

Hybrid functional calculations of defects in materials for optoelectronics and quantum information science

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Acknowledgments:

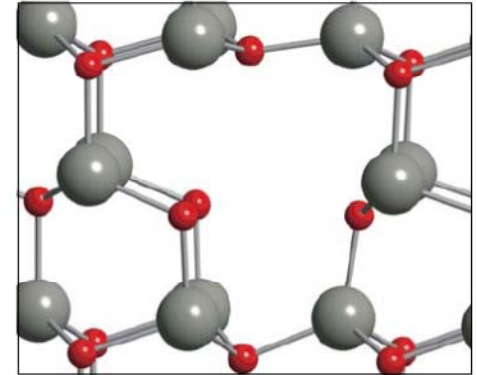
A. Janotti, J. Lyons, J. Varley, J. Weber (UCSB)
P. Rinke (FHI), M. Scheffler (UCSB, FHI Berlin)
G. Kresse (U. Vienna)
C. Freysoldt, J. Neugebauer (MPI Düsseldorf)
NSF, DOE, ARO, XSEDE

MDWS1: Quantum and Atomistic Modeling of Materials Defects
IPAM, UCLA, Los Angeles
October 1-5, 2011

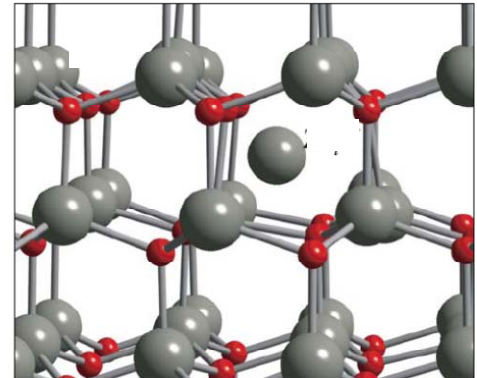
Why study defects?

- **“Defects”**
 - Extended defects
 - » dislocations
 - Point defects:
 - » Native defects
 - » Impurities
- **Defects often determine the properties of materials**
 - Doping and its limitations
 - Device degradation
 - Diffusion
 - » Mediated by point defects

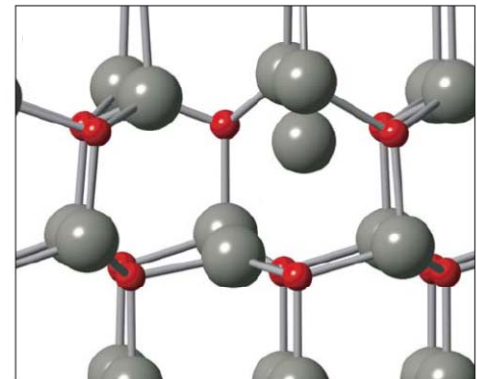
Vacancy



Interstitial



Antisite



Technological significance

- **Semiconductors**

- Achieve higher doping levels
 - » *p*-type doping of AlN would allow UV lasers
- Controllably dope materials *p*-type and *n*-type
 - » Oxides

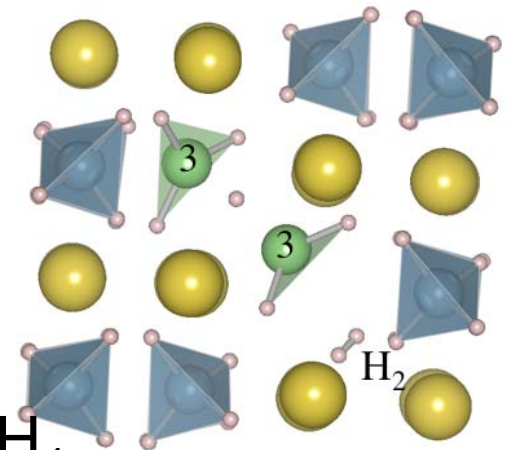
- **Photovoltaics**

- $\text{CuIn}_x\text{Ga}_{1-x}\text{Se}_2$

- **Hydrogen storage materials**

- Kinetics of hydrogen release in NaAlH_4

- **Embrittlement of structural metals**



First-Principles Calculations

- **Solve complex many-body problem for electrons in solid**
 - Nuclei treated classically (Born-Oppenheimer)
- **Replace with a single-particle problem:**

$$\left[-\frac{1}{2}\nabla^2 + v_{\text{KS}}(\mathbf{r}) \right] \psi_i(\mathbf{r}) = \varepsilon_i \psi_i(\mathbf{r})$$

with

$$v_{\text{KS}}(\mathbf{r}) = v_{\text{ext}}(\mathbf{r}) + v_{\text{H}}(\mathbf{r}) + v_{\text{XC}}(\mathbf{r})$$

potential due to nuclei

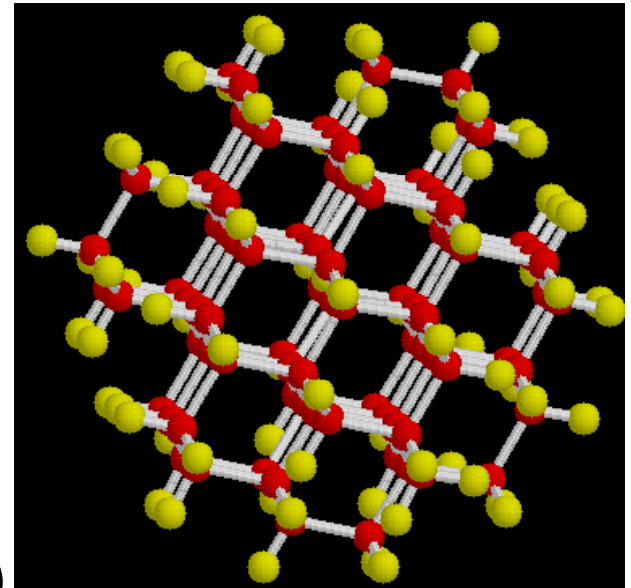
potential due to other
electrons (Hartree)

exchange-correlation potential; captures many-
body effects; functional of charge density

- **Density-functional theory (DFT)**
 - Highly valuable tool for studying defects
 - Very reliable for total energies and atomic structure

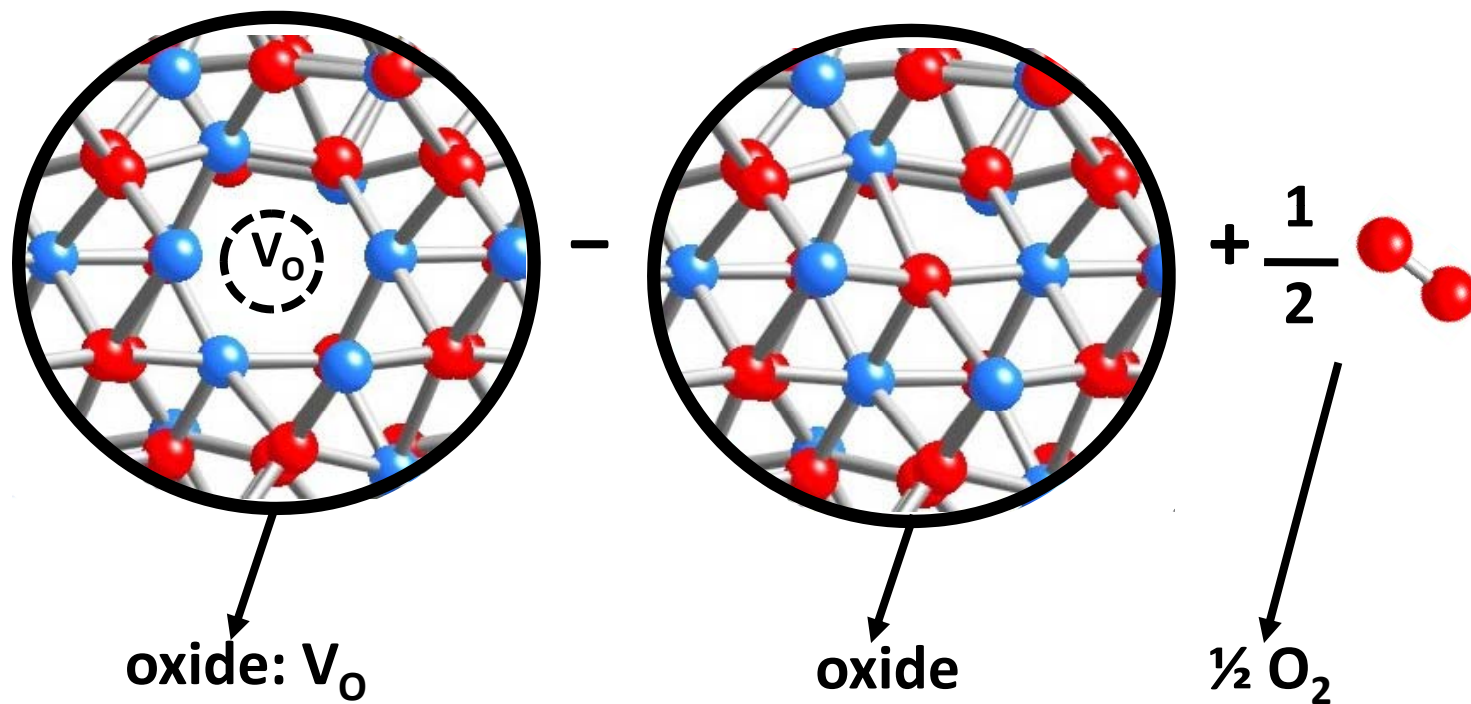
Defect calculations: geometry

- **Green's functions**
 - Implementation difficult
 - Non-intuitive
- **Clusters**
 - Surface effects
 - Not “bulk” (quantum confinement)
- **Supercells: Periodic Boundary Conditions**
 - Need rigorous treatment of electrostatic interactions in dielectric media
 - » C. Freysoldt, J. Neugebauer, and C. G. Van de Walle, Phys. Rev. Lett. **102**, 106402 (2009).



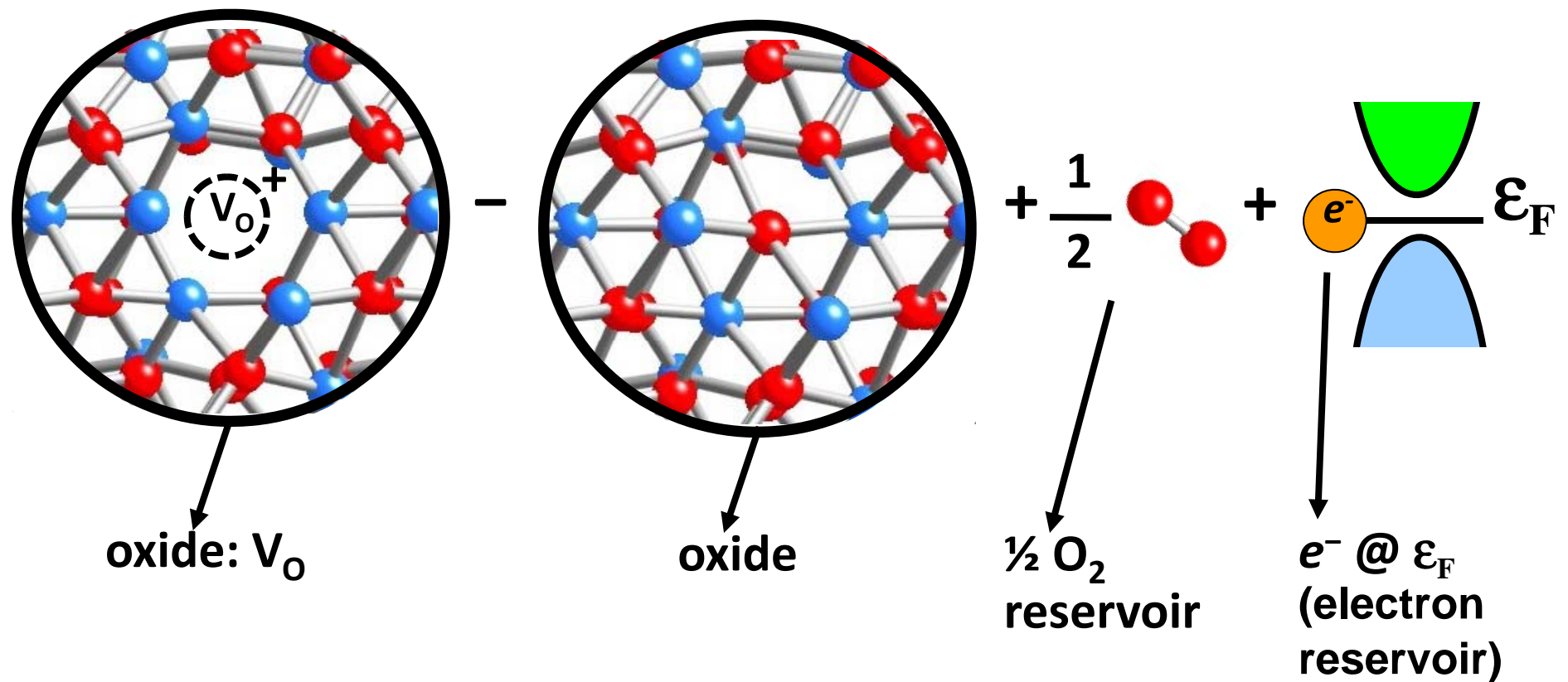
Defect Formation Energies

Example: oxygen vacancy (V_O) in ZnO



Defect Formation Energies

Example: oxygen vacancy (V_O) in ZnO



Formalism

- E_{form} : formation energy

Concentration of defects or impurities:

$$C = N_{\text{sites}} \exp [- E_{\text{form}}/kT]$$

- **Example: oxygen vacancy in ZnO**

$$E_{\text{form}}(V_{\text{O}}^+) = E_{\text{tot}}(V_{\text{O}}^+) - E_{\text{tot}}(\text{bulk}) + \mu_{\text{O}} + E_{\text{F}}$$

