



Inverse methods of Data Analysis in Neutron Scattering

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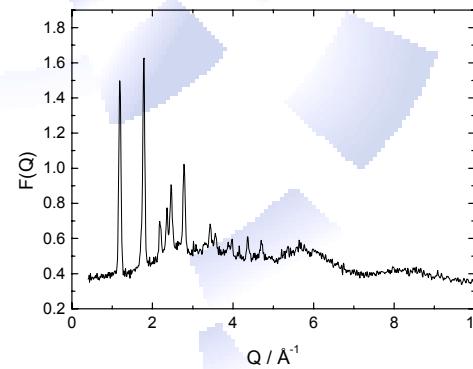


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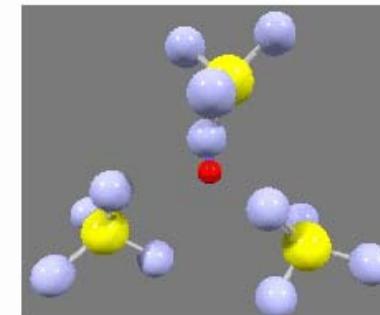
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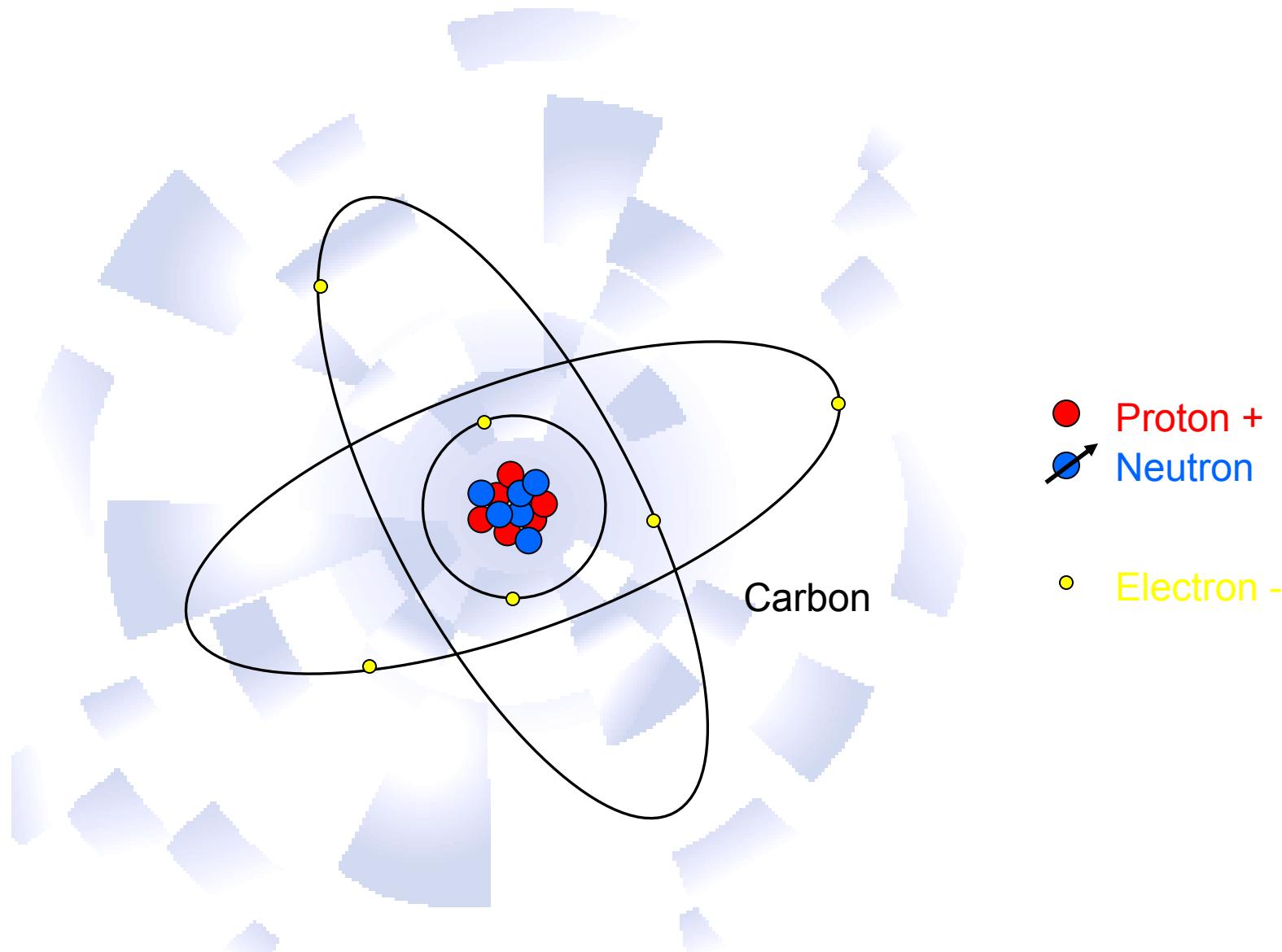
From here



..... to there



What is neutron scattering? A short introduction



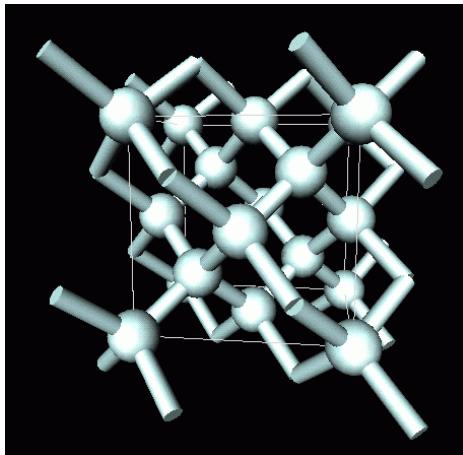


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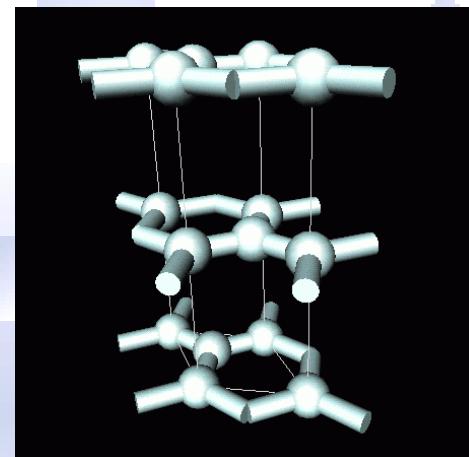


CCLRC

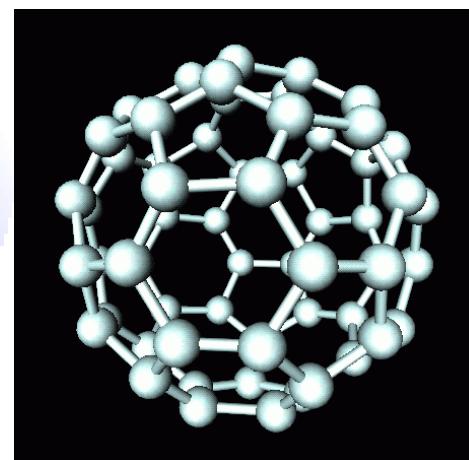
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Diamond

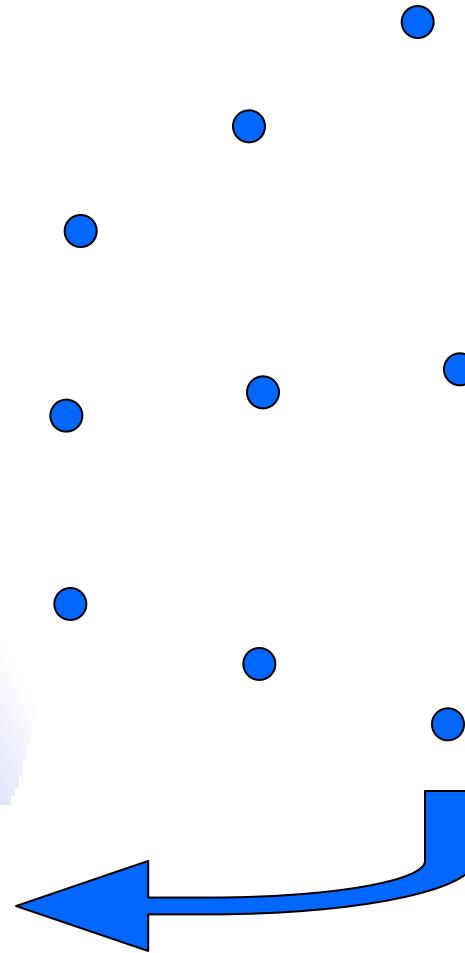
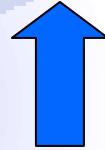
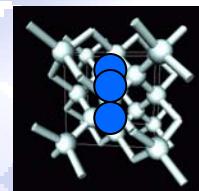
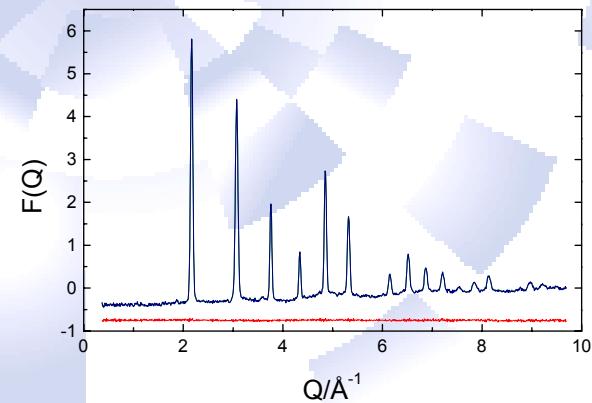


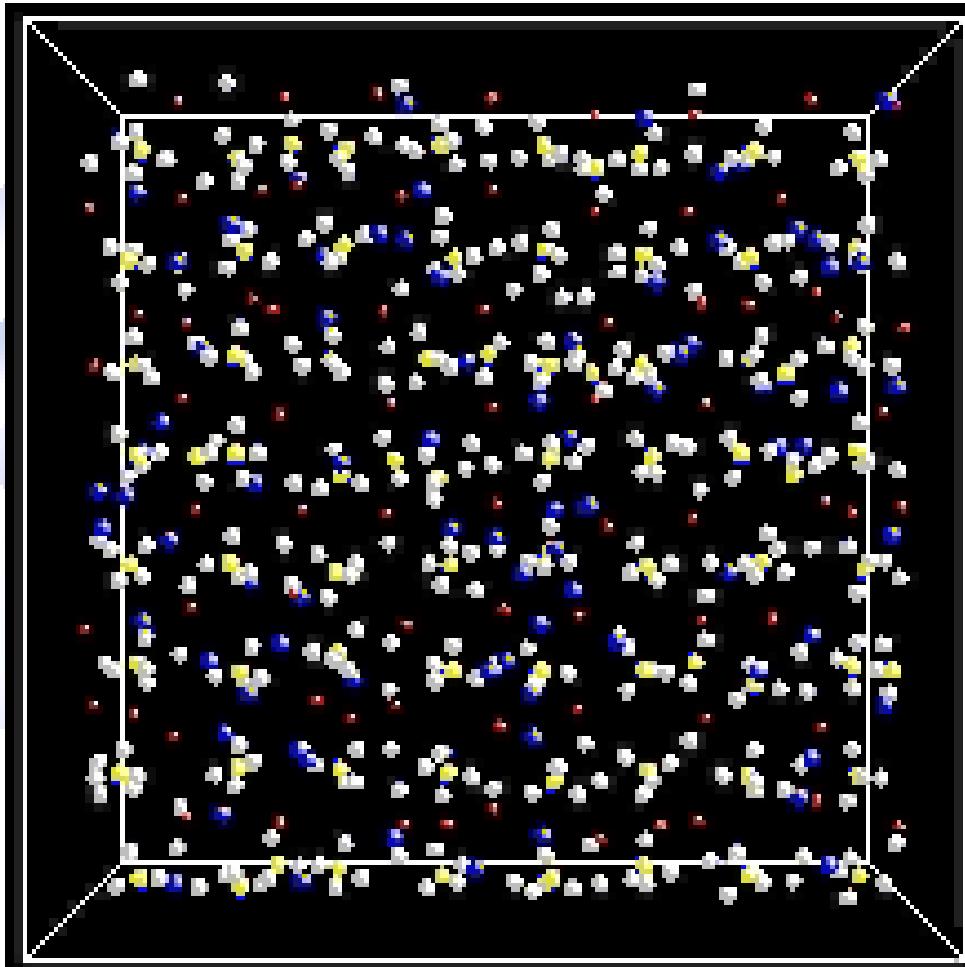
Graphite



Fullerene

Neutron Diffraction



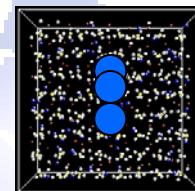
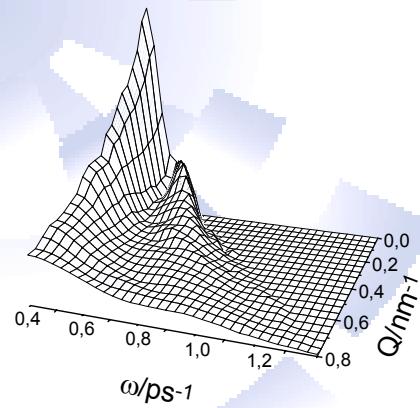


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Inelastic Neutron Scattering



The Nobel Prize in Physics 1994



Clifford G. Shull, MIT, Cambridge, Massachusetts, USA, winner one half of the 1994 Nobel Prize in Physics for development of the neutron diffraction technique.



Shell made use of neutron scattering i.e. of neutrons which change direction without losing energy when they collide with atoms.

Because of the wave nature of neutrons, a diffraction pattern can be measured which indicates where and how the atoms are situated. From the placing of light elements such as hydrogen in metallic hydrides, or hydrogen, carbon and oxygen in organic substances can be determined.

The pattern also shows how atomic dipoles are oriented in magnetic materials, since neutrons are affected by magnetic forces. Shell also made use of this phenomenon in his neutron diffraction technique.



Neutrons see more than X-rays

For many years it was thought that X-rays were the best way to see atoms. But it is now known that neutrons are better. This is because neutrons have no electric charge, so they can penetrate materials that X-rays cannot. They can also penetrate materials that X-rays can't penetrate.

It is known, however, that neutrons do not penetrate the surfaces of the materials they are used to study. So, they must be scattered from the surface to get to the interior of the material. This makes it difficult to use neutrons to study the interior of materials.

Neutrons reveal inner stresses

It is known that neutrons can penetrate deep into materials. This makes it possible to study the interior of materials. But it is also known that neutrons are scattered from the surfaces of materials. This makes it difficult to use neutrons to study the interior of materials.

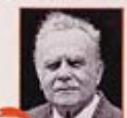
Neutrons show what atoms remember

It is known that neutrons can penetrate deep into materials. This makes it possible to study the interior of materials. But it is also known that neutrons are scattered from the surfaces of materials. This makes it difficult to use neutrons to study the interior of materials.

Neutrons behave as particles and as waves

The Royal Swedish Academy of Sciences has awarded the 1994 Nobel Prize in Physics for pioneering contributions to the development of neutron scattering techniques for studies of condensed matter.

Bertram H. Brockhouse, McMaster University, Hamilton, Ontario, Canada, winner one half of the 1994 Nobel Prize in Physics for the development of neutron spectroscopy.



Brockhouse made use of inelastic scattering i.e. of neutrons, which change both direction and energy when they collide with atoms. They then start or cause atomic oscillations to appear and measure these oscillations in liquid metals. Neutrons can also interact with spin waves in magnets.

With his 3-axis spectrometer Brockhouse measured energies of phonons (atomic vibrations) and magnons (magnetic waves). He also studied how atomic structures in liquids change with time.

Neutrons reveal structure and dynamics

Neutrons show where atoms are

Neutrons bounce against atomic nuclei. They also react with the magnetism of the atoms.

Neutrons show what atoms do

Changes in the energy of the neutrons are first measured in a detector.

Then the neutrons penetrate the sample they want to inspect. They then interact with phonons in the sample. They then penetrate into the sample again. Finally, they penetrate into the sample again. This is called multiple scattering.



Finally, the neutrons are detected in a detector.

...and the signals from several detectors are combined.



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LIGA
VETENSKAPS
AKADEMIE
EN
THE ROYAL SWEDISH ACADEMY OF SCIENCES

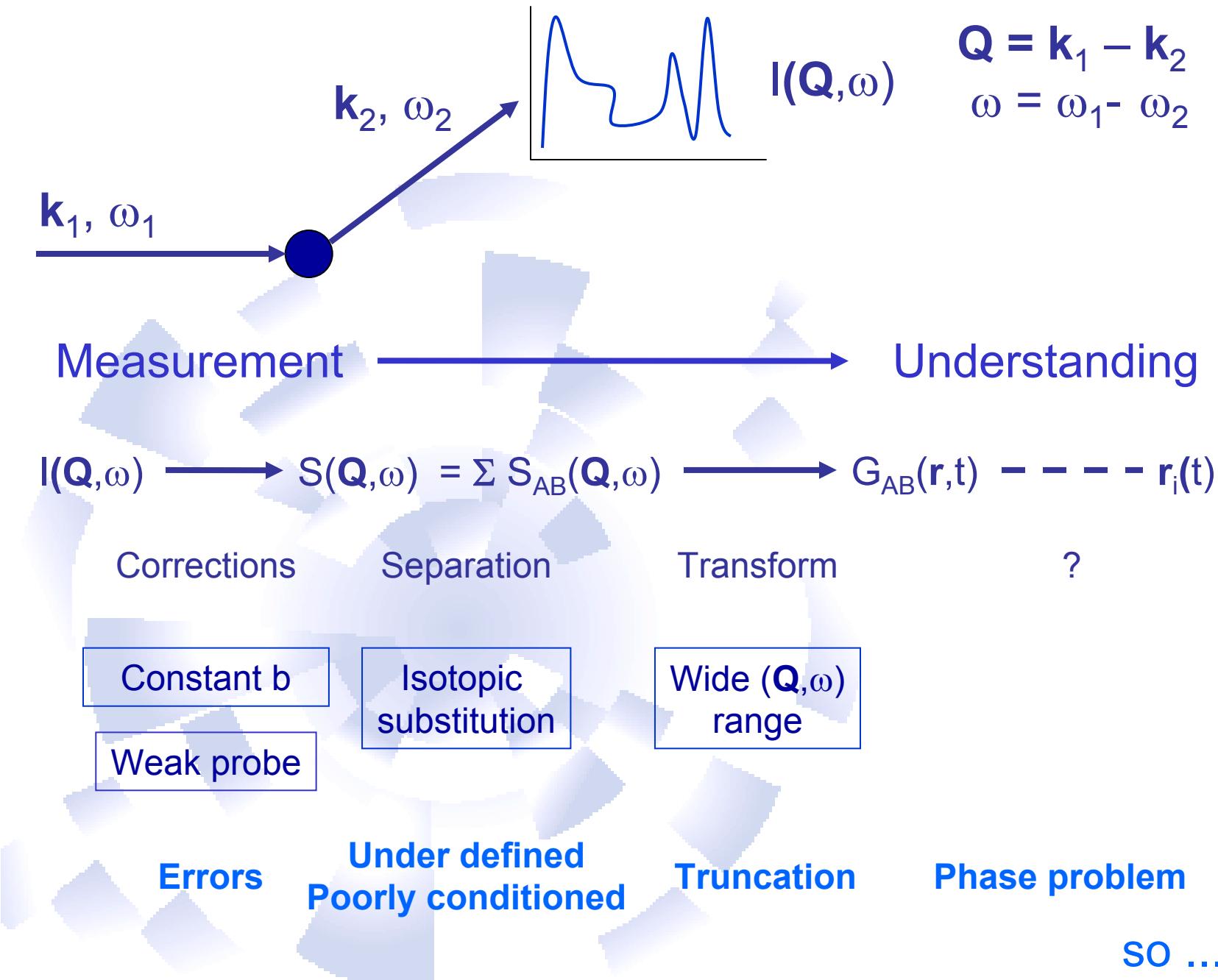
Further reading:
• J. S. Blau and C. G. Shull, *Nature*, 261, 224 (1976).
• J. S. Blau and C. G. Shull, *J. Crystallogr. Struct. Analysis*, 1, 127 (1981).

