

Part I. Implication from QEC.

(Q) "TQO" in 1D?

$$H = \sum h_j, \quad h_j^2 = h_j = h_j^\dagger \\ [h_j, h_{j'}] = 0$$

$$\Pi_{GS} = \prod_j h_j$$

Tool: Q. Info Π_{GS} TFAE, Given $A \subseteq \Lambda$

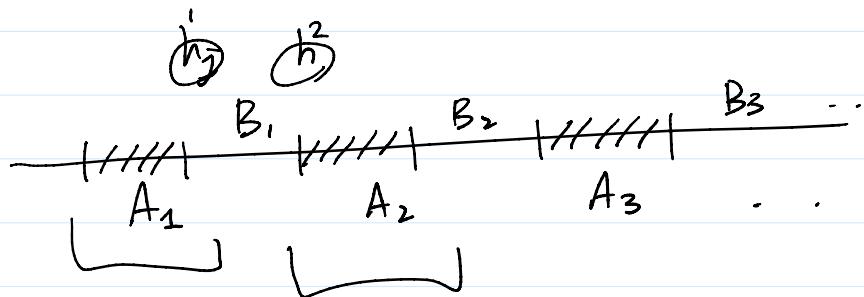
$$1. (TQO) \quad O_A \quad \Pi_{GS} O_A \Pi_{GS} = c(O_A) \Pi_{GS} \quad \checkmark \\ \underbrace{[O_A, \Pi_{GS}]}_0 = 0$$

$$2. (\text{Decoupling}) \quad I_p(A : R) = 0$$

$$3. (\text{Cleaning}) \quad \text{Logical } U \quad (\text{i.e. } [U, \Pi_{GS}] = 0) \\ \exists V^* \text{ s.t. } (V^* - U) \Pi_{GS} = 0$$



arXiv: [\[1610.06169\]](https://arxiv.org/abs/1610.06169) Limits on the storage of quantum information in a volume of space (arxiv.org)



$$\Pi_{GS} \xrightarrow{\quad} O_{A_1} \quad \Pi_{GS} O_{A_2} \xrightarrow{\quad} \Pi_{GS} = c(O_{A_1}) c(O_{A_2}) \Pi_{GS}$$

\Rightarrow If A_1, A_2 are "correctable", $\xrightarrow{\text{so?}} A_1 \cup A_2$.

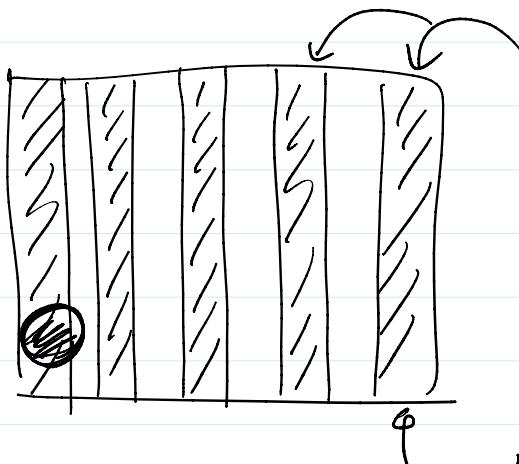
$\bigcup_i A_i$ is corr.

$\bigcup_i B_i$ is corr.

Thm in 1D \nvdash TQO.

This argument adopts an idea from [0810.1983] A no-go theorem for a two-dimensional self-correcting quantum memory based on stabilizer codes (arxiv.org)

2D



$\bigcup A_i$
 $\bigcup B_i$

$$\text{TlGS} \geq \sum_{i=1}^n \langle \tilde{\gamma} \rangle < i$$

n ()

must support
some "logical" op.

How large am TlGS be?



$$\dim \text{TlGS} = n$$

$$- \frac{S_{NR}}{S_N} + S_N + S_R = 0$$

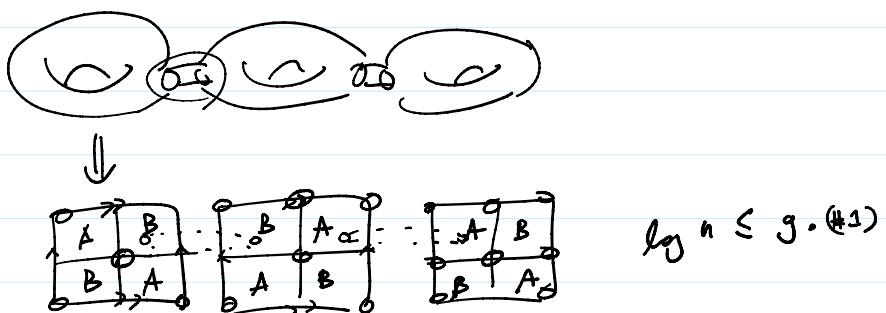
$$- \frac{S_{SR}}{S_N} + S_s + S_R = 0$$

+

2 S - ->

$$\text{pure. } |\Psi\rangle = \sum_i |\Psi_i\rangle^{(i)} / \sqrt{n} \quad \left| \begin{array}{l} \dim \mathcal{H}_{\text{sys}} = n \\ \downarrow \\ S_R = \log n \end{array} \right.$$

$$2S_R = 0 \Rightarrow n=1$$

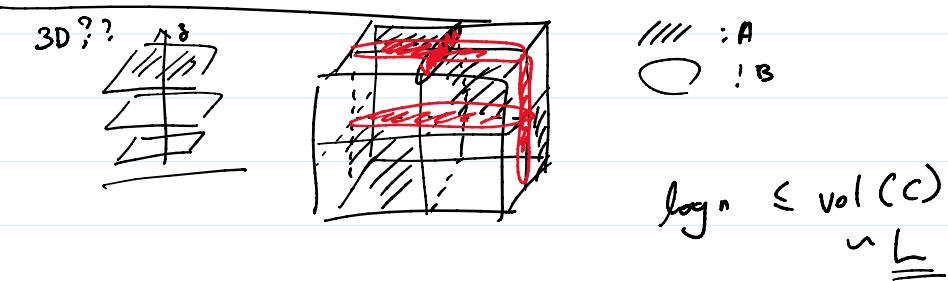


$$S_A + S_R = S_{AR} = S_{BC} \leq S_B + S_C$$

$$+) S_B + S_R = S_{BR} = S_{AC} \leq S_A + S_C$$

$$2S_R \leq 2S_c$$

$$\underline{\log n} \leq \frac{\underline{\text{vol}(C)}}{\underline{\Theta(1)}}$$



[0909.5200] Tradeoffs for reliable quantum information storage in 2D systems (arxiv.org)
[2009.13551] A degeneracy bound for homogeneous topological order (arxiv.org)

Part 2 : Current polys for Pauli Ham.



$$\Lambda = \underline{\mathbb{F}_2[\mathbb{Z}^\oplus]} \ni f = \text{subset in } \Lambda$$

$$\mathbb{F}_2[\underline{x^\pm, y^\pm}]$$



$$\begin{aligned} \textcircled{1} &\rightarrow \begin{pmatrix} 1 + y^+ \\ 1 + x^+ \end{pmatrix} \\ \textcircled{2} &\rightarrow \begin{pmatrix} 1 + y^- \\ 1 + x^- \end{pmatrix} \end{aligned}$$

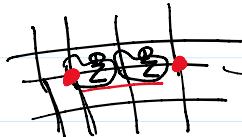
$$\begin{aligned} \mathbb{F}_2[x^\pm, y^\pm] \\ \textcircled{1} \rightarrow x^a y^b \begin{pmatrix} 1 + x^+ \\ 1 + y^- \end{pmatrix} = 0 \\ \textcircled{2} \rightarrow \end{aligned}$$