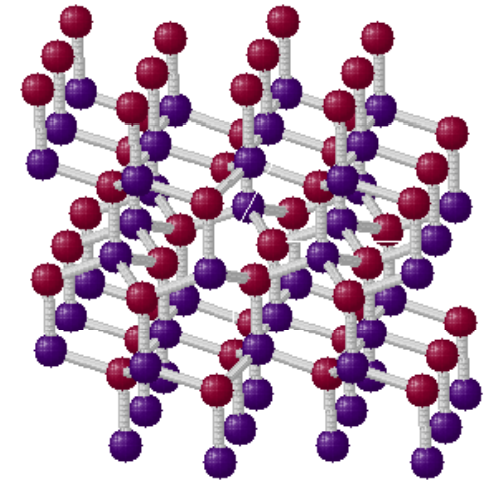
μ_{O} : energy of oxygen in reservoir, i.e., oxygen chemical potential

E_{F} : energy of electron in its reservoir, i.e., the Fermi level

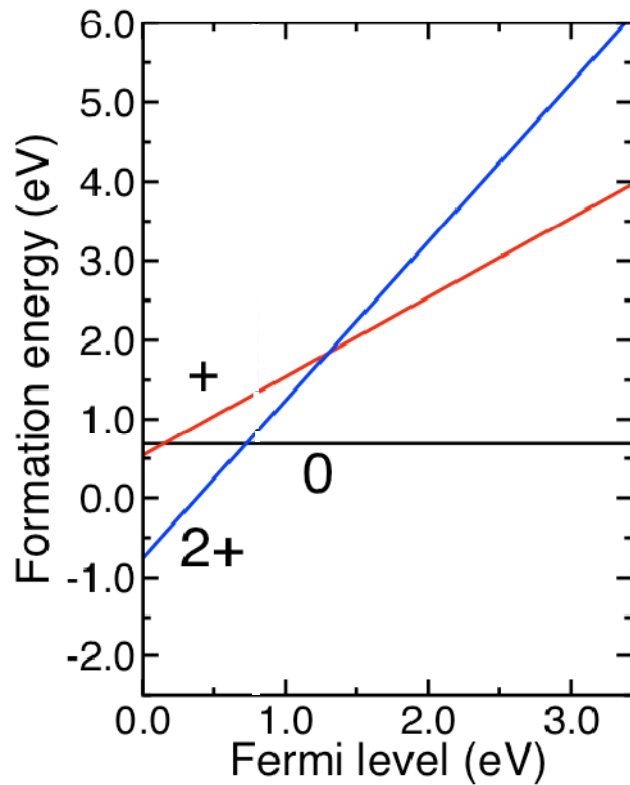
- **General expression**

$$E_{\text{form}}(D^q) = E_{\text{tot}}(D^q) - E_{\text{tot}}(\text{bulk}) + n_i \mu_i + qE_{\text{F}}$$

n_i : number of atoms being exchanged to form the defect

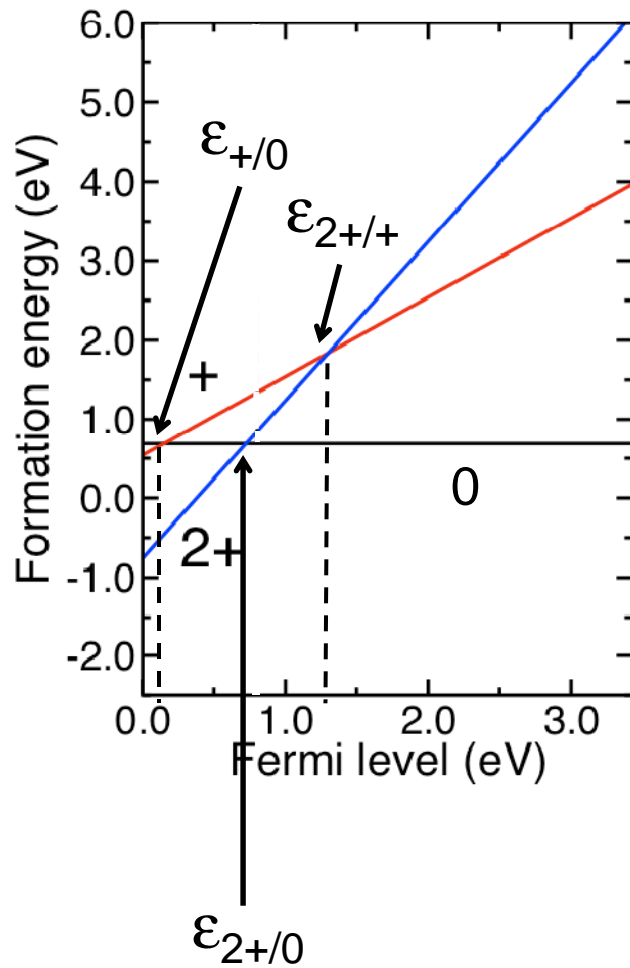


Formation energy

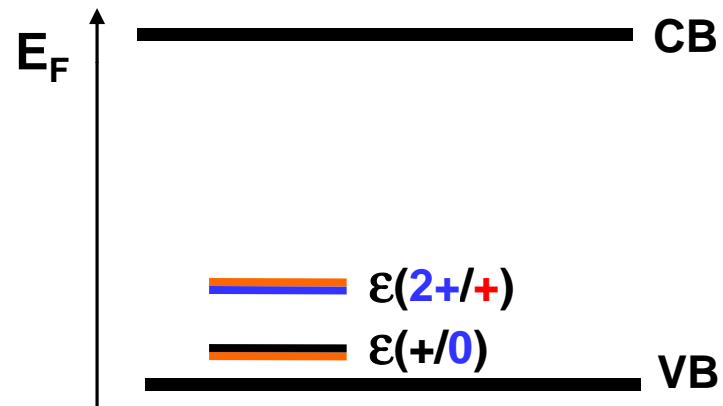


- V_O in ZnO
- DFT
- Local Density Approximation

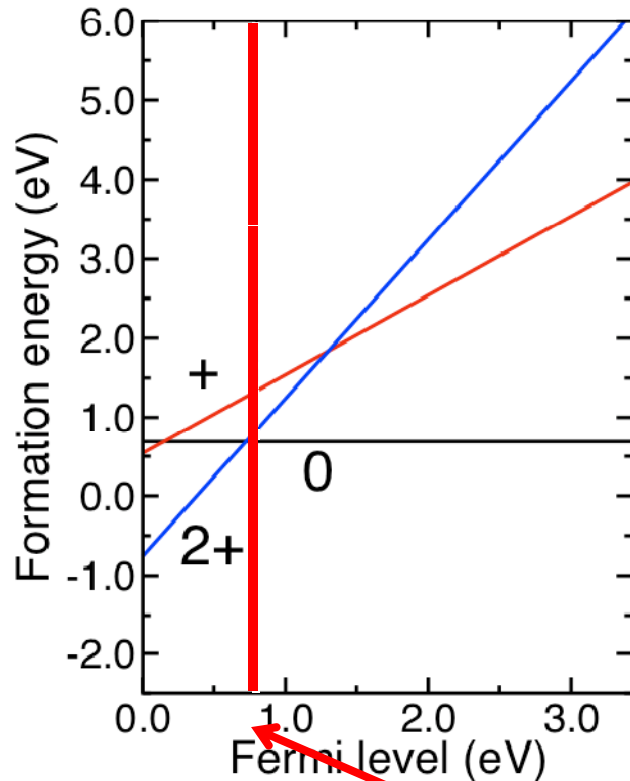
Transition levels



Charge-state transition levels



Issues...



DFT/LDA band gap

- **Band gap problem**

- DFT

- » Local Density

- » Approximation (LDA)

- » Generalized Gradient

- » Approximation (GGA)

- **Affects formation energies and transition levels**

- Even for neutral charge states, if defect-induced Kohn-Sham states are occupied with electrons

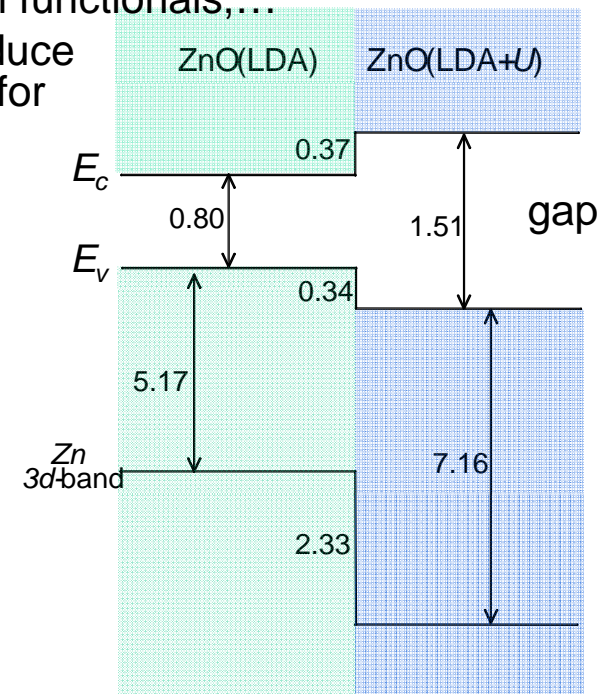
Band-gap corrections

- **Empirical corrections**

- *Ad hoc* corrections
 - » “Scissors operator”
 - » “Modified pseudopotentials”
 - » Issues
 - Hard to control
 - Unintended consequences (indirect vs. direct gaps, ...)
- Extrapolating calculations that yield different gaps
 - » Different plane-wave cutoffs, exchange-correlation functionals, ...
 - » Issue: different choices of parameter not only produce different gap, but also different levels of accuracy for description of defect

- **Physically meaningful corrections**

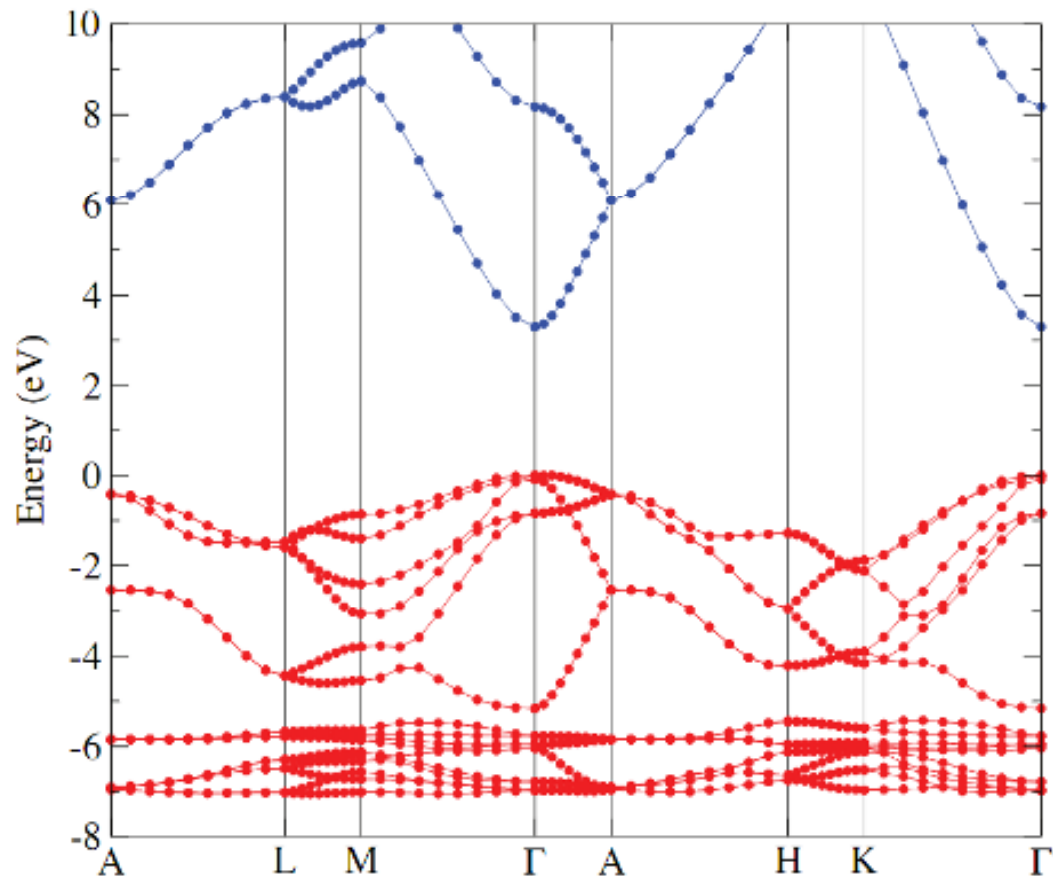
- Self-interaction corrections
- “LDA+ U ” approach
 - » Issues:
 - Determination of U
 - Band gap not fully corrected