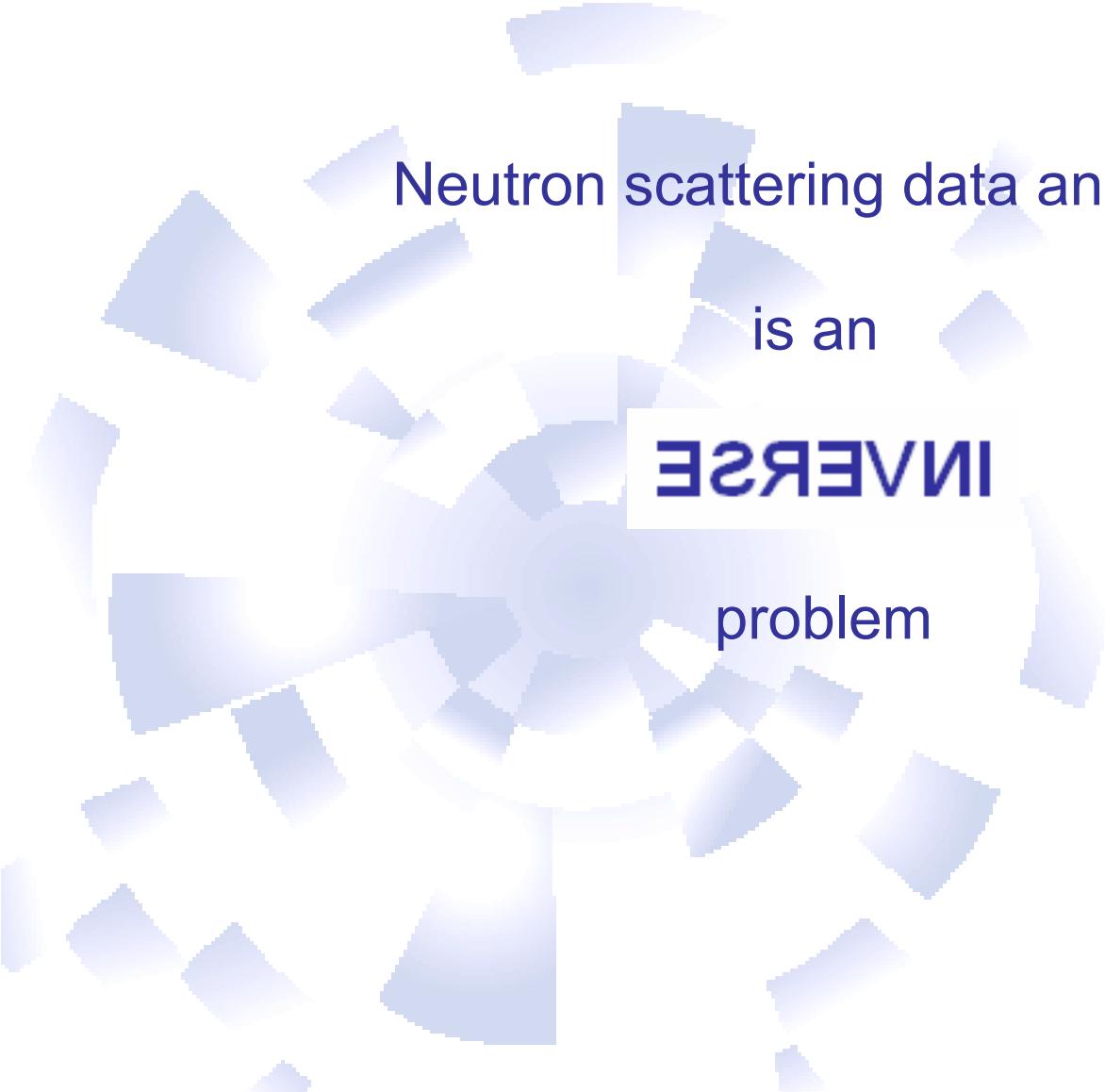
'Neutrons tell you where the atoms are and what the atoms do'

No they don't!

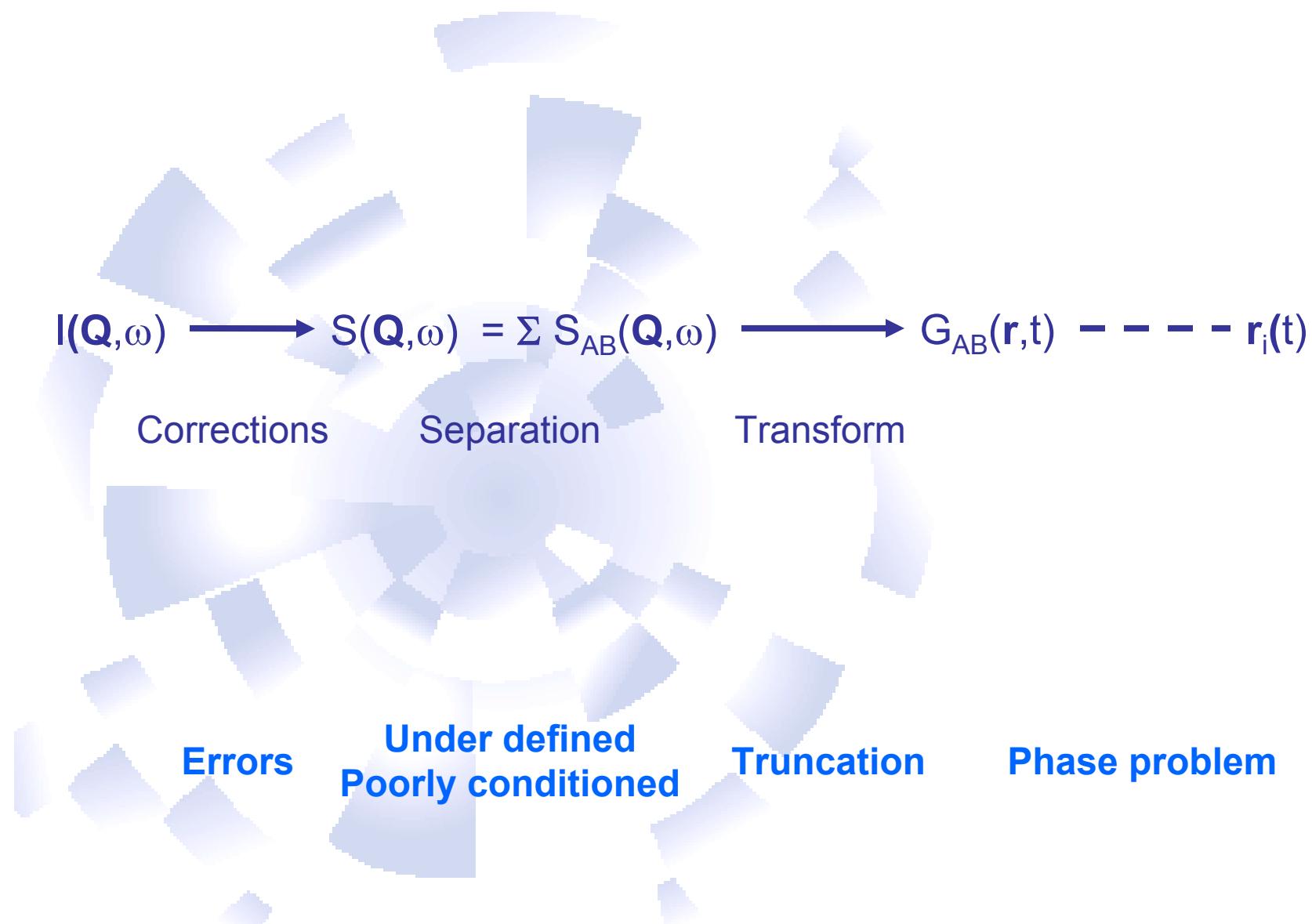
$$\mathbf{Q} = \mathbf{k}_1 - \mathbf{k}_2$$

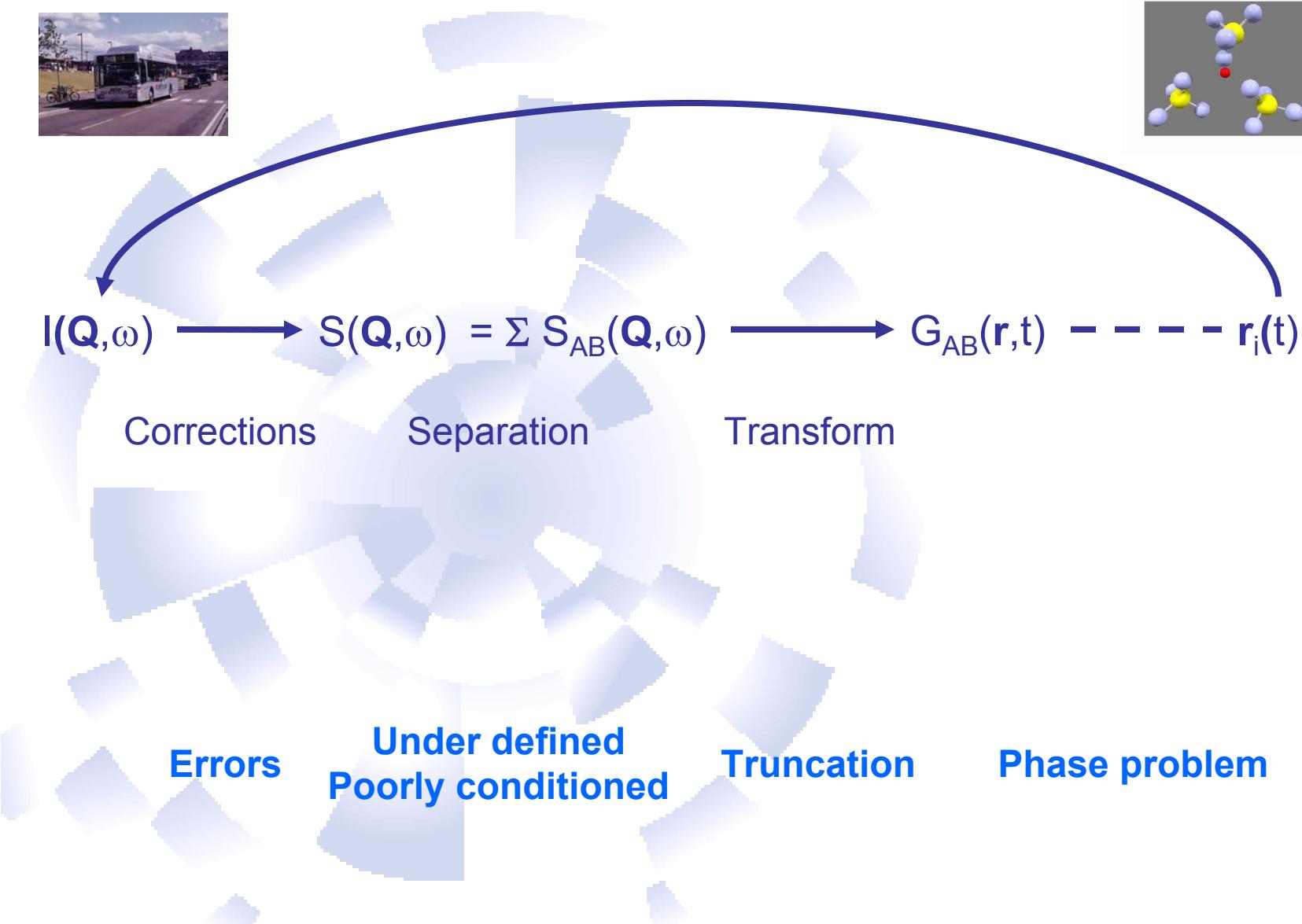
$$\omega = \omega_1 - \omega_2$$





Neutron scattering data analysis
is an
INVERSE
problem





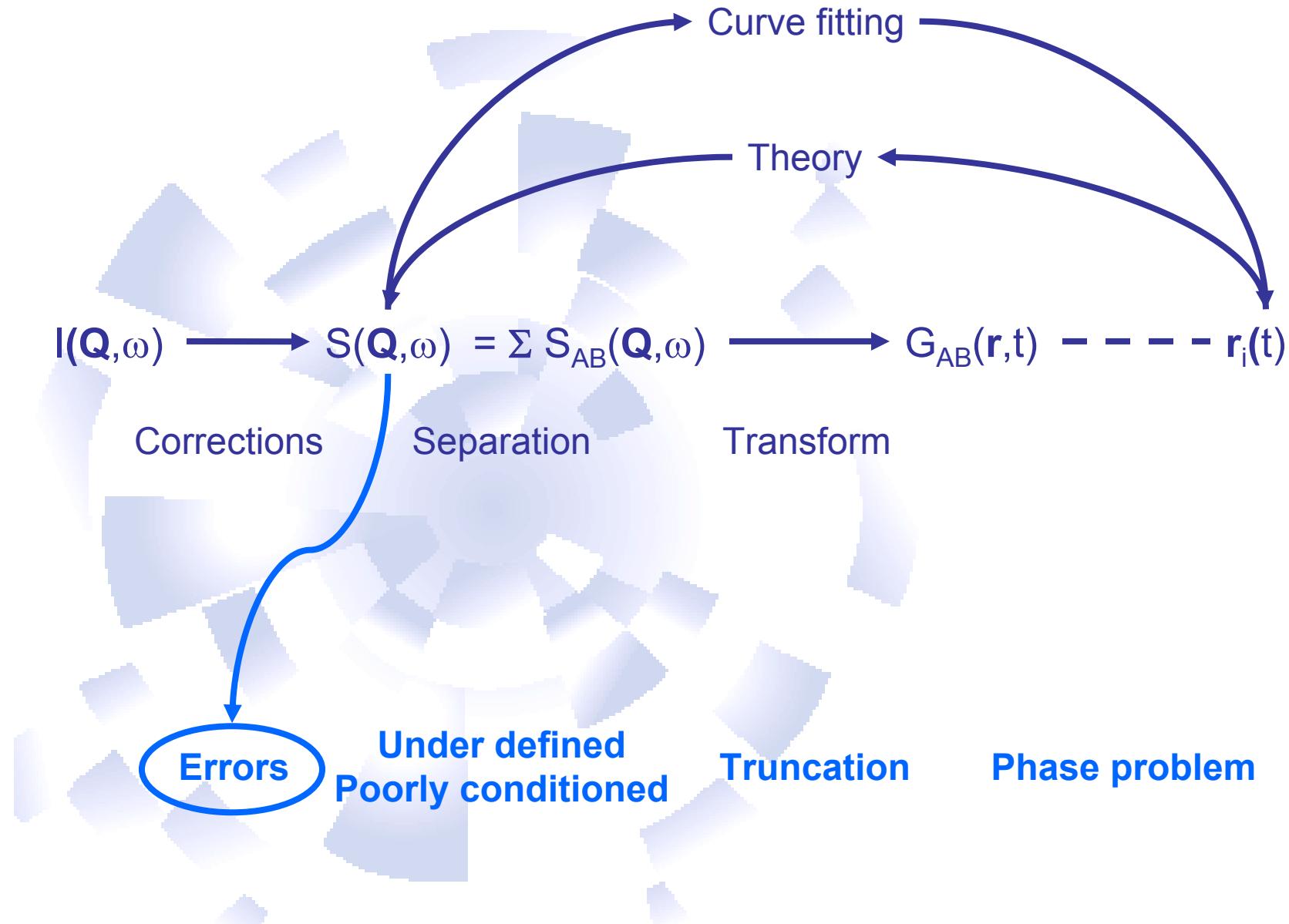
$$S(\mathbf{Q}, \omega) = \frac{1}{2\pi} \int e^{-i\omega t} S(\mathbf{Q}, t) dt \quad S(\mathbf{Q}, t) = \int e^{i\mathbf{Q} \cdot \mathbf{r}} G(\mathbf{r}, t) d\mathbf{r}$$

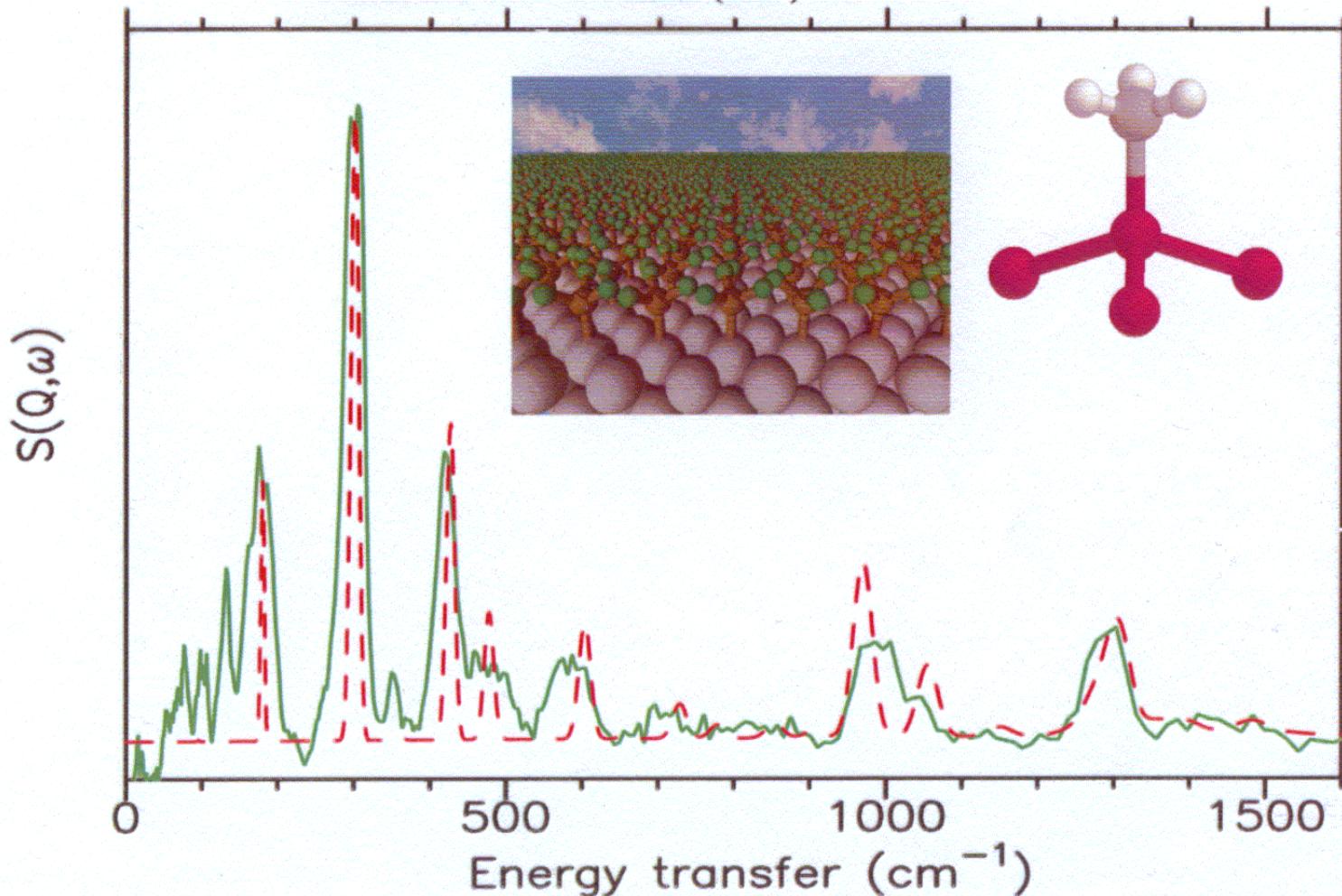
$$S(\mathbf{Q}) = S(\mathbf{Q}, \omega = 0) = \int S(\mathbf{Q}, t) dt = \frac{1}{N} \sum_{j,j'} \int < e^{-i\mathbf{Q} \cdot \mathbf{R}_j(0)} e^{i\mathbf{Q} \cdot \mathbf{R}_{j'}(t)} > dt$$

$$S(\mathbf{Q}) = \frac{1}{N_c} e^{-WQ^2} \left| \sum_k e^{-i\mathbf{Q} \cdot \langle \mathbf{R}_k(t) \rangle} \right|^2$$

$$I(\theta) \propto F(\mathbf{Q}) = \int_{-\infty}^{\infty} S(\mathbf{Q}, \omega) d\omega = S(\mathbf{Q}, t = 0) = \int e^{i\mathbf{Q} \cdot \mathbf{r}} G(\mathbf{r}, t = 0) d\mathbf{r}$$

$$= \frac{1}{N} \sum_{j,j'} < e^{-i\mathbf{Q} \cdot (\mathbf{R}_j(0) - \mathbf{R}_{j'}(0))} > = \frac{1}{N} \left| \sum_j e^{-i\mathbf{Q} \cdot \mathbf{R}_j(0)} \right|$$





$$S(\omega) = \left[A_0 \delta(\omega) + \sum_{j=1}^N A_j \frac{\alpha_j}{\omega^2 + \alpha_j^2} \right] \otimes R(\omega) + \beta(\omega) + \gamma(\omega)$$

What is N ?

