0$ )

There exists a unique regular solution  $u$  for all  $t > 0$ .



## Criterium

Solution exists up to time  $T^*$ , where

$$T^* = \sup\{t < 0 : M(t) < \infty\}.$$

Proof:

$M$  bounded  $\Rightarrow u_x$  bounded  $\Rightarrow$  Can extend solution past  $T^*$

# Asymptotics for $b = 0$ (linear case - no free boundary)

$$(P_0) \quad \begin{cases} p_t = \partial_v [vp] + p_{vv} + N(t)\delta_{v=V_R}, := \mathcal{L}p, & v < 0, t > 0, \\ N(t) = - \left. \frac{\partial p}{\partial v} \right|_{v=0}, & t > 0, \\ p(-\infty, t) = 0, \quad p(0, t) = 0, & t > 0, \\ p(v, 0) = p_I(v), & v < 0. \end{cases}$$

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Known (Caceres-Carrillo-Perthame):

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- Solution decays exponentially fast to equilibrium.
- Entropy methods+Poincaré inequality
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Question: decay.... at what rate **????**

There exists a one parameter family of equilibria

$$p_{\infty}(v) = \begin{cases} e^{-v^2/2} & v \in (-\infty, V_R), \\ \alpha_0 e^{-v^2/2} \int_v^0 e^{v^2/2} dv & v \in (V_R, 0], \end{cases}$$

for

$$\alpha_0 := \left( \int_{V_R}^0 e^{v^2/2} dv \right)^{-1}.$$

# Standard Fokker-Planck equation

- Compute the eigenvalues:  $\mathcal{L}p = \lambda p$ , where

$$p_t = \partial_v [vp] + p_{vv} := \mathcal{L}p$$

- Norm:

$$\|p\|_{L^2_{exp}(\mathbb{R})}^2 := \int_{\mathbb{R}} \left( e^{v^2/4} |p(v)| \right)^2 dv.$$

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- Eigenvalues are

$$\lambda = -n, \quad n \in \mathbb{N}.$$

- Eigenfunctions are given by the Hermite polynomials  $H_n$  as

$$p_n(v) = H_n(v)e^{-v^2/2}.$$

$$(P_0) \quad \begin{cases} p_t = \partial_v [vp] + p_{vv} + N(t)\delta_{v=V_R}, := \mathcal{L}p, & v < 0, t > 0, \\ N(t) = - \left. \frac{\partial p}{\partial v} \right|_{v=0}, & t > 0, \\ p(-\infty, t) = 0, \quad p(0, t) = 0, & t > 0, \\ p(v, 0) = p_I(v), & v < 0. \end{cases}$$



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**Conclusion:** the spectral gap may have nothing to do with the exponential decay to equilibrium !!!

## Tenth **WORKSHOP ON INTERACTIONS BETWEEN DYNAMICAL SYSTEMS AND PARTIAL DIFFERENTIAL EQUATIONS**

Barcelona, May 28 - June 1, 2012

Financial **support** available.

Deadline: March 30, 2012.

- M. Capinski (Krakow): Geometric methods for invariant manifolds in dynamical systems
- M. Gidea (Chicago): Aubry Mather Theory from a Topological Viewpoint
- H. Shahgholian (KTH): Obstacle type free boundaries (Theory and applications)
- E. Valdinoci (Milano): Local phase transition equations

Thank you!!