Band-gap corrections: Hybrid functionals

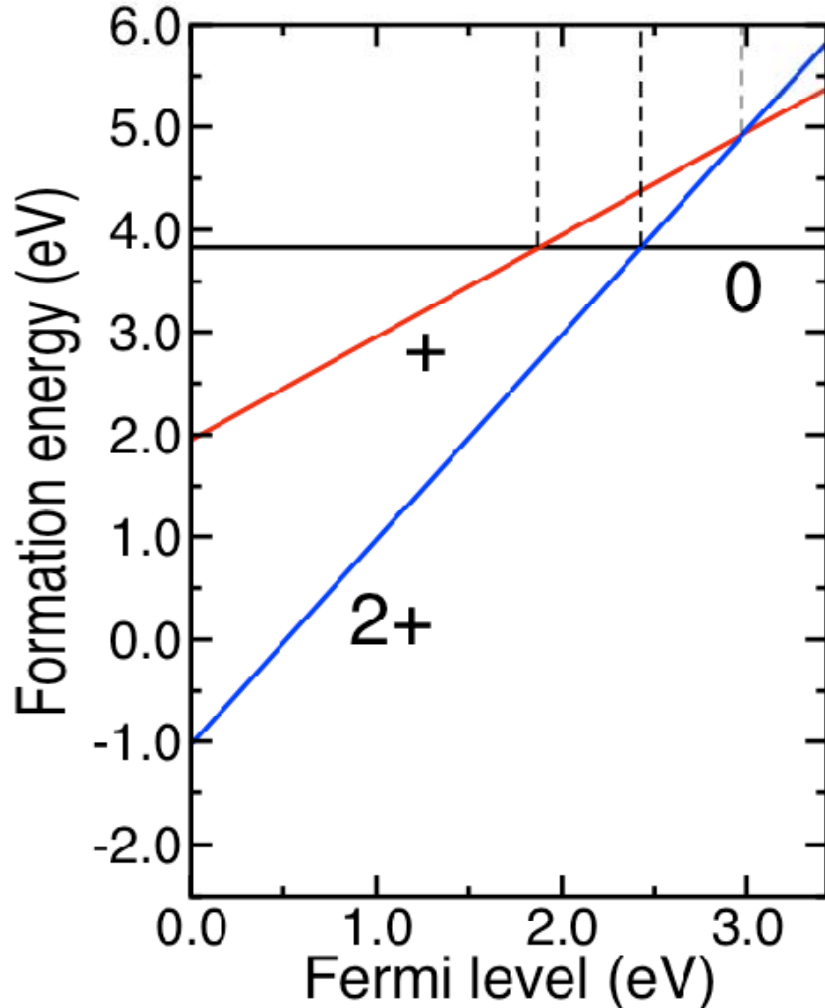
- **Mixing of Hartree-Fock (exact exchange) and GGA**
 - A. D. Becke, J. Chem. Phys. **98**, 1372 (1993).
- **Functionals mix in ~25% of exact exchange**
 - PBE0
 - » J. P. Perdew, K. Burke, and M. Ernzerhof, Phys. Rev. Lett. **77**, 3865 (1996);
J. P. Perdew, M. Ernzerhof, and K. Burke, J. Chem. Phys. **105**, 9982 (1996).
 - HSE
 - » Exact exchange only for short-range interactions
 - » J. Heyd, G. E. Scuseria and M. Ernzerhof, J. Chem. Phys. **118**, 8207 (2003).
- **Obtain very good description of many properties**
 - Structural properties as good as, or better, than with LDA/GGA
 - Band gaps close to experiment
 - Better description of charge localization

HSE calculations for ZnO



		PBE	HSE	Exp.
ZnO	a (\AA)	3.28	3.23	3.25
	c (\AA)	5.28	5.24	5.21
	E_g (eV)	0.73	3.35	3.40
	ΔH_f (eV)	-2.93	-3.66	-3.64

Example: V_O in ZnO



Value of formation energy :

- High formation energy
- Unlikely to form
 - » Experiment: electron spin resonance detects no oxygen vacancies in as-grown ZnO

Transition levels:

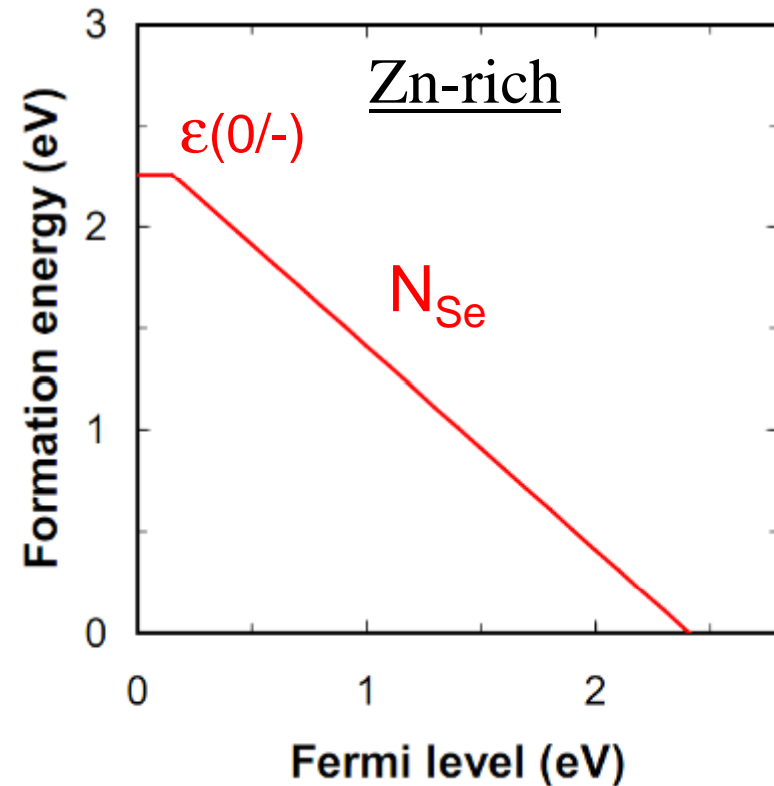
- Not a shallow donor!
- Not a source of conductivity
- ➔ Unintentional conductivity due to impurities

***p*-type doping of ZnO**

- **Numerous reports of *p*-type ZnO**
 - Using N and other group-V impurities (Sb, As, P)
- **Reliability? Reproducibility?**
- **No *pn* junctions**
- **Nitrogen is often regarded as most suitable hole dopant**
 - Shallow acceptor in ZnSe

Test case: N_{Se} in ZnSe

- Similar semiconductor
- N_{Se} known shallow acceptor in ZnSe
- Effective p -type dopant
- Results agrees well with experimental values
 - Theory: $E_A = 150$ meV
 - Exp.: $E_A = 110$ -130 meV



	PBE	HSE	Exp.
ZnSe a (Å)	5.73	5.69	5.67
ZnSe E_g (eV)	1.19	2.80	2.82
ΔH_f (eV)	-1.43	-1.74	-1.69

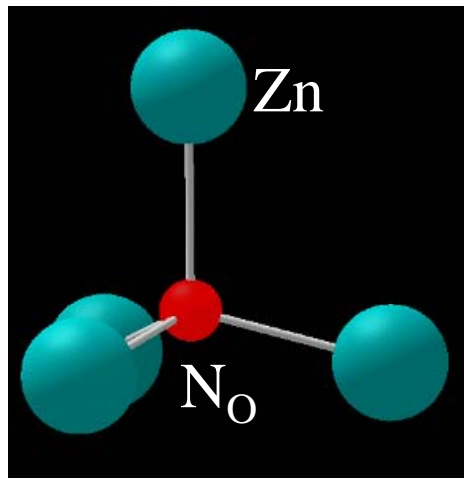
Nitrogen acceptor in ZnO

Deep acceptor

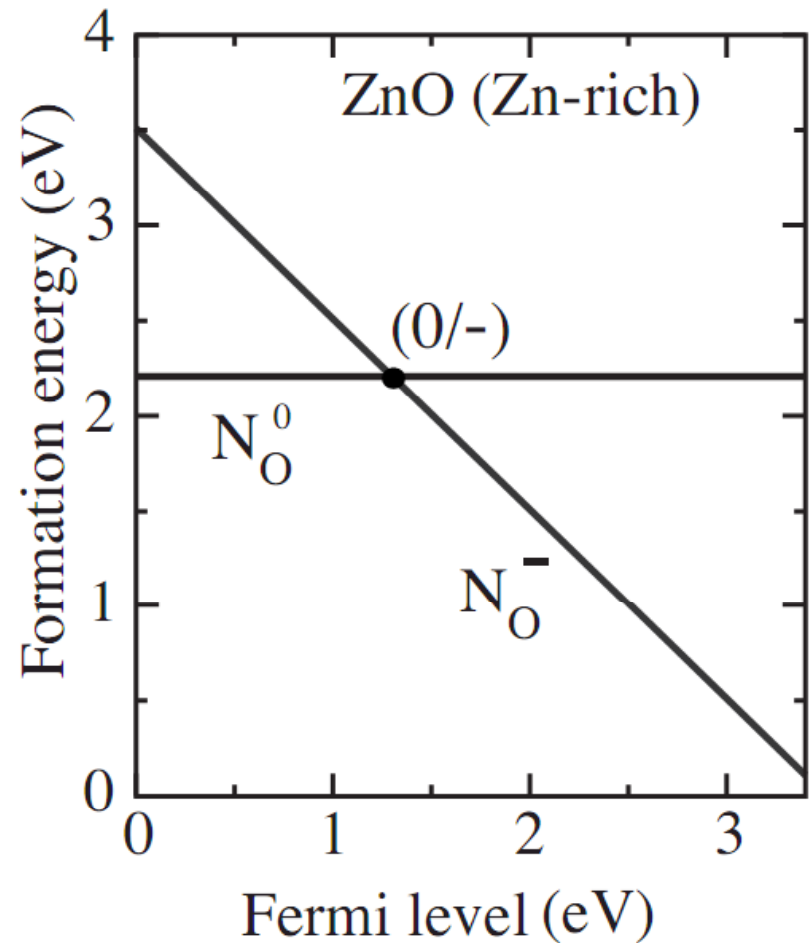
Large ionization energy: 1.3 eV

Off-center relaxation

Localized wavefunction



Axial bond = 2.12 Å
Planar bonds = 1.94 Å



J. L. Lyons, A. Janotti, and C. G. Van de Walle, Appl. Phys. Lett. **95**, 252105 (2009).

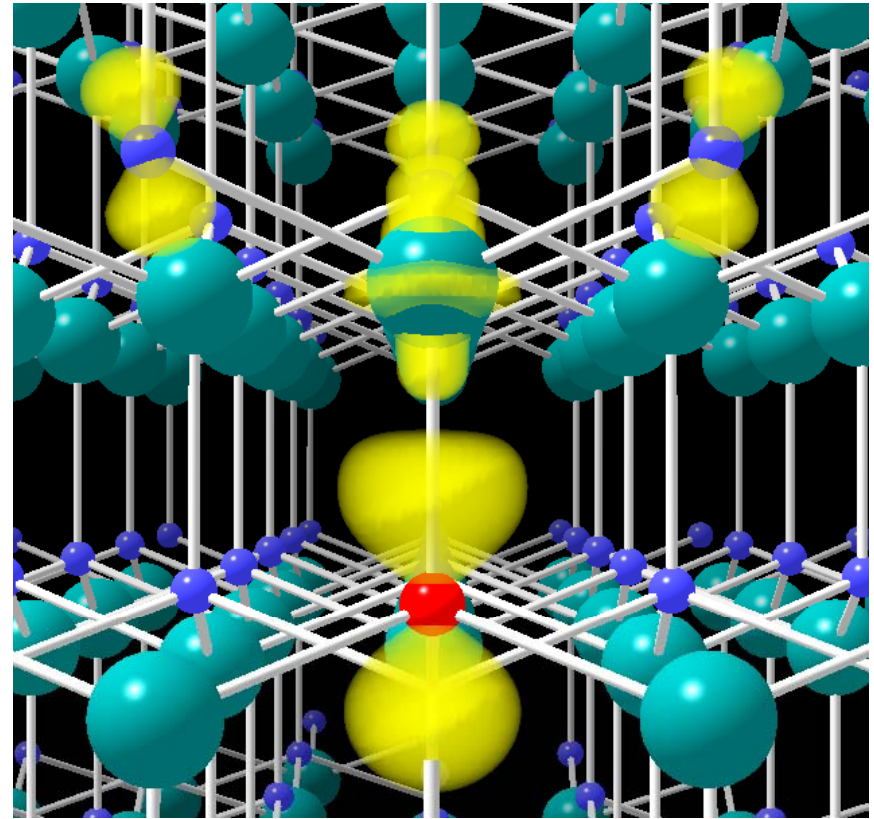
N_o acceptor: experiment

Localization on N atom
Directed towards axial zinc neighbor
Agrees with EPR signals

W.E. Carlos, E.R. Glaser, and D.C. Look, *Physica B* **308**, 976 (2001).
N. Y. Garces *et al.*, *Appl. Phys. Lett.* **80**, 1334 (2002).

LDA does *not* produce this localized wave function!
Need **hybrid functional!**

“Why nitrogen cannot lead to *p*-type conductivity in ZnO”, J. L. Lyons, A. Janotti, and C. G. Van de Walle, *Appl. Phys. Lett.* **95**, 252105 (2009).

