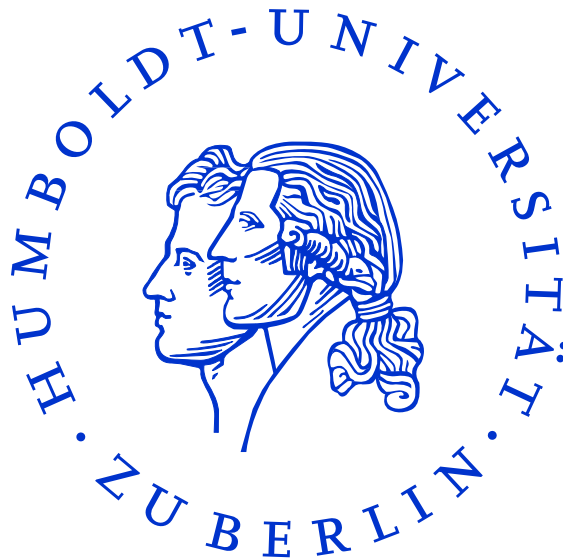


# Dynamical systems perturbed by big jumps: meta-stability and non-local search\*

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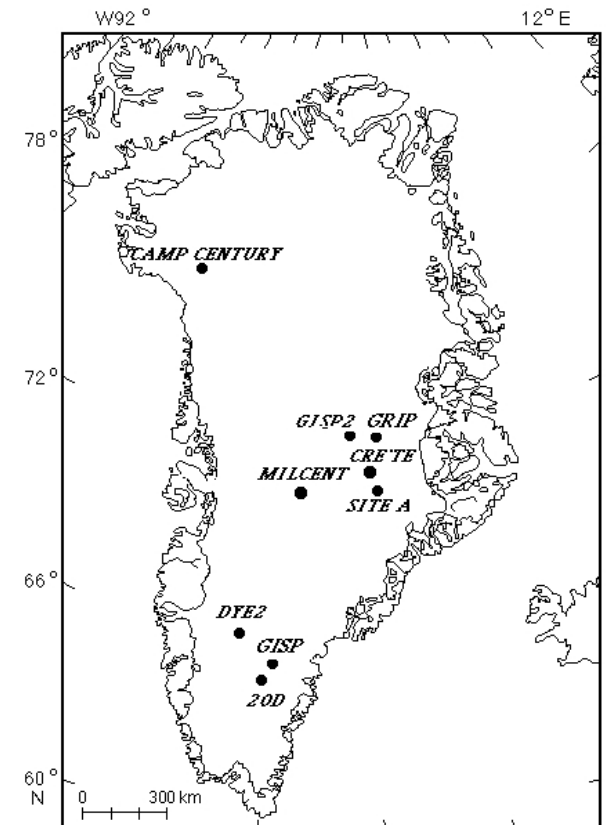
\*Supported by the DFG research projekt *Stochastic Dynamics of Climate States*

# 1. Motivation

Greenland ice-core data allows to reconstruct Earth's climate up to 200.000 years before present.

International projects:

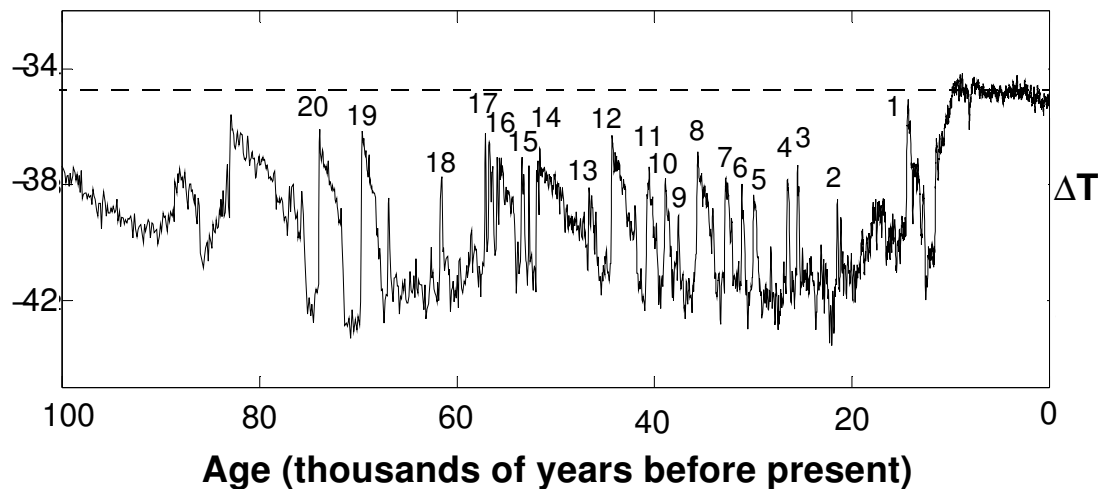
GRIP (3028 m), GISP2 (3053.44 m), NGRIP (3084.99 m)



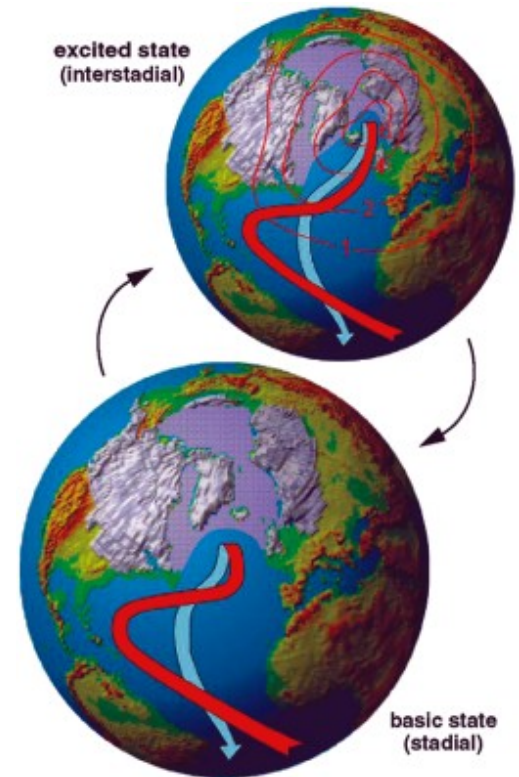
## 2. Paleo proxy data. Dansgaard–Oeschger events

Paleo data proxies: oxygen isotopes, dust, volcanic markers etc.

Global climate during the last glacial ( $\sim 120\,000 - 10\,000$  b. p.) has experienced at least 20 abrupt and large–amplitude shifts (Dansgaard–Oeschger events).



