

Introduction to Numerical Relativity

- Textbooks: • M. Alcubierre "Intro to 3+1 NR" 2008
- T. Baumgarte & S. Shapiro: "Numerical Relativity" 2010
- "NR: Starting from scratch" 2021
- M. Shibata: "Numerical Relativity" 2015
- [also online lecture notes, eg, Gourgoulhon gr-qc/0703035; LRRS]

• What is numerical relativity?

↳ solving EEs (or extensions) in 3+1 dims, typically using HPC to model the nonlinear regime of gravity

• Science cases

GW source modelling:



$v/c \ll 1, \Phi \ll 1$   
post-Newtonian

$v/c \lesssim 1$   
NR  
nonlinear

- ringdown  
- perturbation theory  
- NR

Ingredients

1) theoretical model

TODAY EEs in vacuum in 4D asympt. flat.

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 0$$

2) Spacetime decomposition

→ space, time - dep. explicit ←

→ kinematic evol.

3) 3+1 decomposition of field eqs.

→ dynamics (hyperbolic PDEs)

→ constraint (elliptic PDEs)

←  
 $\mathbb{F}(4,1)$   
 x tensor  
 tutorial

4) initial data

$$(g_{\mu\nu}, \partial_t g_{\mu\nu})|_{t=0}$$

→ phys. system under consideration  
 (e.g. single or binary BHs)

→ solve constraints  
 (4 coupled, elliptic PDEs)

5) Evol. eqs → well-posed initial value problems

6) Gauge choices, treatment of BH

7) Observables

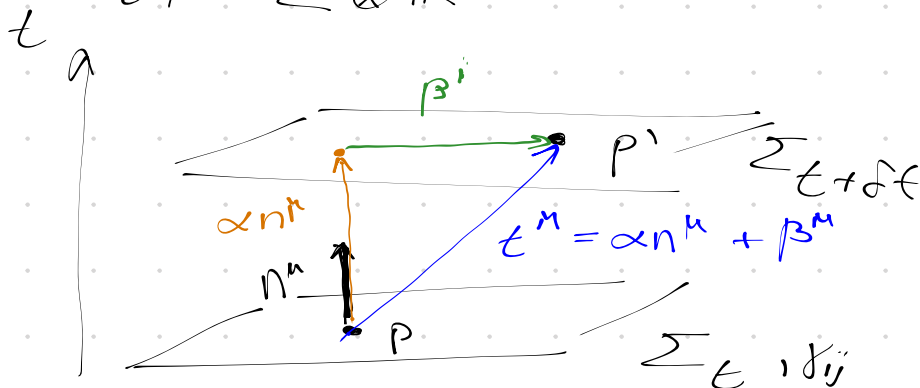
→ wave extraction (Newman-Penrose form.)  
 metric perturbations

→ apparent horizon

## 2) 3+1 of spacetime

↳ foliate 4D manifold  $(\mathcal{M}, g)$  into  
 spacelike hypersurfaces  $(\Sigma, \gamma)$  parametrized  
 by  $t \in \mathbb{R}$

$$\mathcal{M} = \Sigma \otimes \mathbb{R}$$



$\mu, \nu, \dots = 0, \dots, 3$   
 $i, j, \dots = 1, 2, 3$   
 $x^\mu$  - coords

↳  $n^\mu$  - normal, timelike vector  $n^\mu n_\mu = -1$

↳ on  $\Sigma$ : induced, spatial metric  $\gamma_{ij}$   
 measure proper distances  $dl^2 = \gamma_{ij} dx^i dx^j$

Lapse  $\alpha$ : proper time between hypersurfaces  
measured by observer moving along  $n^\mu$

L shift  $\beta^i$ : relative velocity between normal  
observer and lines of const. spat. coords.

by construction  $\beta^\mu n_\mu = 0$

L metric  $g_{\mu\nu}$  in terms of  $(\gamma_{ij}, \alpha, \beta^i)$

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

$$= -(\alpha^2 - \beta^i \beta_i) dt^2$$

$$+ 2 \gamma_{ij} \beta^i dt dx^j + \gamma_{ij} dx^i dx^j$$

note:  $\gamma_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$

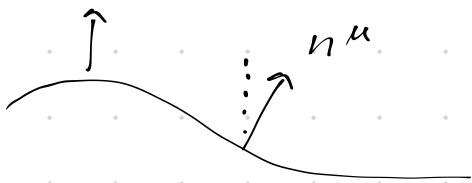
•  $\gamma$  defines a projection op.

