

## Practical Implementations: The Nuts and Bolts of DFT Part II



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# Outline

- Part I: The Self Consistency Cycle (SCF)
  - Density mixing
  - Metallic systems: Fermi level broadening
  - Preconditioning
- Part II: Structure optimization
  - Local structure optimization
  - "Forces" in density functional theory
  - Vibrations in the harmonic approximation
- Part III: Beyond conventional CPUs
  - Calculating with Graphical Processing Units (GPUs)

## Recap: Electron Density

All ground state properties are related to the electron distribution [1]:

 $\Psi(\mathbf{r}_1,\mathbf{r}_2,\ldots,\mathbf{r}_n) \Longleftrightarrow n(\mathbf{r})$ 

• Map electron density on effective one-particle orbitals with energy  $\varepsilon_i[2]$ :

$$n(\mathbf{r}) \Longrightarrow \sum_{i}^{\infty} f_i \left| \phi_i(\mathbf{r}) \right|^2$$

 $H[n]\phi_i(\mathbf{r}) = \epsilon_i \phi_i(\mathbf{r})$ 

"Occupation numbers"

$$f_i = \begin{cases} 0 \dots 2 & \text{not spin-polarized} \\ 0 \dots 1 & \text{spin-polarized} \end{cases}$$

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[1] P. Hohenberg, W. Kohn, Phys Rev. (1964), B864
[2]: W. Kohn, L.J. Sham, Phys. Rev. (1965), A1133

## Fractional Occupation Numbers

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- Justified as a statistical average
- Practical necessity: Brillouin zone integrals for metallic systems (later in this talk)
- Formal necessity: Physics of metals needs statistical mechanics
   Can lead to unexpected occupations for pure state systems





 $n(\mathbf{r}) \Longrightarrow \sum f_i |\phi_i(\mathbf{r})|^2$ 

#### Recap: Kohn-Sham Equations

 $H[n]\phi_i(\mathbf{r}) = \epsilon_i \phi_i(\mathbf{r})$ 

$$\left(-\frac{1}{2}\nabla^2 + \int_{\Omega} \frac{n(\mathbf{r})}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + V_{\rm xc}(\mathbf{r}) + V_{\rm ext}(\mathbf{r})\right) \phi_i(\mathbf{r}) = \epsilon_i \phi_i(\mathbf{r})$$

kinetic contribution

electron-electron interaction

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electron-ion interaction external fields

...

- We want to determine all solutions  $\Phi_i$  such that
  - the total energy E[n] is minimized
  - H[n] is consistent with  $\{\Phi_i\}$

#### Self Consistent Field Method (S.C.F.)

What kind of initial guess?



## The Initial Guess

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Random Numbers ?

- mainly in plane-wave based methods
- \* easy to implement
- no chemical motivation (unbiased)
- no chemical motivation (slow)
- Completely impractical for an all electron code (density is dominated by core electrons close to the nuclei)

$$n^0(\mathbf{r}_i) = RND([0\dots 1])$$

$$\int_{\Omega} n^0(\mathbf{r}) d\mathbf{r} = N_e$$



## The Initial Guess

Superposition of atomic densities (used in FHI-aims):

$$n^{\mathrm{at}}(\mathbf{r}) = \sum_{i} f_i^{\mathrm{occ}} |\chi_i(\mathbf{r})|^2$$

2

$$n^{0}(\mathbf{r}) = \sum_{I}^{N_{\mathrm{at}}} n_{I}^{\mathrm{at}}(\mathbf{r} - \mathbf{R}_{I})$$



- Extended Hückel Theory [1] (in Quantum Chemistry codes)
- not implemented in FHI-aims
  - Linear Combination of Atomic Orbitals

$$\phi_i^0(\mathbf{r}) = \sum_i c_{ij} \chi_j(\mathbf{r} - \mathbf{R}_j)$$

Hamiltonian [2,3]:

 $H_{ii}$ : parameterized atomic valence ionization energies

$$H_{ij} = KS_{ij} \frac{H_{ii} + H_{jj}}{2}$$

- Solve set of linear equations  $\hat{H}\mathbf{c}_i = \epsilon_i \hat{S}\mathbf{c}_i$
- Improved flexibility, specific orbitals can be populated  $n^{0}(\mathbf{r}) = \sum_{i} f_{i}^{\text{occ}} |\phi_{i}^{0}(\mathbf{r})|^{2}$

[1] R. Hoffmann, J Chem. Phys (**1963**), 1397 [2] R. S. Mulliken, J. Chem. Phys. (**1946**) 497 [3] M. Wolfsberg and L. Helmholtz, J. Chem. Phys (**1952**), 837

## Importance of Initial Guess

- Different minima might exist, leading to qualitative different properties
- e.g. O<sub>2</sub>: singlet vs. triplet
- symmetric solutions can be higher in energy



## Density Update

Evaluate new KS density from current KS orbitals

 $n_{\rm KS}^j(\mathbf{r}) = \sum_i f_i |\phi_i^j(\mathbf{r})|^2$ 

Naive way: Replace old density by new density

 $n^j(\mathbf{r}) = n^j_{\mathrm{KS}}(\mathbf{r})$ 

no convergence
 oscillating total energy



#### Simple Linear Mixing

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Define difference density

$$\Delta n(\mathbf{r}) = n_{\rm KS}^j(\mathbf{r}) - n^{j-1}(\mathbf{r})$$

Take only a fraction of this (underrelaxation method):

 $n^{j}(\mathbf{r}) = n^{j-1}(\mathbf{r}) + \alpha \Delta n(\mathbf{r})$ 

 too high α: unstable SCF behaviour
 too small α:

slow convergence

optimal α is
 system dependent



Pulay mixing (Direct inversion of iterative subspace)

the residual density: change between KS density and input density

 $R^{j}(\mathbf{r}) = n^{j}(\mathbf{r}) - n^{j}_{\mathrm{KS}}(\mathbf{r})$ 

Try to predict R<sup>j+1</sup> from N previous residues

$$R^{j+1}(\mathbf{r}) = \sum_{i=0}^{N-1} \beta_i R^{j-i}(\mathbf{r})$$

• find  $\beta_i$  by minimizing

$$F = ||R^{j+1}(\mathbf{r})||^2 - \lambda \left(1 - \sum_{i=0}^N \beta_i\right)$$

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[1] P. Pulay, Chem. Phys. Lett. 73, 393 (1980)

Pulay mixing (Direct inversion of iterative subspace)

Construct "optimal density":

$$n^{\mathrm{opt}}(\mathbf{r}) = n^{j}_{\mathrm{KS}} + R^{j}(\mathbf{r})$$

be conservative

$$n^{j}(\mathbf{r}) = (1 - \alpha)n^{j-1}(\mathbf{r}) + \alpha n^{\text{opt}}(\mathbf{r})$$

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increase of convergence speed
can be nearly independent of α
for insulating systems
α = 0.2 - 0.6
for metallic systems
α can be 0.05



### Metallic systems

- Metallic systems are more demanding in convergence
- Band crossing at the Fermi level
- Bad convergence of band structure energy in reciprocal space:

$$E_{\text{Band}} = \sum_{n} \frac{1}{\Omega} \int_{\Omega} \epsilon_{n\mathbf{k}} \Theta(\epsilon_{n\mathbf{k}} - \epsilon_{\text{Fermi}}) d\mathbf{k}$$

We operate with a finite k-mesh

$$\frac{1}{\Omega} \int\limits_{\Omega} d\mathbf{k} \to \sum_{N_k} w_{\mathbf{k}}$$

for metals : T=oK occupancies would jump discontinuously from 1 to 0 at the Fermi level

## Metallic systems

Fermi surfaces = collection of k-points with energy ε<sub>Fermi</sub>

related to Fermi level crossing of bands (structure-rich Fermi surface[1])