Pauli Z

$$\begin{pmatrix} 1+x \\ 0 \end{pmatrix}$$

$$\textcircled{2} = \sigma^+ = \begin{pmatrix} 1+x & 1+y \\ 0 & 1-x-y \end{pmatrix}$$

$$= 1+2x+x^2 = 1+x^2$$



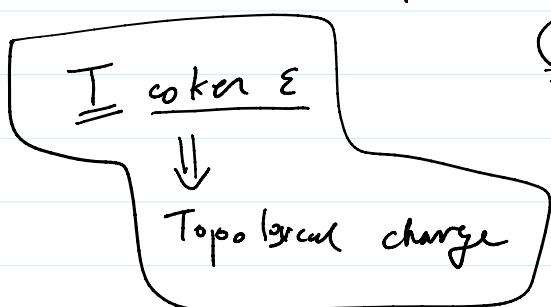
[1204.1063] Commuting Pauli Hamiltonians as maps between free modules (arxiv.org)

$$\begin{array}{ccc} R & \xleftarrow{\varepsilon} & R^2 \\ \text{position of } + & & \text{position of } \oplus \text{ and } \ominus \\ \hline R & = & 0 \end{array}$$

$$R = \mathbb{F}_2[x^\pm, y^\pm]$$

$$\underline{\text{coker } \varepsilon} = \frac{\overset{\circ}{R}}{\underline{\text{im } \varepsilon}} = \mathbb{F}_2[x^\pm, y^\pm] / (1+x, 1+y) \cong \mathbb{F}_2$$

$$\underline{\text{torsion }} T(M) = \{ z \in M \mid \exists r \in R \setminus \text{zero divisor } r z = 0 \}$$



$$\frac{(1+x)}{11} e^{\text{coker } \varepsilon} = 0$$

$$(1+x) \equiv 0$$

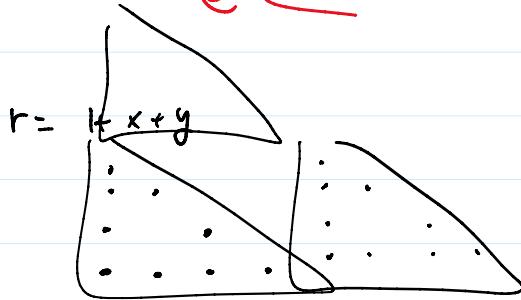
If \exists binomial zero-divisor,
 $x^a + y^b$

$$\left(\frac{x^a + y^b}{r} \right) e = \varepsilon(p)$$

$$(a+b)^{2^m} = a^{2^m} + b^{2^m}$$

$$\varepsilon(r^{2^m-1}p) = r^2 e \quad r^{2^{m-1}} \quad r^{2^{m-2}} \quad r^4 \cdot r^2 \cdot r \cdot p$$

$(\dots)(\dots)(\dots)(\dots) e$



Part 3 : Fractions

$$\varepsilon = \left(\frac{1+x+y+z}{1+xw+yz+zx}, \frac{1+xw+yz+zx}{1+x+y+z} \right)$$

Claim: $\not\exists$ binomial zero-divisor on $\text{dom } \varepsilon$

Pf)

$$\left\{ \begin{array}{l} 1+x+y+z=0 \\ 1+xw+yz+zx=0 \end{array} \right.$$

F_4

$$\left\{ \begin{array}{l} x = 1+t \\ y = 1+\omega t \\ z = 1+\omega^2 t \end{array} \right. \quad \left(\begin{array}{l} \omega^2 + \omega + 1 = 0 \\ \omega \leftrightarrow \omega^2 \end{array} \right)$$

$$(1+x+y+z) = (1)(1+x+y+z) + \dots + (1)(x+y+z) + \dots$$

$$(1 + x^a y^b z^c) = u \underbrace{(1 + x + y + z)}_{\begin{matrix} \downarrow \\ 11 \\ 0 \end{matrix}} + v \underbrace{(x + y + z^2 + zx)}_{\begin{matrix} \downarrow \\ 11 \\ 0 \end{matrix}}$$

$$(1 + t)^a (1 + \omega t)^b (1 + \omega^2 t)^c = 1$$

[1305.6973] Lattice quantum codes and exotic topological phases of matter (arxiv.org) contains this proof of no mobile excitation.

[1101.1962] Local stabilizer codes in three dimensions without string logical operators (arxiv.org) where this model is written first has more elementary argument.