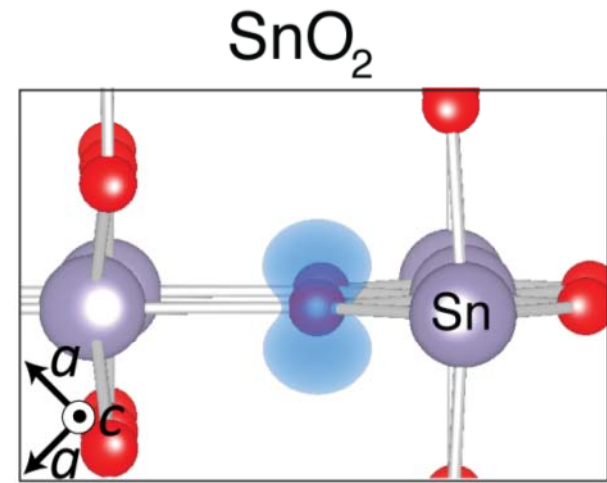


Spin density of nitrogen state
(Isosurface is 5% of maximum density)

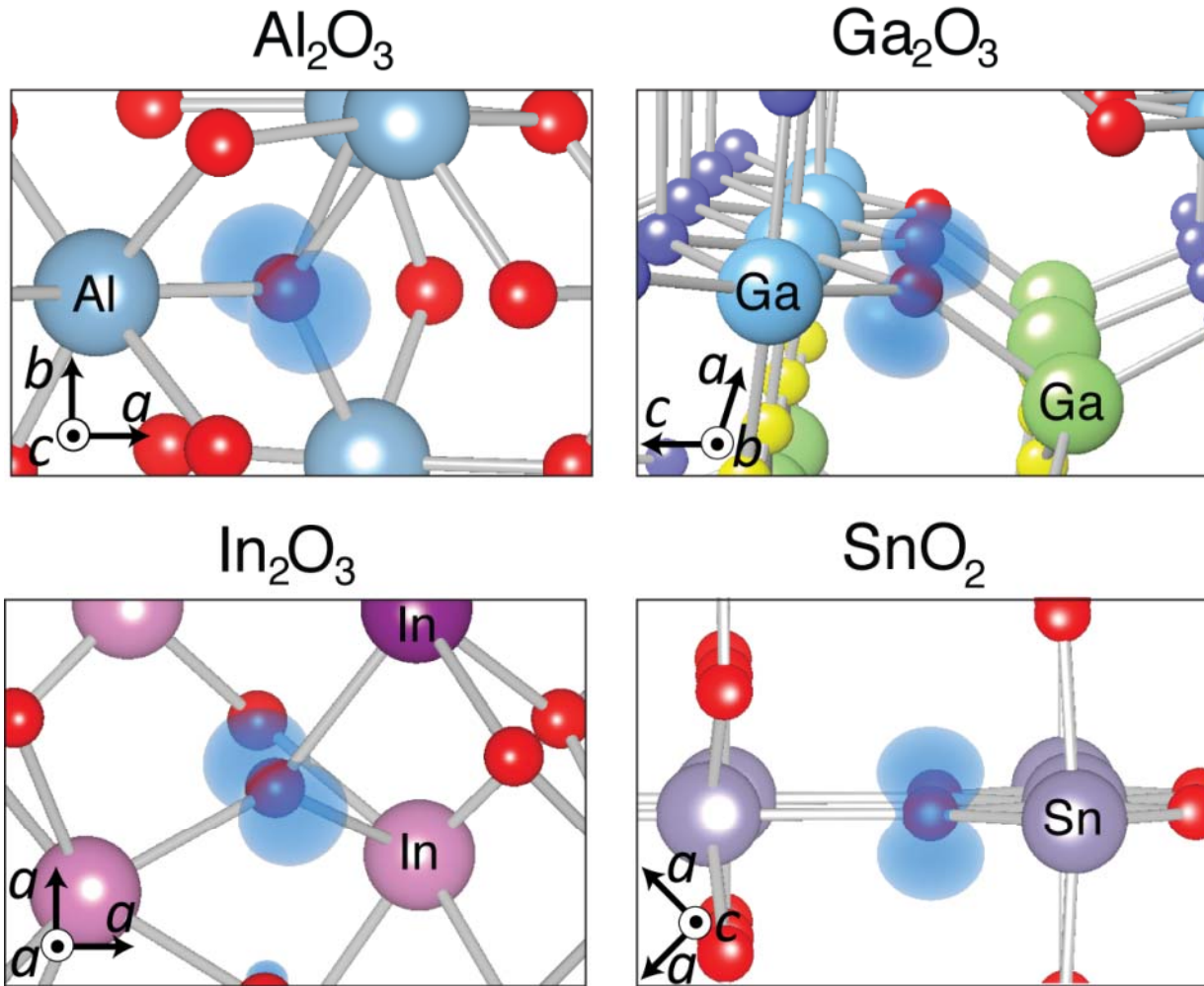
Conductivity: *p* type

Broader problem: self-trapping

- In ZnO, holes are trapped in a localized state bound to acceptor impurities
 - Rather than a delocalized, extended effective mass state
- Localization of holes occurs in many oxides
- Polarons
 - Presence of a charge in the valence band leads to large distortion of the lattice
 - Hole becomes trapped
 - “Self-trapping”
- Major obstacle for *p*-type doping!
 - Even if can get holes in the valence band, they will have very low mobility



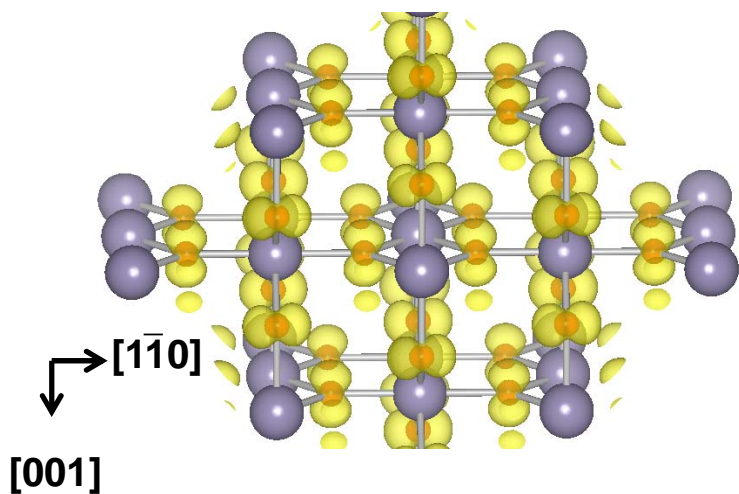
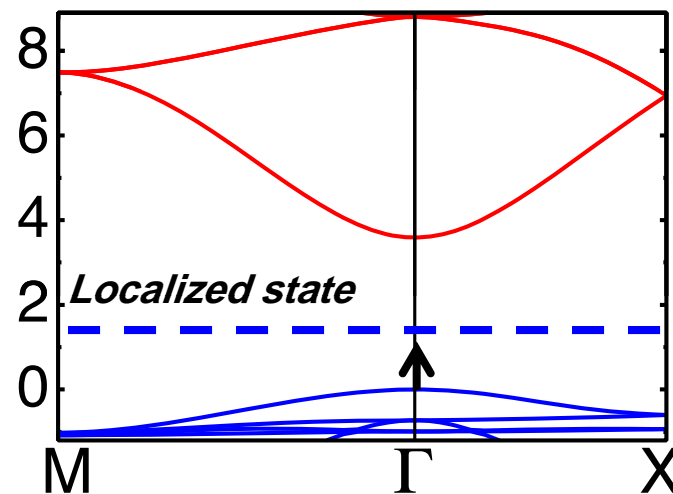
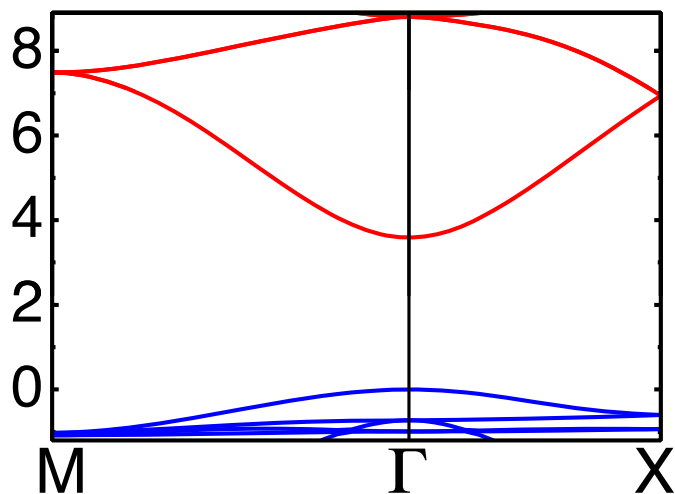
Self-trapped holes



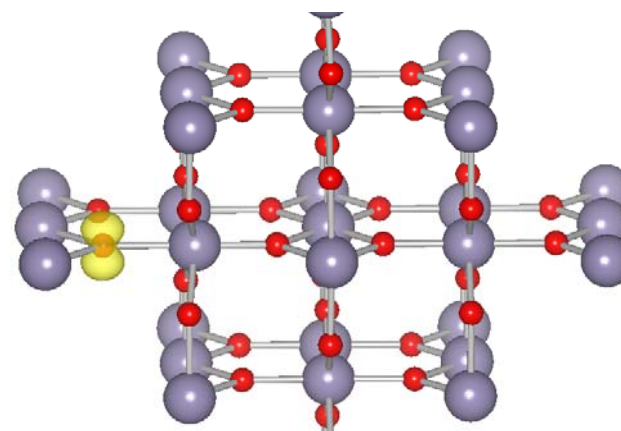
J. Varley, A. Janotti, C. Franchini, and C. G. Van de Walle,
Phys. Rev. B **85**, 081109(R) (2012).

Signatures of self-trapped holes

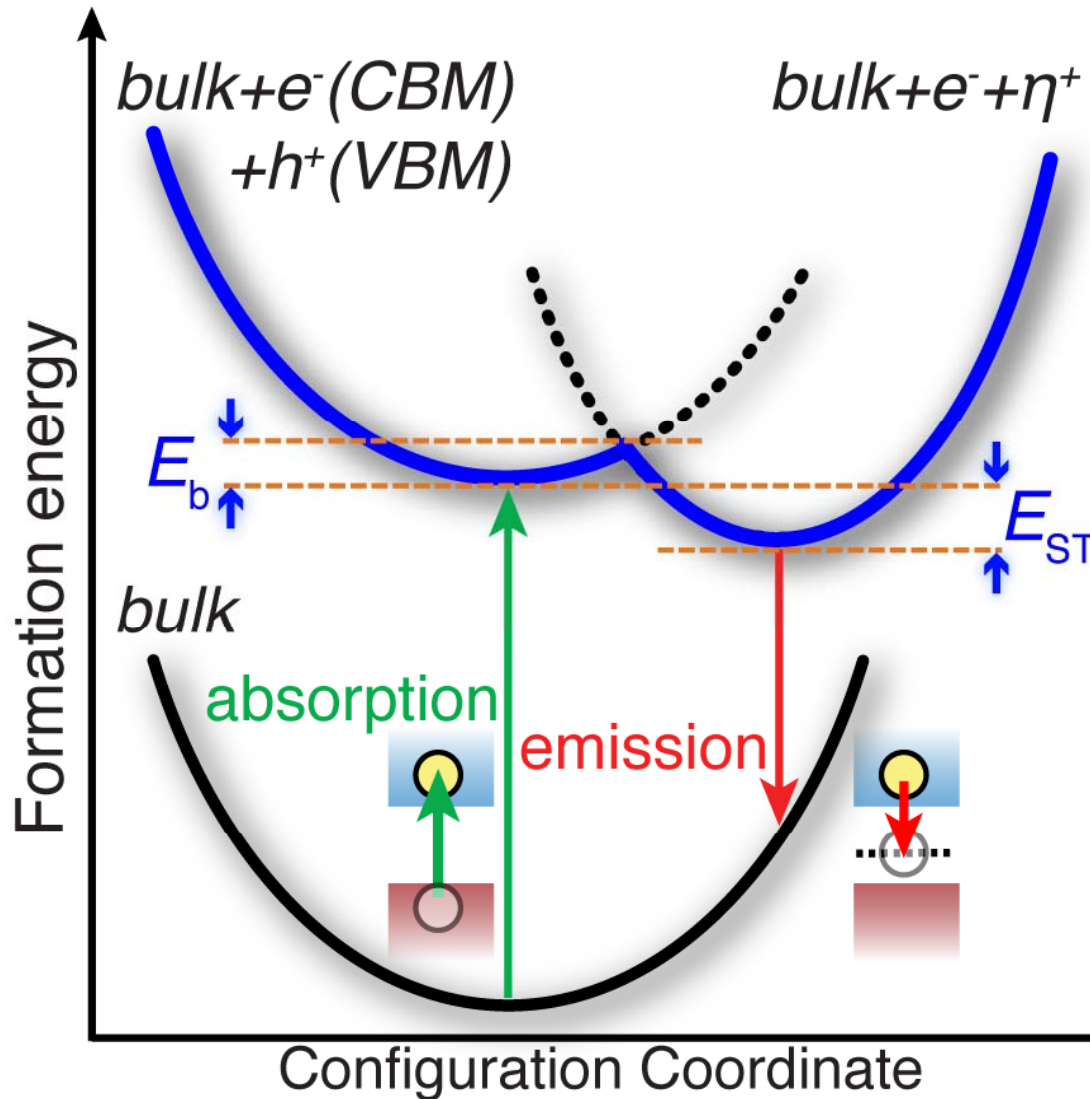
SnO₂



Self-trapping



Energy barriers & optical signatures



* Need **hybrid functional!** LDA/GGA do not produce localization!

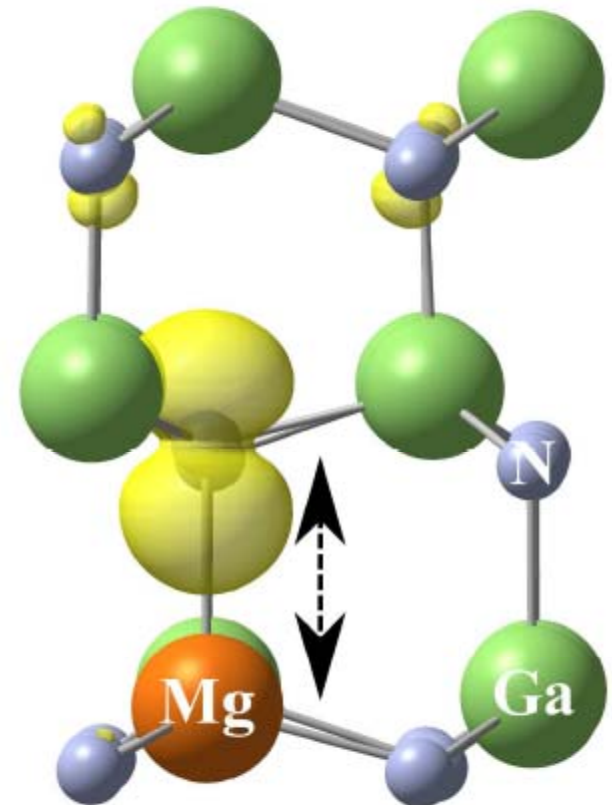
* Explains broad luminescence peaks observed in many oxide semiconductors

* Tendency to form small hole polarons is an obstacle to *p*-type conductivity

J. Varley, A. Janotti, C. Franchini, and C. G. Van de Walle,
Phys. Rev. B **85**, 081109(R) (2012)

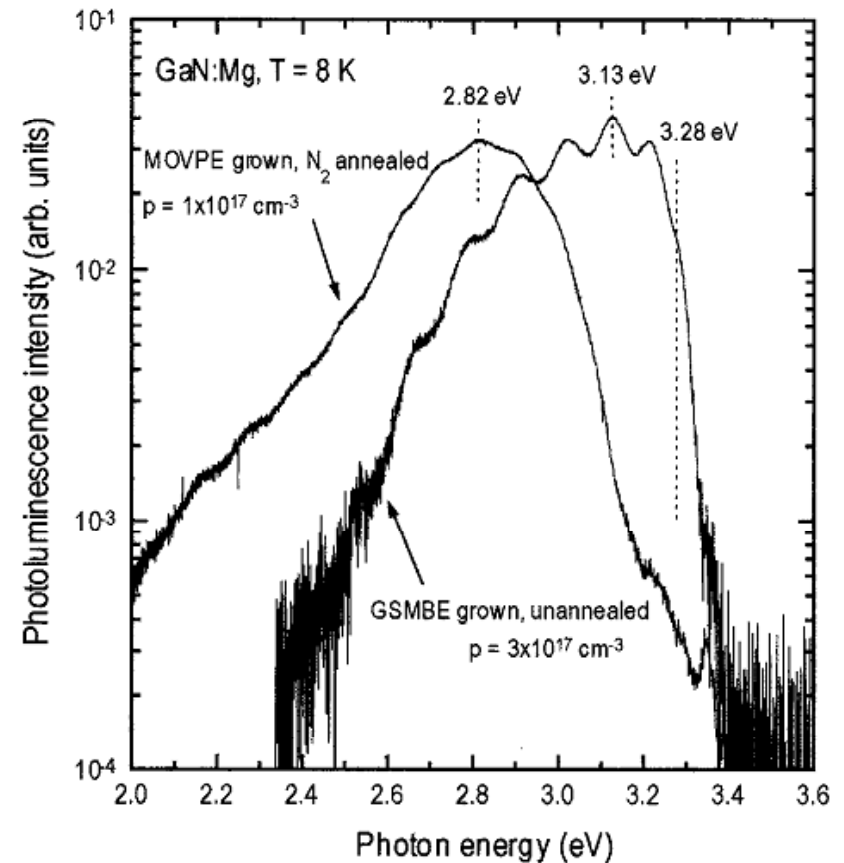
Magnesium acceptor in GaN

- Mg incorporates on Ga site
- Only known effective p -type dopant for GaN
- p -type doping is still limiting factor for devices
- 150-300 meV ionization energy
 - Few percent ionization of acceptors
- **Optical behavior?**



Mg-related optical signals in GaN

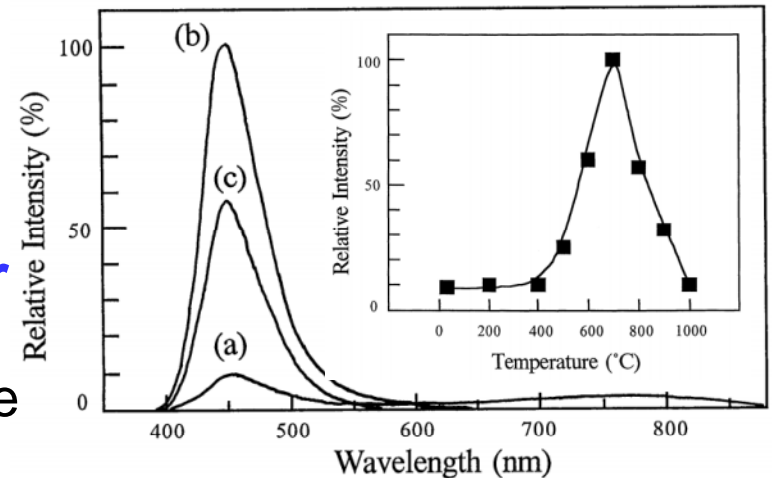
- Two photoluminescence (PL) signals observed in GaN:Mg
- 3.27 eV “donor-acceptor pair” (DAP)
 - Acceptor “A1”
B. Monemar *et al.*, Phys. Rev. Lett. 102, 235501 (2009).
 - Unannealed, Mg-doped GaN
- ~2.8 eV blue luminescence
 - Annealed, p-type GaN:Mg



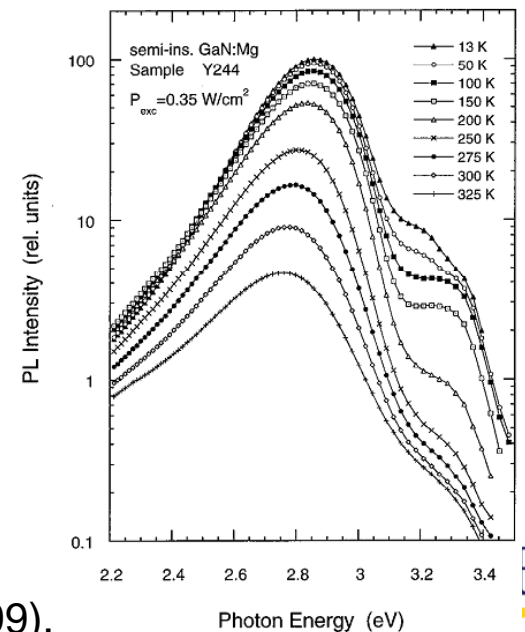
C. Leroux *et al.*, JAP 86, 3726 (1999).

2.8 eV Blue Luminescence (BL)

- Strong increase in intensity after post-growth anneal
- Previous proposals: deep donor
 - But: no deep donor signal from optically detected magnetic resonance (ODMR)
 - Instead: Mg acceptor signal...
E. R. Glaser *et al.*, Phys. Rev. B 65, 085312 (2002)



S. Nakamura *et al.*, JJAP 31, L139 (1992).

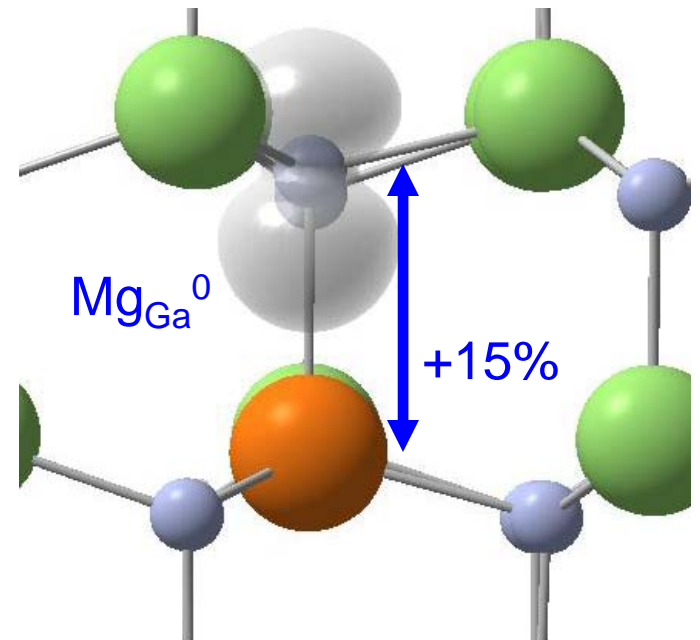
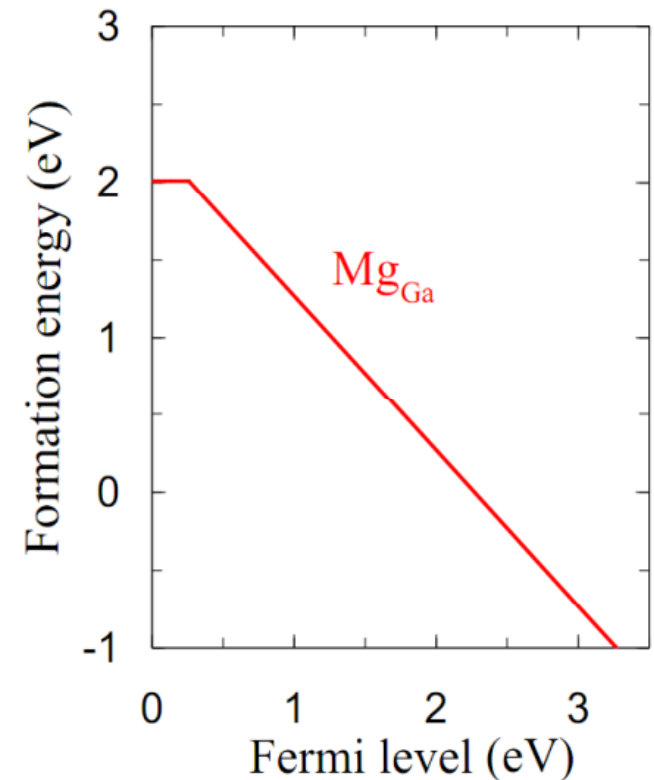


M. A. Reshchikov *et al.*, PRB 59, 13176 (1999).

Results: Mg_{Ga} in GaN

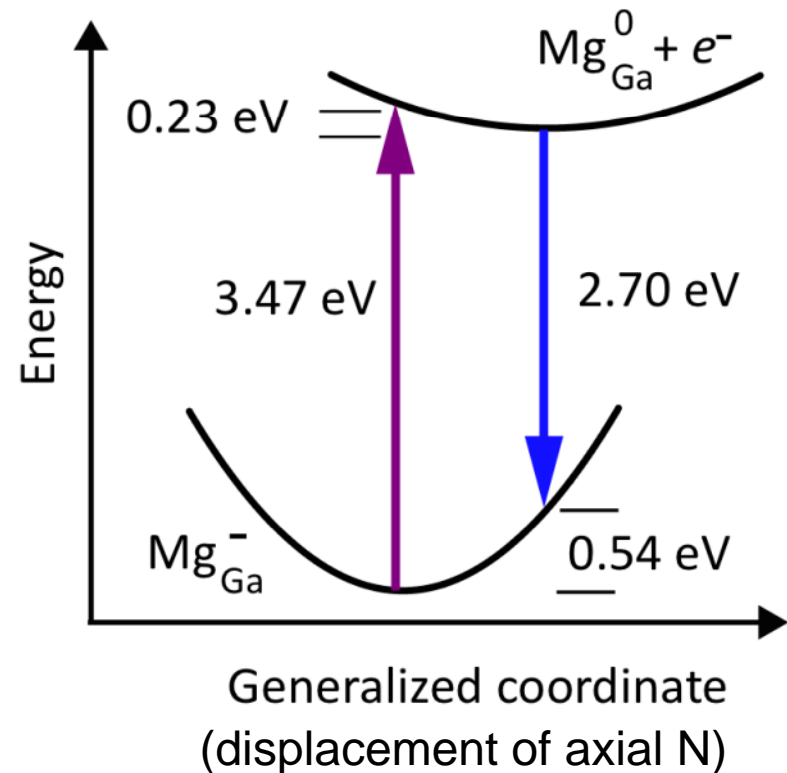
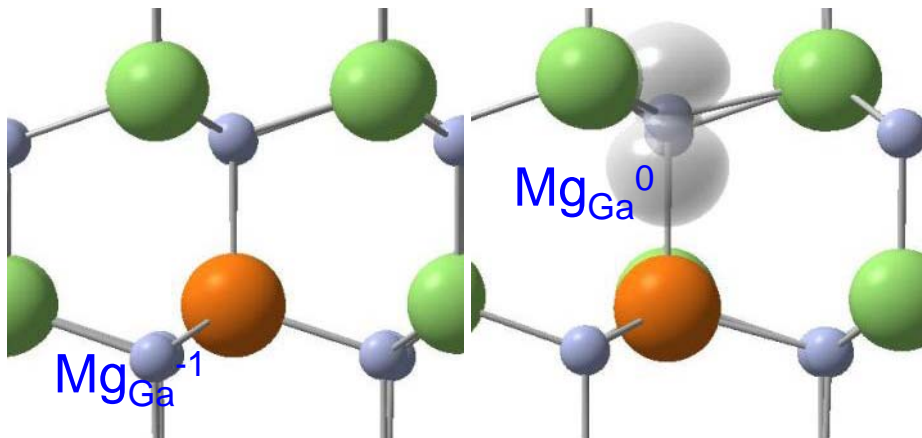
- Mg_{Ga} : acceptor, ionization energy of 260 meV
- Hole for Mg_{Ga}^0 is localized on nitrogen neighbor
- Large distortion in neutral charge state
 - Axial Mg-N bond increases by 15%
 - See also Lany and Zunger, Appl. Phys. Lett. 96, 142114 (2010).

Spin density of Mg_{Ga}^0 .
Isosurface shows the localized hole.



Results: optical signals of Mg_{Ga}

- Large relaxations
→ broad lines
- Electron recombining (from CBM or shallow donor) with Mg_{Ga}^0 leads to blue luminescence!



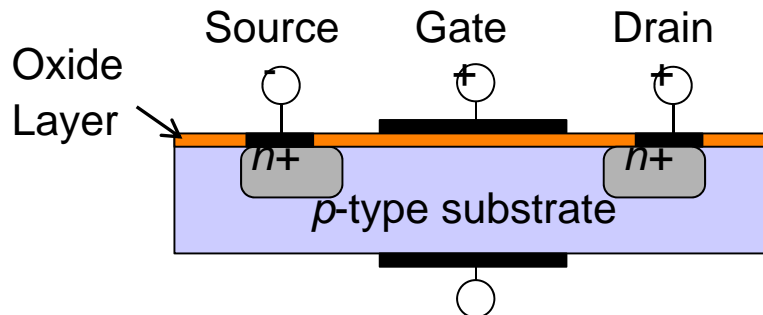
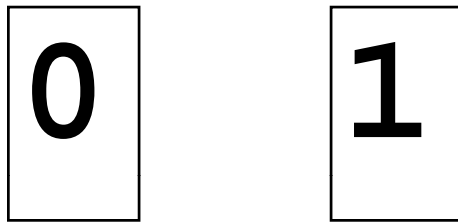
J. L. Lyons, A. Janotti, and C. G. Van de Walle,
Phys. Rev. Lett. **108**, 156403 (2012)

A 3D molecular model of a crystal lattice. The atoms are represented by blue and black spheres connected by white rods. A red laser beam is directed at a specific atom in the lattice, which is highlighted with a glowing purple and green aura. The background is a dark, textured surface.

Quantum Computing with Defects

Paradigm shift

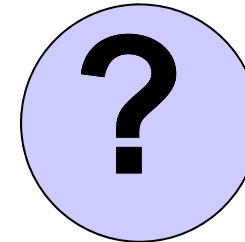
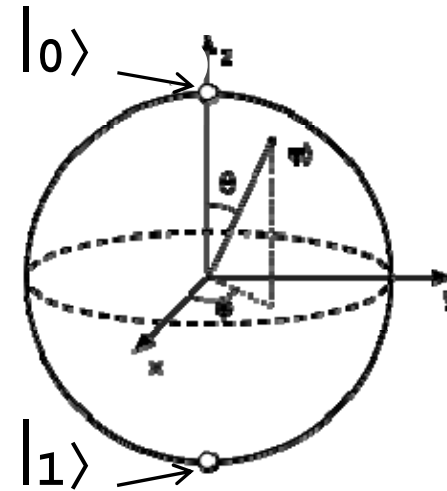
Classical Computing



Si CMOS (Transistors)

- Discrete states
- On (1) or Off (0) states
- Classical bits (0 or 1)

Quantum Computing



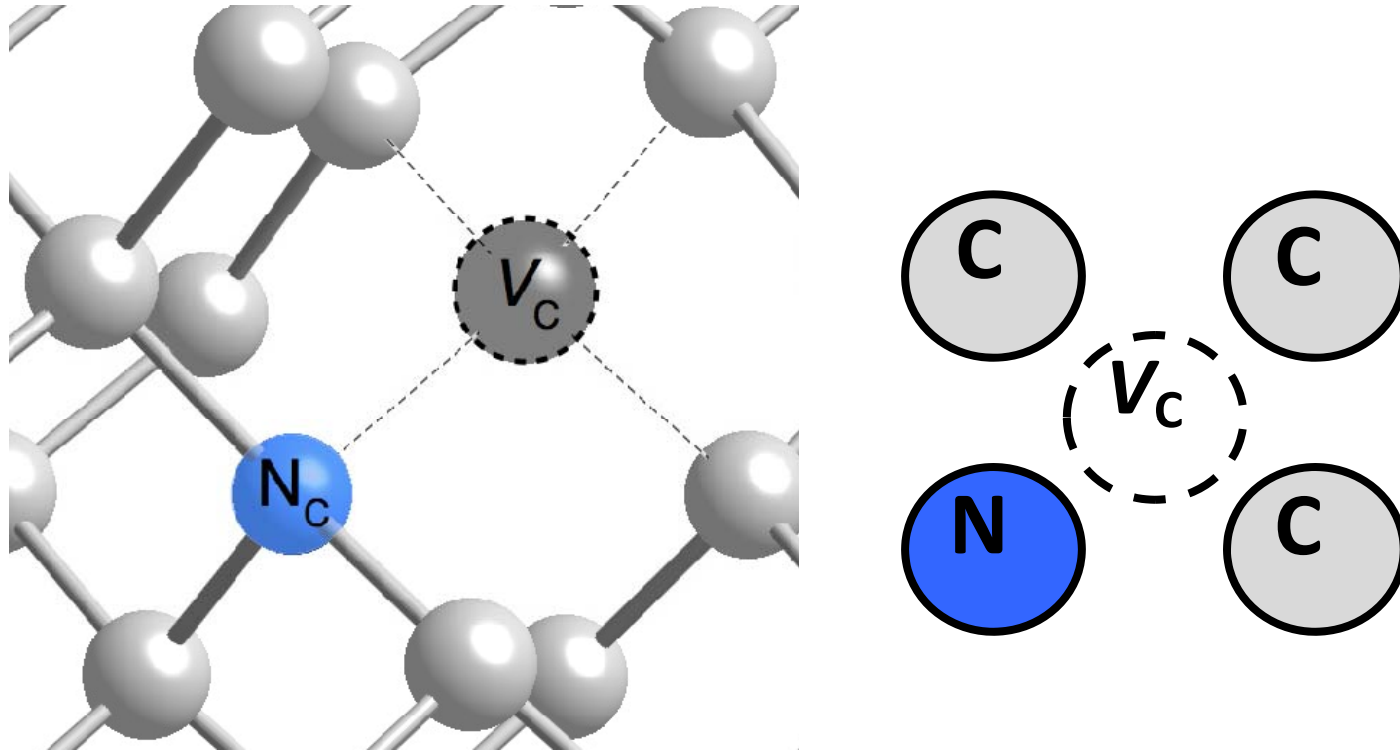
Qubits (Quantum bits)

- Quantum states
- Superposition of $|0\rangle$ and $|1\rangle$
- Bloch sphere

Qubit Requirements

- **Isolated**
 - design atomic-like state
- **Addressable**
 - can initialize, manipulate, and read out the quantum state
- **Engineerable**
 - create qubits in controlled and systematic manner

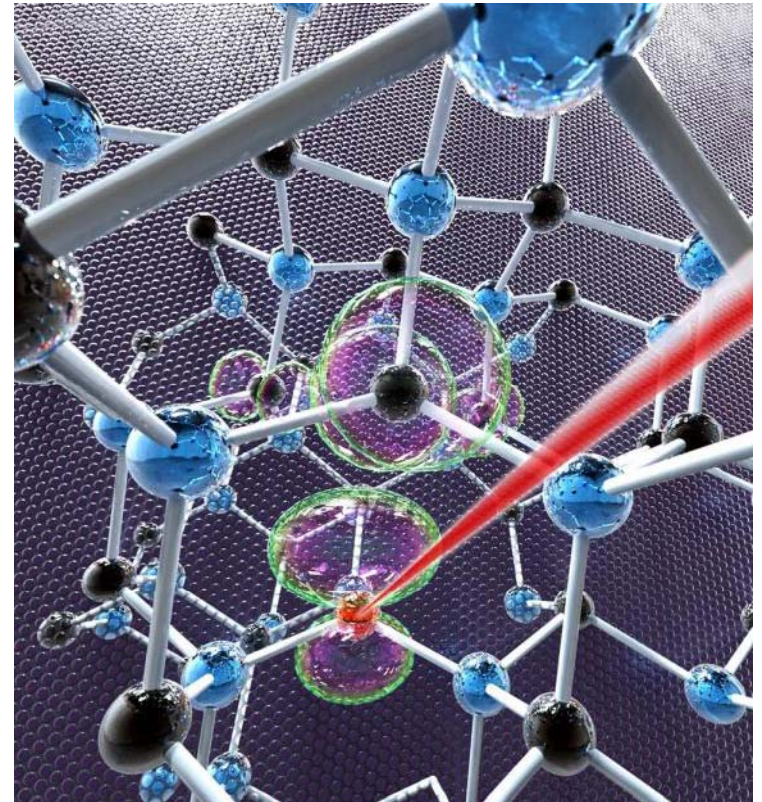
The NV center in Diamond



- A “defect” with many desirable properties
 - Room temperature qubit
 - Optically addressable single spin
 - Millisecond spin coherence times

Other Defects with Similar Properties?

- **Diamond: difficult to grow and process**
- **Do similar defects exist in a more technologically mature material?**
- **If so, how do we find them?**
- **Propose a set of physical criteria for host material and defect**
- **Develop a framework for testing criteria**
- **Illustrate criteria by example**



– vacancies and complexes in other semiconductors

Quantum computing with defects, J. R. Weber, et al., Proc. Natl. Acad. Sci. 107, 8515-8518 (2010).

Important Host Criteria

conduction band

_____ e_x, e_y
_____ $a_1(2)$

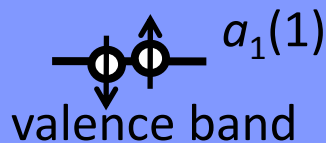
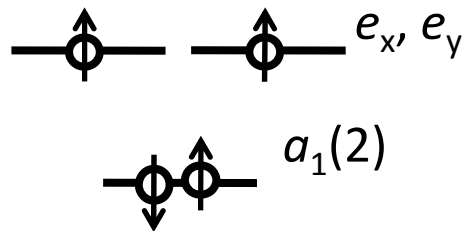
_____ $a_1(1)$
valence band

- **Wide band gap**
 - Avoid coupling to bulk states
- **Small spin-orbit coupling**
 - Avoid unwanted spin flips
- **Constituent elements with no net nuclear spin (if possible)**
 - Avoid spin bath effects
- **High quality crystals/thin films**
 - Experimentally feasible to address individual defects

Important Defect Criteria

conduction band

6 electrons ($S=1$; $m_s=+1,0,-1$)

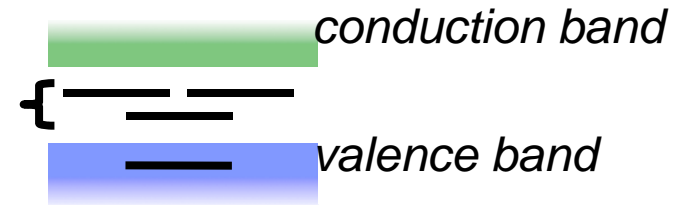


valence band

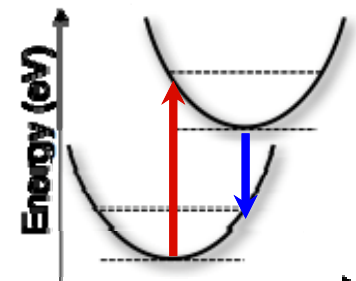
- **Defect states well localized within band gap**
 - Avoid coupling to bulk states
- **Splitting between defect levels not too small**
 - Avoid thermal coupling
- **Spin state capable of behaving as qubit**
 - Exploit spin sublevels
 - Form qubit states

Calculations Give Insight Into:

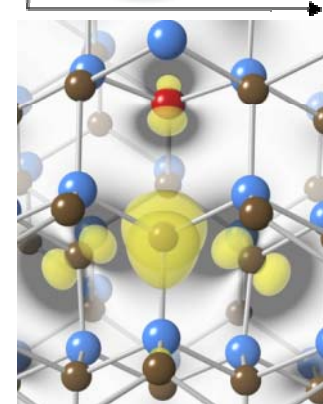
- Location of defect levels relative to band edges
single-particle states



- Excitation properties associated with the defect
configuration-coordinate diagrams

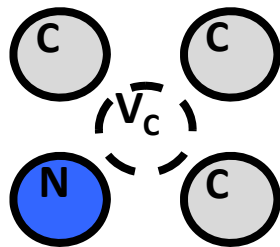
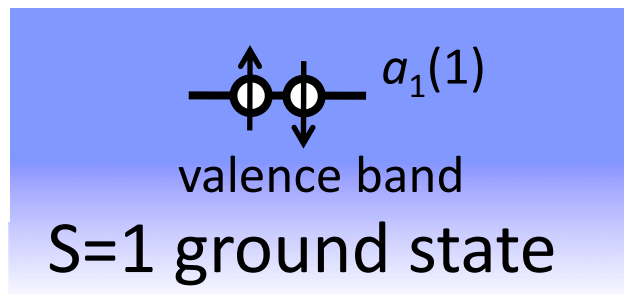
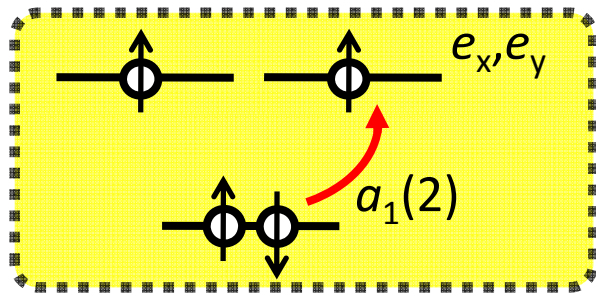


- **Defect stability: spin states, charge states, and binding energies**
 - formation energies



Calculated NV⁻¹ Defect Levels

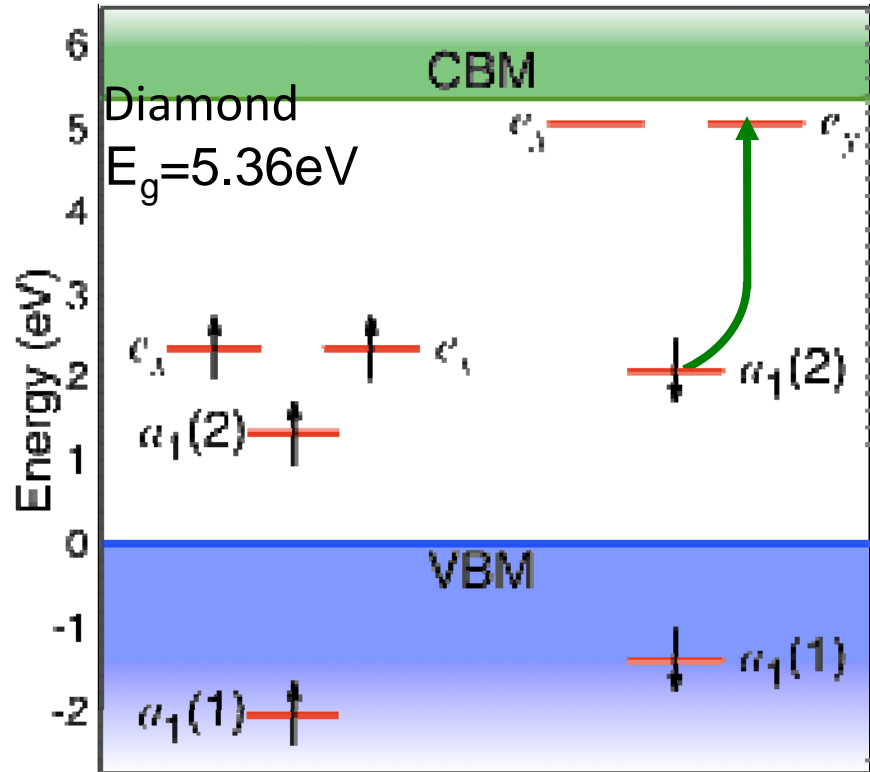
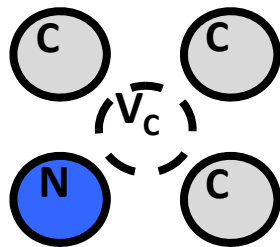
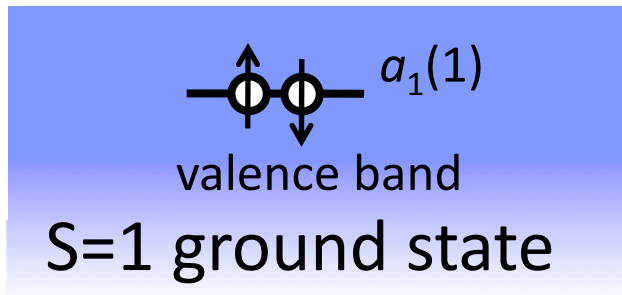
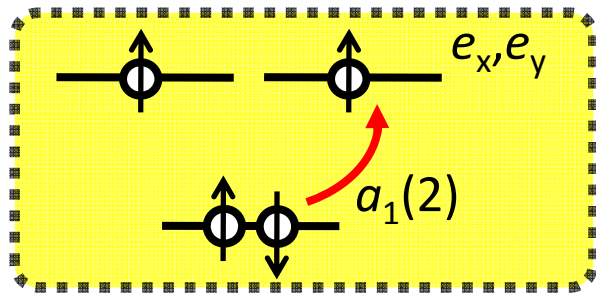
Single-particle picture



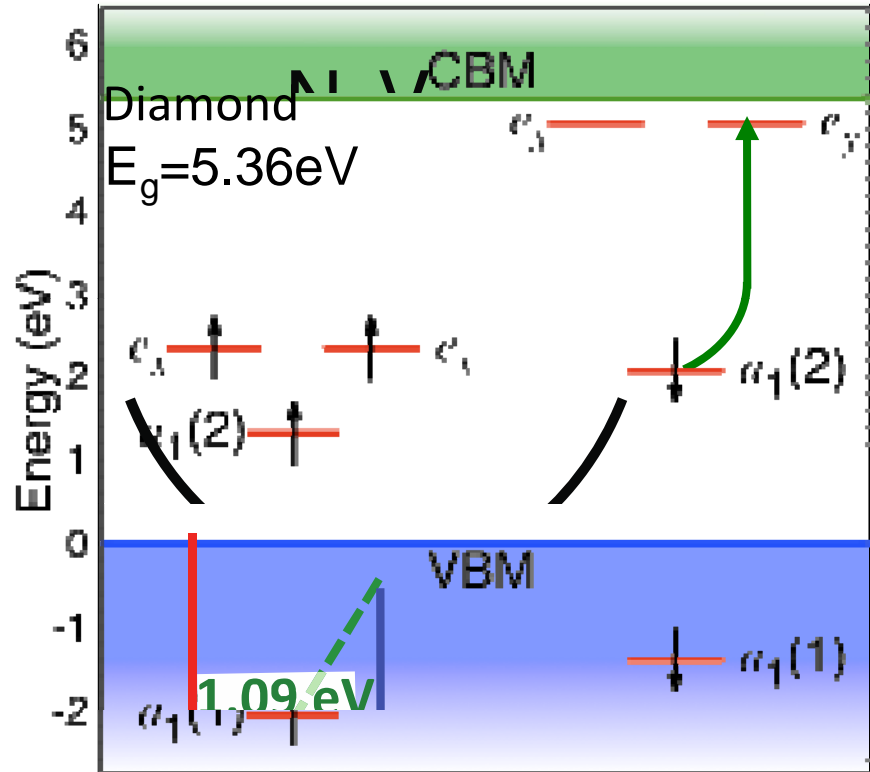
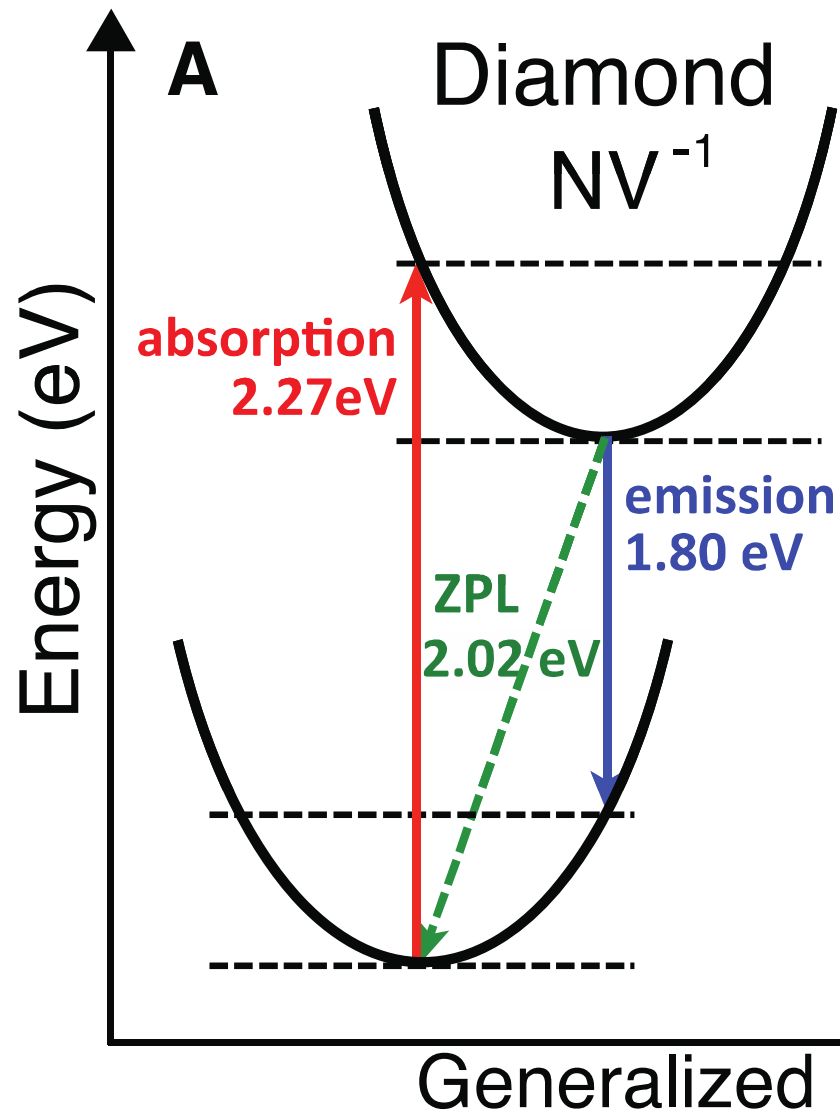
- Spin polarized calculation
 - Spin-majority and minority channels
- Excitation of $a_1(2)$ spin-minority electron
- Promoting $a_1(2)$ electron to higher lying empty state

Calculated NV⁻¹ Defect Levels

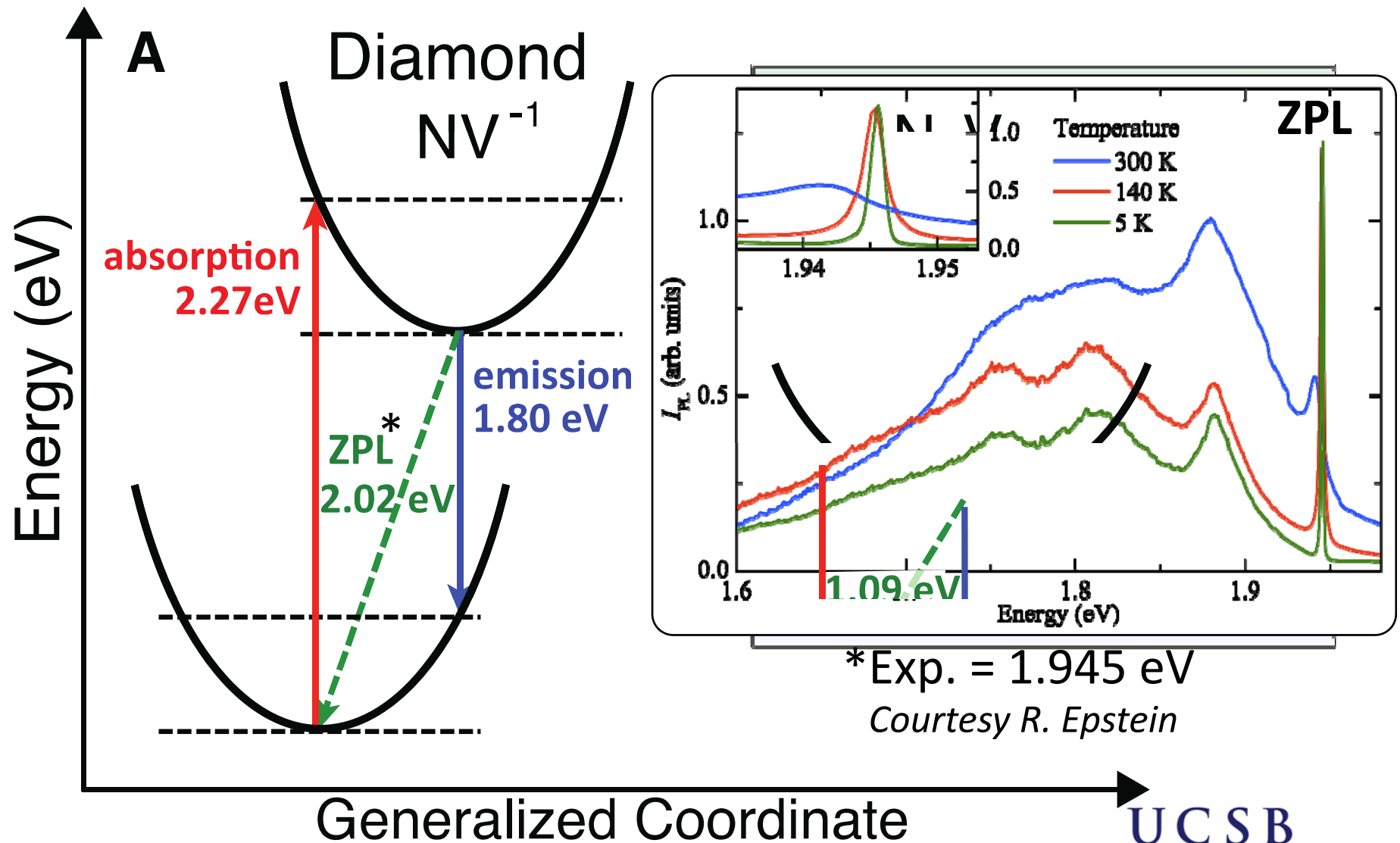
Single-particle picture



Calculated NV⁻¹ Defect Levels

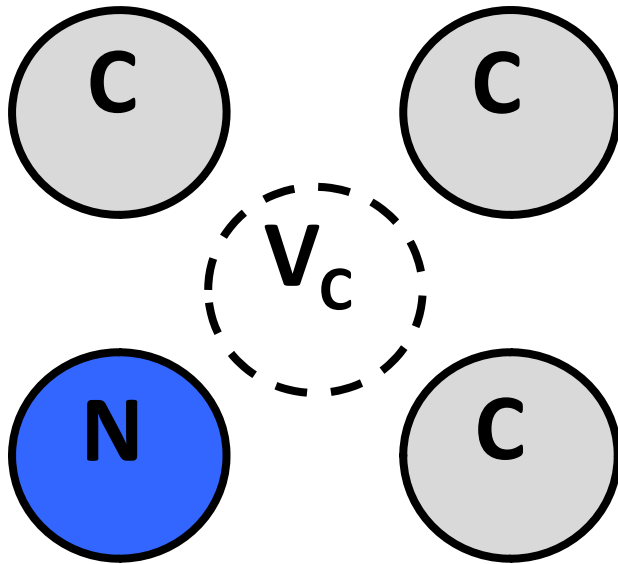


Calculated NV⁻¹ Defect Levels



What About Other Defects?

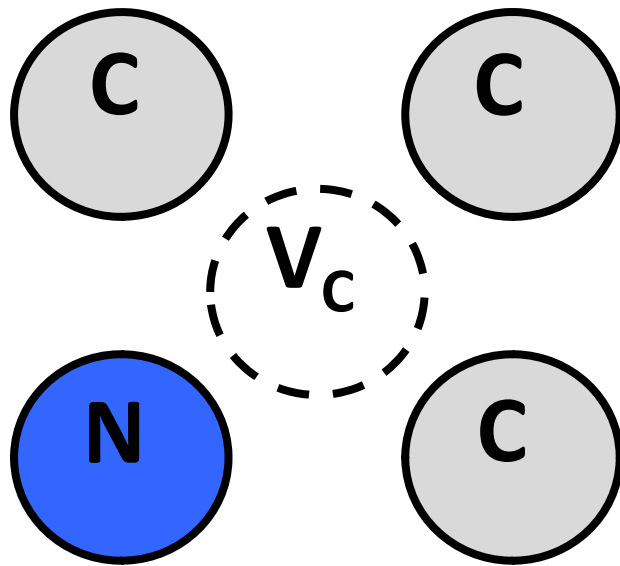
NV: Diamond



$$E_g = 5.36 \text{ eV}$$

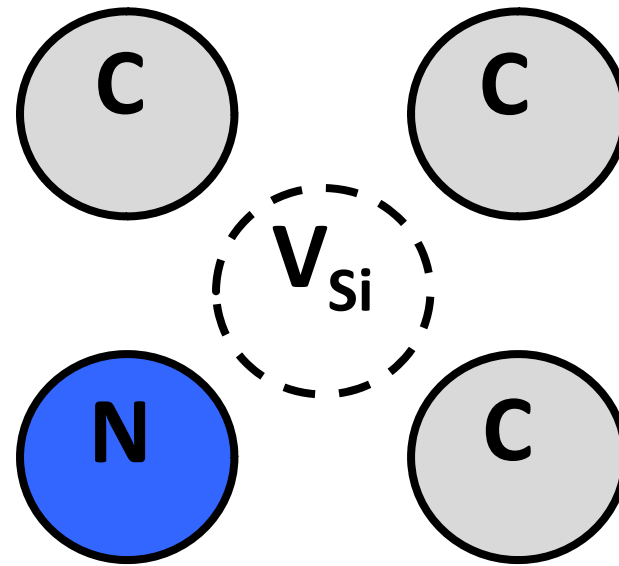
What About Other Defects?

NV: Diamond



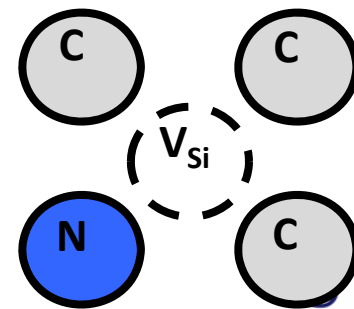
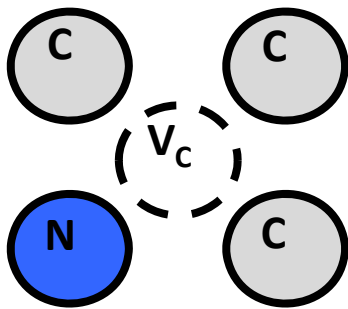
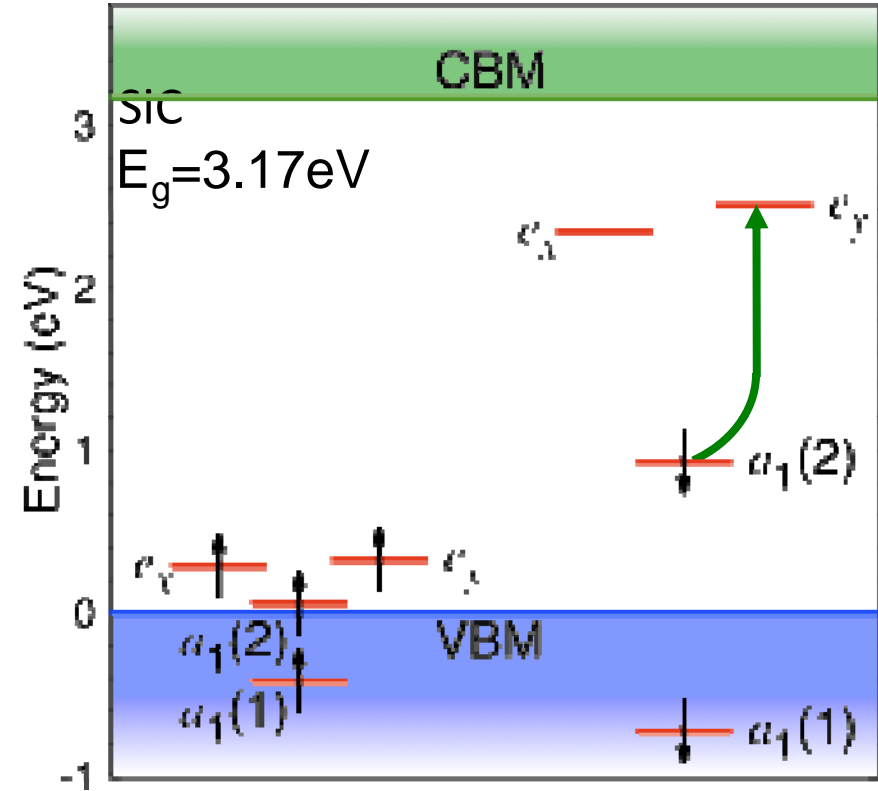
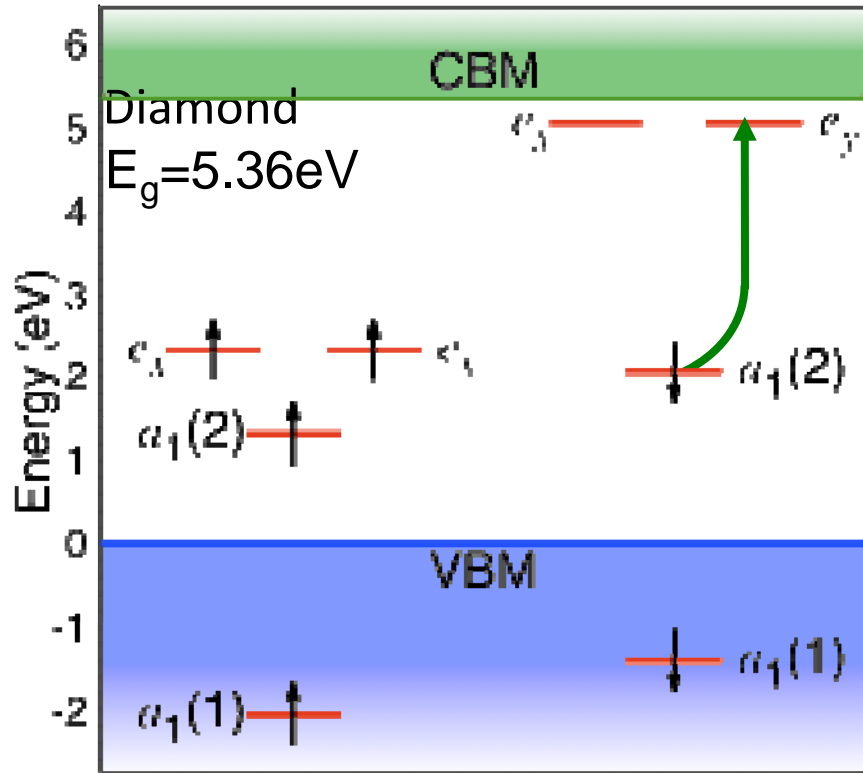
$$E_g = 5.36 \text{ eV}$$

$N_C V_{Si}$: SiC (4H)

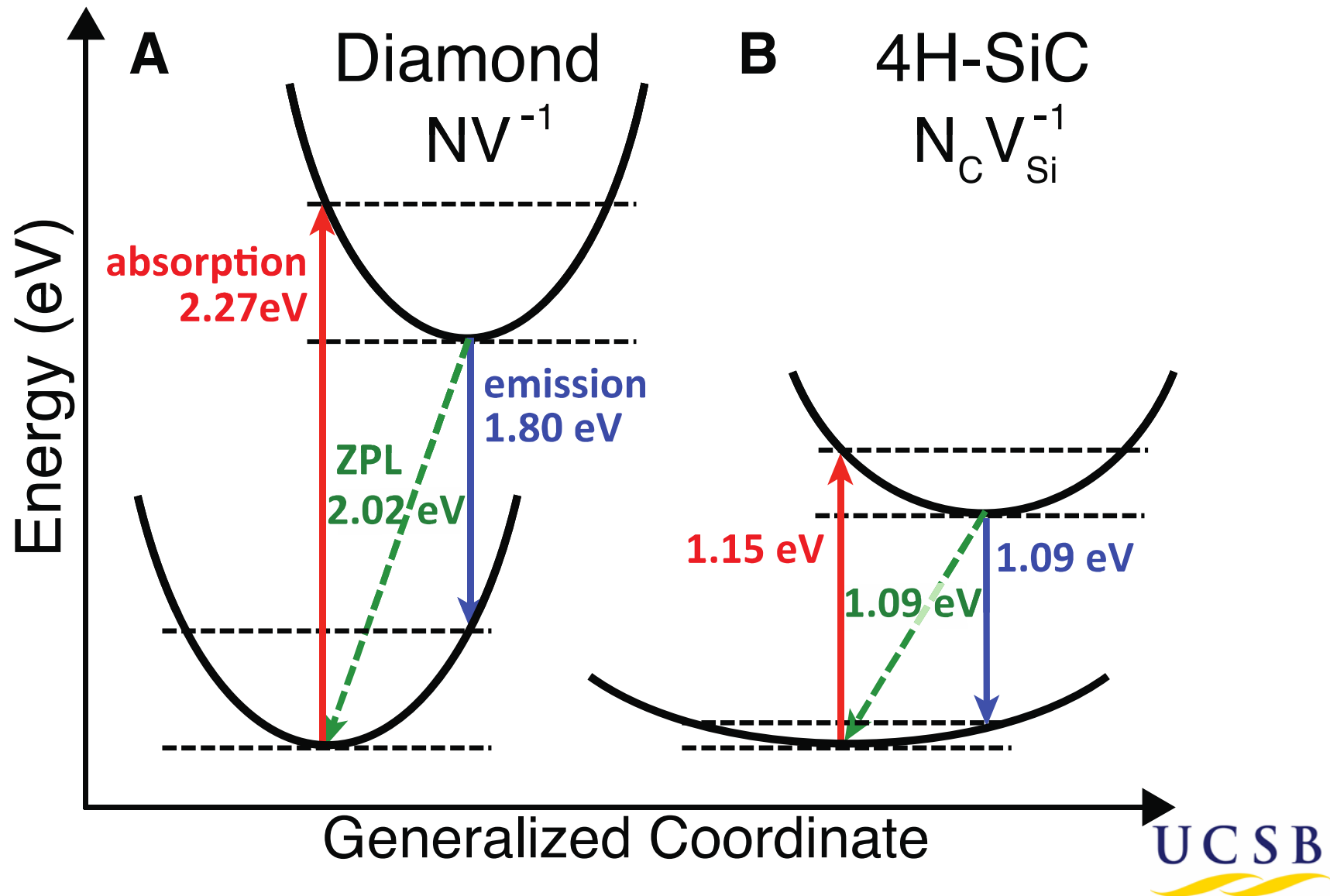


$$E_g = 3.17 \text{ eV}$$

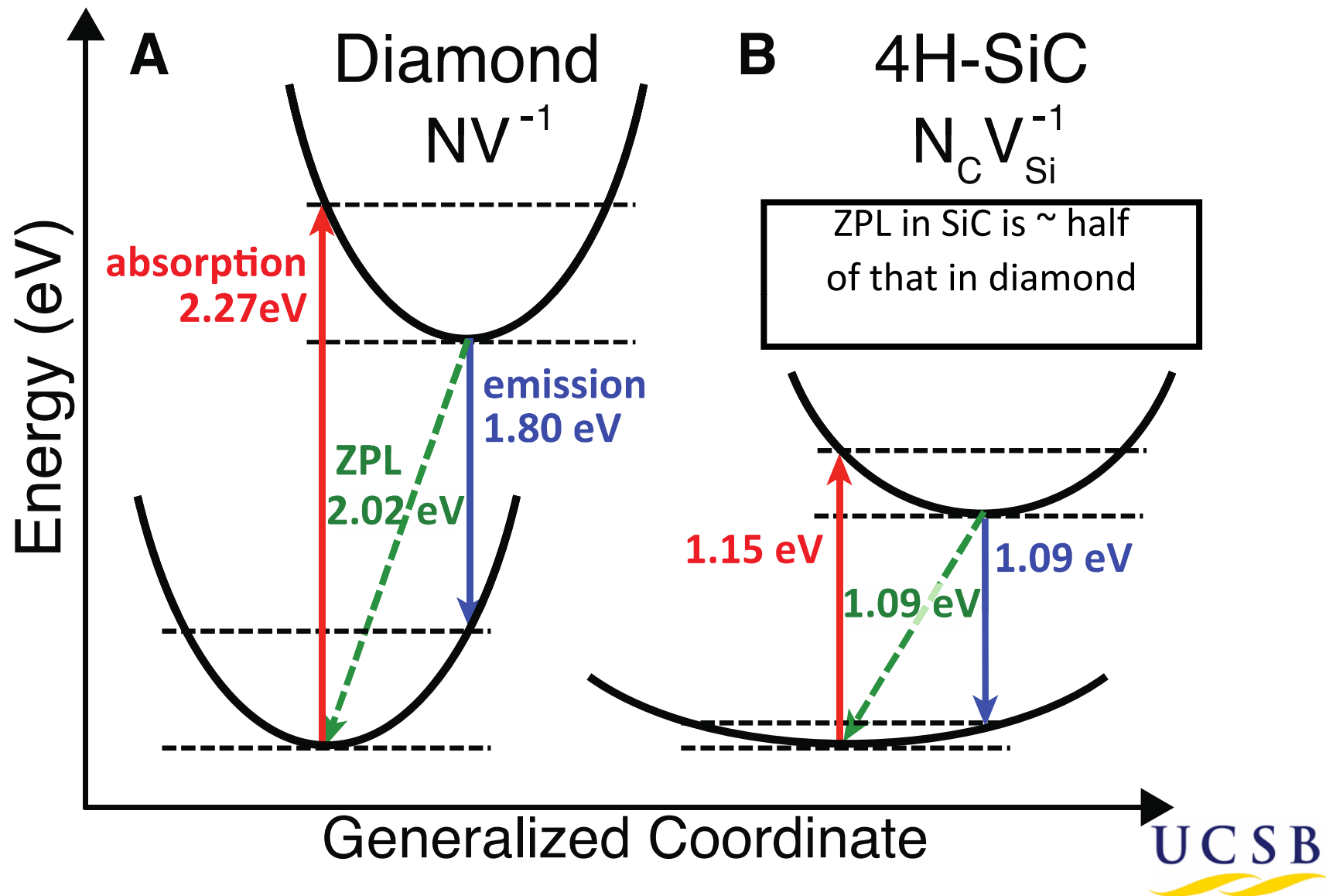
Calculated NV⁻¹ Defect Levels



Calculated NV⁻¹ Excitation Energies



Calculated NV⁻¹ Excitation Energies



Summary

- **Understanding of point defects**
- **Examples where hybrid functionals are crucial for correct physics**
- **Point defects not responsible for unintentional conductivity**
 - Unintentional incorporation of impurities
- **Problems achieving p -type conductivity in wide-band-gap materials**
- **Transport limited by polarons**
- **Defects can be useful!**
 - Qubits for quantum computing