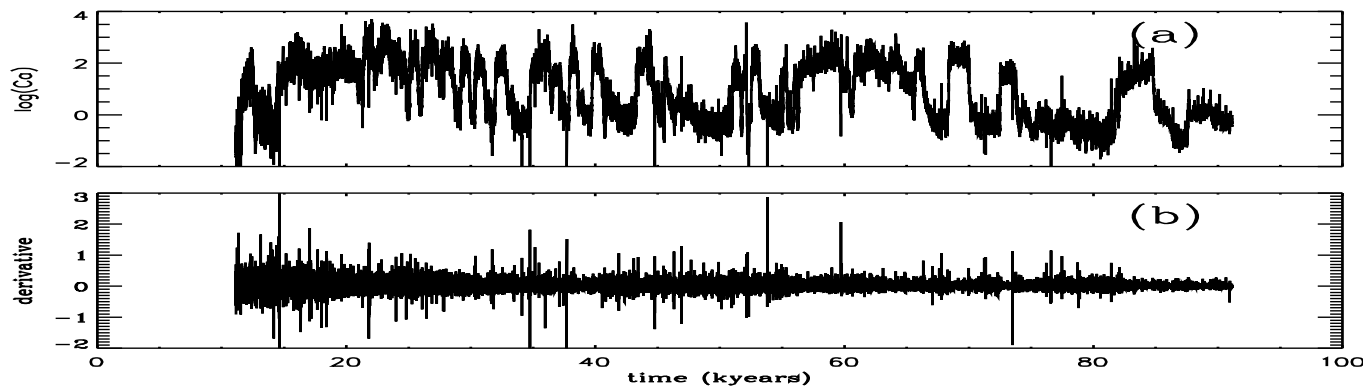
- rapid warming by 5 – 10 °C within at most a few decades
- plateau phase with slow cooling lasting several centuries
- rapid drop to cold stadial conditions



Simulations: Ganopolsky&Rahmstorf,  
Potsdam Institute for Climate Impact  
Research

### 3. Paleo proxy data: detailed look.

The calcium (Ca) signal from the GRIP ice-core:  
about 80,000 data points from 11 kyr to 91 kyr before present.



Typical interjump time: 1000 – 2000 years, mean waiting time  $\sim 1470$  years

#### What triggers the transitions?

Langevin equation for climate dynamics

$$dX(t) = -U'(X(t)) dt + \varepsilon dL(t)$$

$U$  – double-well potential, wells correspond to the climate states.

P. Ditlevsen (*Geophys. Res. Lett.* 1999): histogram analysis of the noise.

Noise  $L$  has  $\alpha$ -stable component with  $\alpha \approx 1.75$ .

## 4. The random perturbation $L$ . Lévy processes

$L$  is a Lévy process:

- independent increments  
 $L(s) \perp L(t) - L(s)$
- stationary increments  
 $L(t) - L(s) \stackrel{d}{=} L(t - s)$
- $L(0) = 0$
- stochastically continuous (no fixed jumps)
- right-continuous with left limits

$B$  is a Brownian motion:

- independent increments
- stationary increments
- $B(0) = 0$
- $B(t) - B(s) \sim \mathcal{N}(0, t - s)$
- continuous paths

Description via Fourier transform:  $\varphi_t(\lambda) = \mathbf{E}e^{i\lambda L(t)}$

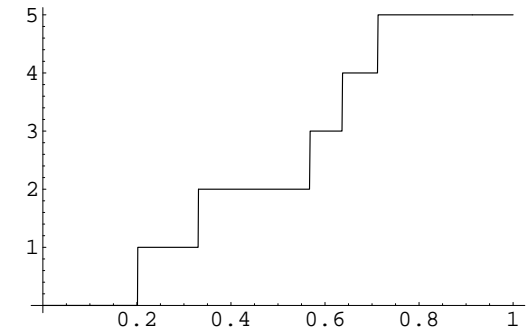
$$L(t + s) = L(t) + (L(t + s) - L(t)) \stackrel{d}{=} L(t) + L'(s)$$

$$\varphi_{t+s}(\lambda) = \varphi_s(\lambda) \cdot \varphi_t(\lambda) \quad \Rightarrow \quad \varphi_t(\lambda) = \mathbf{E}e^{i\lambda L(t)} = e^{t\Phi(\lambda)}$$

## 5. Examples

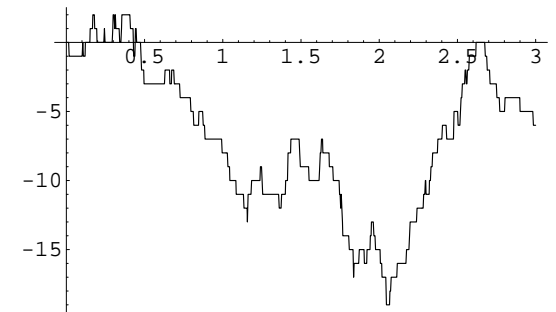
**Poisson process**, intensity  $c > 0$

$$\begin{aligned} \mathbf{E}e^{i\lambda L(t)} &= \sum_{k=1}^{\infty} e^{i\lambda k} e^{-ct} \frac{(ct)^k}{k!} = e^{c(e^{i\lambda} - 1)} \\ &= \exp \left[ t \int (e^{i\lambda y} - 1) c \delta_1(dy) \right] \end{aligned}$$



**Compound Poisson process**,  $c > 0$ , jumps  $\sim \sigma$

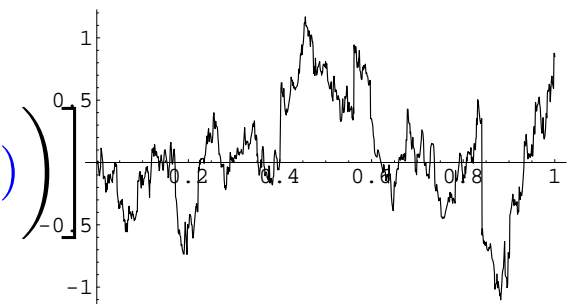
$$\mathbf{E}e^{i\lambda L(t)} = \exp \left[ t \int (e^{i\lambda y} - 1) c \sigma(dy) \right]$$



**General case**, Lévy–Hinchin Formula

$$\mathbf{E}e^{i\lambda L(t)} = \exp \left[ t \left( -\frac{d}{2} \lambda^2 + i\mu\lambda + \int (e^{i\lambda y} - 1 - i\lambda y \mathbb{I}\{|y| \leq 1\}) \nu(dy) \right) \right]$$

$$d \geq 0, \mu \in \mathbb{R}, \nu(\{0\}) = 0, \int (y^2 \wedge 1) \nu(dy) < \infty$$



## 6. $\alpha$ -stable Lévy processes (Lévy flights)

Jump measure  $\nu(y) = \frac{1}{|y|^{1+\alpha}}$ ,  $\alpha \in (0, 2)$ .

$$\mathbf{E}e^{i\lambda L(t)} = \exp\left(t \int (e^{i\lambda y} - 1 - i\lambda y \mathbb{I}\{|y| \leq 1\}) \frac{dy}{|y|^{1+\alpha}}\right)$$

Markov process with a generator

$$Af(x) = \int [f(x+y) - f(x) - yf'(x)\mathbb{I}\{|y| \leq 1\}] \frac{dy}{|y|^{1+\alpha}},$$

$$Af(x) = -(-\Delta)^{\alpha/2} f(x) \quad \text{— fractional Laplacian}$$

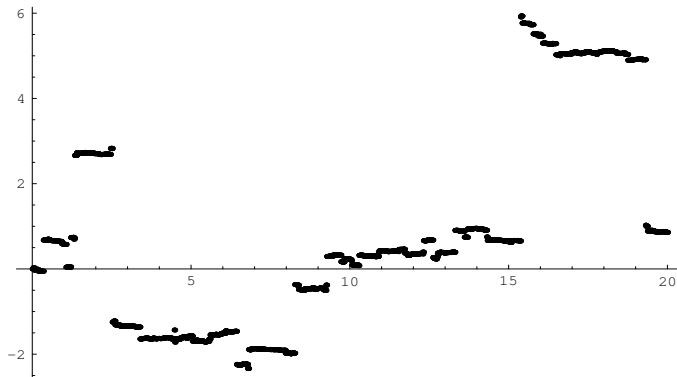
$L$  is a **purely jump** process.