$$\gamma^M{}_\nu = \delta^M{}_\nu + n^M n_\nu$$

• a vector  $V^M = \underbrace{d} n^M + \underbrace{V}^M$

$$d = -n_\nu \underbrace{V}^\nu \quad V^i = \gamma^i{}_\nu \underbrace{V}^\nu$$

L extrinsic curvature  $K_{\mu\nu}$



$$K_{\mu\nu} = - \gamma^k{}_\mu \nabla_k n_\nu$$

sign convention

one can show:  $K_{\mu\nu} = -\frac{1}{2\alpha} \mathcal{L}_n \gamma_{\mu\nu}$

$$\Rightarrow \boxed{(\partial_t - \mathcal{L}_\beta) \gamma_{ij} = -2\alpha K_{ij}} \quad \text{kinematic eq.}$$

### 3) 3+1 split of field equations

↳ dynamic

(i) Gauss-Codazzi relations

$$\rightarrow {}^{(4)}R_{\mu\nu\sigma\rho} \leftrightarrow {}^{(3)}R_{\mu\nu\sigma\rho}$$

$$\nabla_\mu$$

$$D_\mu = (\mathcal{L}_\beta D)_\mu$$

On Monday 20 Sept.  
 xtensor for decomposition  
 • download notebook  
 • install xtensor  
 (xact.es)

(ii) projection of energy-momentum tensor

$$T_{\mu\nu} : \quad -\rho = T_{\mu\nu} n^\mu n^\nu \quad (\text{energy density})$$

$$-j_\mu = -T_{\rho\nu} j_\mu^\rho n^\nu \quad (\text{energy-mom. flux})$$

$$-S_{\mu\nu} = T_{\rho\sigma} j_\mu^\rho j_\nu^\sigma \quad (\text{spatial stress tensor})$$

(iii) decomposition of EEs:

• Hamiltonian constraint  $\mathcal{H} = 2 G_{\mu\nu} n^\mu n^\nu$

$$= 0$$

$$= {}^{(3)}R - K_{ij} K^{ij} + K^2$$

• Momentum constraint

$$\mathcal{M}_i = -j_i^\mu n^\nu G_{\mu\nu} = D_j^j K_{ij} - D_i K = 0$$

Evolution Eqs:

$$j_\mu^\rho j_\nu^\sigma G_{\rho\sigma} = 0 \Rightarrow \mathcal{L}_\beta K_{ij}$$

dynamical  $\partial_t K_{ij} = \mathcal{L}_\beta K_{ij} - D_i D_j \alpha$

$$+ \alpha [{}^{(3)}R_{ij} + K K_{ij} - 2 K_{ik} K^k_j]$$

kinematic:  $\partial_t \delta_{ij} = \mathcal{L}_\beta \delta_{ij} - 2\alpha K_{ij}$

Note: structure:

$$\partial_t u \approx F[\partial_i \partial_j u, \partial_i u, u]$$

vector of  
evol vars

in practise: adopt Method of Lines

→ rhs  $(F(\partial_i \partial_j u, \partial_i u, u))$   
↳ choose a type discretization  
for spatial [finite differences,  
spectral method]

→ lhs:  $\partial_t u$

↳ choose time integrator (e.g. Runge-Kutta  
4<sup>th</sup> order)

↳ note: adopt a free evolution scheme, i.e.,

- solve constraints only for initial data
- monitor constraints during evolution  
(check that they remain satisfied)
- use BI to show that constraints  
remain satisfied if • satisfied @  $t=0$

• system evolves  
according to evol.  
system.

Well-posed initial value formulation

Def: A system of PDEs

$$\begin{cases} \partial_t f = A^p \partial_p f + Bf \\ f(t=0) = g \end{cases}$$

is said to be well-posed IVP  
if there exists a unique solution  
that depends continuously on smooth ID

in particular, a system is well-posed, if  
 $\forall k = \text{const}, a = \text{const}$ . s.t. for all  $t$   
 we have

$$\|f(t, \cdot)\| \leq k e^{a \cdot t} \|f(t=0, \cdot)\|$$

1st person: compare to wave-eqn

V1: 4D  
 $E\bar{E}S$

$$R_{\mu\nu} = 0$$

$$= \gamma^{\kappa\lambda} \frac{\partial^2}{\partial x^\kappa \partial x^\lambda} g_{\mu\nu}$$

$$+ \gamma^{\kappa\lambda} \frac{\partial^2}{\partial x^\mu \partial x^\nu} g_{\kappa\lambda}$$

$$+ \gamma^{\kappa\lambda} \frac{\partial^2}{\partial x^\mu \partial x^\nu} g_{\kappa\lambda}$$

$$+ \text{l.o.t.}$$

scalar  
 wave eqn  
 $\square \phi = \gamma^{\mu\nu} \frac{\partial^2}{\partial x^\mu \partial x^\nu} \phi = 0$

harmonic gauge  
 $\square x^\mu = 0$

V2: 3+1  $E\bar{E}S$

$$\partial_t \gamma_{ij} \approx -K_{ij}$$

$$\partial_t K_{ij} \approx -\partial_i \partial_j \alpha$$

$$+ \alpha R_{ij} + \text{l.o.t.}$$

$$\approx -\partial_i \partial_j \alpha$$

$$+ \alpha \left[ \gamma^{lm} \frac{\partial^2}{\partial x^i \partial x^j} \gamma_{lm} + \gamma^{lm} \frac{\partial^2}{\partial x^i \partial x^j} \gamma_{lm} + \dots \right] + \text{l.o.t.}$$

Scalar wave eq  
 $\square \phi = 0$   
 introduce  $\pi = -\partial_t \phi$   
 $\partial_t \phi = -\pi$   
 $\partial_t \pi = -\Delta \phi$   
 $= -\gamma^{ij} \partial_i \partial_j \phi$

cause for illposedness

"cure": Baumgarte-Shapiro (1999) -  
 Shibata-Nakamura (1995) (BSSN)  
 for  $n=3$

vars:  $W = f^{-1/6}$   
 $\tilde{\gamma}_{ij} = W^2 \gamma_{ij}$  (s.t.  $\tilde{f} = 1$ )  
 $K = \gamma^{ij} \dot{\gamma}_{ij}$ ,  $\tilde{A}_{ij} = W^2 A_{ij} = W^2 (K_{ij} - \frac{1}{3} \gamma_{ij} K)$   
 $\tilde{\Gamma}_{rel}^i = \gamma^{kj} \dot{\gamma}_{kj}$ ,  $\tilde{\Gamma}_{rel}^i = -\dot{\gamma}_{ij} \tilde{\gamma}^{ij}$

# Gauge choices, i.e., $(\alpha, \beta^i)$

note: simple is not always best  
 (e.g.  $\alpha = 1, \beta^i = 0 \Rightarrow$  reach BH sing. in finite time)

Wishlist: • avoid reaching singularity

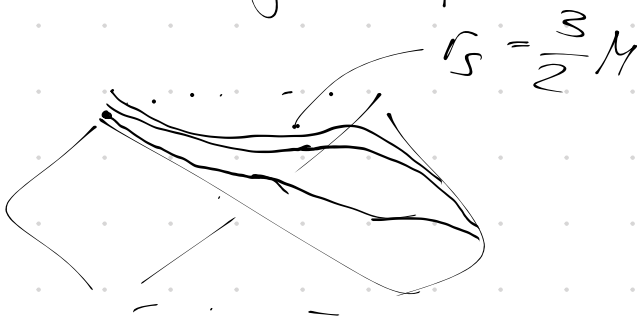
- evol eqs + gauge cond. form a well-posed PDE system
- easy to implement spec., avoid elliptic eqs (comput. expensive) use evol. eqs.

common choices: puncture gauge

$\Gamma$ -driver slicing

$$\left| \frac{\partial \alpha}{\partial t} = -2\alpha(K - K_0) \right|$$

- $\alpha \rightarrow 0$  near singularity
- "singularity avoidance"
- spac. hypersurfaces cannot be arbitrarily close to singularity



$\Gamma$ -driver shift

$$\frac{\partial \beta^i}{\partial t} = \beta^j \Gamma^i_{jk} - \gamma^i_j \beta^j$$

$\approx \partial_j^2 f^{ij}$

advection

## Initial data

Goal: prescribe  $(g_{ij}, K_{ij})|_{t=0}$

↳ Count: specify  $12^{\text{ind.}}$  components

↳ solve constraints: specify 4 ind. comp.

↳ "free" choice for  $\mathcal{S}$  <sup>remaining</sup>:

- motivated by physical system

- simplify

• conformal decomposition:

$$\begin{aligned} g_{ij} &= \chi^4 \overset{\uparrow}{\hat{g}}_{ij} \\ K_{ij} &= \chi^{-2} \overset{\uparrow}{A}_{ij} + \frac{1}{3} \chi^4 \overset{\uparrow}{K} \end{aligned}$$

• Example: single BH:

• time-symmetry:  $K_{ij} = 0$   
 $\Rightarrow \mathcal{M}_i = 0 \quad \checkmark$

• conformal flatness, i.e.,  $\overset{\uparrow}{\hat{g}}_{ij} = \gamma_{ij}$   
 $\overset{\uparrow}{\Delta} = \overset{\uparrow}{\Delta}_{\text{fl}}$

• asymp. flatness, i.e.,  $\left( \lim_{r \rightarrow \infty} \chi = 1 \right)$   
 $\Rightarrow \left[ \mathcal{H} = \overset{\uparrow}{\Delta}_{\text{fl}} \chi \right]$

• simplest non-trivial solution:

$$\chi = 1 + \frac{R}{2r}$$

identify  $R = \frac{M}{2} \rightarrow \chi = 1 + \frac{M}{2r}$

$$\rightarrow ds^2 = -\alpha^2 dt^2 + \left(1 + \frac{M}{2r}\right)^4 \gamma_{ij} dx^i dx^j$$

Schwarzschild in isotropic coord. rat's

Note: superpose for headon (zero momenta)

• seed solution for BBH, together with Bowen-York extrinsic cur