How big can we justifiably make N on
the basis of the data?

D S Sivia and C J Carlile J. Chem. Phys. **96** 171 1992

Bayes theorem

$$\text{prob}(A_0, A_j, \alpha_j | N, d) \propto \text{prob}(d | A_0, A_j, \alpha_j, N) \times \text{prob}(A_0, A_j, \alpha_j | N)$$

Posterior

Likelihood function

Prior

If we assume a uniform prior (i.e. we don't know anything about the answer beforehand) and independent data (uncorrelated errors) then the most probable 'posterior' is just the 'best fit'.

But to fit we need to know N ...

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$$\text{prob}(N|d) \propto \text{prob}(d|N) \times \text{prob}(N)$$



$$\text{prob}(d|N) \propto \int \dots \int \text{prob}(d, A_0, A_j, \alpha_j | N) \partial A_0 \partial^N A_j \partial^N \alpha_j$$

$$\text{prob}(d|N) \propto \int \dots \int \text{prob}(d | A_0, A_j, \alpha_j, N) \times \text{prob}(A_0, A_j, \alpha_j | N) \partial A_0 \partial^N A_j \partial^N \alpha_j$$

If we assume no knowledge of N then

$$\text{prob}(N|d) \propto \text{prob}(d|N)$$

so

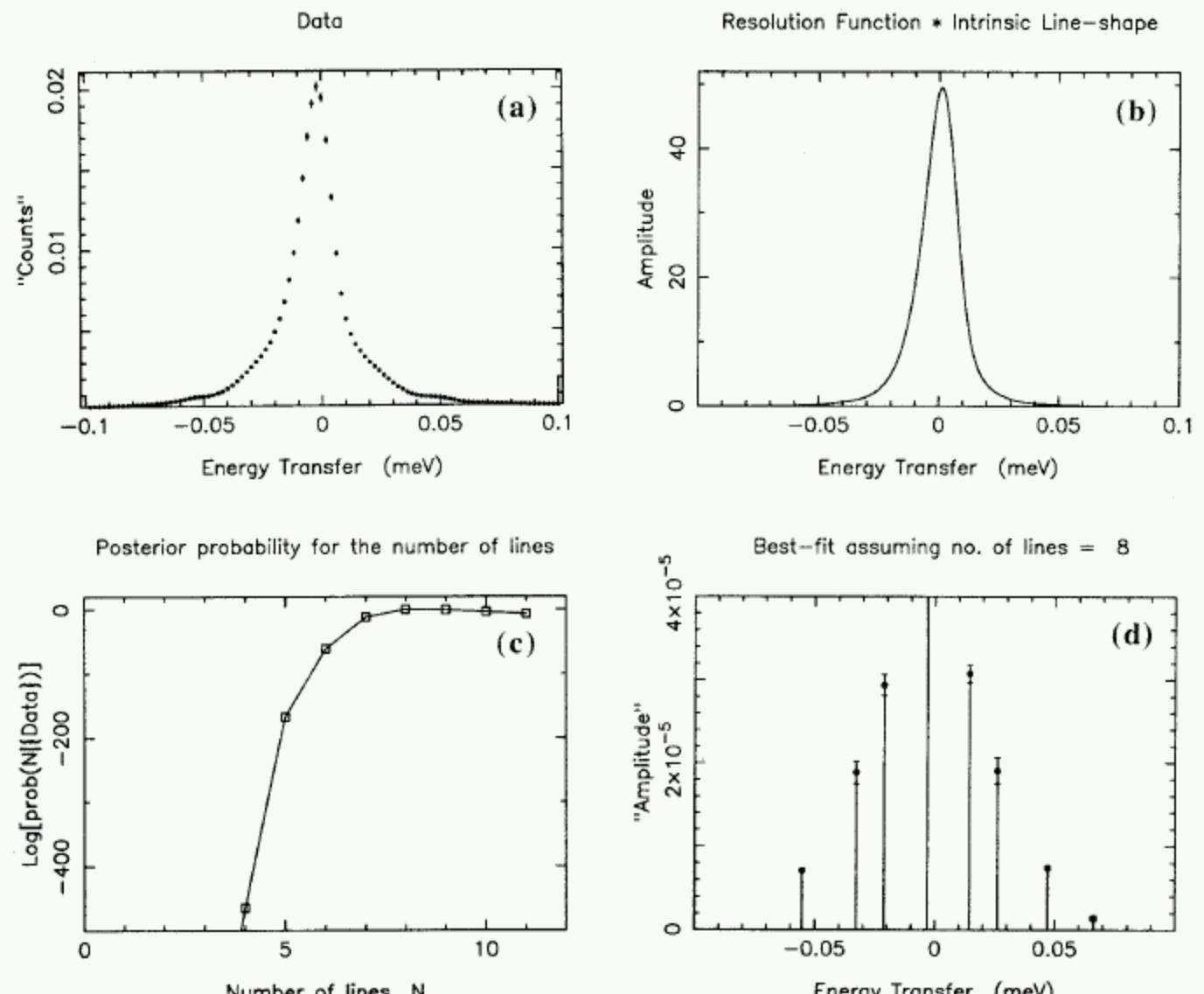
$$\text{prob}(N) = \text{uniform}$$

so calculate the integral above and then maximise

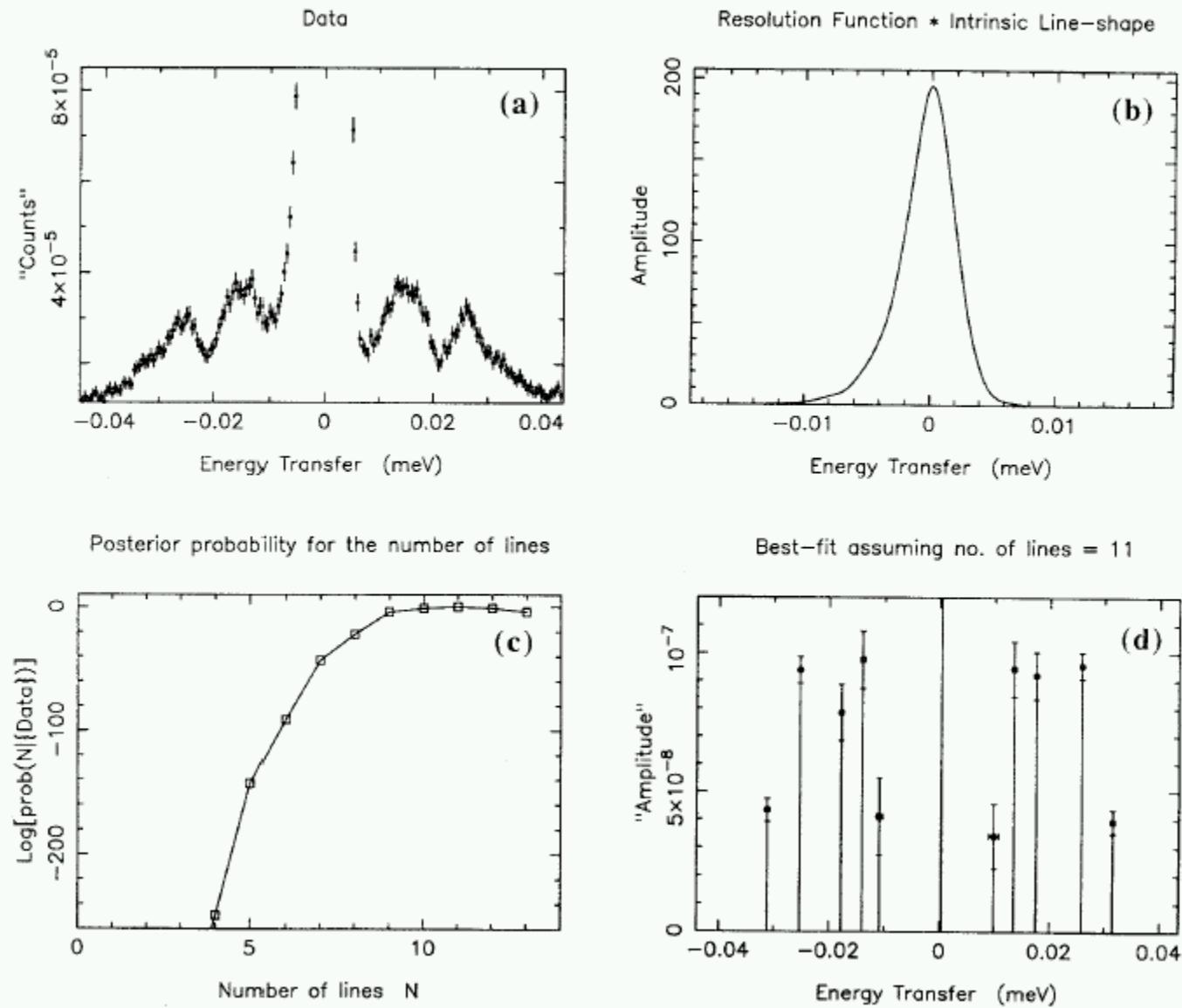
$$\text{prob}(N|d)$$

D S Sivia and C J Carlile J. Chem. Phys. **96** 171 1992

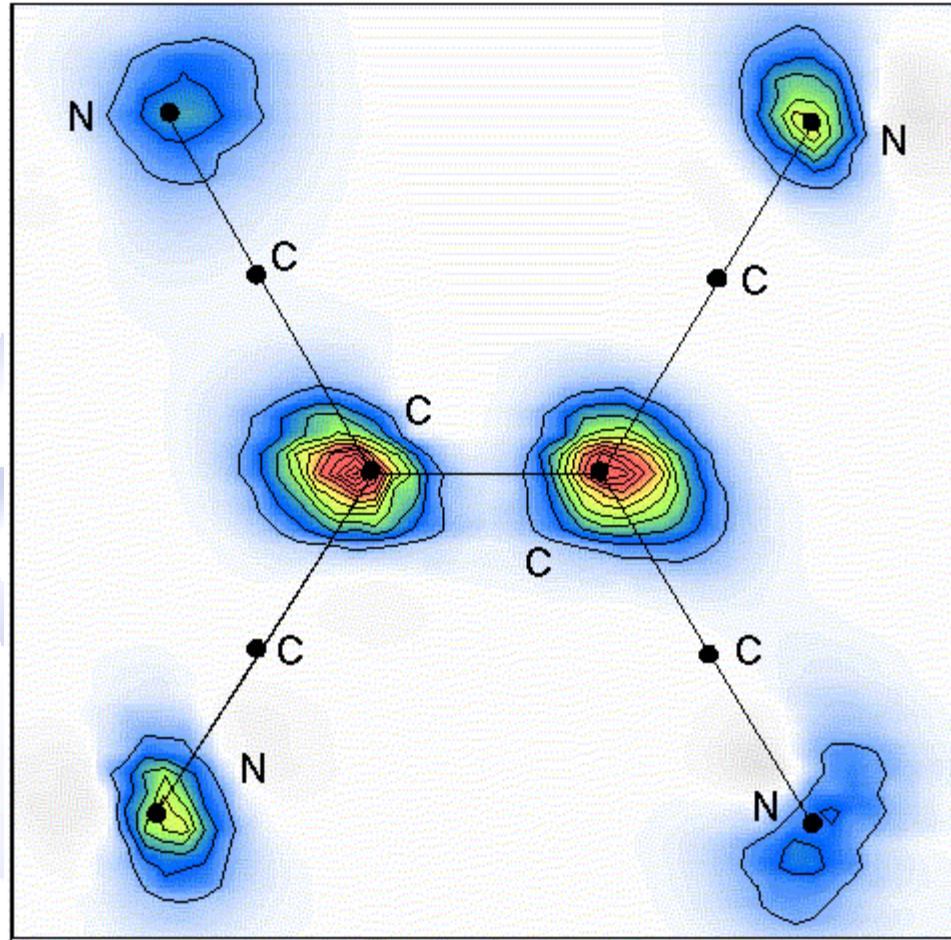




D S Sivia and C J Carlile J. Chem. Phys. **96** 171 1992

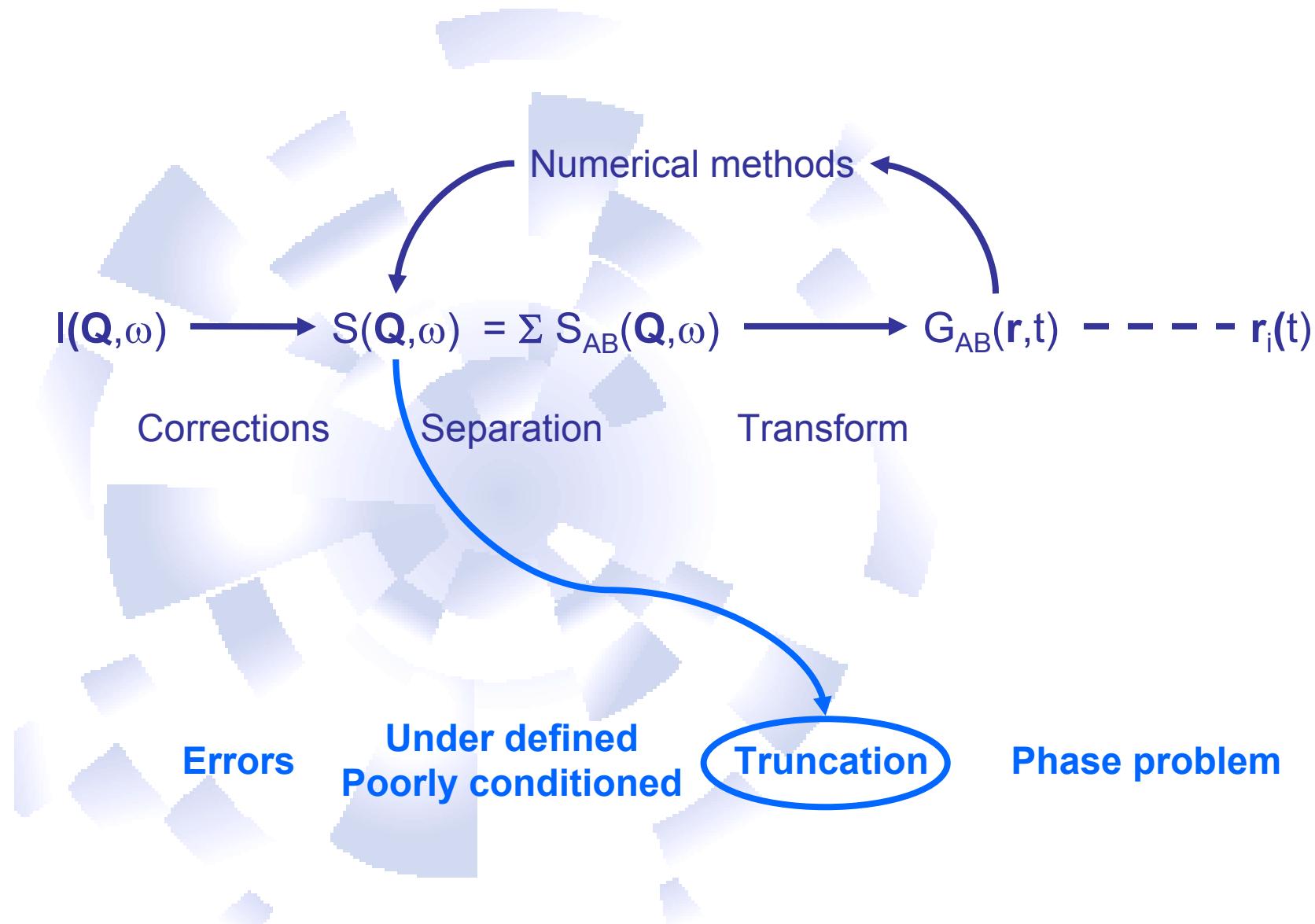


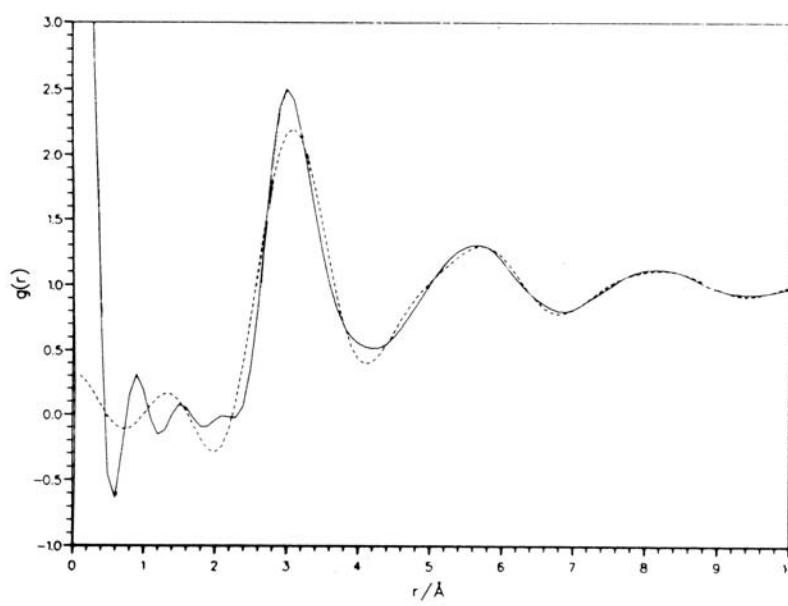
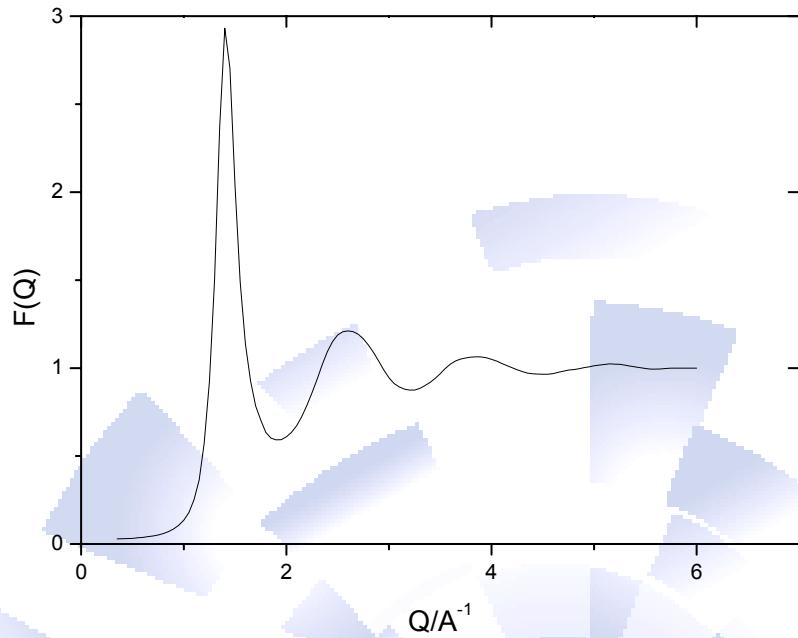
D S Sivia and C J Carlile J. Chem. Phys. **96** 171 1992



What not to do

Choose the prior to get the answer you want





Truncation
(limited data range)



‘Smoothing’

$$n(r) = 4\pi r^2 \rho g(r)$$

$$H = -\sum n(r) \ln(n(r)/p(r))$$

‘Flattest’ (Max. Ent.)