"level switching" of partially filled orbitals during SCF possible [1] <u>http://www.phys.ufl.edu/fermisurface/periodic\_table.html</u>

#### Fermi level broadening

• Stabilize integration by replacing  $\Theta$  with a smoother function

$$f_{\text{Fermi-Dirac}}\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) = \frac{1}{\exp\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) + 1}$$
 [1]

$$f_{\text{Gaussian}}\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) = \frac{1}{2} erfc\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) \quad [2]$$

$$f_{\text{Methfessel-Paxton}}\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) = \frac{1}{2} erfc\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) + \sum_{i=1}^{n} A_i H_{2i-1}\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right) \exp\left(-\left(\frac{\epsilon - \epsilon_{\text{Fermi}}}{\sigma}\right)^2\right)$$
(3)

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[1] Mermin, N. D. (1965), Phys.Rev., 137(5A), A1441-A1443
[2] Fu, C. L., & Ho, K. M. (1983), Physical Review B, 28, 5480-5486.
[3] Methfessel, M., & Paxton, A. T. (1989) Physical Review B, 40(6), 3616-3621

### Fermi Level Broadening

- electronic smearing avoids discontinuity at the Fermi level
- can be related to electronic temperature in Fermi-Dirac method



### Convergence

- Step function approach fails to converge
- Electronic broadening crucial to speed up the convergence



#### Free Energy

- Total energy depends on variations in  $f_i$
- Optimize free energy instead  $F = E \sigma S(\sigma)$
- Backextrapolation possible  $\sigma$  –
- Extrapolation Error << 1 meV</p>
- Applicable for metals, but not, in general, for atoms or small molecules with discrete electronic levels



Kresse, G., & Furthmüller, J. (1996). Computational Materials Science, 6(1), 15–50 M. Weinert and J.W. Davenport. Phys. Rev. B. 45 (1992) 13709. R.M. Wentzcovitch, J.L. Martins and P.B. Allen, Phys. Rev B 45 (1992) 11372

#### Surfaces, Interfaces, Defects, Thin Films

- e.g. Mo(100) surface
  - difficult electronic structure



- Long range density fluctuations in Surfaces, Interfaces, Defects, and Thin Films can lead to charge sloshing [1]
- Let mixing depend on the density  $n^{j}(\mathbf{r}) = n^{j-1}(\mathbf{r}) + \hat{G}\Delta n(\mathbf{r})$
- Idea: Treat density in Fourier space and damp long range density fluctuations

$$\Delta n(\mathbf{G}) = \frac{1}{\Omega} \int \Delta n(\mathbf{r}) e^{-\mathbf{G} \cdot \mathbf{r}} d\mathbf{r}$$

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[1] Kerker, G. P. (1981). Physical Review B, 23(6), 3082-3084

## Kerker Preconditioning

Based on Thomas Fermi Theory in reciprocal space[1,2]:

$$\hat{G}(\mathbf{G}) = \beta \frac{|\mathbf{G}|^2}{|\mathbf{G}|^2 + \lambda^2}$$

β: mixing parameter
λ: screening constant



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[1] M. Manninen, R. Nieminen, and P. Hautojavri, Phys. Rev. B 12, 4012 (1975) [2] Kerker, G. P. (1981). Physical Review B, 23(6), 3082–3084

## Summary SCF

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the SCF cycle

(

*initial guess* 



• approaches of density mixing  $n^{j}(\mathbf{r}) = (1 - \alpha)n^{j-1}(\mathbf{r}) + \alpha n^{\text{opt}}(\mathbf{r})$ 

Fermi level broadening

• preconditioner 
$$\hat{G}(\mathbf{G}) = \beta \frac{|\mathbf{G}|^2}{|\mathbf{G}|^2 + \lambda^2}$$





## Finding Stable Structures

- "The space of possible structures is endless, as well as the number of methods to find them":
  - Stochastical or Monte Carlo Methods (Talk: Peter Kratzer)
  - Molecular Dynamics (Talk: Luca Ghiringhelli, MauroMaggioni)
  - