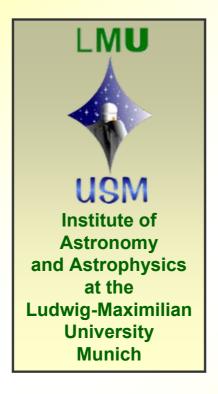
# Consistent Models of the Expanding Atmospheres of Hot Star



May 19, 2005 Adalbert W.A. Pauldrach

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# Consistent Models of the Expanding Atmospheres of Hot Star

# **Theory - Models**

Observations Motivation

non-equilibrium thermodynamics (NLTE) radiation hydrodynamics

**Boltzmann - equation** 

$$\frac{\partial f}{\partial t} + \underline{\mathbf{v}} \cdot \nabla f + \frac{1}{m} \underline{\mathbf{F}} \cdot \nabla_{\mathbf{v}} f = \left[ \frac{\partial f}{\partial t} \right]_{\text{ic}}$$

time development of the distribution functions

Application
Spectral Diagnostics

<mark>sub-group</mark>

Low Mass Stars

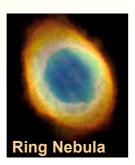
0.4 - 1.44 M<sub>☉</sub>

CSPN (SN Ia)

luminous,10<sup>4</sup>L<sub>⊙</sub>

radiation, UV/EUV

rare, evolution



Hot Stars

T<sub>eff</sub> = 15000 - 100000 K

show wind features

**Expanding Atmospheres** 

characterized by

high radiation energy density

sub-group

Massive Stars

20 - 150 M<sub>o</sub>



O/B (WR SN II)

young,10<sup>6</sup> yr

luminous,10<sup>6</sup>L<sub>⊙</sub>

rare, IMF

radiation, UV/EUV

10-4  ${
m M_{\odot}/yr} \geq \dot{{
m M}} {
m v_{\infty}}$ 

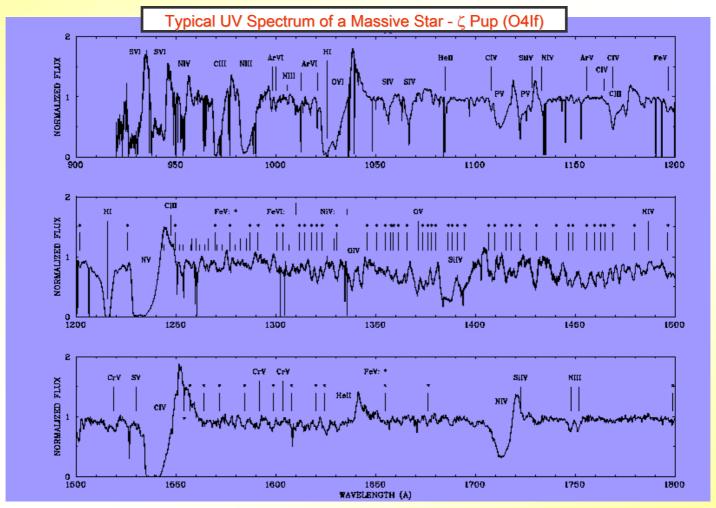
ightarrow  $m \dot{M}~v_{\infty} \leq 4000~km/s$ 

In this talk focus on the status of UV Spectral Diagnostics!

## Hot Stars: Observations

# Emergent UV Spectra of O-type Stars are Old Fashioned!

Walborn, Nichols-Bohlin, Panek, 1985, NASA Ref. Publ. 1155; Morton and Underhill, 1977, Ap. J. Suppl. 33, 83



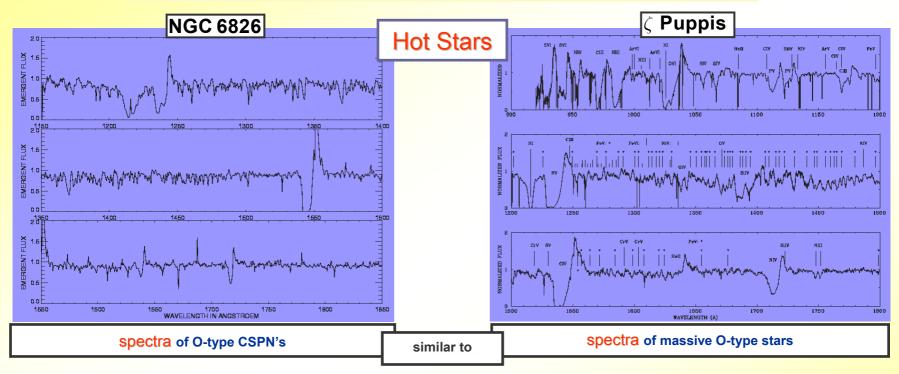
Signatures of P-Cygni lines

**Hundreds of strongly wind contaminated lines** 

SVI,CIII,NIII,OVI,PV,NV,OIV,OV,SiIV,CIV,HEII,NIV FeV, NiV, FeIV, FeVI, CrV, ArV, ArVI

## Hot Stars: Old Fashioned or Trendy?

# Realistic Synthetic UV Spectra of O-type Stars A Diagnostic Tool with great Astrophysical Potential!?

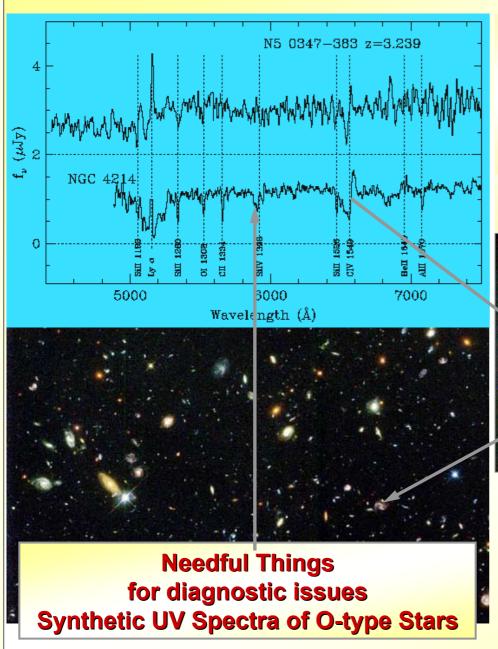


Spectra contain information about stellar and wind parameters and abundances

Spectra available for > 25 years
but most work done concentrated on qualitative results and arguments

Powerful Tool for modeling realistic synthetic spectra for diagnostic issues is still in development!

## Motivation: Galaxies and Galaxy Clusters are Trendy !



#### Corresponding Spectrum of a Starburst Galaxy at z=3

Keck spectrum of a z > 3 galaxy, chosen to illustrate the variety of features encountered.

For comparison a recent HST spectrum of the central starburst region in the Wolf-Rayet galaxy NGC4214 is also shown.