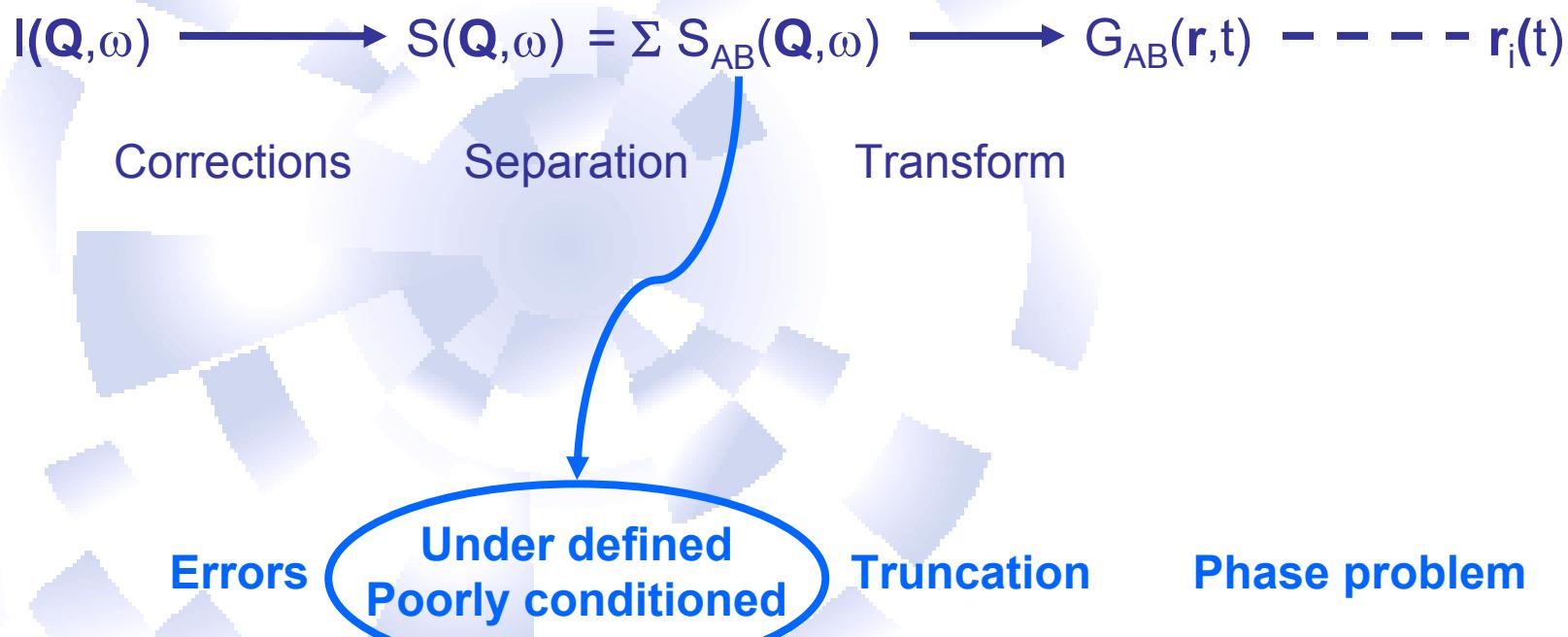
$$H' = \sum w(r) d''(r)$$

$$d(r) = n(r) - p(r)$$

‘Least noisy’

$$I = \int (1 + d'^2(r))^{\frac{1}{2}} dr$$

‘Shortest line’



$$F^{(k)}(Q) = \sum_{\alpha} \sum_{\beta} c_{\alpha\beta}^{(k)} (A_{\alpha\beta}(Q) - 1)$$

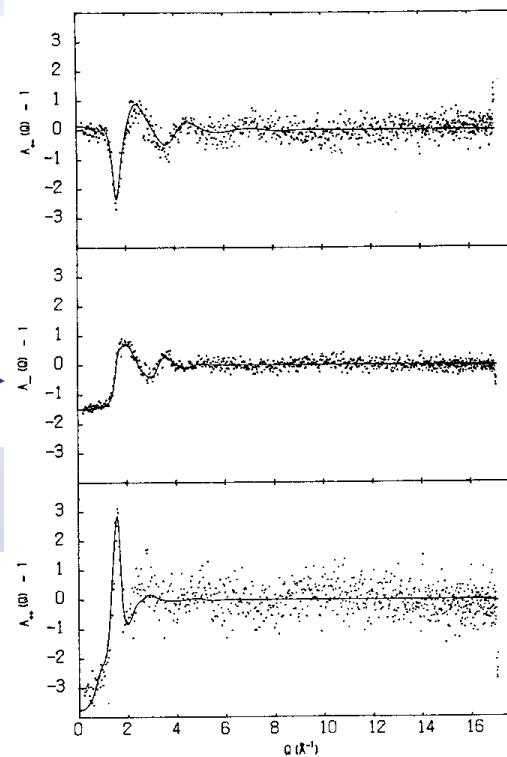
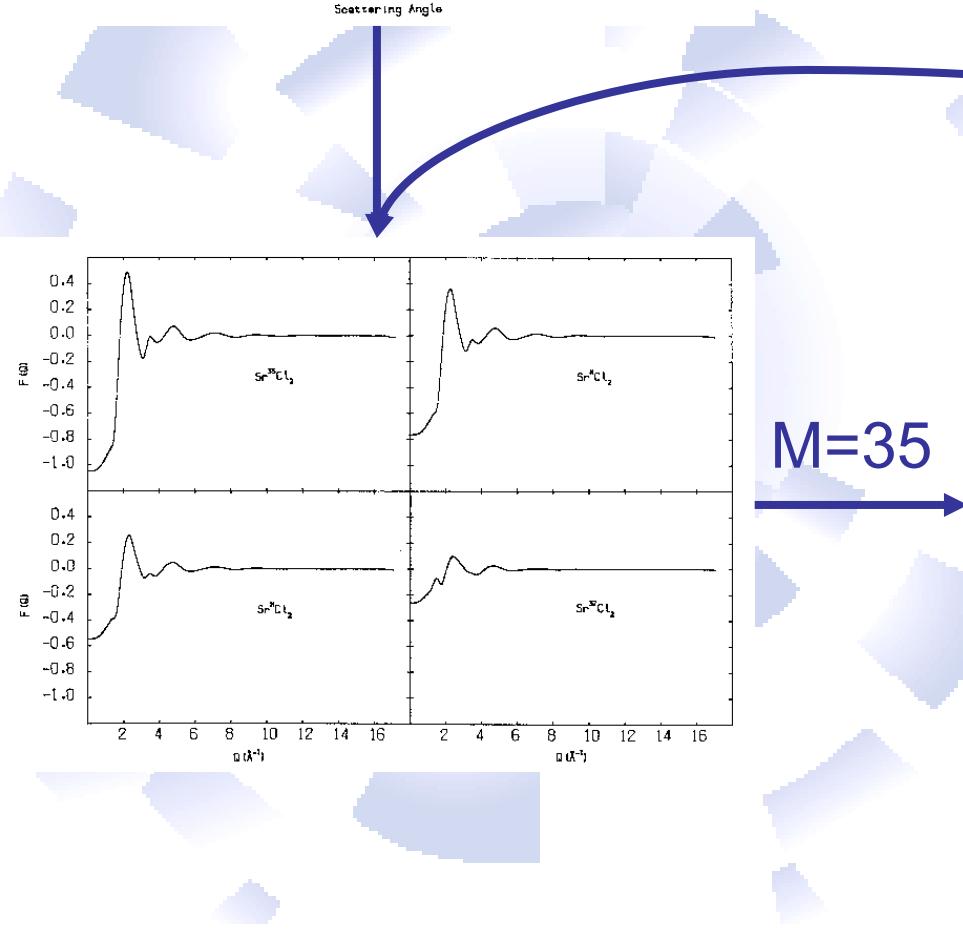
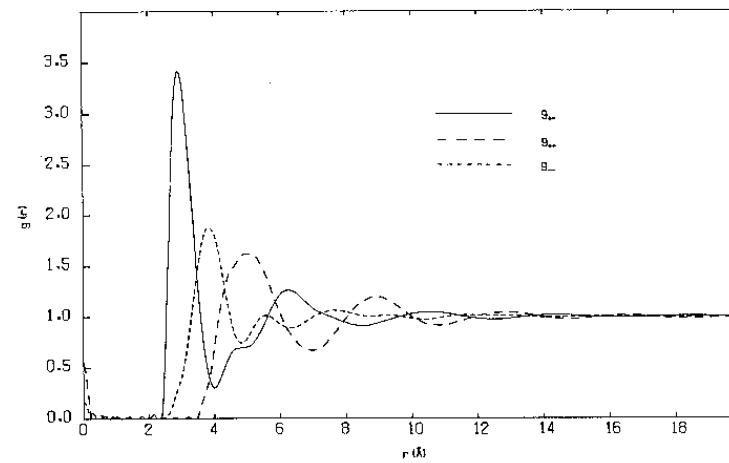
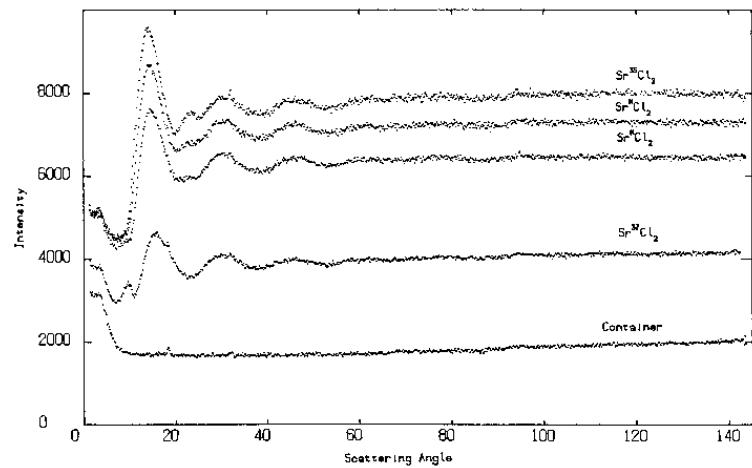
$$\vec{F} = \mathbf{C} \vec{A}$$

$$\vec{A} = \mathbf{C}^{-1} \vec{F}$$

Make a set of measurements with varying coefficients $c_{\alpha\beta}^{(k)}$
(isotopic substitution, anomalous scattering, EXAFS) and solve

C is poorly conditioned (nearly singular), so small errors in F
lead to large errors in A

The Turing M conditioning number is a measure of the
conditioning of matrix \mathbf{C}



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0.136538 0.200896 0.168601
0.754507 0.657018 0.704565
0.641932 0.726614 0.68932

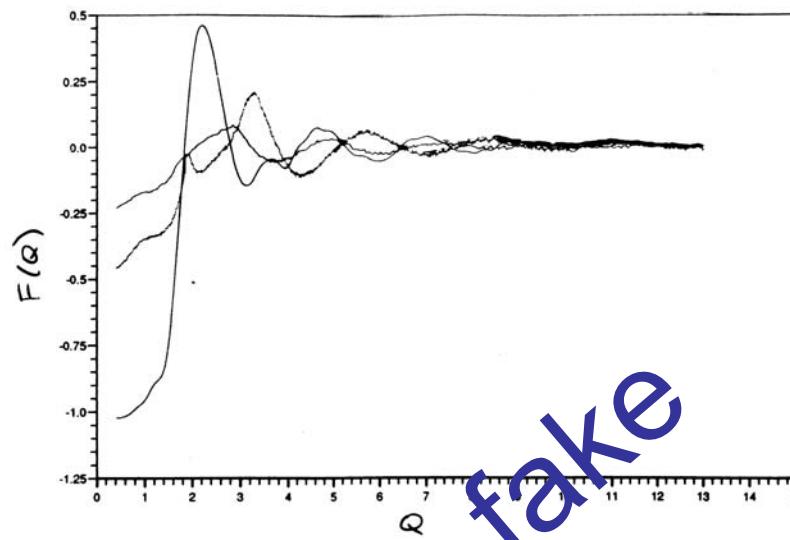
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187.0622 38.92702 -85.54158763
-348.8769 -82.06621 170.6639011

841.7145

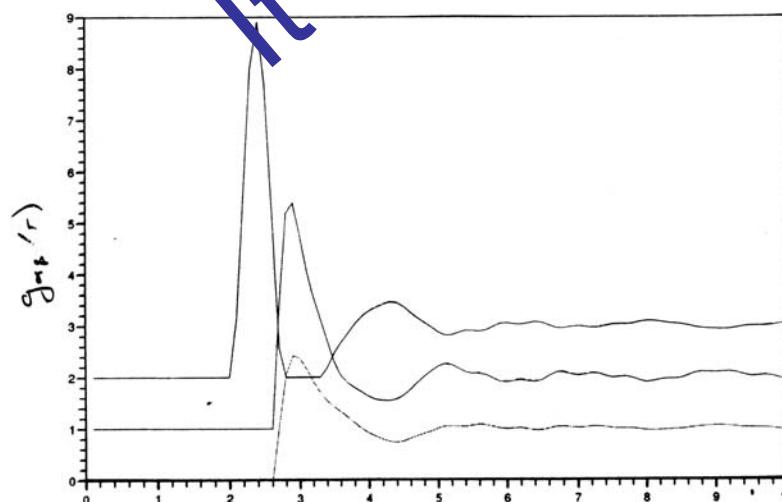


ISIS

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PLOTTED: 13-AUG-1992 14:03:32 by process Robert
FILE: f13.dat



TITLE: ~~g(r)~~
PLOTTED: 13-AUG-1992 14:04:08 by process Robert
FILE: g12.dat

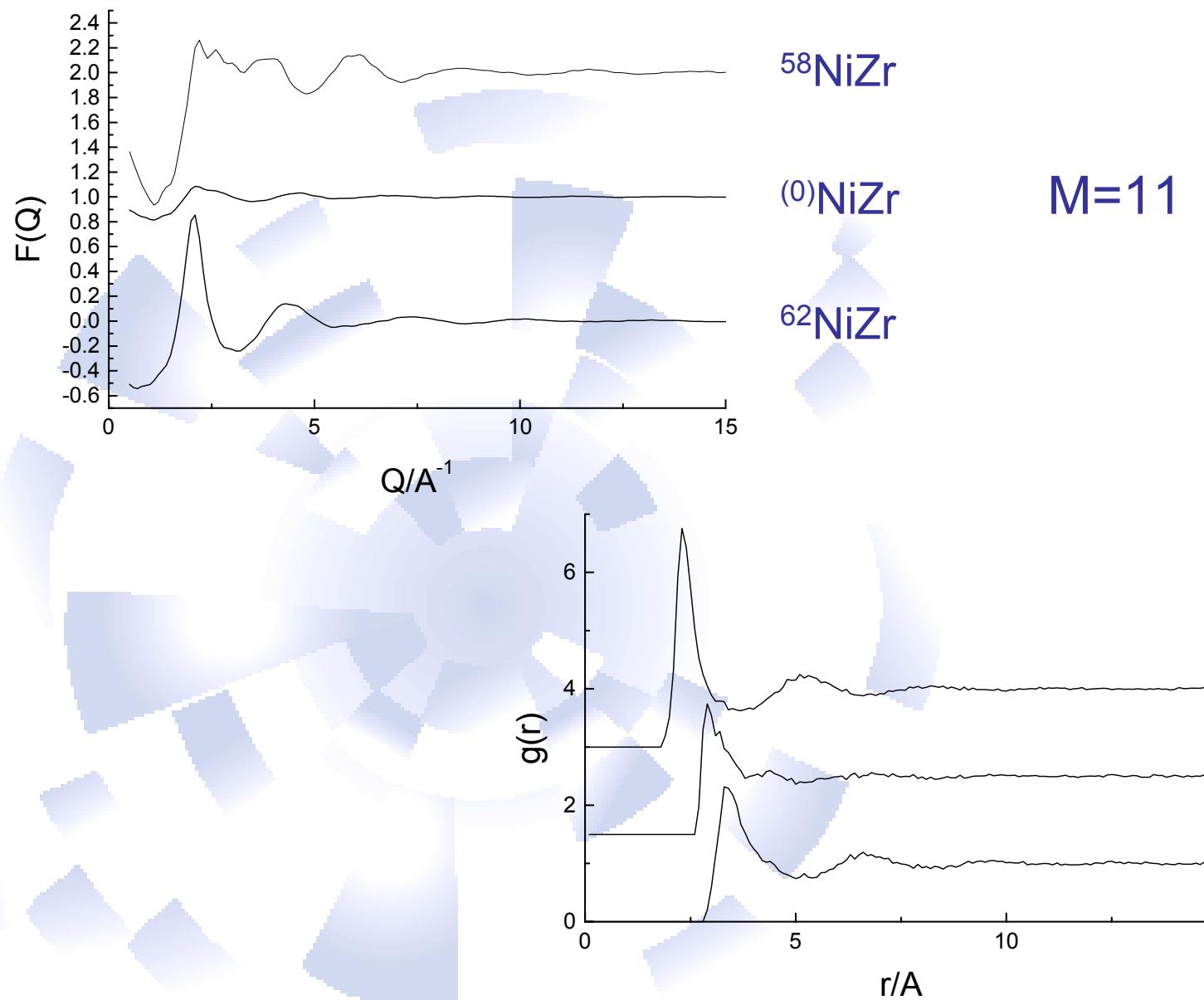


CCLRC

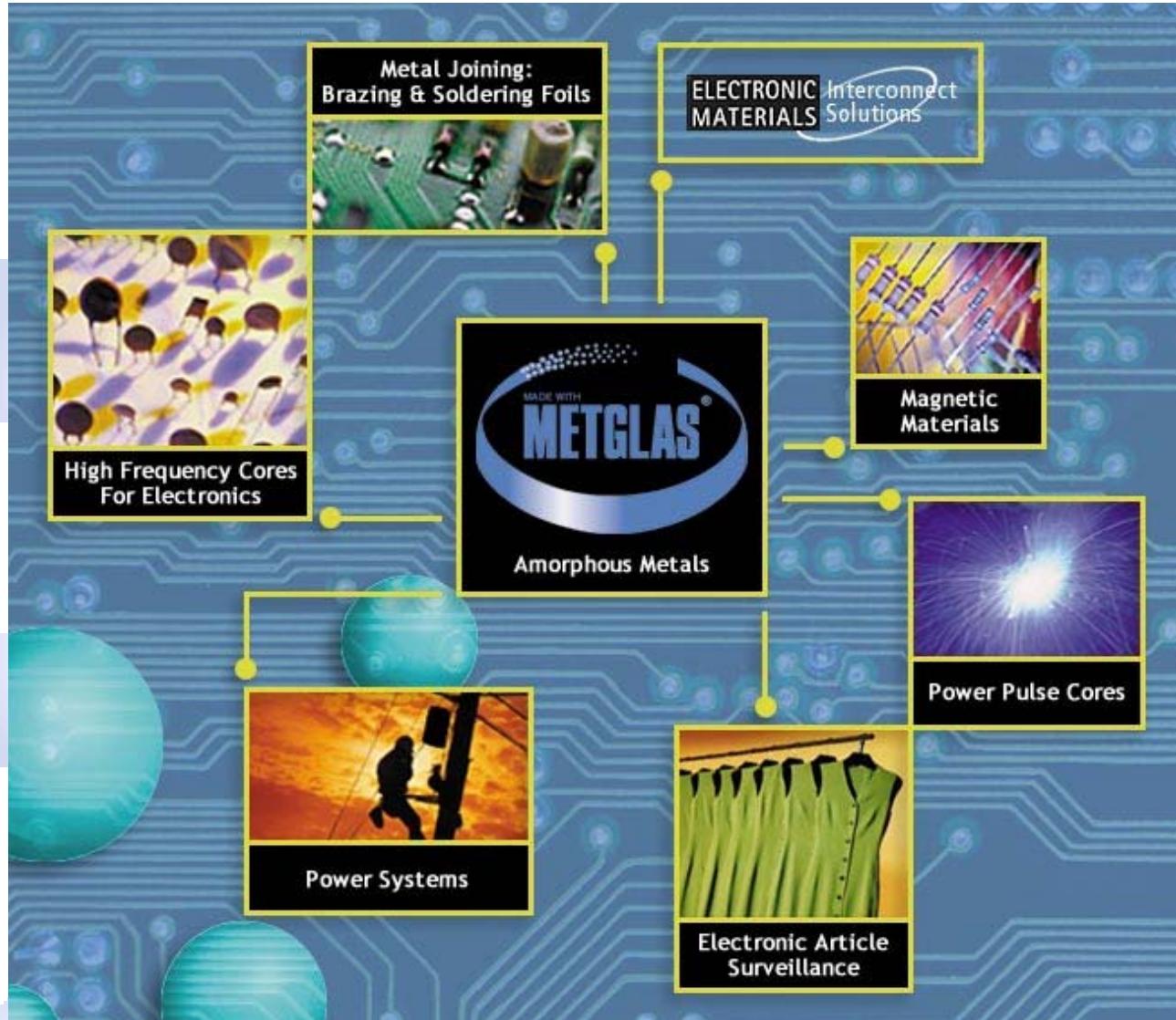
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Metallic glass: NiZr

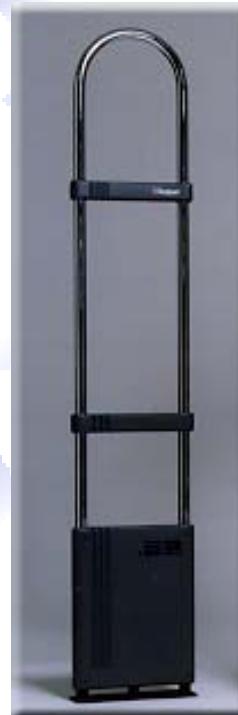


Metallic glasses



ISIS

Metallic glasses



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Metallic glass: CuZr



$^{63}\text{CuZr}$

$(63+65)/2\text{CuZr}$

$^{65}\text{CuZr}$

M=286



$$c^{(k)} = \begin{pmatrix} c_{11} \\ c_{12} \\ \cdot \\ \cdot \\ c_{nn} \end{pmatrix} \quad \hat{c}^{(k)} = c^{(k)} / \sum c_{\alpha\beta}^2$$

$$\theta_{jk} = \cos^{-1}(\hat{c}^{(j)} \cdot \hat{c}^{(k)})$$

θ_{jk} is a measure of the relative information content
in $F^{(j)}(Q)$ and $F^{(k)}(Q)$

(90° is good!)

NiZr neutron + Ni/Zr EXAFS

CuZr neutron + Cu/Zr EXAFS



NiZr neutron isotopes

CuZr neutron isotopes



$$I(Q, \omega) \longrightarrow S(Q, \omega) = \sum S_{AB}(Q, \omega) \longrightarrow G_{AB}(r, t) - - - r_i(t)$$

Corrections

Separation

Transform

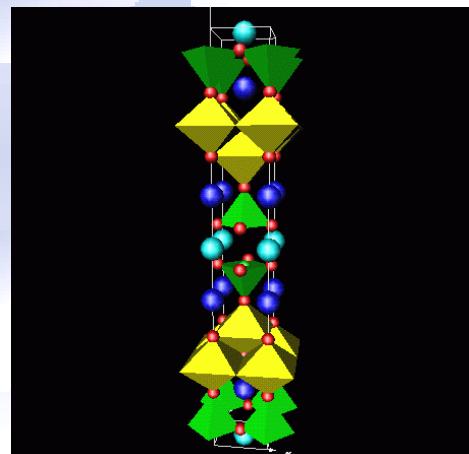
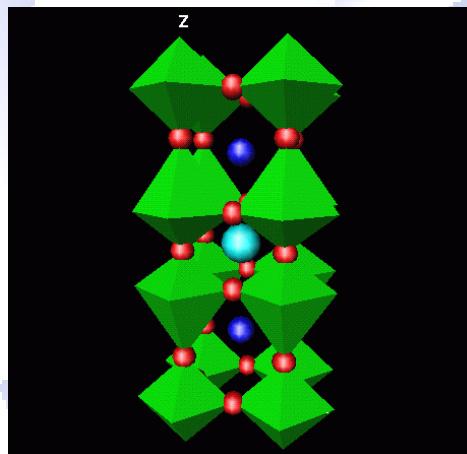
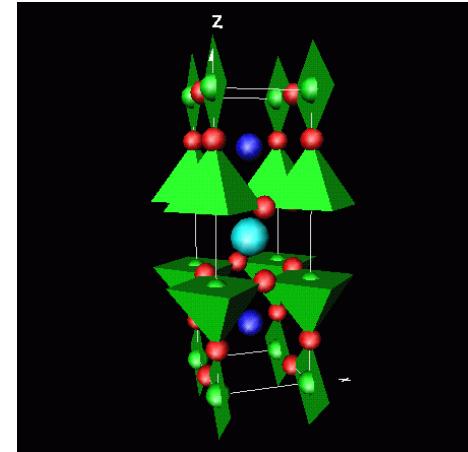
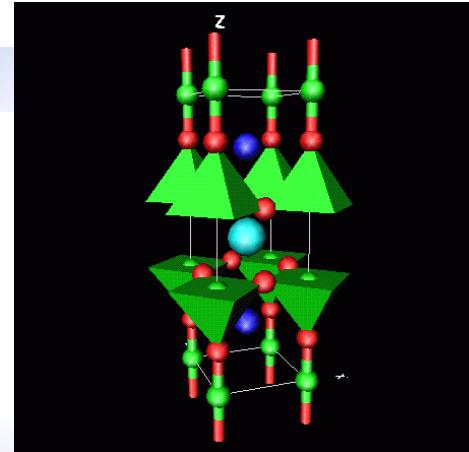
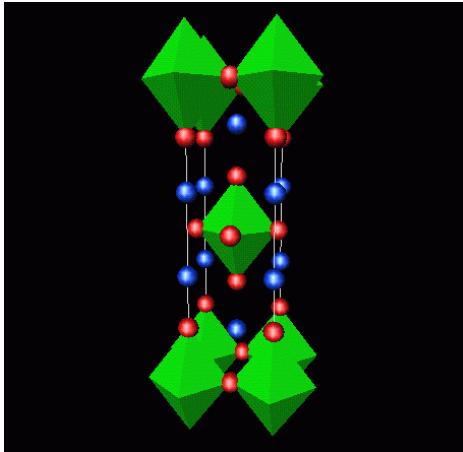
Errors

**Under defined
Poorly conditioned**

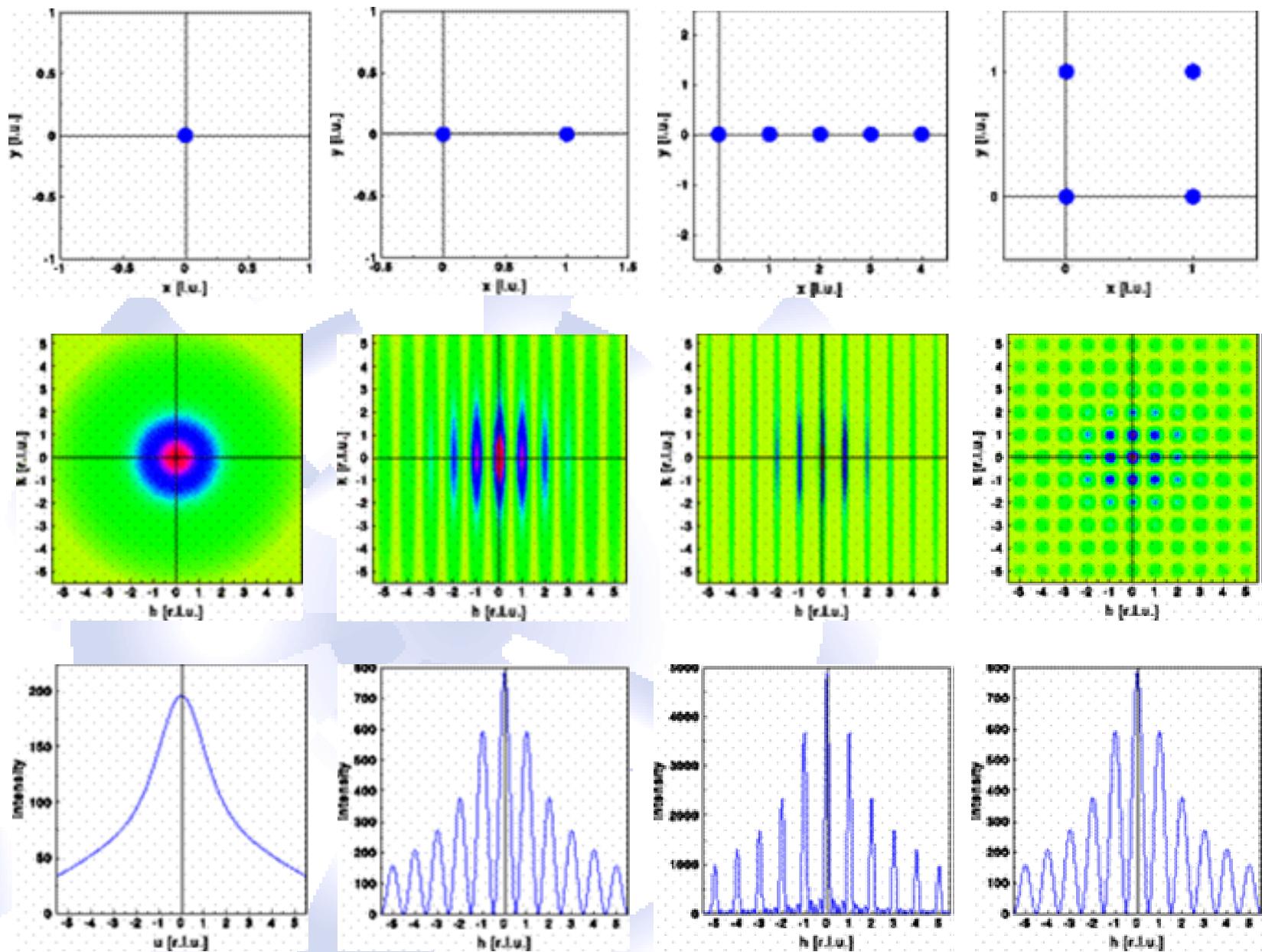
Truncation

Phase problem

so ...



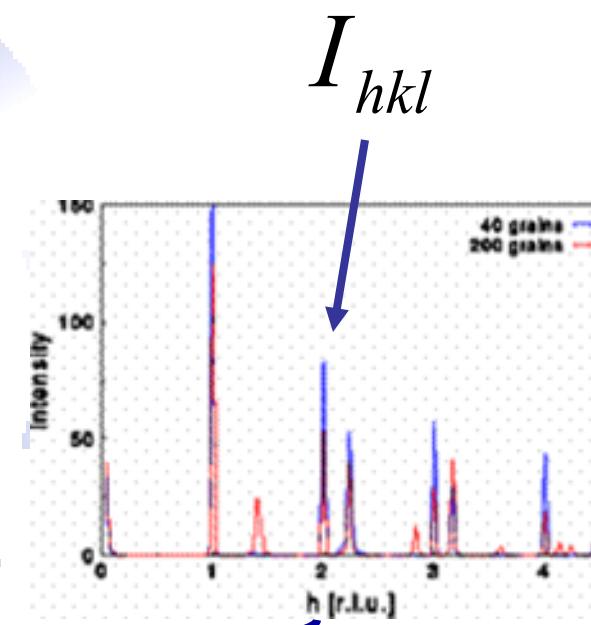
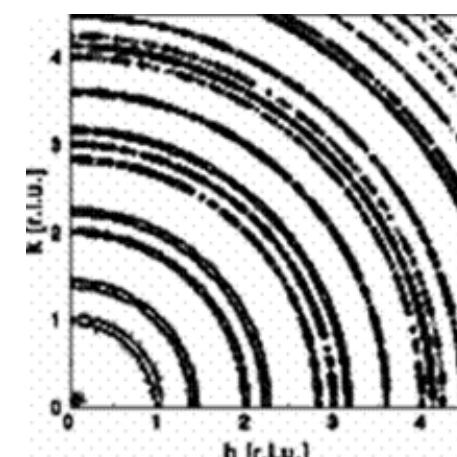
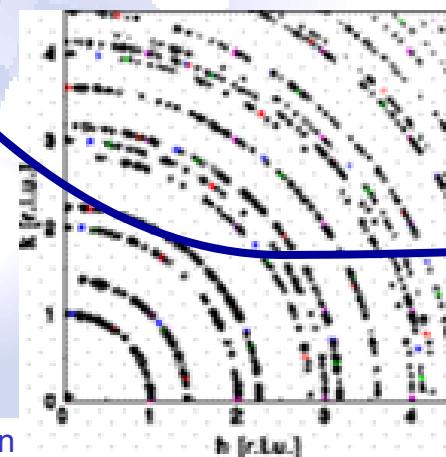
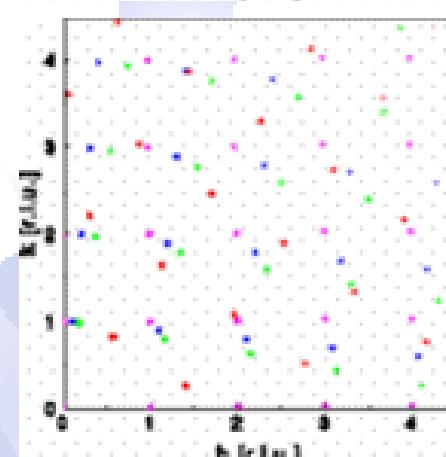
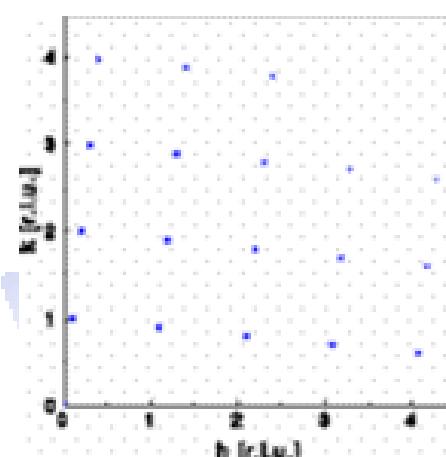
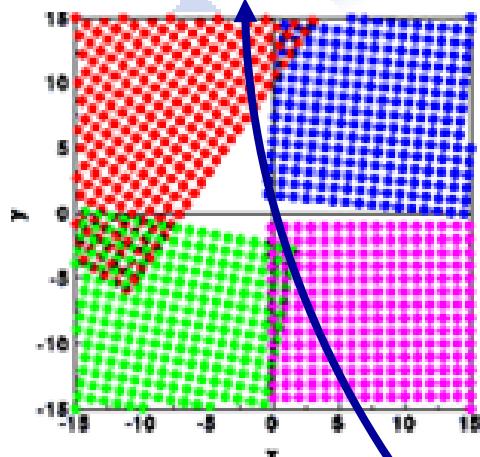
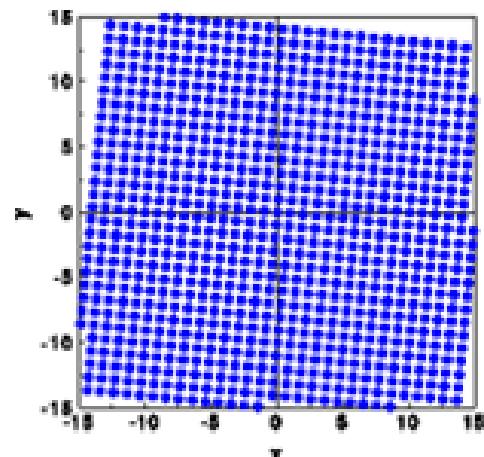
High temperature
superconductors



Figures courtesy of Thomas Proffen



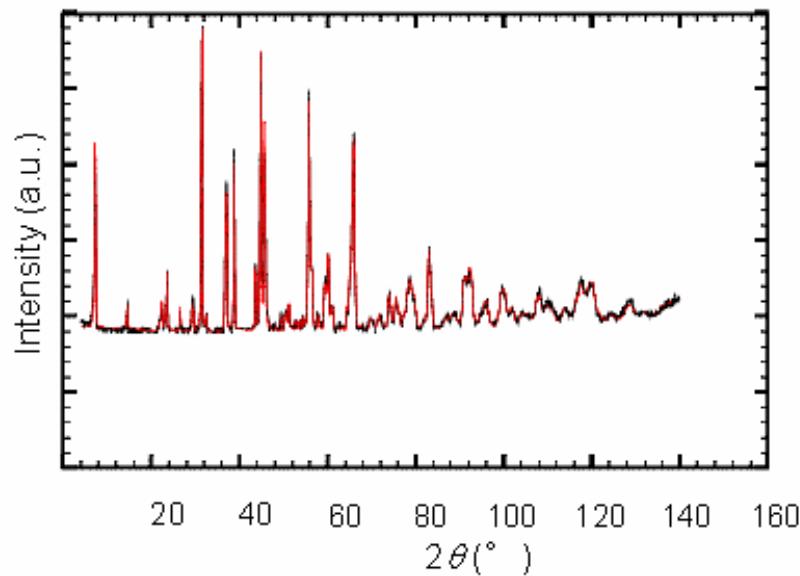
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Figures courtesy of Thomas Proffen

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$$I_{hkl} \propto |F_{hkl}|^2$$



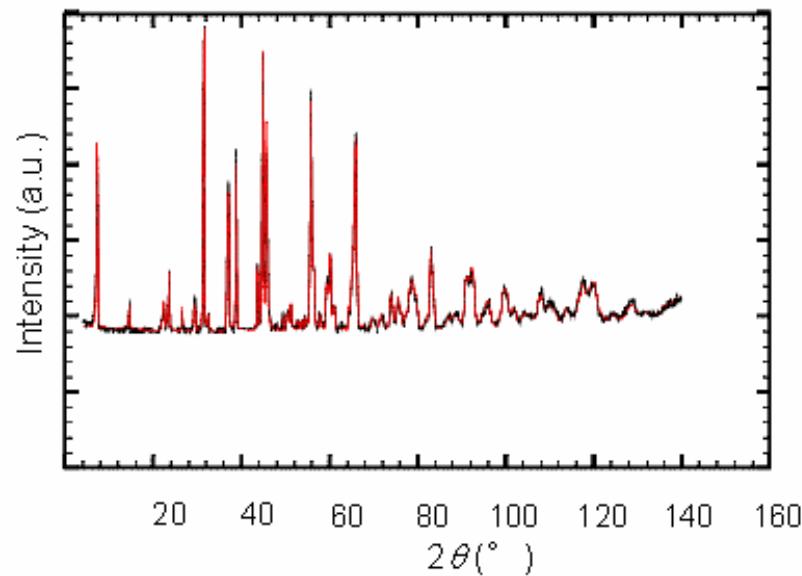
$$F_{hkl} = \sum_{j=1}^{N_j} b_j \exp(2\pi i(hx_j + ky_j + lz_j))$$

$$F_{hkl} = \sum_{j=1}^{N_j} b_j \exp(2\pi i \mathbf{r}_j \cdot \mathbf{s}_{hkl})$$

$$\mathbf{r}_j = x_j \mathbf{a} + y_j \mathbf{b} + z_j \mathbf{c}$$

$$\mathbf{s}_{hkl} = h \mathbf{a}^* + k \mathbf{b}^* + l \mathbf{c}^*$$

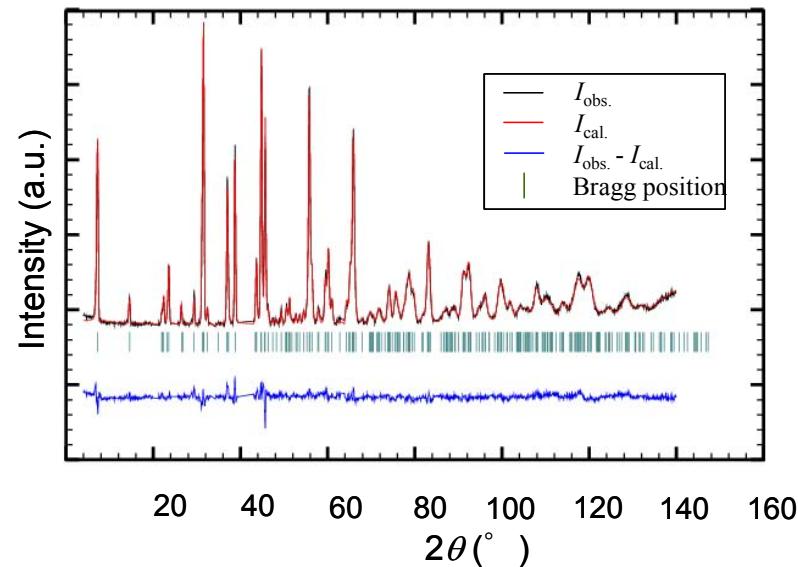
Rietveld refinement (parametric modeling)



$$y_i^c = s \sum_{hkl} L_{hkl} |F_{hkl}|^2 \phi(2\theta_i - 2\theta_{hkl}) P_{hkl} A + y_i^b$$

$$S_y = \sum_i w_i (y_i - y_i^c)^2$$

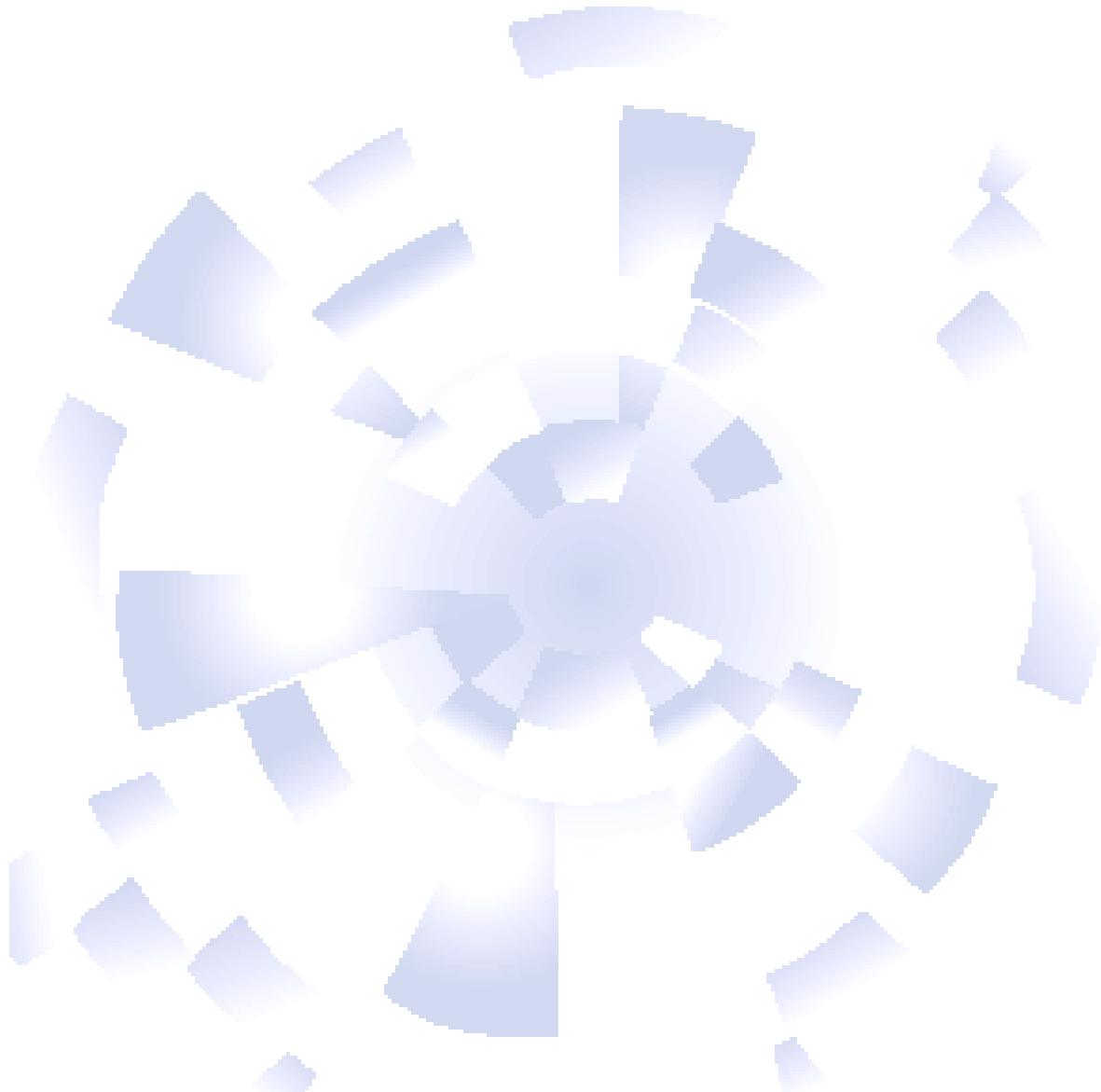
Name	x	y	z	B	occ.	Mult
Yb	0.50000	0.50000	0.50000	0.244	1.000	1
Sr	0.50000	0.50000	0.18448	0.971	0.452	2
Ba	0.50000	0.50000	0.18448	0.971	0.548	2
Cu1	0.00000	0.00000	0.00000	0.335	1.000	1
Cu2	0.00000	0.00000	0.35540	0.095	1.000	2
O1	0.00000	0.50000	0.00000	0.100	0.914	1
O2	0.50000	0.00000	0.37812	0.448	1.000	2
O3	0.00000	0.50000	0.38058	0.583	1.000	2
O4	0.00000	0.00000	0.16120	0.895	1.000	2
O5	0.50000	0.00000	0.00000	0.100	0.013	1
Cell parameters	:	3.78672	3.85536	11.58989		
		90.00000	90.00000	90.00000		
Overall scale factor	:	5.364356990	0.041989010			
Eta(p-v) or m(p-vii)	:	0.32626	0.01554			
Overall tem. factor	:	0.00000	0.00000			
Halfwidth parameters	:	1.64226	-0.93251	0.22901		
Asymmetry parameters	:	0.08420	0.02110			
GLOBAL PARAMETERS						
Zero-point:	-0.0016	0.0028				
Background Polynomial Parameters						
757.356	-10.6976	0.243093	-0.211296E-02	0.664176E-05		



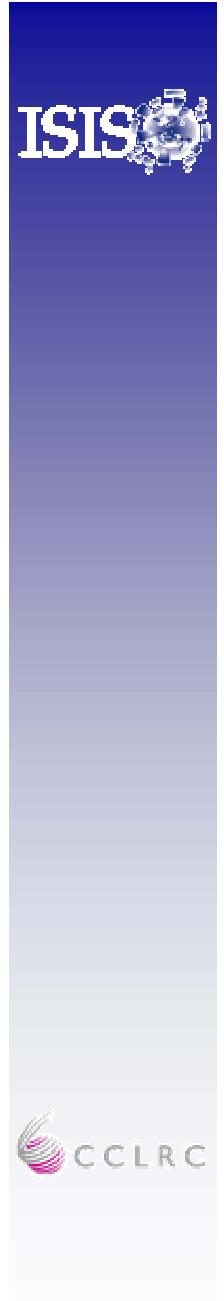
$$R_p = 4.17, R_{wp} = 5.47, \chi^2 = 3.68$$

Constrained Rietveld refinement

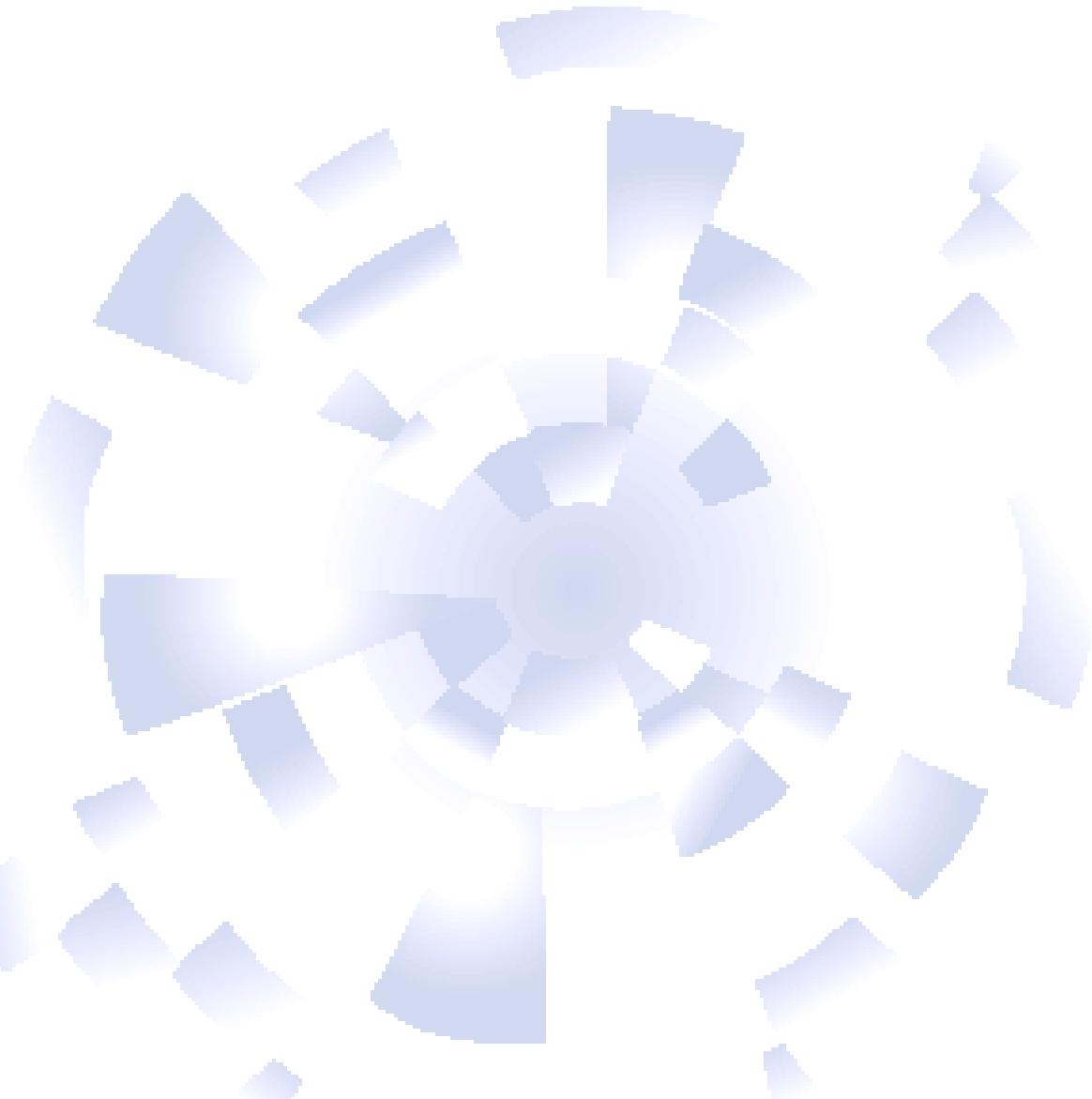
Simulated annealing (Monte Carlo)



$$\sum_i w_i (y_i - y_i^c)^2$$



Simulated annealing (Monte Carlo)

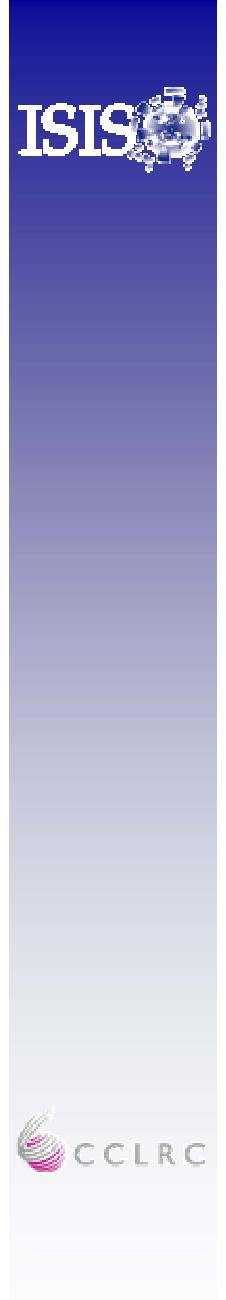
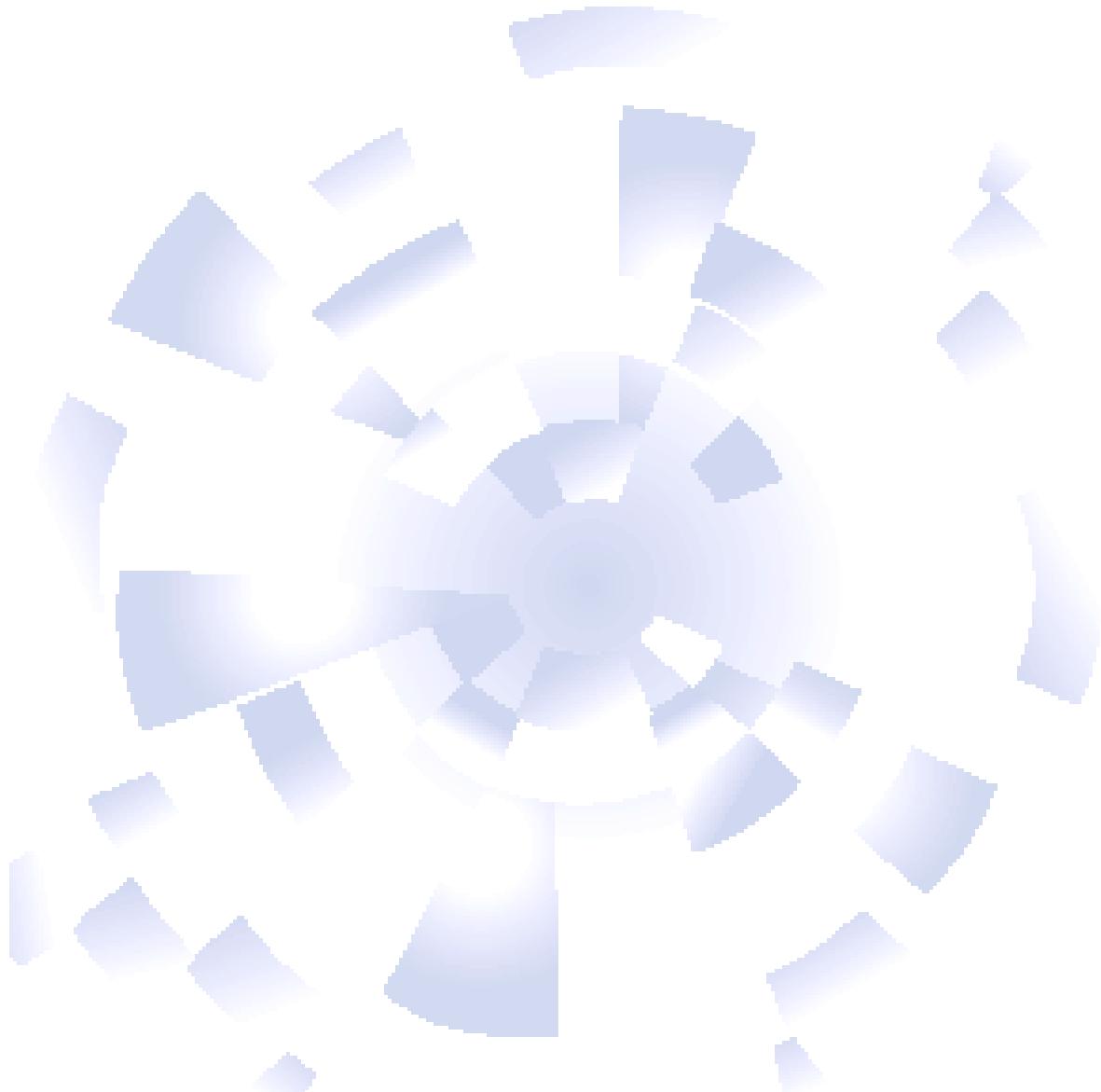


$$\exp\left(-\sum_i w_i (y_i - y_i^c)^2\right)$$



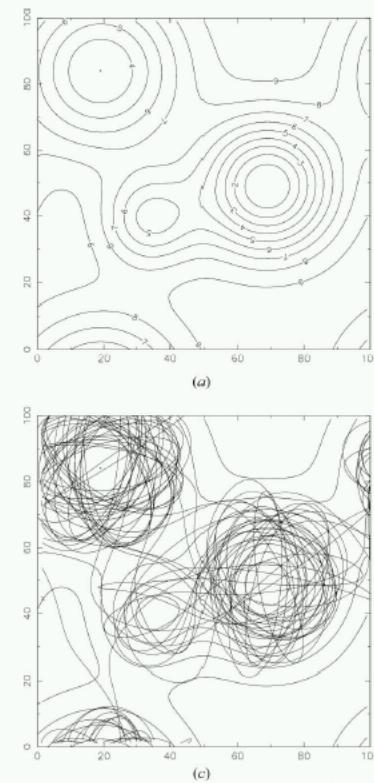
Simulated annealing (Monte Carlo)

$$\exp\left(-\sum_i (y_i - y_i^c)^2 / T\right)$$



Simulated annealing (Monte Carlo)

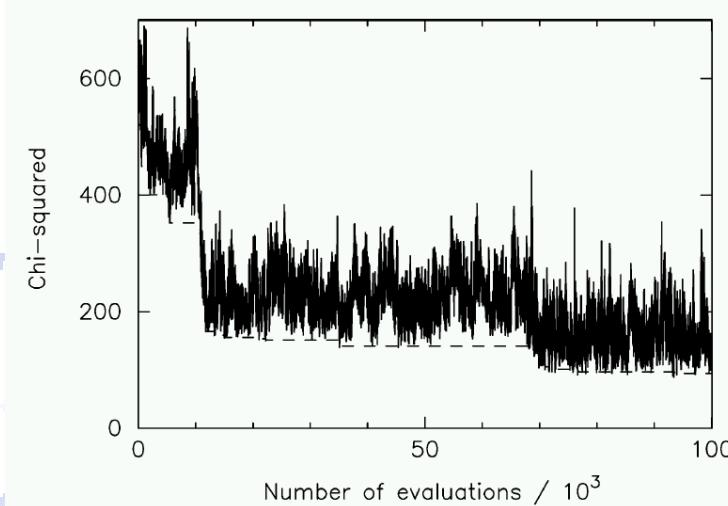
Hybrid Monte Carlo (molecular dynamics)



$$\exp\left(-\sum_i (y_i - y_i^c)^2 / T\right)$$

$$H(t) = \frac{1}{2} \sum_{i=1}^N m_i v_i^2(t) + \chi^2(\mathbf{r}(t))$$

$$\exp(-(E_m - E_0)/T)$$



Johnston, David, Markvardsen, Shankland. Acta Cryst A **58** 441 2002



$\alpha\text{-AgI}$

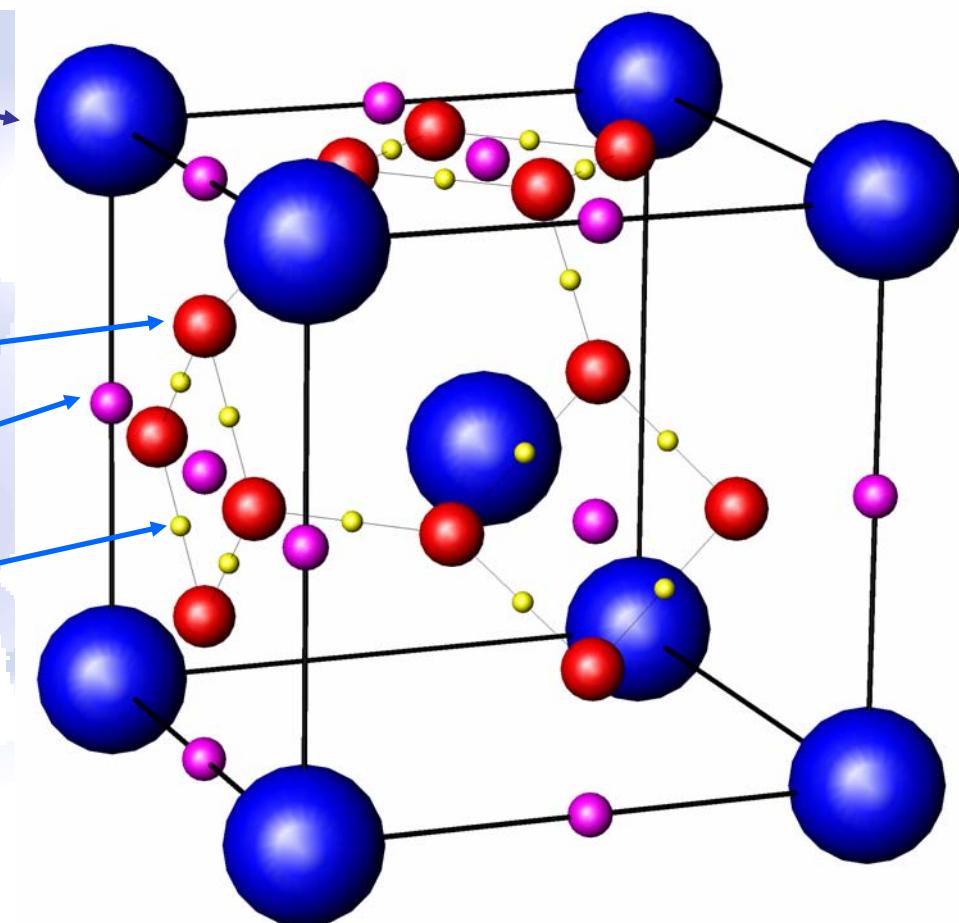
b.c.c. I⁻

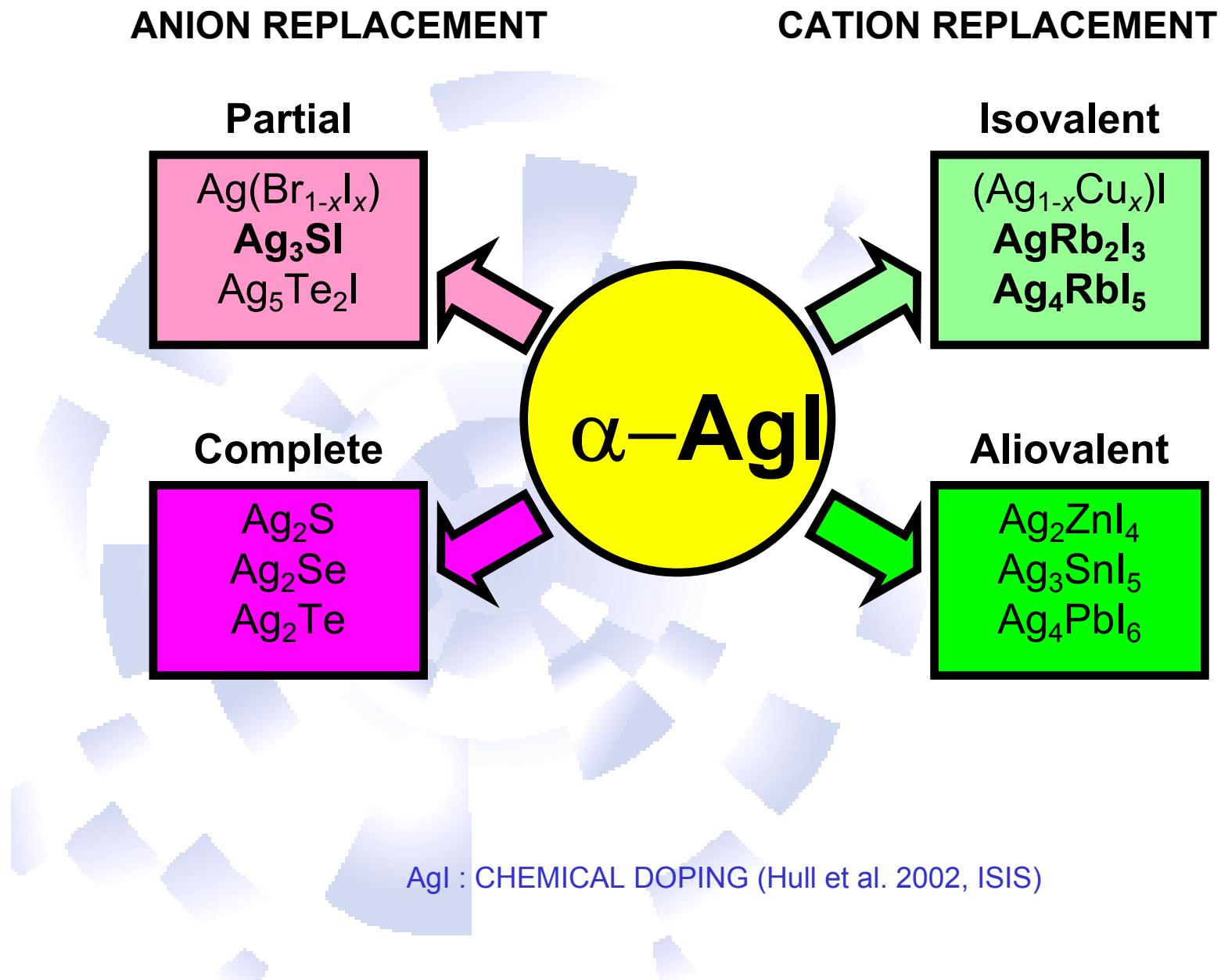
Ag⁺ sites.....

12 × tetrahedral

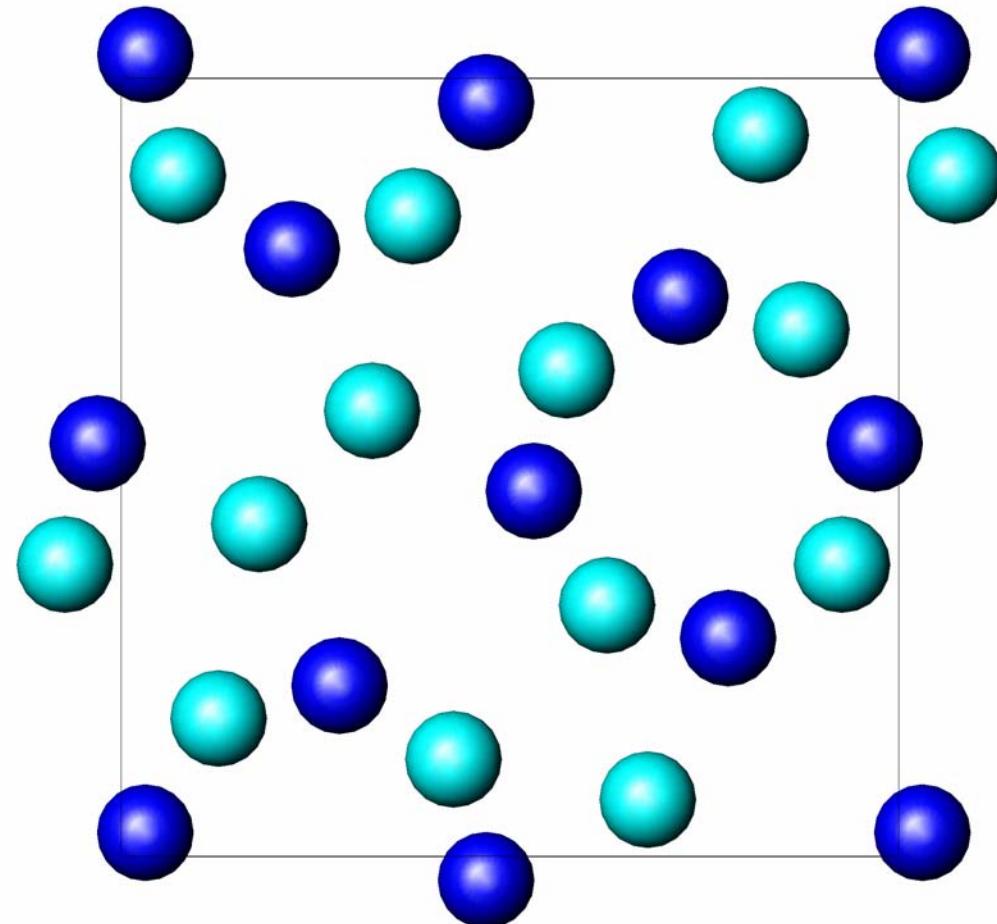
6 × octahedral

24 × trigonal





Ag₄RbI₅ : IODINE SUBSTRUCTURE

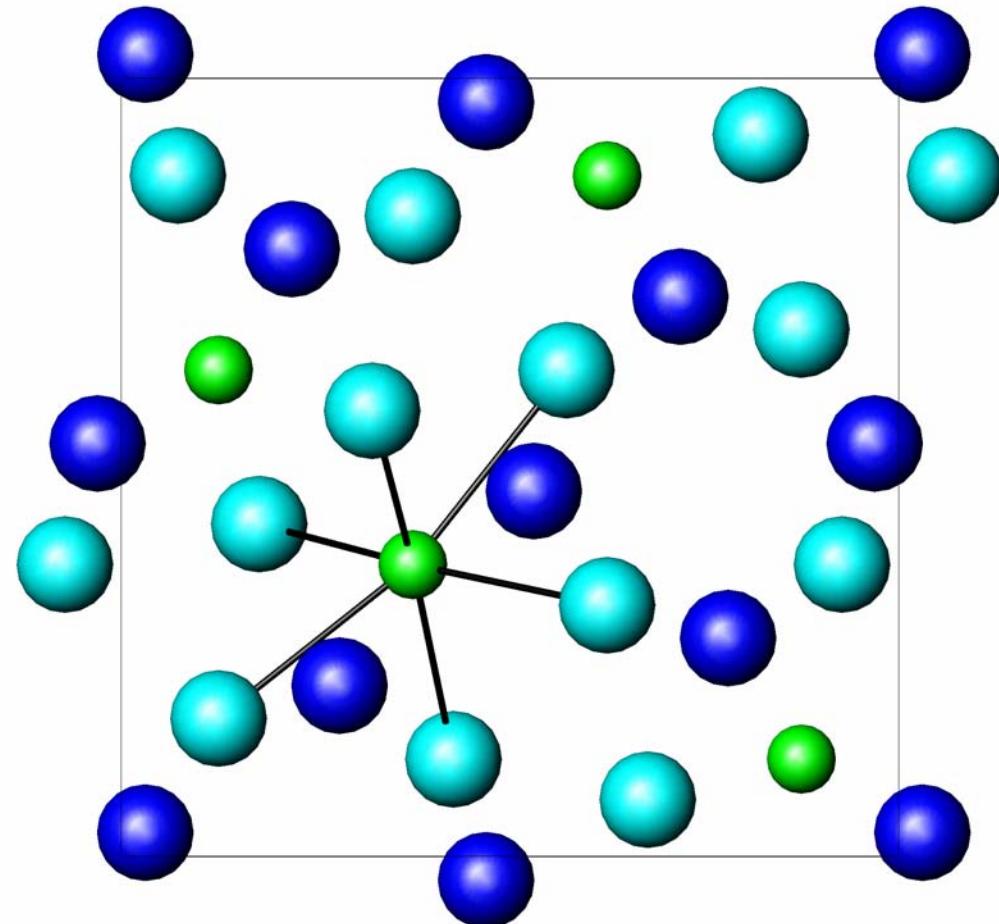


P 4₁32
a ~11.24 Å

I1 in 8(c)
 x, x, x
 $x \sim 0.031$

I2 in 12(d)
 $\frac{1}{8}, y, \frac{1}{4} + y$
 $y \sim 0.177$

AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

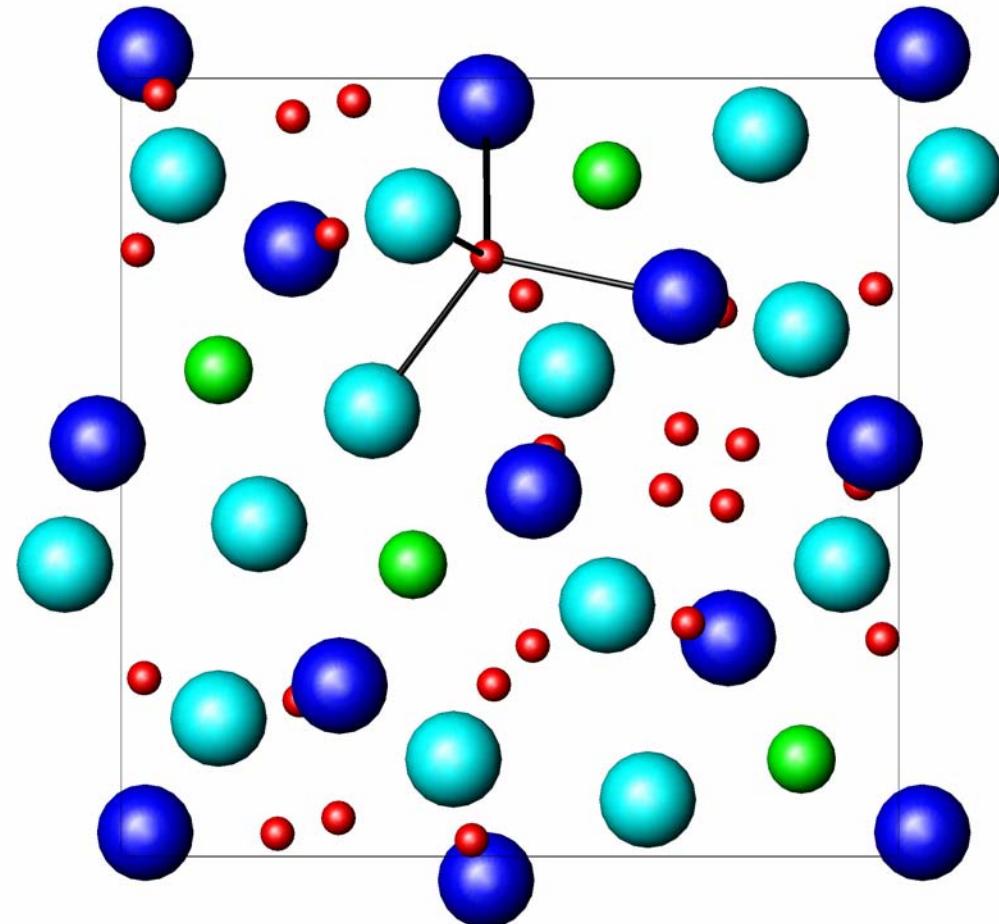
Ag_4RbI_5 : Rb POSITIONS

P 4₁32
a ~11.24 Å

Rb in 4(a)
 $\frac{3}{8}, \frac{3}{8}, \frac{3}{8}$

AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

Ag_4RbI_5 : Ag1 SITES



AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

$P\ 4_132$
 $a \sim 11.24\text{\AA}$

Ag1 in 24(e)

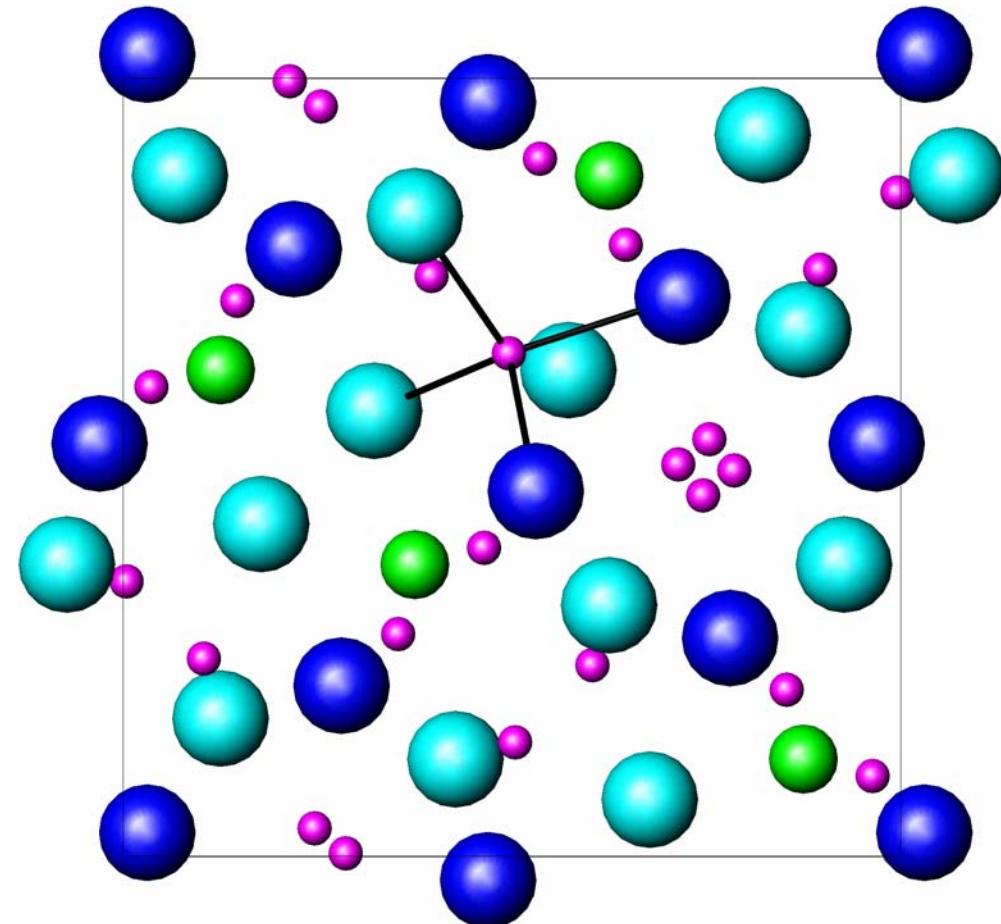
x,y,z

$x \sim 0.531$

$y \sim 0.272$

$z \sim 0.806$



Ag_4RbI_5 : Ag2 SITES

AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

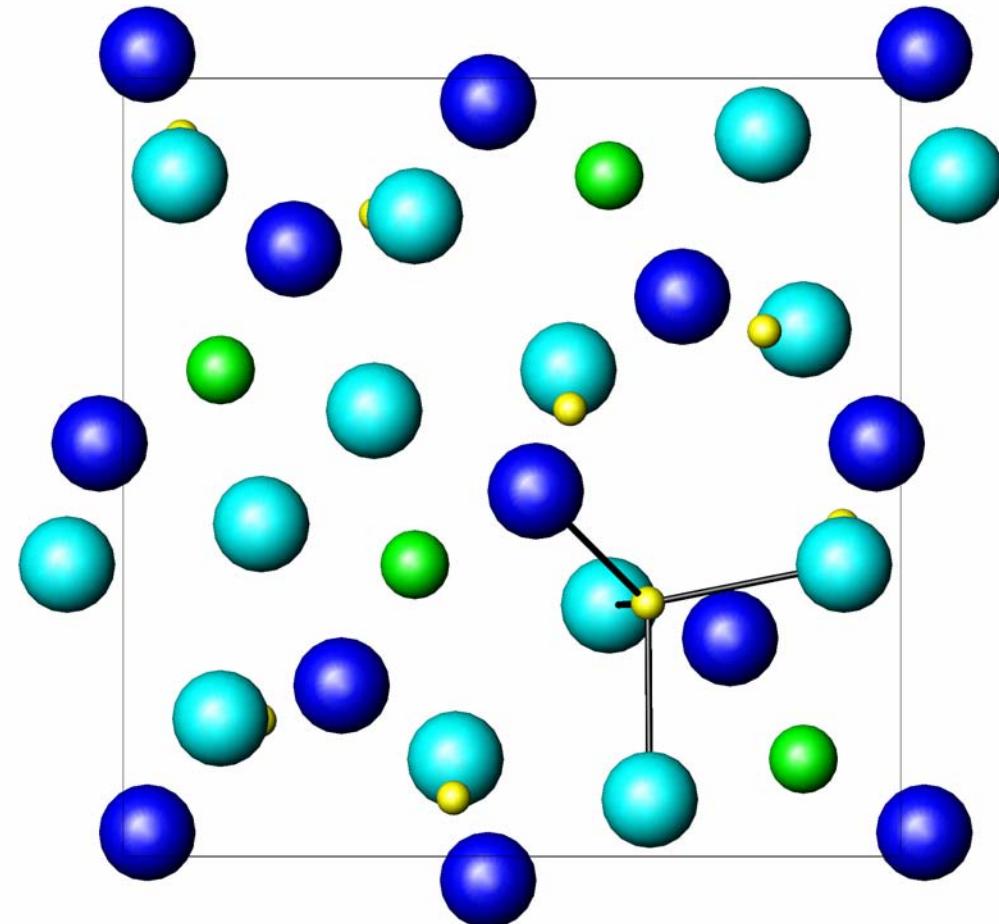
P 4₁32
 $a \sim 11.24\text{\AA}$

Ag2 in 24(e)

x, y, z
 $x \sim 0.993$
 $y \sim 0.855$
 $z \sim 0.206$

Ag_4RbI_5 : Ag3 SITES

ISIS



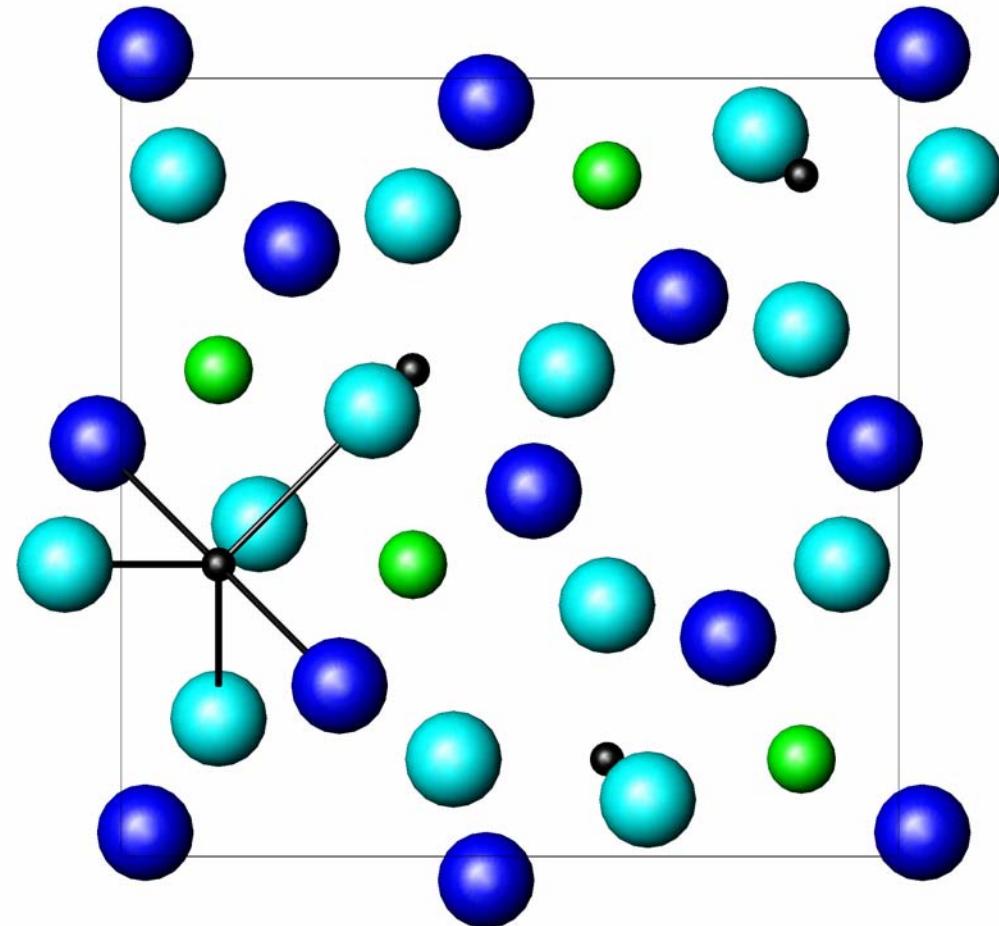
$\text{P} 4_132$
 $a \sim 11.24 \text{\AA}$

Ag3 in 8(c)
 x, x, x
 $x \sim 0.177$

AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

CCLRC

Ag_4RbI_5 : Ag4 SITES



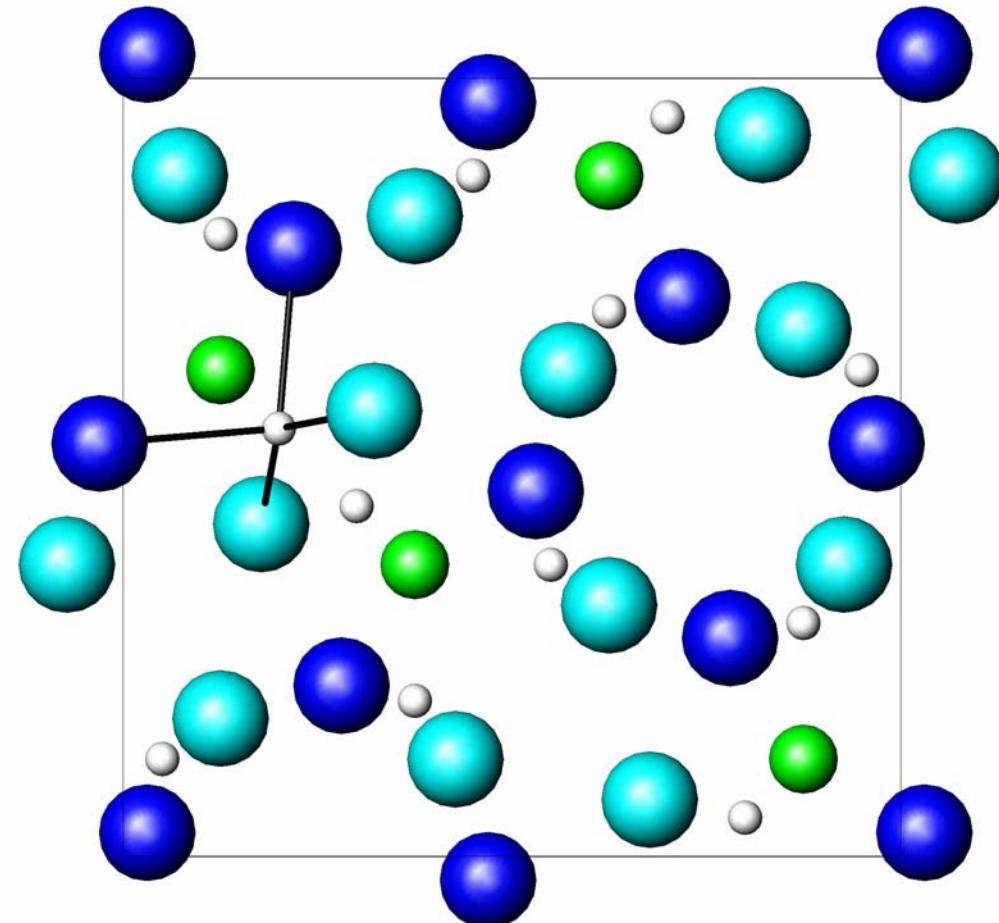
AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

$P 4_1 32$
 $a \sim 11.24 \text{\AA}$

Ag4 in 4(b)
 $7/8, 7/8, 7/8$



Ag_4RbI_5 : Ag5 SITES

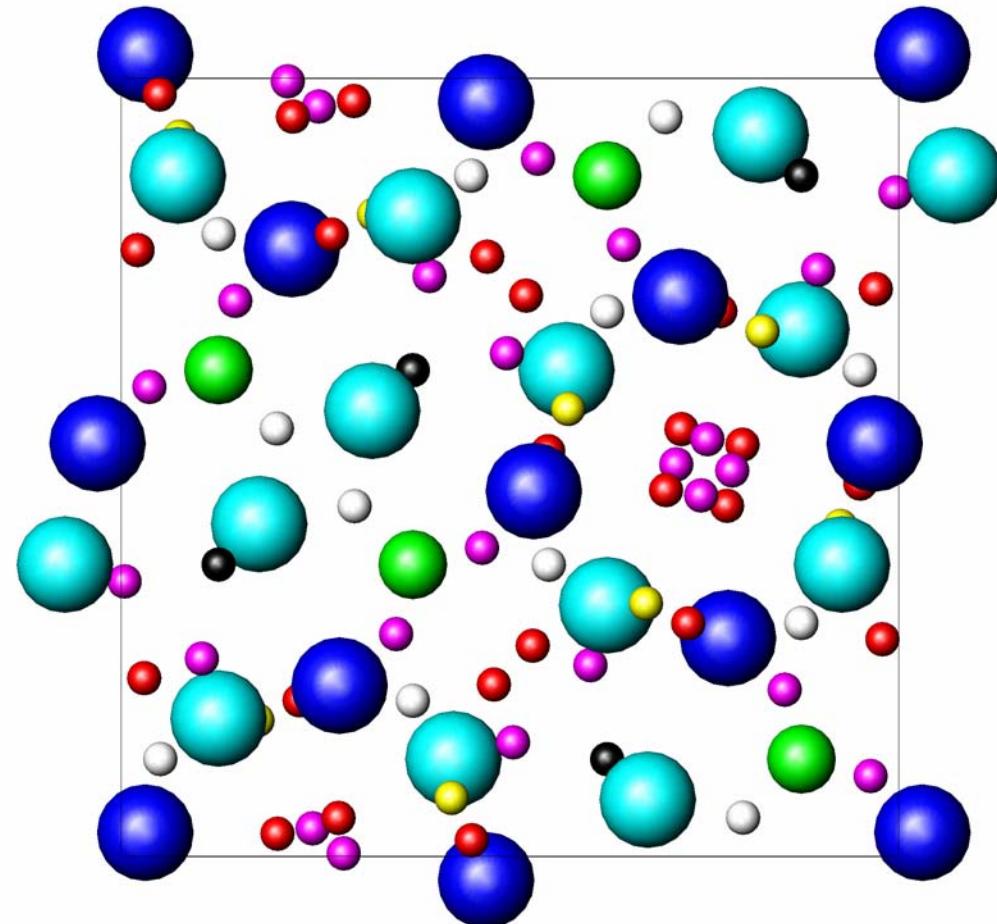


$P\ 4_132$
 $a \sim 11.24\text{\AA}$

Ag5 in $12(d)$
 $\frac{1}{8}, y, \frac{1}{4} + y$
 $y \sim 0.764$

AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

Ag₄RbI₅ : Ag SITES



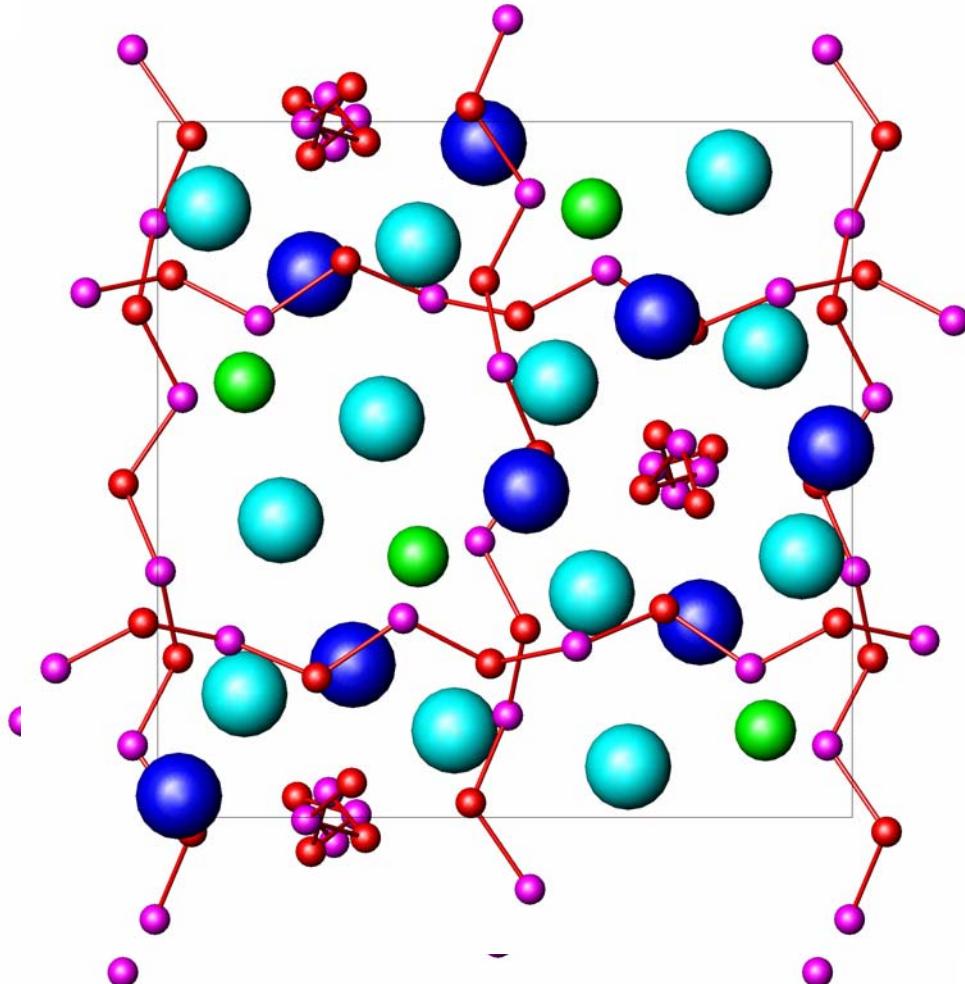
P 4₁32
a ~11.24 Å

20 × Ag⁺
per unit cell

How are they
distributed
over the 72
sites?

AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

Ag₄RbI₅ : CONDUCTION PATHWAYS



AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

Find a preferential occupancy of the **Ag1** and **Ag2** sites.

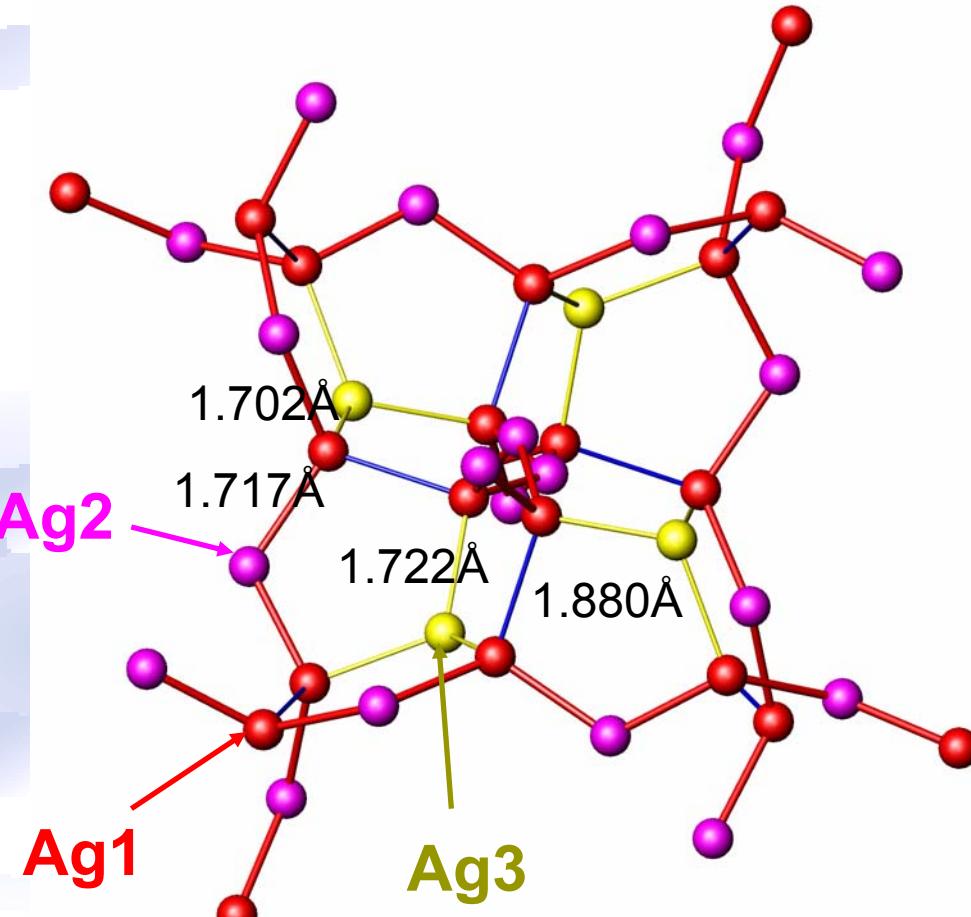
- Ag⁺ hop between pairs of these sites in <001> directions.
- conduction of Ag⁺ occurs along one-dimensional channels.

Ag_4RbI_5 : DIFFUSION BETWEEN CHANNELS

There are two plausible routes for Ag^+ to hop between channels.

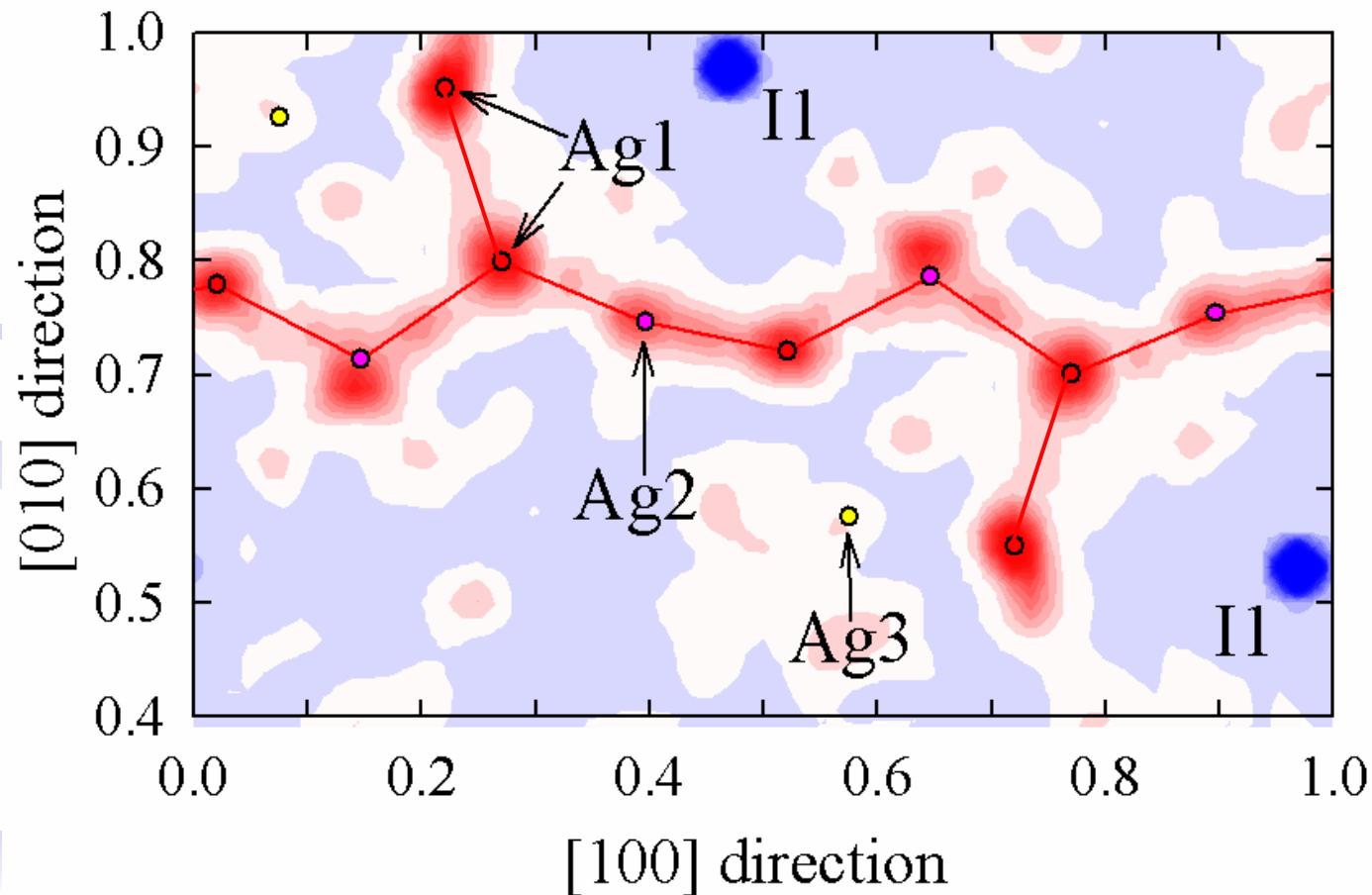
- Direct hop :
 $\text{Ag1} \rightarrow \text{Ag1}$

- Indirect hop :
 $\text{Ag1} \rightarrow \text{Ag3} \rightarrow \text{Ag1}$



AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)

Ag_4RbI_5 : MAXIMUM ENTROPY FOURIER MAP

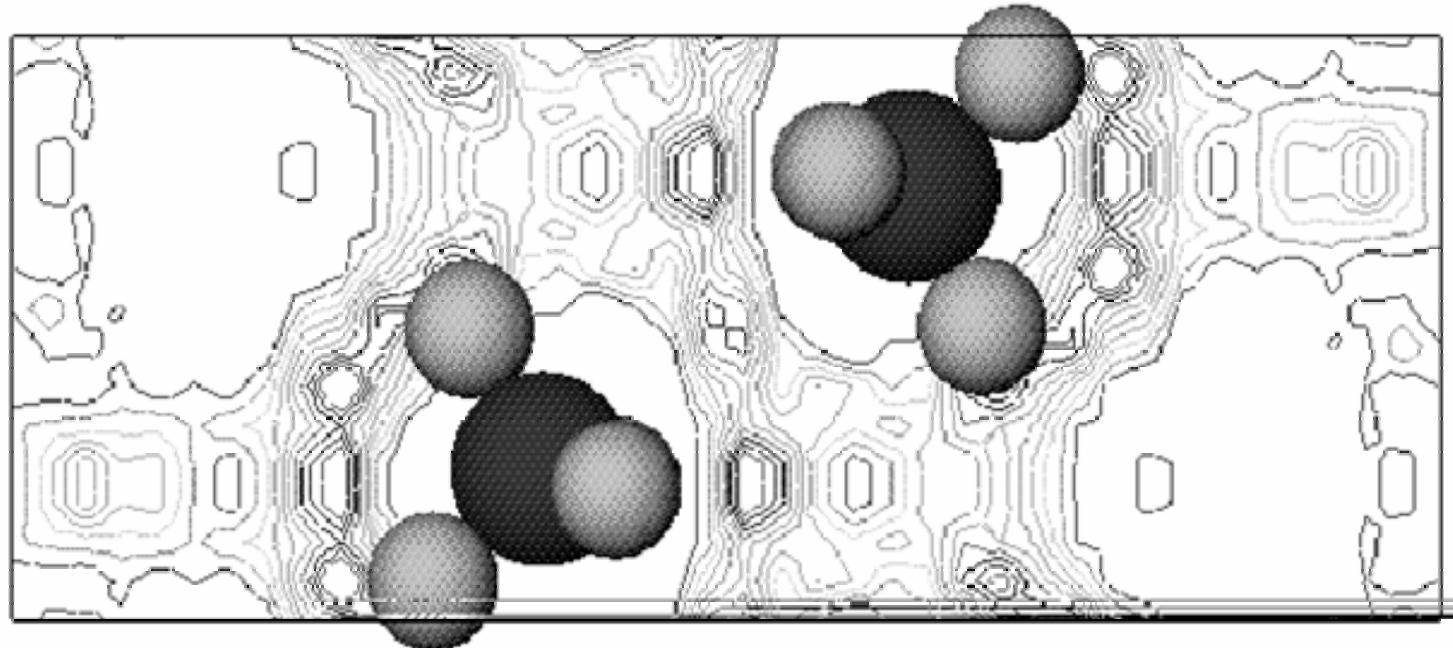


AgI : CHEMICAL DOPING (Hull et al. 2002, ISIS)



UCLA
12/11/2003

Atomic (structural) modeling



All inverse problems have solutions

Some inverse solutions have problems