Control of jumps:  $J(t) = L(t) - L(t-)$  the jump size at time  $t > 0$ .

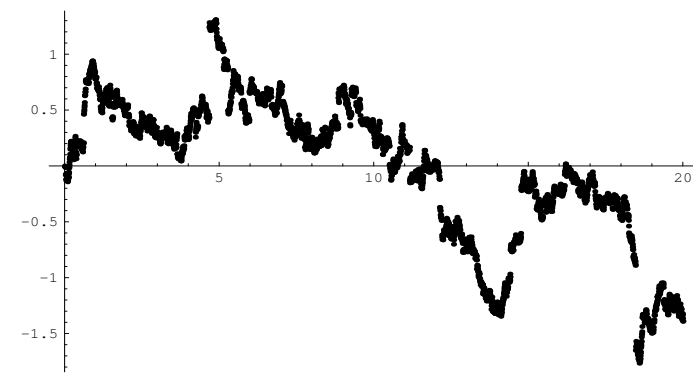
$$N(t, A) = \#\{s \in (0, t] : J(s) \in A\} \sim \mathbf{Poisson} \left( t \int_A \frac{dy}{|y|^{1+\alpha}} \right)$$

$$\text{For example: } \#\{s \in (0, t] : J(s) \geq 1\} \sim \mathbf{Poisson} \left( t \int_1^\infty \frac{dy}{|y|^{1+\alpha}} \right)$$

## 7. $\alpha$ -stable Lévy processes II



$$\alpha = 0.75$$



$$\alpha = 1.75$$

**Length** of sample paths: finite for  $\alpha \in (0, 1)$  and infinite for  $\alpha \in [1, 2)$ .

**Self-similarity:**  $(L(ct))_{t \geq 0} = (c^{1/\alpha} L(t))_{t \geq 0}$ . Hurst parameter  $\mathbb{H} = \frac{1}{\alpha} \geq \frac{1}{2}$ .

**Heavy tails:**  $\mathbb{P}(L(t) \geq u) \approx \frac{c}{u^\alpha}, u \rightarrow \infty$ .

Explicit form of the **Fourier transform**  $\mathbb{E}e^{i\lambda L(t)} = e^{-c(\alpha)|\lambda|^\alpha t}$ ,

$\alpha = 1$       Cauchy process       $\frac{1}{\pi} \frac{1}{1+x^2}$

$\alpha = 2$       Brownian motion       $\frac{1}{\sqrt{2\pi}} e^{-\frac{x^2}{2}}$



## 8. Object of study. Simple system with Lévy perturbation

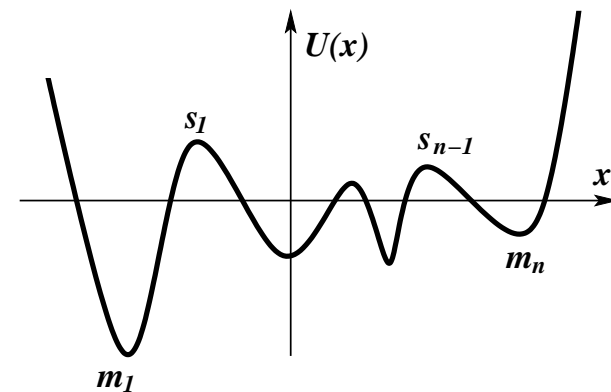
Small noise ( $\varepsilon \downarrow 0$ ) asymptotics of solutions of SDE

$$X_x^\varepsilon(t) = x - \int_0^t U'(X_x^\varepsilon(s)) ds + \varepsilon L(t), \quad \varepsilon \downarrow 0.$$

- $L$  —  $\alpha$ -stable (symmetric) Lévy process (maybe + Brownian motion)

### Regular $n$ -well potential

(smooth,  $U''(m_i), U''(s_i) \neq 0$ ,  
 $|U'(x)| > |x|^{1+\delta}$ ,  $x \rightarrow \pm\infty$ )

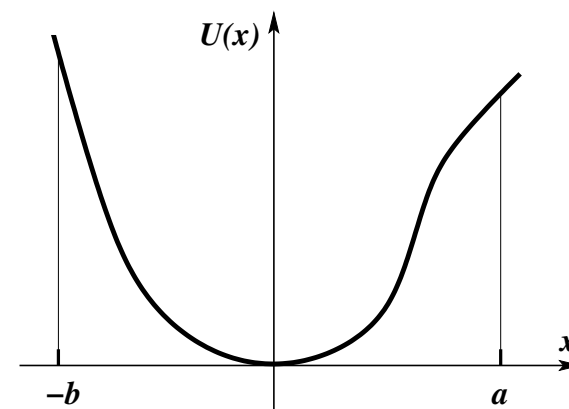


### **Meta-stable behaviour**

Transitions between the wells

### Regular one-well potential

(smooth,  $U''(0) > 0$ )



### **Exit time:**

$$\sigma_x(\varepsilon) = \inf\{t \geq 0 : X_x^\varepsilon(t) \notin [-b, a]\}$$

$a, b < \infty$  ( $b = \infty$ )

P. Imkeller & I. Pavlyukevich

*Stoch. Proc. Appl.* 116, 2006; *J. Phys. A: Math. Gen.* 39, 2006; *ESAIM: P&S*, 2008

## 9. What is known

**Freidlin, Wentzell** (Random Perturbations of Dynamical Systems, 1979):

*Gaussian perturbations:*

$$\hat{X}_x^\varepsilon(t) = x - \int_0^t U'(\hat{X}_x^\varepsilon(s)) ds + \varepsilon W(t), \quad \varepsilon \downarrow 0.$$

Perturbation:  $\varepsilon W$  — standard Brownian motion of small amplitude (in  $\mathbb{R}^d$ ).

*Locally infinitely divisible perturbations* leading to Gaussian  $\varepsilon L^\varepsilon$  with jump part

$$A^\varepsilon f(x) = \int_{\mathbb{R}^d \setminus \{0\}} \left[ f(x + \varepsilon y) - f(x) - \langle \varepsilon y, \nabla f(x) \rangle \right] \frac{\nu(x, dy)}{\varepsilon}$$

Lévy measure  $\nu$  has all exponential moments.

For example,  $\varepsilon L^\varepsilon(t) = \varepsilon \pi^{1/\varepsilon}(t) - t$ ,  $\pi$  — standard Poisson process.

*Effects of heavy tails*

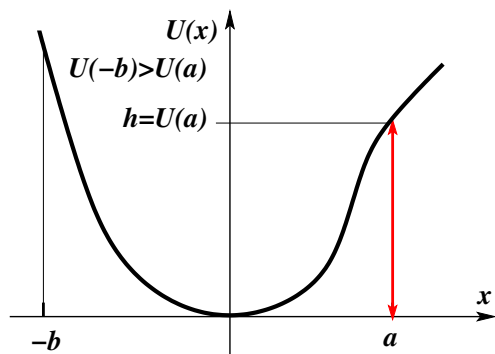
**Godovanchuk** (Theor. Prob. Appl. 26, 1982):

‘Large deviations’ for Markov processes with heavy jumps.

**Samorodnitsky, Grigoriou** (Stoch. Proc. Appl. 105, 69–97, 2003):

Tails of solutions of SDEs driven by Lévy processes with power tails.

# 10. First exit: Gaussian vs. Lévy



$$\hat{X}_x^\varepsilon(t) = x - \int_0^t U'(\hat{X}_x^\varepsilon(s)) ds + \varepsilon W(t)$$

Large deviations (Freidlin–Wentzell):

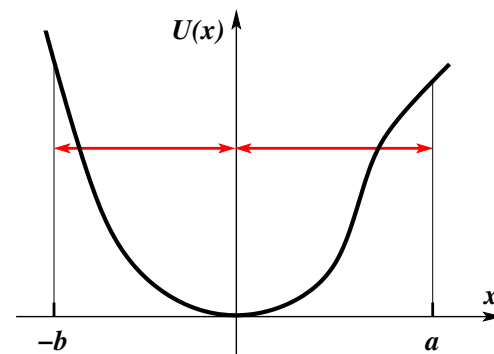
$$\mathbf{P}(e^{(2h-\delta)/\varepsilon^2} < \hat{\sigma}_x(\varepsilon) < e^{(2h+\delta)/\varepsilon^2}) \rightarrow 1$$

Mean exit time (Kramers, Day, Bovier):

$$\mathbf{E}\hat{\sigma}_x(\varepsilon) \approx \frac{\varepsilon\sqrt{\pi}}{U'(a)\sqrt{U''(0)}} \exp\left(\frac{2h}{\varepsilon^2}\right)$$

Exponential exit (Day, Bovier)

$$\mathbf{P}\left(\frac{\hat{\sigma}_x(\varepsilon)}{\mathbf{E}\hat{\sigma}_x(\varepsilon)} > u\right) \rightarrow \exp(-u)$$



$$X_x^\varepsilon(t) = x - \int_0^t U'(X_x^\varepsilon(s)) ds + \varepsilon L(t)$$

$$\mathbf{P}\left(\frac{1}{\varepsilon^{\alpha-\delta}} < \sigma_x(\varepsilon) < \frac{1}{\varepsilon^{\alpha+\delta}}\right) \rightarrow 1$$

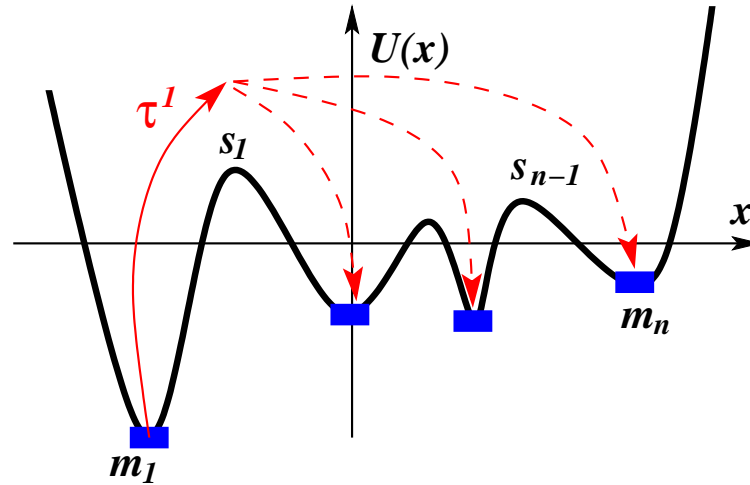
$$\mathbf{E}\sigma_x(\varepsilon) \approx \frac{\alpha}{\varepsilon^\alpha} \left[\frac{1}{a^\alpha} + \frac{1}{b^\alpha}\right]^{-1}$$

$$\mathbf{P}\left(\frac{\sigma_x(\varepsilon)}{\mathbf{E}\sigma_x(\varepsilon)} > u\right) \rightarrow \exp(-u)$$

# 11. Transitions

For small  $\Delta > 0$  denote  $\mathbf{B}_i = \{y : |y - m_i| \leq \Delta\}$  and

$$\tau_x^i(\varepsilon) = \inf\{t \geq 0 : X_x^\varepsilon(t) \in \cup_{j \neq i} B_j\}$$



**Theorem 1.** For  $x \in B_i$  the following holds as  $\varepsilon \rightarrow 0$ :

$$\mathbf{P}(\varepsilon^\alpha \tau_x^i(\varepsilon) > u) \rightarrow e^{-q_i u},$$

$$\mathbf{E} \tau_x^i(\varepsilon) \approx \frac{1}{\varepsilon^\alpha q_i},$$

$$q_i = \int_{\mathbb{R} \setminus (s_{i-1}, s_i)} \frac{dy}{|y - m_i|^{1+\alpha}},$$

$$\mathbf{P}(X^\varepsilon(\tau_x^i(\varepsilon)) \in B_j) \rightarrow \frac{q_{ij}}{q_i},$$

$$q_{ij} = \int_{(s_{j-1}, s_j)} \frac{dy}{|y - m_i|^{1+\alpha}}, \quad i \neq j.$$

## 12. Meta–stability

**Theorem 2.** *Let  $x \in (s_{j-1}, s_j)$  and  $t > 0$ . Then*

$$X_x^\varepsilon \left( \frac{t}{\varepsilon^\alpha} \right) \xrightarrow{d} Y_{m_j}(t),$$

*where  $Y$  is a Markov chain on  $\{m_1, \dots, m_n\}$  with a generator  $Q = (q_{ij})$ ,  $q_{ii} = -q_i$ .*

**Remark.**  $Y$  has the invariant measure

$$\pi(dy) = \sum_{j=1}^n \pi_j \delta_{m_j}(dy),$$

$$\pi_j > 0,$$

where  $Q^T \pi = 0$ .

### 13. Meta-stable behaviour. Gaussian case

$$\hat{X}_x^\varepsilon(t) = x - \int_0^t U'(\hat{X}_x^\varepsilon(s)) ds + \varepsilon W(t)$$

Different life times for each well:  $\mathbf{E}\tau_x^i(\varepsilon) \sim \exp(V_i/\varepsilon^2)$

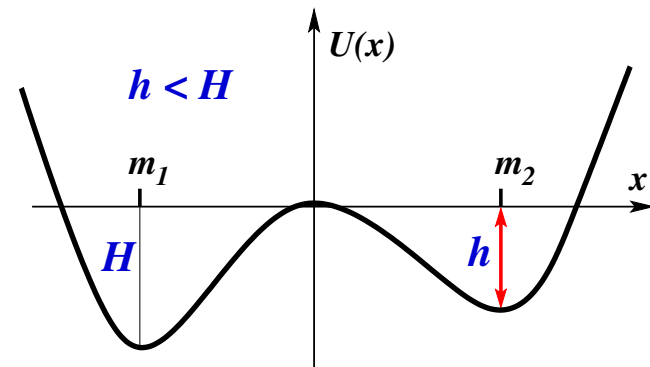
$\Rightarrow n - 1$  critical exponents  $0 = \lambda_0 < \lambda_1 < \dots < \lambda_{n-1}$  such that for  $\Lambda \in (\lambda_i, \lambda_{i+1})$  the process  $\hat{X}_x^\varepsilon \left( te^{\Lambda/\varepsilon^2} \right)$  converges to  $\mu(\Lambda, x)$

**Double-well potential:** critical time scale:

$$\lambda_\varepsilon \approx \frac{2\pi}{\sqrt{|U''(0)|U''(m_2)|}} \exp(2h/\varepsilon^2)$$

Convergence of fin. dim. distrib.:  
(Kipnis, Newman; Mathieu)

$$\hat{X}^\varepsilon(\lambda_\varepsilon t) \rightarrow \hat{Y}(t)$$



Generator of  $\hat{Y}$

$$\begin{pmatrix} 0 & 0 \\ 1 & -1 \end{pmatrix}, \text{ and } \hat{Y}(0) = \begin{cases} m_1, & \text{if } x < 0, \\ m_2, & \text{if } x > 0. \end{cases}$$

## 14. Doubts and questions

$$X_x^\varepsilon(t) = x - \int_0^t U'(X_x^\varepsilon(s)) ds + \varepsilon L(t), \quad \varepsilon \downarrow 0.$$

- Why  $\alpha$ -**stable** (non-Gaussian, heavy-tail) noise?
- Why **small** noise?

## 15. A natural system with small noise

Is there a Gaussian system with *a priori* small noise?

Simulated annealing (Kirkpatrick et al., Geman&Geman, Černý, ~ 1985):

$$\hat{Z}(t) = z - \int_0^t U'(\hat{Z}(s)) ds + \int_0^t \sigma(s) dW(s),$$

$\sigma(t)$  — ‘temperature’,  $\sigma(t) \rightarrow 0$  as  $t \rightarrow \infty$ .

ANNEALING: gradual cooling of steel (copper, glass etc.) in order to induce softness, to relieve internal stresses, to refine the crystalline structure.



Carbon steel:



before annealing

fully annealed at 920 °C





## 16. Gaussian diffusion: long-time behaviour

Stochastic optimisation: look for a **global minimum**  $m^*$  of  $U$ .

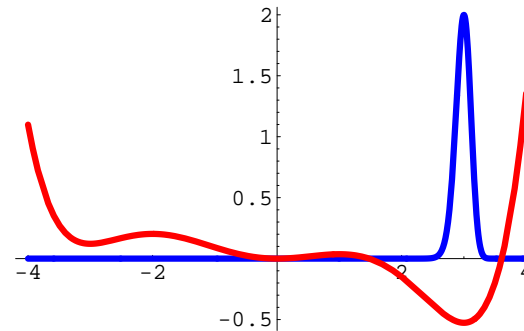
$$\hat{X}(t) = x - \int_0^t U'(\hat{X}(s)) ds + \varepsilon W(t)$$

Generator  $A_\varepsilon f = \frac{\varepsilon^2}{2} \Delta f - U' f'$

Invariant measure:  $A_\varepsilon^* \mu = 0$

$$\mu_\varepsilon(dx) = c(\varepsilon) e^{-2U(x)/\varepsilon^2} dx$$

$$\mu_\varepsilon \Rightarrow \delta_{m^*}, \quad \varepsilon \rightarrow 0$$



The spectrum  $\{-\Lambda_k^\varepsilon, \Psi_k^\varepsilon\}_{k \geq 0}$  of  $A_\varepsilon$  is discrete.

### SPECTRAL GAP

$$\Lambda_0^\varepsilon = 0, \quad \Psi_0^\varepsilon(x) = 1$$

$$\Lambda_1^\varepsilon \sim \exp(-\Theta/\varepsilon^2) \quad (\text{Friedman, Day, Bovier et al.})$$

$$\mathbf{E}_x f(\hat{X}^\varepsilon(t)) = \langle f, 1 \rangle_{L^2(\mu_\varepsilon)} 1(x) + e^{-\Lambda_1^\varepsilon t} \langle f, \Psi_1^\varepsilon \rangle_{L^2(\mu_\varepsilon)} \Psi_1^\varepsilon(x) + \dots$$

$$\text{Law}(\hat{X}^\varepsilon(t)) \rightarrow \mu_\varepsilon, \quad t \rightarrow \infty$$

## 17. Gaussian simulated annealing

$$\hat{Z}_z(t) = z - \int_0^t U'(\hat{Z}_z(s)) ds + \int_0^t \sigma(s) dW(s), \quad \sigma(t) = \left( \frac{\theta}{\ln(\lambda + t)} \right)^{1/2}.$$

**INTUITION:**  $\hat{Z}(t) \approx \hat{X}^{\sigma(t)}(t)$

**“FOURIER EXPANSION”:**

$$\mathbf{E}_{0,z} f(\hat{Z}(t)) \approx \langle f, 1 \rangle_{L^2(\mu_{\sigma(t)})} 1 + e^{-\Lambda_1^{\sigma(t)} t} \langle f, \Psi_1^{\sigma(t)} \rangle_{L^2(\mu_{\sigma(t)})} \Psi_1^{\sigma(t)} + \dots$$

**CONVERGENCE:**

$$\mu_{\sigma(t)} \Rightarrow \delta_{m^*},$$

$$\Lambda_1^{\sigma(t)} t \sim e^{-\Theta/\sigma(t)^2} t = \frac{t}{(t + \lambda)^{\Theta/\theta}} \rightarrow \begin{cases} \infty, & \theta > \Theta, \\ 0, & \theta < \Theta, \end{cases} \quad t \rightarrow +\infty,$$

$$\theta > \Theta \quad \Rightarrow \quad Z(t) \rightarrow m^*.$$

This works, see Chiang&Hwang&Sheu, Holley&Kusuoka&Stroock and others

**HOWEVER: very slow cooling, local search, unknown critical rate  $\Theta$ ...**

## 18. Modifications: big jumps

Szu and Hartley (*Phys. Lett. A* 122, 1987): “fast simulated annealing”

Discrete schema with Cauchy jumps  $\xi_k$

$$X_{kh} = X_{(k-1)h} + U'(X_{(k-1)h})h + \frac{\xi_k}{(k-1)h}$$

$$\Delta U = U(X_{kh}) - U(X_{(k-1)h})$$

Metropolis acceptance probability:

Accept always if  $\Delta U \leq 0$

Accept with probability  $\exp(-\Delta U \cdot kh)$

if  $\Delta U > 0$

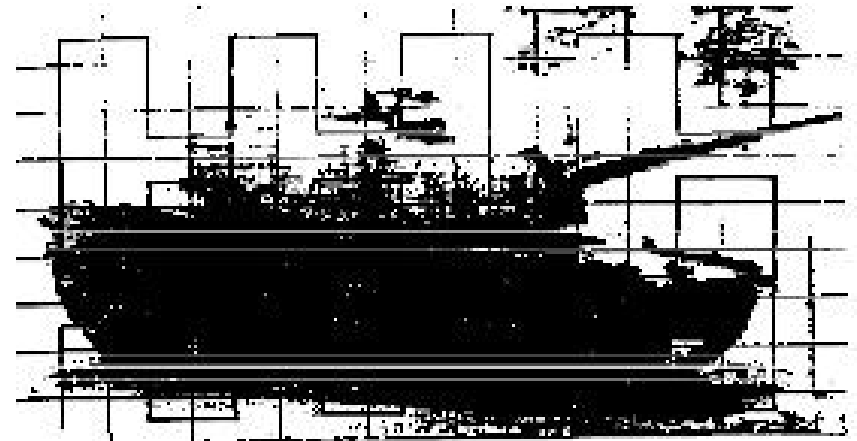
Non-local search

Fast algebraic cooling

Convergence ???

Probably works ???

Applications to image recognition  $\longrightarrow$



## 19. Simulated annealing with $\alpha$ -stable processes

Analogously to Gaussian simulated annealing:

$$Z(t) = z - \int_0^t U'(Z(s)) ds + \int_0^t \frac{dL(s)}{(\lambda + s)^\theta}, \quad \lambda > 0, \theta > 0.$$

Convergence of  $Z(t)$  as  $t \rightarrow \infty$ ?

Time-homogeneous case:

$$X^\varepsilon \left( \frac{t}{\varepsilon^\alpha} \right) \rightarrow Y(t)$$

$Y$  — Markov chain on  $\{m_1, \dots, m_n\}$  with generator  $Q = (q_{ij})$

$$q_{ij} = \int_{(s_{j-1}, s_j)} \frac{dy}{|y - m_i|^{1+\alpha}}, \quad q_{ii} = -q_i = - \int_{\mathbb{R} \setminus (s_{i-1}, s_i)} \frac{dy}{|y - m_i|^{1+\alpha}}$$

Law( $Y(t)$ )  $\rightarrow \pi$ ,  $Q^* \pi = 0$

$$Z(t) \approx X^{1/(\lambda+t)^\theta}(t) \approx Y \left( \frac{t}{(t + \lambda)^{\alpha\theta}} \right) \Rightarrow \begin{cases} \pi, & \alpha\theta < 1, \\ \text{no convergence,} & \alpha\theta > 1. \end{cases}$$

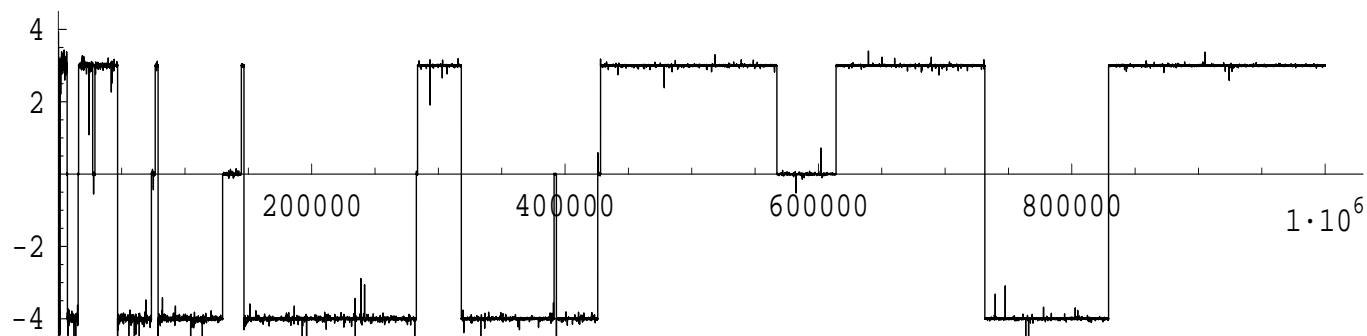
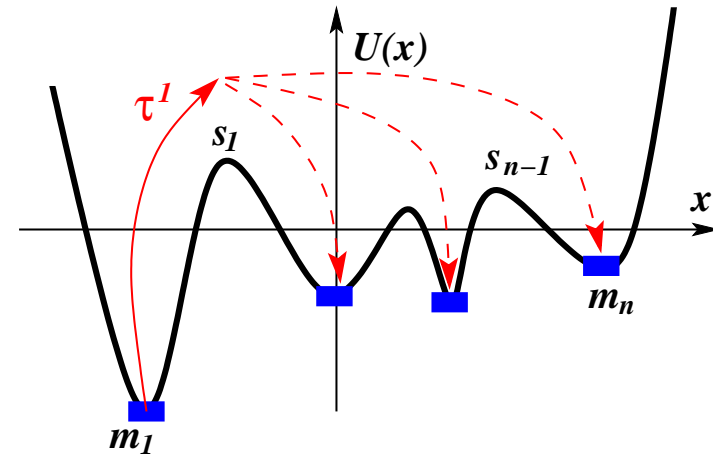
## Slow cooling: convergence

**Theorem 3. [slow cooling]** *Let  $\alpha\theta < 1$ . Then*

$$\mathbf{E}\tau^i(\lambda) \approx \frac{\lambda^{\alpha\theta}}{q_i},$$

$$\mathbf{P}(Z(\tau^i(\lambda)) \in B_j) \rightarrow \frac{q_{ij}}{q_i}, \quad \lambda \rightarrow \infty,$$

$$Z(t) \Rightarrow \pi, \quad t \rightarrow \infty.$$

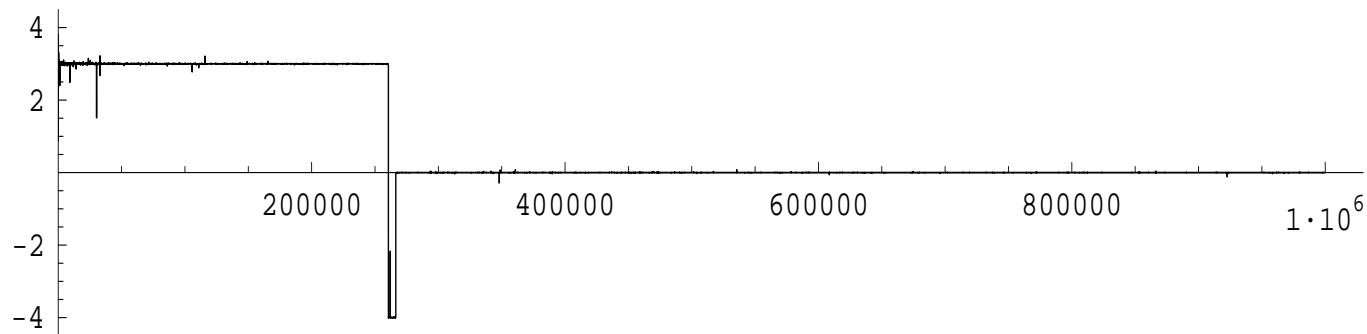


## 21. Fast cooling: trapping

**Theorem 4. [trapping]** *Let  $\alpha\theta > 1$ . Then*

$$\mathbf{P}(\tau^i(\lambda) < \infty) = \mathcal{O}(\lambda^{1-\alpha\theta}), \quad \lambda \rightarrow \infty,$$

$$\mathbf{E}\tau^i(\lambda) = \infty.$$



## 22. Gaussian vs. $\alpha$ -stable simulated annealing

Simulated annealing with  $\alpha$ -stable process:

- allows to determine the measure  $\pi$ , i.e. the **sizes** of potential wells
- the search is **non-local**
- polynomially **fast** cooling

**BUT:** how to find the **global** minimum of  $U$ ?

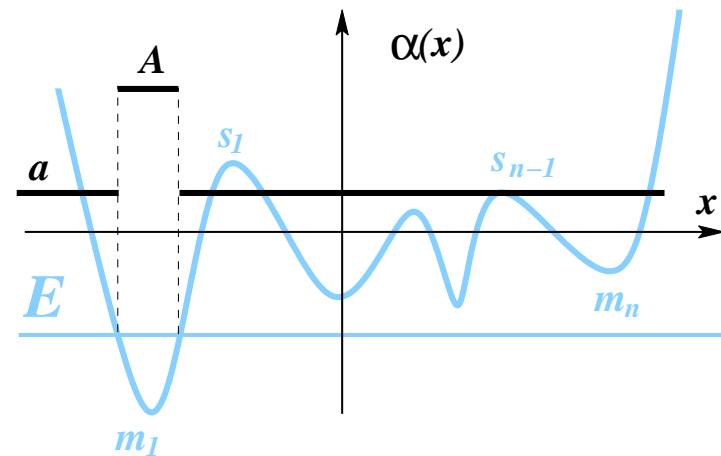
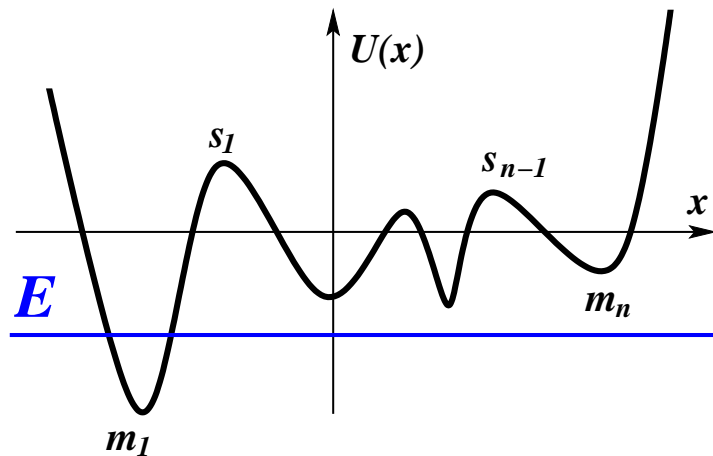
## 23. Simulated annealing with *stable-like* processes

In progress: I.P., *J. Comp. Phys.*, 2007

$$V(t) = v - \int_0^t U'(V(s-)) ds + \int_0^t \frac{dH(V(s-), s)}{(\lambda + s)^\theta}, \quad \lambda > 0, \theta > 0$$

$H$  a *stable-like process* with the jump measure

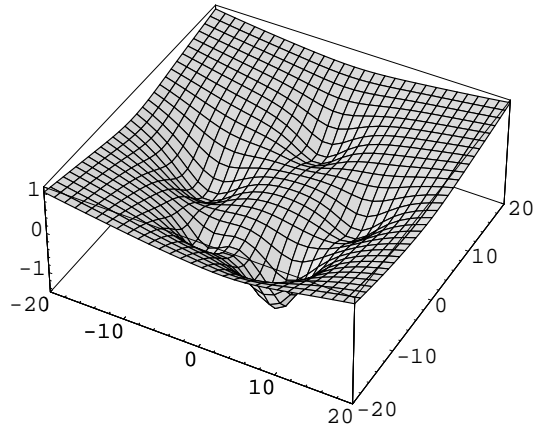
$$\nu_x(dy) = \frac{dy}{|y|^{1+\alpha(x)}}$$



$$\alpha(x) = \begin{cases} A, & U(x) \leq E, & \text{trapping,} \\ a, & U(x) > E, & \text{exit,} \end{cases} \quad 0 < a < A < 2, \quad a\theta < 1 < A\theta.$$



## 24. An example of non-local search

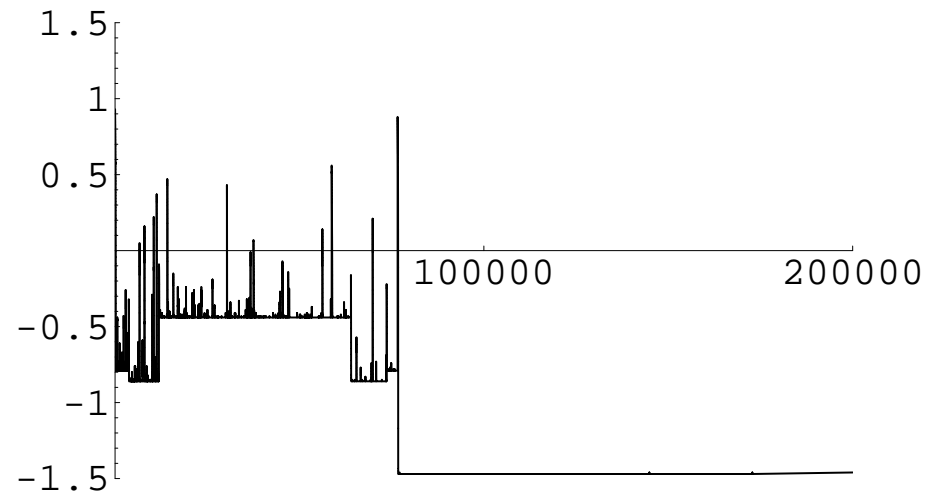
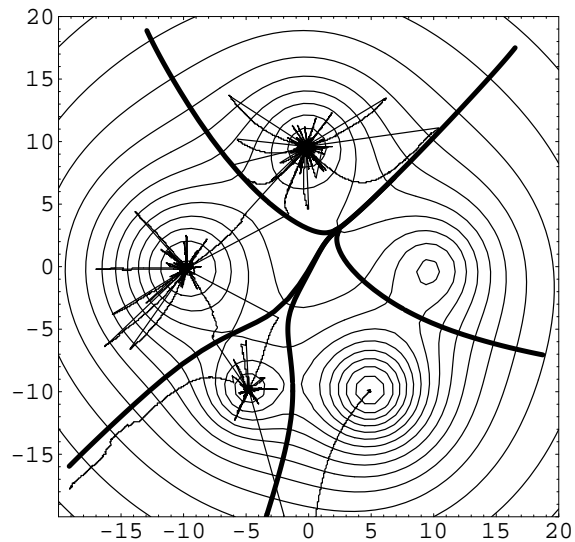


The lowest minimum:  $U(4.9, -9.9) = -1.46$

The second lowest minimum:  $U(-9.7, -0.1) = -0.85$

$$\alpha(x) = \begin{cases} 1.8, & \text{if } U(x) \leq -1, \\ 1.1, & \text{if } U(x) > -1, \end{cases} \quad \theta = 0.75.$$

Look for a local minimum:  $U(m) \leq -1$ .

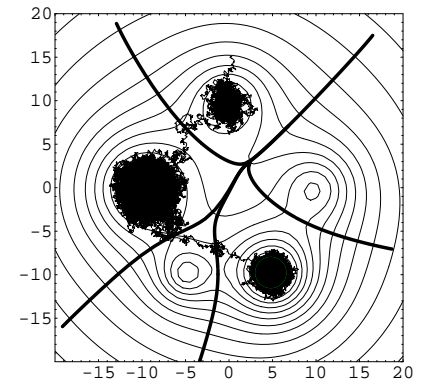
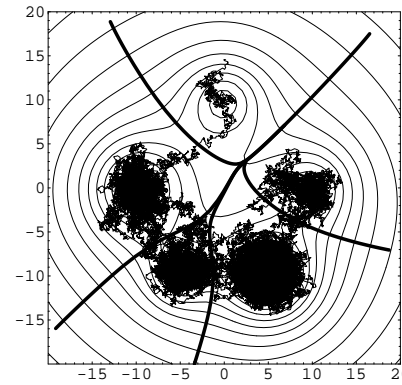
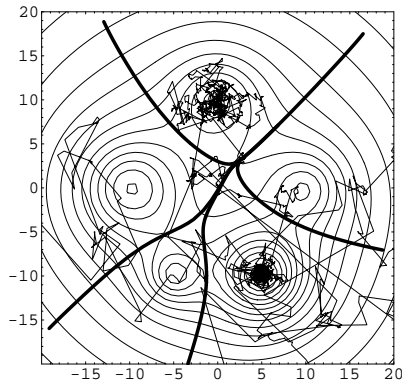
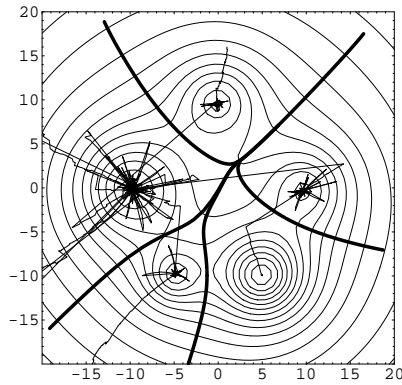


$$V_k = V_{k-1} - \nabla U(V_{k-1})h - \frac{\xi_k^h(V_{k-1})}{(\lambda + (k-1)h)^\theta}$$

$$0 \leq k \leq 2 \cdot 10^6, \lambda = 10^4, V_0 \in [-20, 20]^2, h = 0.1$$

**Success ratio: 96 out of 100.**

# 25. Numerical example in $\mathbb{R}^2$



**Stable-like simulated annealing**

$\alpha = 1.1, A = 1.8, \theta = 0.75$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$1 \cdot 10^2$	175458	61	0.0004	90	0.0004
$5 \cdot 10^2$	<b>93273</b>	<b>75</b>	<b>0.0004</b>	<b>96</b>	<b>0.0004</b>
$1 \cdot 10^3$	135081	62	0.0004	93	0.0004
$5 \cdot 10^3$	148972	60	0.0004	93	0.0004
$1 \cdot 10^4$	264070	47	0.0004	85	0.0004

**Fast simulated annealing**

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^1$	586029	67	0.0005	69	0.0004
$5 \cdot 10^1$	5787	97	0.0019	98	0.0006
$1 \cdot 10^2$	<b>1887</b>	<b>100</b>	<b>0.0035</b>	<b>100</b>	<b>0.0008</b>
$5 \cdot 10^2$	4477	99	0.0294	100	0.0038
$1 \cdot 10^3$	7552	70	0.0600	100	0.0112

**Gaussian SDE**

$\theta = 3$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^3$	8749	62	0.1211	76	0.1086
$10^4$	<b>12768</b>	<b>71</b>	<b>0.1170</b>	<b>79</b>	<b>0.1109</b>
$10^5$	34961	61	0.1204	81	0.1154
$10^6$	37262	66	0.0888	86	0.0923

**Gaussian simulated annealing**

$\theta = 3$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^3$	<b>81943</b>	<b>67</b>	<b>0.0786</b>	<b>93</b>	<b>0.0804</b>
$10^4$	113516	62	0.0715	94	0.0756
$10^5$	301035	44	0.0685	79	0.0643
$10^6$	573990	50	0.0662	63	0.0633

## 26. Numerical example in $\mathbb{R}^4$ . Shekel's function

$$S_{10,4}(\mathbf{y}) = - \sum_{i=1}^{10} \frac{1}{c_i + \|\mathbf{y} - \mathbf{a}_i\|^2}, \quad \mathbf{y} \in \mathbb{R}^4.$$

### Stable-like simulated annealing

$\alpha = 1.2, A = 1.9, \theta = 0.6$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^2$	50198	39	0.3598	92	0.0887
$5 \cdot 10^2$	46768	52	0.2456	92	0.0866
$1 \cdot 10^3$	<b>43951</b>	<b>63</b>	<b>0.2072</b>	<b>95</b>	<b>0.0810</b>
$5 \cdot 10^3$	63960	56	0.0848	95	0.0595
$1 \cdot 10^4$	77239	53	0.0572	97	0.0515

### Fast simulated annealing

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
50	612902	67	0.1385	70	0.0179
100	354639	81	0.2499	82	0.0324
<b>500</b>	<b>57214</b>	<b>7</b>	<b>0.8834</b>	<b>99</b>	<b>0.1470</b>
1000	99846	0	—	99	0.2644
1500	139922	0	—	100	0.4101

### Gaussian SDE

$\theta = 10$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^3$	37089	1	1.5448	7	1.7751
$10^4$	26749	10	1.5667	18	1.6223
$10^5$	24657	52	1.5839	45	1.6339
$10^6$	<b>26611</b>	<b>59</b>	<b>1.5631</b>	<b>67</b>	<b>1.3583</b>
$10^7$	53874	54	1.4889	69	1.5545

### Gaussian simulated annealing

$\theta = 10$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^3$	16889	69	1.4155	79	1.2417
$10^4$	<b>18722</b>	<b>78</b>	<b>1.2062</b>	<b>90</b>	<b>1.3376</b>
$10^5$	59266	68	1.0551	89	1.1888
$10^6$	183003	74	0.9401	84	1.0355
$10^7$	536875	67	1.1161	72	1.0520

## 27. Numerical example in $\mathbb{R}^6$ . Hartman's function

$$H_{4,6}(\mathbf{y}) = - \sum_{i=1}^4 c_i \exp \left( - \sum_{j=1}^6 a_{ij} (y_i - p_{ij})^2 \right), \quad \mathbf{y} \in \mathbb{R}^6.$$

### Stable-like simulated annealing

$\alpha = 1.5, A = 1.9, \theta = 0.6$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
100	20070	13	0.0295	84	0.0044
250	18492	18	0.0257	84	0.0043
<b>500</b>	<b>12841</b>	<b>43</b>	<b>0.0264</b>	<b>80</b>	<b>0.0049</b>
1000	18624	42	0.0155	83	0.0042
5000	29495	73	0.0043	83	0.0025

### Fast simulated annealing

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
10	1302701	35	0.0044	35	0.0001
50	1525483	17	0.0323	23	0.0029
100	1451393	2	0.0428	26	0.0049
500	1348318	0	—	11	0.0323
1000	865728	0	—	7	0.0388

### Gaussian SDE

$\theta = 6$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^6$	360969	0	—	0	—
$10^7$	191639	0	—	0	—
$10^8$	597533	0	—	0	—

### Gaussian simulated annealing

$\theta = 6$

$\lambda$	$\langle N_{\text{first}} \rangle$	$k = 5 \cdot 10^4$		$k = 5 \cdot 10^5$	
		$N_k$	$\Delta_k$	$N_k$	$\Delta_k$
$10^5$	668062	0	—	0	—
$10^6$	193789	0	—	0	—
$10^7$	466592	0	—	1	0.0430

# Pictures

**Page 1** [http://www.ngdc.noaa.gov/paleo/globalwarming/gallery/icecore\\_4.jpg](http://www.ngdc.noaa.gov/paleo/globalwarming/gallery/icecore_4.jpg)

[http://ess.geology.ufl.edu/ess/Notes/Paleoclimatology/Paleoclimate\\_Slides/greenland.gif](http://ess.geology.ufl.edu/ess/Notes/Paleoclimatology/Paleoclimate_Slides/greenland.gif)

**Page 2** From: Ganopolski, A. and Rahmstorf, S., Abrupt glacial climate changes due to stochastic resonance. *Physical Review Letters* 88(3), 038501+, 2002.

**Page 3** From: Ditlevsen, P. D., Observation of  $\alpha$ -stable noise induced millennial climate changes from an ice record, *Geophysical Research Letters* 26(10), 1441–1444, 1999.

**Page 15** <http://char.txa.cornell.edu/media/metal/ann.gif>

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**Page 18** From: H. Szu, Automated fault recognition by image correlation neural network technique. *IEEE Transactions on Industrial Electronics*, 40(2):197–208, 1993.

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