Note the characteristic P-Cygni lines – CIV, SilV Steidel et al., 1996, Ap.J.L. 462, L17

Distance:~12 Mpc
located in the constellation of Coma Berenices.
© Hubble Space Telescope



Starburst Ring

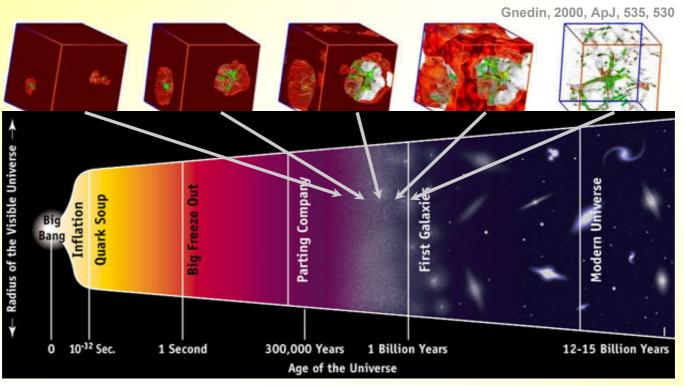
Dominated by Massive Stars

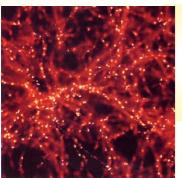
## Motivation: Reionization of the universe

#### **End of Cosmic Dark Ages**

absence of Gunn-Peterson trough in the spectra of high-redshift quasars implies that the universe was reionized at a redshift of z ~ 6

Fan et al., 2000, AJ, 120, 1167





# Population III stars (Z<10<sup>-3</sup>) are important for the reionization

Carr et al.,1984,ApJ 277, 445

but

ionization history of the IGM depends crucially on the IMF of the First Stars

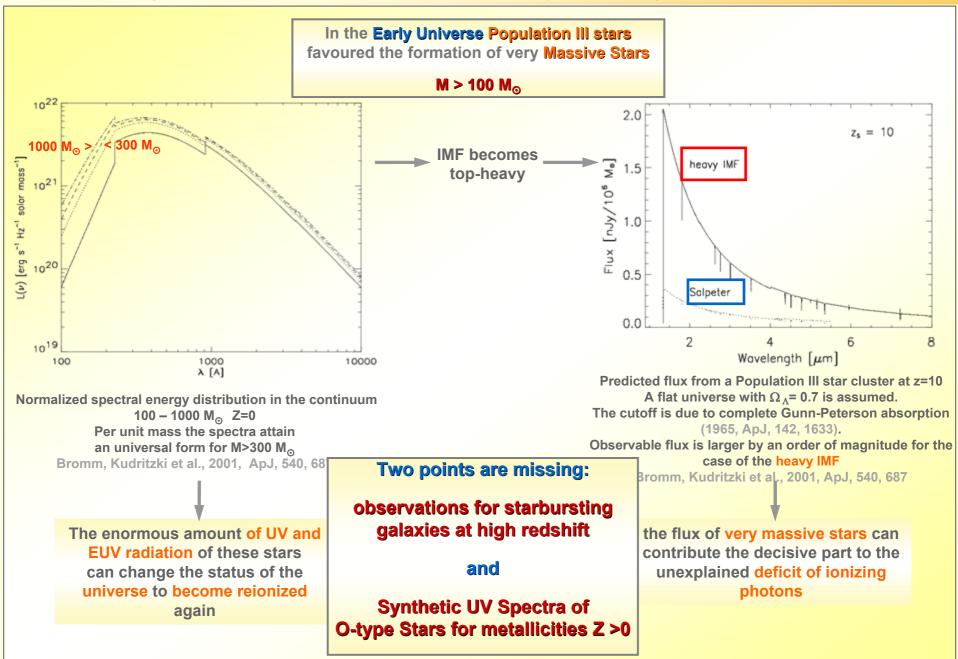
Bromm et al., 1999, ApJ, 527, L5

Numerical simulation of the first stars in the universe

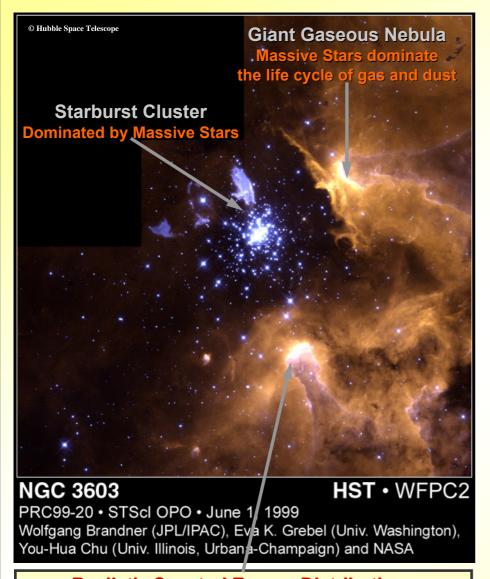
Davé, Katz, Weinberg, 2002, ApJ, 579, 23



## First generations of Stars: Ionization Efficiency of a Heavy Initial Mass Function



## Motivation: The impact of massive stars on their Environment



Realistic Spectral Energy Distributions
of massive stars and stellar clusters are required
to analyze the excited HII regions

#### **Giant Galactic Nebula NGC 3603**

various stages of the life cycle of stars are captured

© Hubble Space Telescope

Near the center of the view is a Starburst Cluster dominated by young, early O-type stars

The Interaction of ionizing radiation with cold molecular-hydrogen cloud material results in the giant gaseous nebula structures to the right of the cluster.

In the Carina spiral arm of the Galaxy at a distance of about 6 kpc field 2.5 x 2.5 arcmin<sup>2</sup> (~ 6 x 6 pc<sup>2</sup>)

# Massive stars dominate the life cycle of gas and dust in several types of galaxies

#### Interaction leads to

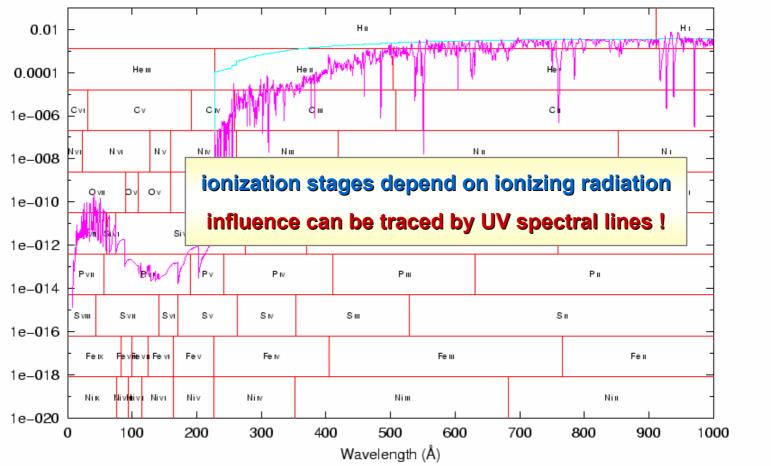
chemical enrichment of the ISM large amount of momentum and energy input into the ISM basis for star formation

## **Ionizing flux of ζ Puppis (O4If)**



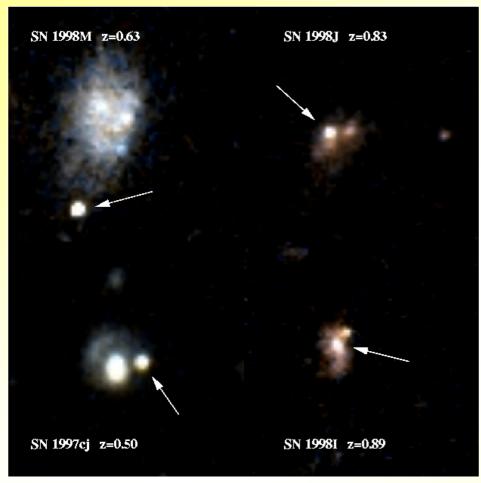
## **UV Spectra of O Stars:**

test involves hundreds of spectral signatures of different ionization stages



## Motivation: SNIa as distance indicators

### SNIa are Standard Candles!



Observations: 4 High-z Supernovae imaged with HST <a href="http://cfa-www.harvard.edu/cfa/oir/Research/supernova/highz/figures/index.html">http://cfa-www.harvard.edu/cfa/oir/Research/supernova/highz/figures/index.html</a>.

**Are SNIa Standard Candles?** 

## surprising result:

Distant SNe la appear fainter than a standard candle in an empty Friedmann model of the universe!

# SNIa-luminosity distances

indicate

accelerated expansion!

## **Are SNIa Standard Candles**

independently of age?
Or is there some evolution of
SNIa-luminosity
with age?

## Spectroscopy

is a powerful tool to search for differences!

### **SN 1992A**

at early phases ~ 5 days after B maximum

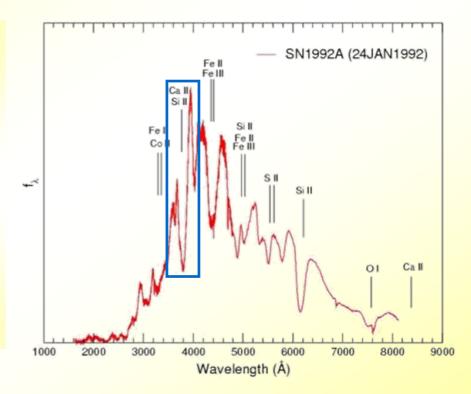
## Early epochs ( < 2 weeks after maximum):

No hints for H and He lines

Prominent absorption features of mainly intermediate mass elements embedded in a non-thermal pseudo continuum (SiII, OI, SII, CaII, MgII...)

→ "photospheric epoch"

Characteristic P-Cygni line profiles blue-shifted absorption red-shifted emission broad lines → high ejecta velocities



HST spectrum of 24. Jan. 1992 from Kirshner et al., 1993, ApJ 415, 589

Are SNe la standard candles in a cosmological sense?

**Required:** 

Realistic Models and Synthetic Spectra of Type la Supernovae

## Motivation: Conclusion

# Hot Stars play an important role in Astrophysics

# they dominate the physical conditions of galaxies and the life cycle of gas and dust

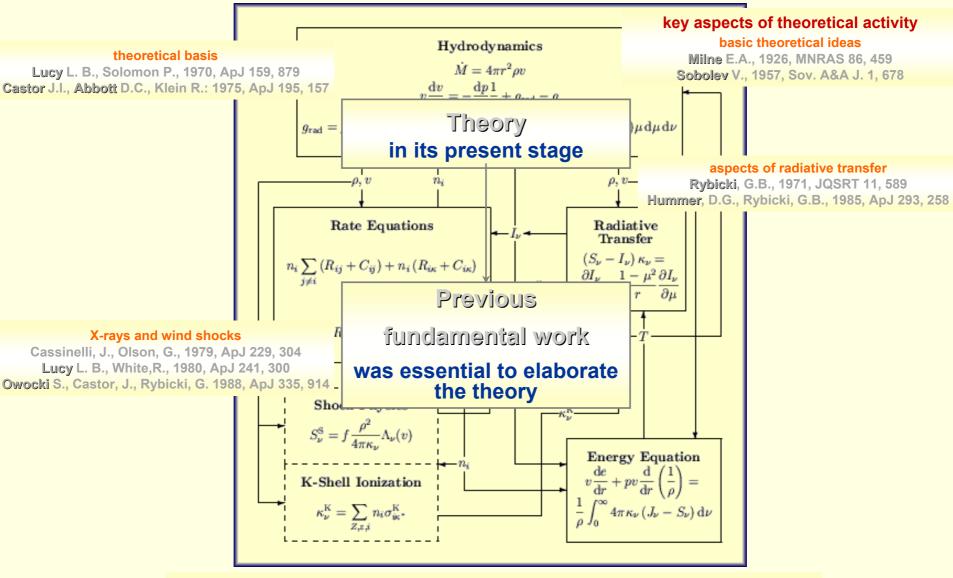
chemical enrichment of the ISM - significant impact on chemical evolution large amount of momentum and energy input into the ISM - major dynamical role

relevance to cosmological issues

Realistic Synthetic UV Spectra of Hot Stars

are still relevant for current astronomical research





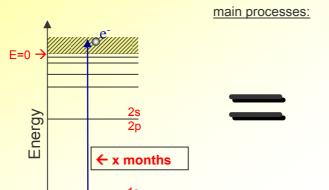
#### basis of our theoretical framework

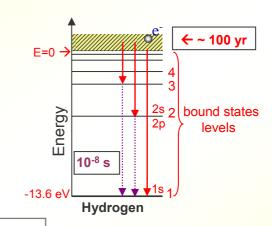
Pauldrach, A.W.A., Puls, J., Kudritzki, R.P.,1986, A&A,164,86; Pauldrach, A.W.A., 1987, A&A, 183, 295
Puls, J., Pauldrach, A.W.A., 1990, PASPC 7, 203; Pauldrach, A.W.A. et al., 1994, A&A, 283, 525
Feldmeier, A. et al., 1997, A&A 320, 899; Pauldrach, A.W.A. et al., 1998, ASPCS 131, 258
Pauldrach, A.W.A., Hoffmann, T.L., Lennon, M., 2001, A&A, 375, 161

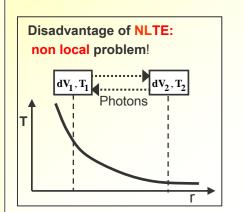
# Theory: Motivation

#### in GNs radiative processes are not balanced in detail!

→ thermodynamic equilibrium can not be established!







-13.6 eV

Hydrogen

just ground states are occupied

→no thermodynamic equilibrium (TE)

instead

Non-Equilibrium Thermodynamics (NLTE)

described by a well defined mixture of macroscopic and microscopic physics

Transport Theory
Thermodynamics

**Quantum Mechanics** 

#### Advantage of TE:

In **TE** state of the gas completely described by

2 macroscopic variables

→ local problem!

$$n_i^* = n_i(\rho, T)$$

 $n_i^* = \text{occupation numbers of bound levels}$   $(\text{number of particles per } cm^3)$ 

 $n_i$  follows from **Boltzmann statistics** 

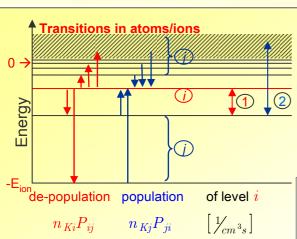
$$n_1 4\pi \int\limits_{\nu_{1k}}^{\infty} \frac{\alpha_{1k}(\nu)}{h\nu} J_{\nu} d\nu = n_P n_e \alpha_R$$
 ionization equation

$$\left[\frac{1}{r^2}\frac{\partial}{\partial r}\right]r^2J_{\nu} = -n_1\alpha_{1k}J_{\nu}(+\eta_{\nu})$$

the equation of transfer in spherical symmetry

**Integro-Differential Equation** 

# Theory: Motivation

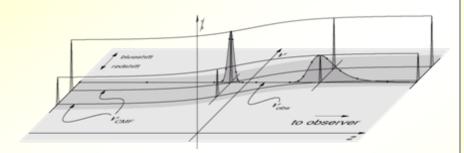


bound-bound transitions excitation and de-excitation radiative/collisional rates

 $R_{ij}$  ,  $C_{ij}$ 

bound-free transitions ionization and recombination processes

R<sub>ik</sub>, C<sub>ik</sub>

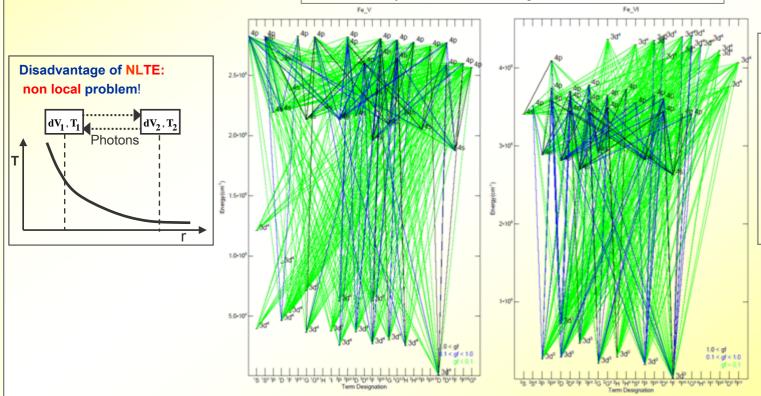


0-th moment of distribution function  $f_{Ki}$  for each element K, level i:

$$\underline{\nabla} \cdot (n_{Ki}\underline{\mathbf{u}}) = \int_{-\infty}^{+\infty} \left[ \frac{\partial f_{Ki}}{\partial t} \right]_{ic} d^3 \mathbf{v} = \sum_{j \neq i} n_{Kj} P_{ji} - \sum_{j \neq i} n_{Ki} P_{ij}$$

yields net rate of changes for each K, i

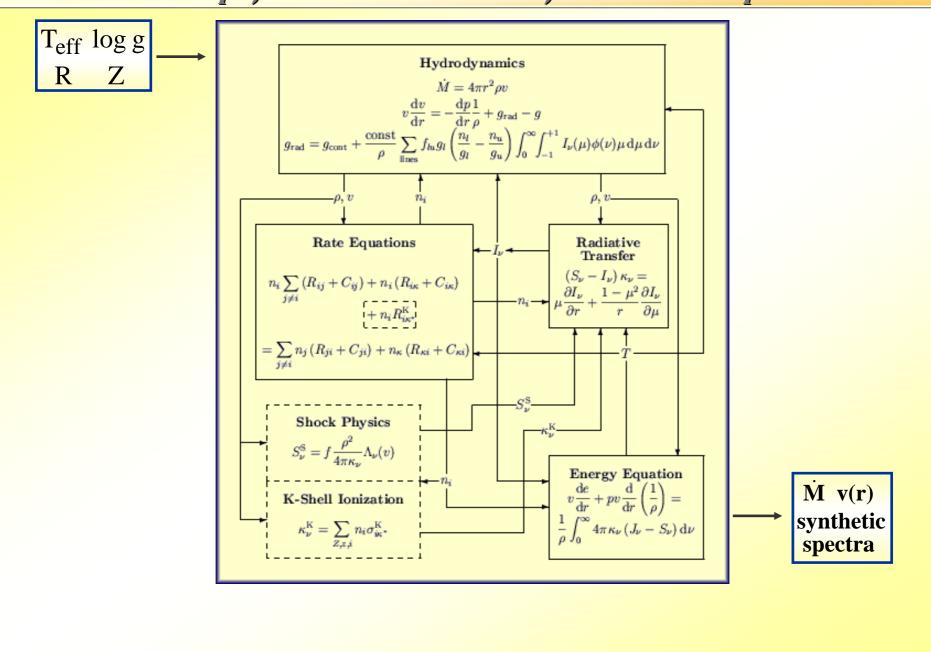
Nonequilibrium thermodynamics is a rather esoteric field of physics that we would do well to avoid!

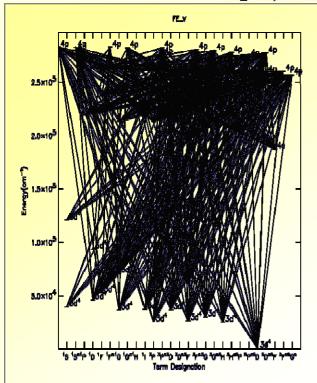


# Non-Local Thermodynamic Equilibrium (NLTE)

described by a well defined
mixture of
macroscopic and
microscopic physics
Transport Theory

Thermodynamics
Quantum Mechanics





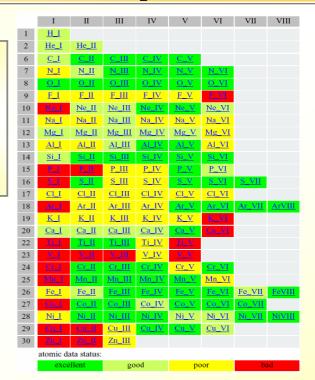
## well-elaborated

atomic models are required in order to avoid

GIGO

## **Rate Equations**

$$n_i \sum_{i \neq j} (R_{ij} + C_{ij}) + n_i (R_{ik} + C_{ik})$$
  
=  $\sum_{i \neq j} n_j (R_{ji} + C_{ji}) + n_k (R_{ki} - C_{ki})$ 



#### **NLTE Models** require atomic data for

lines, collisions, ionization, recombination

#### Essential for occupation numbers, line blocking, line force

Accurate atomic models have been included

26 elements

149 ionization stages

**5,000** levels ( + **100,000** )

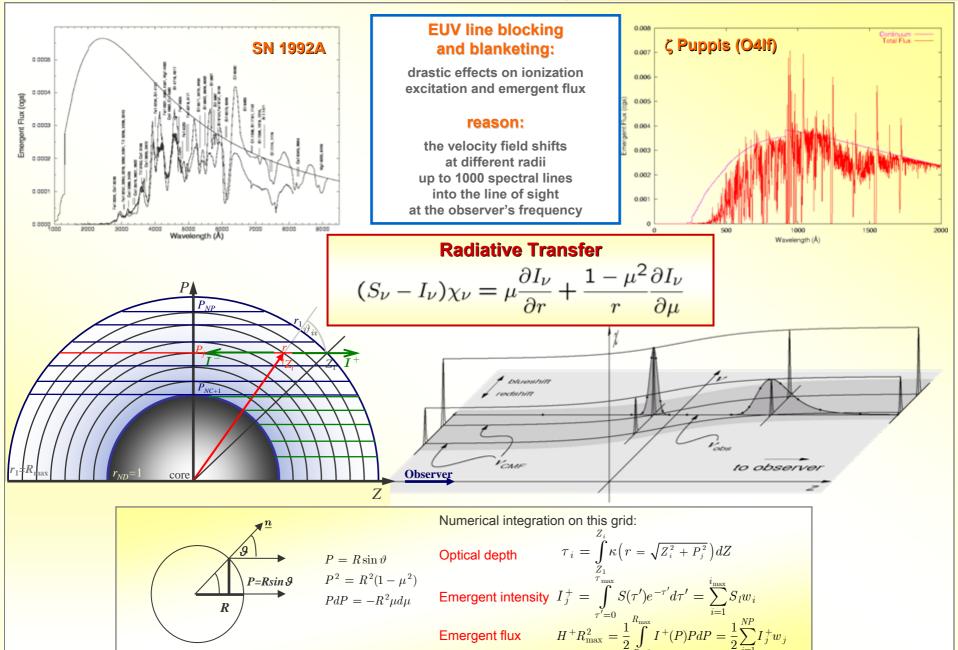
20,000 diel. rec. transitions

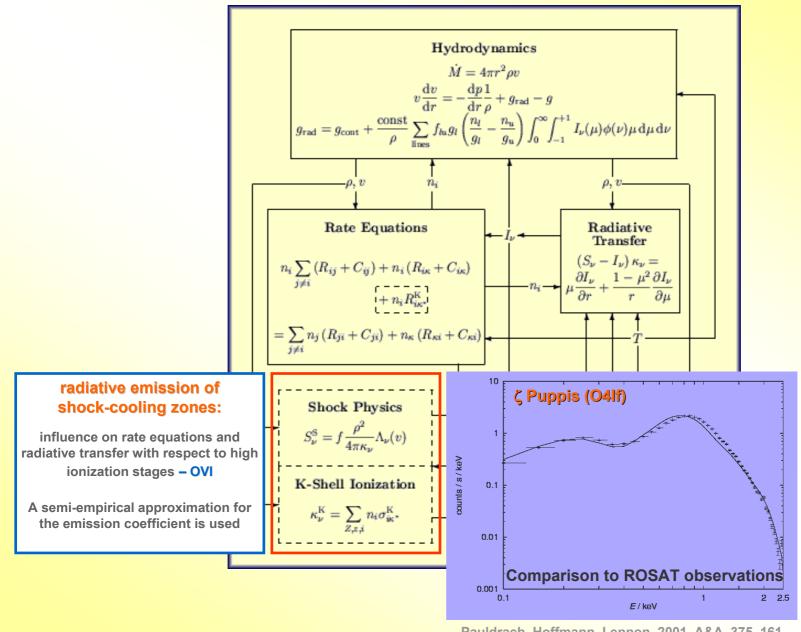
4 10<sup>6</sup> b-b line transitions

Auger-ionization

recently improved models are based on Superstructure

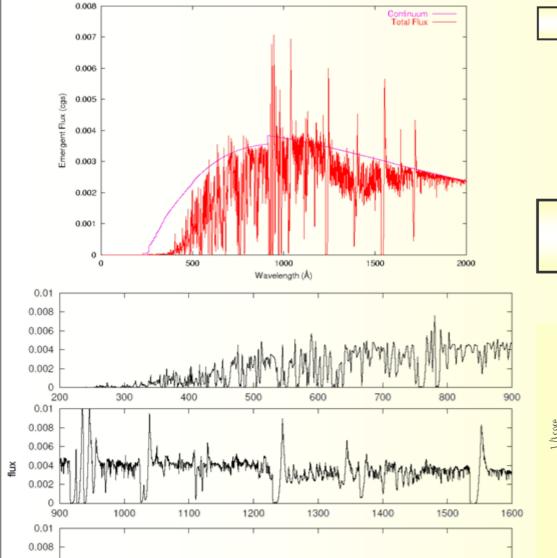
Eisner et al., 1974, CPC 8,270





Pauldrach, Hoffmann, Lennon, 2001, A&A, 375, 161

# Solution of the system > Tool for spectrum synthesis

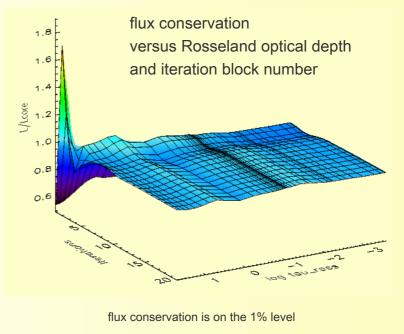


O supergiant model

$$\begin{array}{ll} \mathbf{T_{eff}} &= 45000 \ \mathbf{K} \\ \mathbf{log} \ \mathbf{g} &= 3.6 \\ \mathbf{R} &= 18 \ \mathbf{R_{\odot}} \end{array}$$

### Are the models realistic?

Strategy for quantitative spectral UV analysis



1600

0.006 0.004 0.002

1800

1900

wavelength

2000

2100

2200

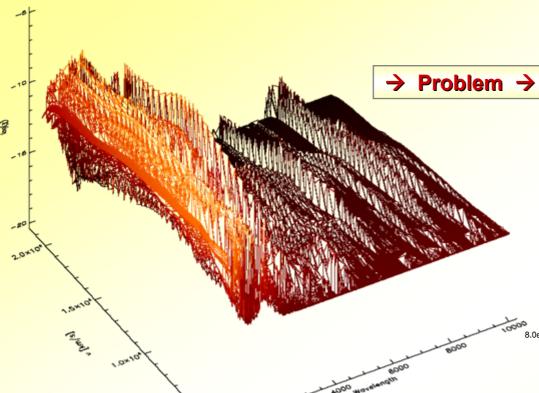
2300

1700

## Results: Synthetic and observed spectra of Supernovae of Type Ia

#### $T_{\rm eff} = 14000 \ {\rm K}$ **SN 1992A Comparison to observations** $R = 9900 R_{\odot}$ at early phases ~ 5 days after B maximum → verification of explosion models $Z = Z_{w_7}$ **Constraints for** Models based on W7 explosion model **Abundances** $v_{\infty} = 22000 \text{ km/s}$ averaged chemical composition above the photosphere **Velocity-structure** Nomoto et al, 1984, ApJ. 286, 644 $L = 3.4 \cdot 10^9 L_{\odot}$ **Density-structure** The present method is already on a quantitative level! Fe II, S II Fe II, Co II SN 1994E SN1992A (24.Jan.1992, t~2d) Model (42W) 45.0 6.0e-14 -Fell F, [erg cm<sup>-2</sup>s<sup>-1</sup>A<sup>-</sup>] Log [F.] (in arbitrary units) 4.0e-14 01 43.0 2.0e-14 W7 23d $\eta = 1.3$ 42.0 4000 6000 8000 10000 12000 3000 7000 8000 1000 2000 10000 Wavelength (Angstroms) Wavelength (Å) Pauldrach et al., 1996, A&A 312, 525 Nugent et al., 1997, ApJ. 485, 812 Sauer and Pauldrach, 2002, Nuclear Astrophysics, MPA/P13

## Results: Synthetic and observed spectra of Supernovae of Type Ia



Diffusion approximation - standard for inner boundary :

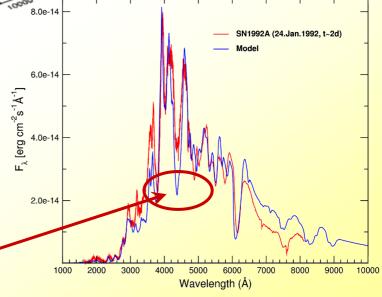
$$I_{\tau_{\text{max}}}^{+} = (B_{\nu} + 3\mu H_{\nu}^{0})(\tau_{\text{max}})$$

$$H_{\nu}^{0}(R) = -\frac{1}{3} \frac{dB_{\nu}}{\kappa_{\nu} dr}$$

Logarithm of the total opacity versus velocity and wavelength of a SN la model

Sauer and Pauldrach, 2005, A&A

Main features of the spectrum
are fairly well reproduced
especially in the blue and UV range
However,
Some features are "too deep"



HST spectrum of 24. Jan. 1992 from Kirshner et al., 1993, ApJ 415, 589 Model Sauer and Pauldrach, 2002, Nuclear Astrophysics, MPA/P13

## Modified diffusion approximation for Supernovae of Type Ia

#### Modified diffusion approximation for inner boundary:

$$I_{\tau_{\text{max}}}^{+} = (J_{\nu} + 3\mu H_{\nu}^{0})(\tau_{\text{max}})$$

after some calculation

$$\Rightarrow H_{\nu}^{0}(R) = -\frac{1}{3} \left( \frac{dB_{\nu}}{\kappa_{\nu} dr} + C_{\nu}^{*}(\tau_{\nu}) \right) - \frac{df_{\nu}}{\kappa_{\nu} dr} \left[ J_{\nu} \left( \frac{3\tau_{\nu}}{\kappa_{\nu} R} + 1 \right) + \tau_{\nu} \left( \frac{dB_{\nu}}{\kappa_{\nu} dr} + C_{\nu}^{*}(\tau_{\nu}) \right) \right]$$

and

$$J_{
u}({ au}_{
u}) = B_{
u}({ au}_{
u}) + rac{C_{
u}}{qf(
u, { au})r^2} e^{-\sqrt{eta_{
u}/q^2 f_{
u}} s_{
u}}$$

with

$$C_{\nu}^{*}(\tau_{\nu}) = \frac{B_{\nu} - J_{\nu}}{f_{\nu}} \left( \frac{3f_{\nu} - 1}{r\kappa_{\nu}} + \frac{df_{\nu}}{\kappa_{\nu}dr} \right) + \sqrt{\beta_{\nu}/q^{2}f_{\nu}} \frac{C_{\nu}}{f_{\nu}r^{2}} e^{-\sqrt{\beta_{\nu}/q^{2}f_{\nu}} s_{\nu}}$$

and

$$C_{\nu} = \frac{-\tilde{B}_{\nu}(\tau_{0\nu}) + j_{\nu}(\tau_{0\nu}) \frac{d[qf_{\nu}(\tau)\tilde{B}_{\nu}(\tau)]}{ds_{\nu}} \Big|_{\tau_{0\nu}}}{(1/qf_{\nu})_{\tau_{0\nu}} + j_{\nu}(\tau_{0\nu}) \sqrt{\beta_{\nu}/q^{2}f_{\nu}}} e^{\sqrt{\beta_{\nu}/q^{2}f_{\nu}}s_{0\nu}}$$

where  $j_{\nu}(\tau_{0\nu}) = \frac{J_{\nu}(\tau_{0\nu})}{H_{\nu}(\tau_{0\nu})} \sim 2$ 

with the Eddington factor  $f_{\nu}=K_{\nu}/J_{\nu}$  and the sphericality factor  $r^2q(r)=\exp\left(\int\limits_1^r \frac{3f(r')-1}{r'f(r')}dr'\right)$  with the Thomson opacity:  $\kappa_e=n_e\sigma_e$ ,  $\eta_e=\kappa_e J$  and  $\beta_{\nu}(\tau_{\nu})=\frac{\kappa_{\nu}}{\kappa_{\nu}+\kappa_{\nu}}$ 

and  $S = \frac{\eta}{\kappa + \kappa} + \frac{\kappa_e}{\kappa + \kappa} \frac{\eta_e}{\kappa} = S^t + \kappa_{Th} J$ 

 $\mathrm{and}\, ds_\nu = -q\kappa_\nu dr, s_\nu = \int\limits_r^{r_{\mathrm{max}}} q\kappa_\nu dr \ \ \mathrm{and}\, d\tau_\nu = -\kappa_\nu dZ, \tau_\nu = \int\limits_Z^{z_{\mathrm{max}}} \kappa_\nu dZ$ 

moments of the radiation field

$$\text{Oth:} J = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) d\mu \text{ , } \tilde{J} = r^2 J \text{ 1st:} H = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu d\mu \text{ , } \tilde{H} = r^2 H \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2} \int\limits_{\mu = -1}^{+1} I(\mu) \mu^2 d\mu \text{ , } \tilde{K} = r^2 K \text{ 2nd:} K = \frac{1}{2}$$

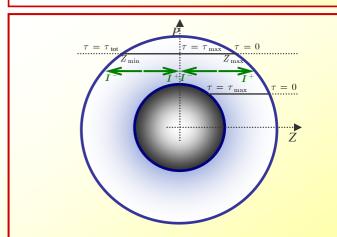
Diffusion approximation - standard for inner boundary :

$$I_{\tau_{\text{max}}}^{+} = (B_{\nu} + 3\mu H_{\nu}^{0})(\tau_{\text{max}})$$
$$H_{\nu}^{0}(R) = -\frac{1}{3} \frac{dB_{\nu}}{\kappa_{\nu} dr}$$

Modifications to the inner boundary that allow deviations of the radiation field from thermal equilibrium conditions is considered.

The dominating contribution from Thomson scattering

is explicitly taken into account



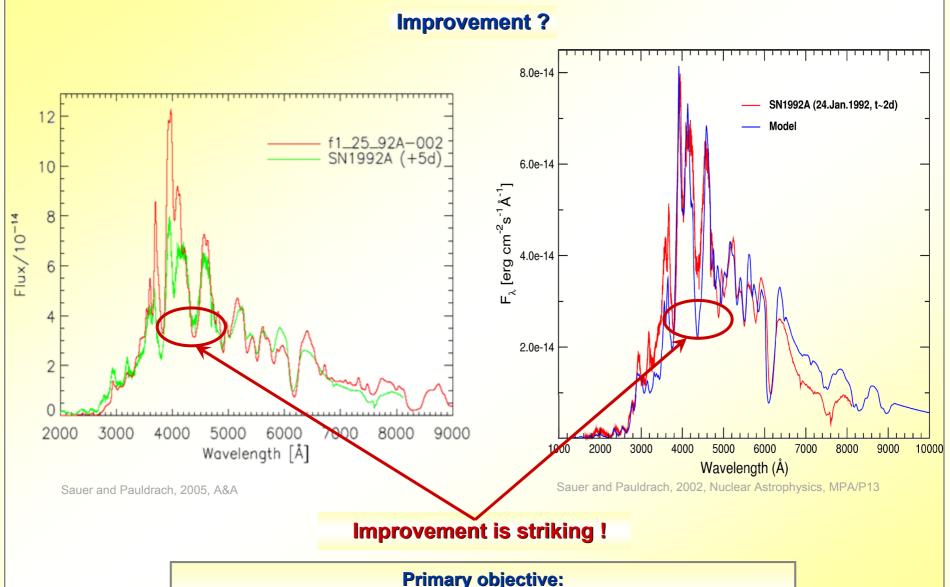
Equation of radiative transfer

$$\frac{d^{2}(I^{+} + I^{-})(\tau)}{d\tau^{2}} = (I^{+} + I^{-})(\tau) - S_{\nu}(\tau)$$

iterated with moment equation of radiative transfer

$$\frac{d^2(qf\tilde{J})}{ds^2} = \frac{1}{q}(\tilde{J} - \tilde{S})$$

## Results: Synthetic and observed spectra of Supernovae of Type Ia

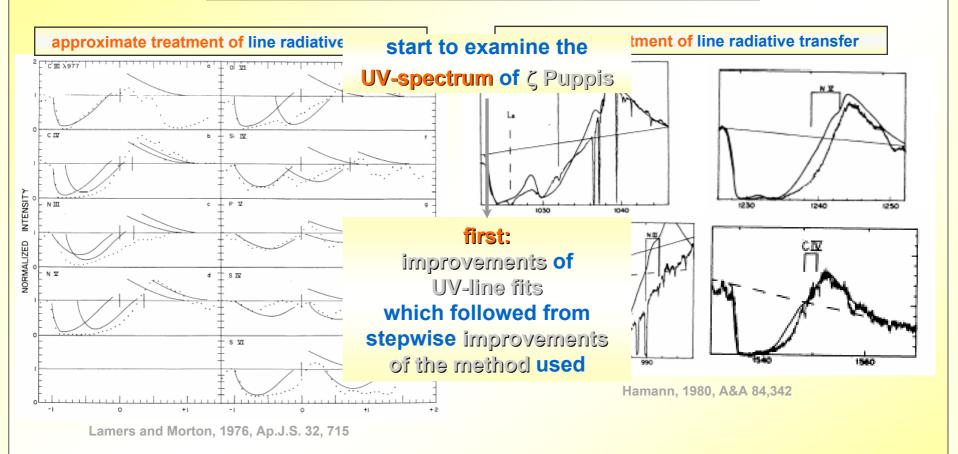


## **Primary objective:**

search for spectral differences between local and distant SNe la! Are SNe la standard candles in a cosmological sense?

## Empirical wind models for the O-supergiant & Puppis

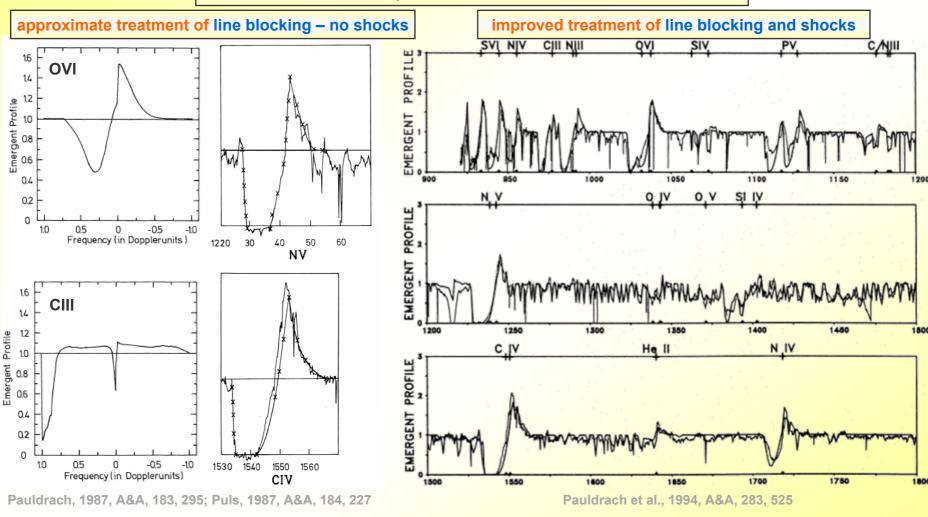
dynamical structure, occupation numbers, and temperature structure are assumed



A more consistent treatment of expanding atmospheres is required for a detailed analysis

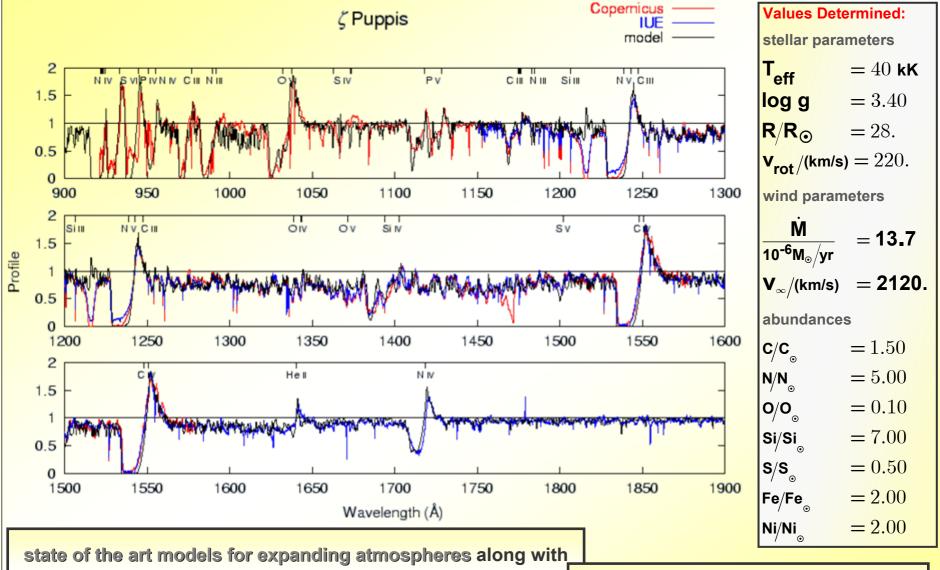
# Consistent wind models of the O-supergiant & Puppis

line blanketing neglected - approximated temperature structure approximate treatment of line blocking simple treatment of shocks



A more realistic treatment of expanding atmospheres is required for a detailed analysis

# Detailed analysis of the hot O-supergiant & Puppis - final step

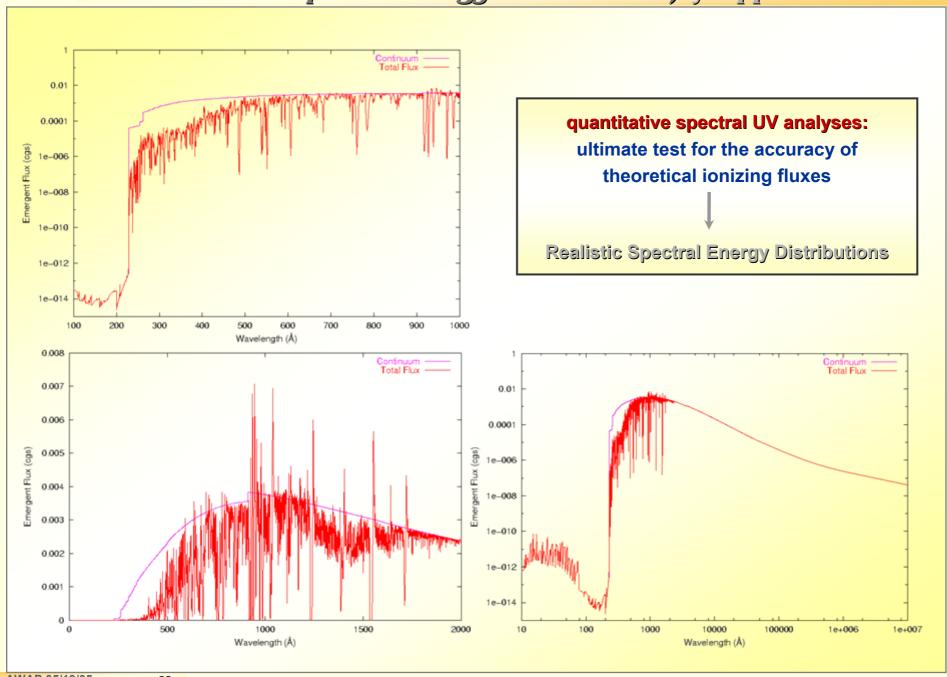


state of the art models for expanding atmospheres along with spectrum synthesis techniques allow the determination of stellar parameters, wind parameters, and abundances

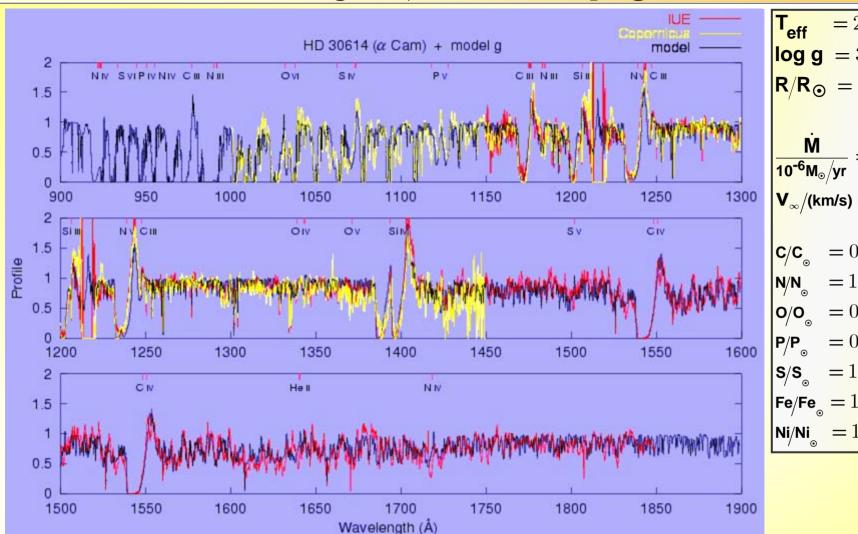
present method of quantitative spectral UV analyses of hot stars leads to realistic models!

Pauldrach, 2003, Reviews in Modern Astronomy, Vol. 16

# Realistic Spectral Energy Distribution of & Puppis



# Detailed analysis of the cool O-supergiant a Came



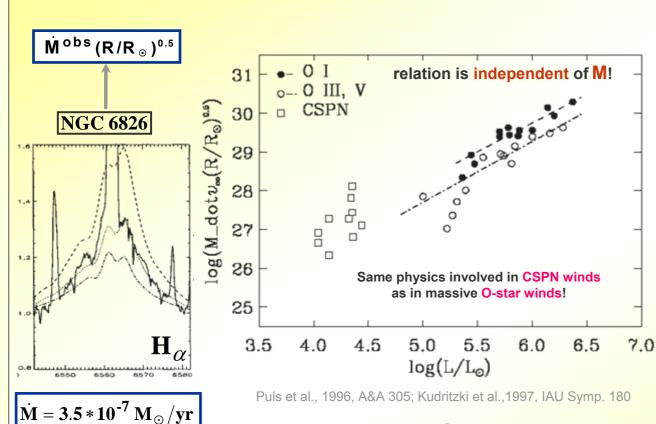
= 29 kK= 3.00 $R/R_{\odot} = 30$  $V_{\infty}/(\text{km/s}) = 1500.$ = 0.05=1.00= 0.30=0.05=1.00 ${
m Fe/Fe}_{\odot}=1.50$ =1.50

Spectrum synthesis technique allows the determination of T<sub>eff</sub> (within 1000 K), log g (within 0.05), R (within 1.0), and the abundances (within a range of 20%)

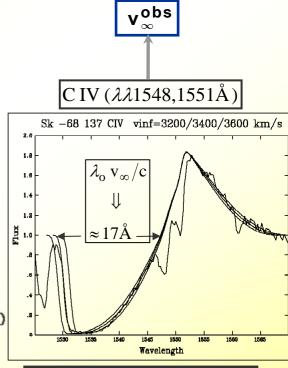
## Wind momentum luminosity relation

winds driven by radiation → mechanical momentum of wind flow mostly a function of photon momentum

 $\dot{M}v_{\infty}R^{1/2} \propto L^{3/2}$ 



independent tool for measuring extragalactic distances up to Virgo and Fornax

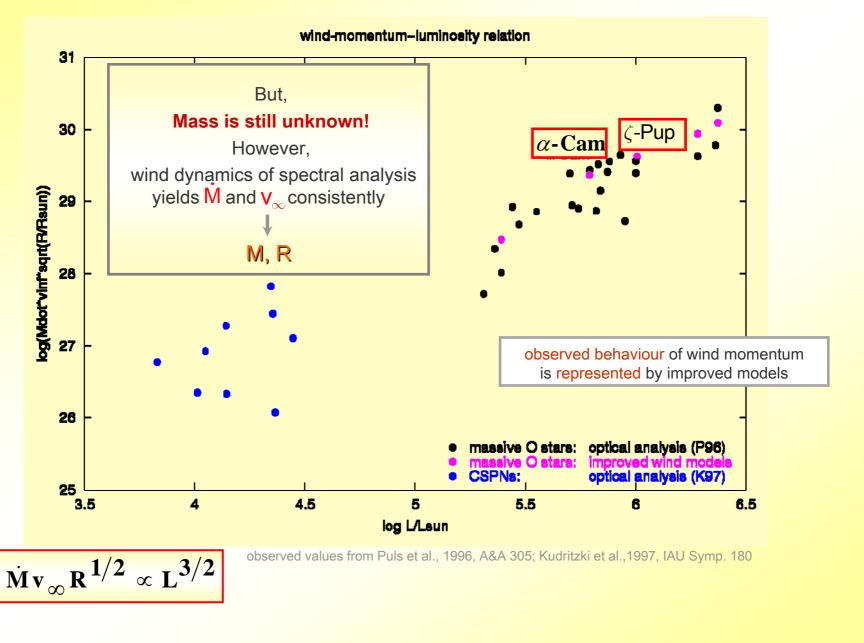


v<sub>∞</sub> can be measured directly

 $\dot{M} = 2.5 * 10^{-7} M_{\odot}/yr$ 

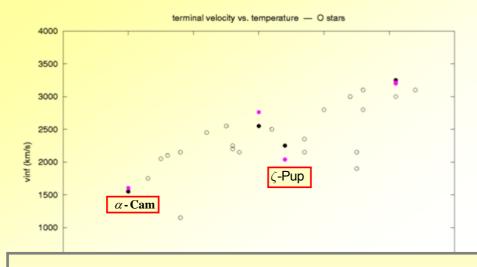
 $\dot{M} = 1.9 * 10^{-7} M_{\odot}/yr$ 

# Wind momentum luminosity relation of massive O-stars

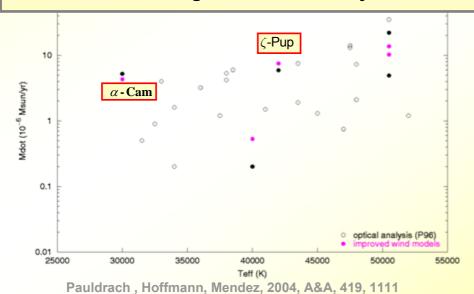


Pauldrach, Hoffmann, Mendez, 2004, A&A, 419, 1111

# Wind properties of massive O-stars



Realistic models are characterized by quantitative spectral UV analysis calculated along with consistent dynamics



$$\mathbf{v}_{\infty} \propto \mathbf{v}_{\mathbf{esc}} = \left(\frac{2\mathbf{GM}}{\mathbf{R}}(1-\Gamma)\right)^{1/2}$$

predicted terminal velocities
of improved models
agree within 10% with observed values

predicted mass-loss-rates of improved models agree within a factor of 2 with observed optical values

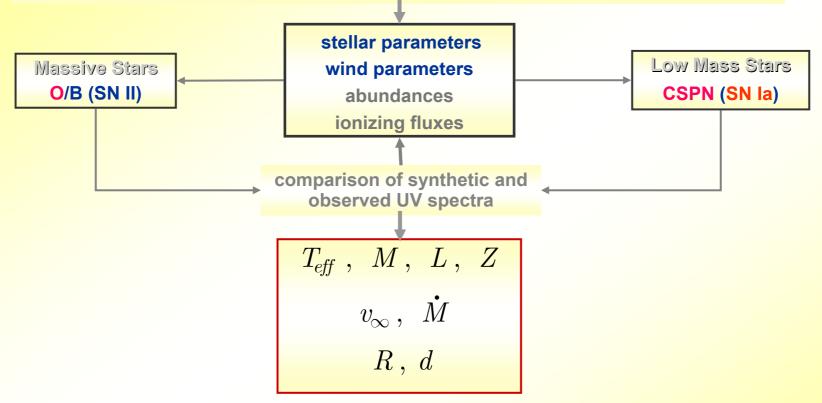
 $\dot{\mathbf{M}} \propto \mathbf{L}$ 

M,R

## Conclusions

Consistent state of the art models for expanding atmospheres of hot stars

a diagnostic tool\* with great astrophysical potential?!



Analysis of integrated synthetic spectra of galaxies based on population synthesis is feasible stellar content, IMF, SFR, Z, burst age

<sup>\*</sup> download of the program package -including an easy to use interface- is possible from http://www.usm.uni-muenchen.de/people/adi/adi.html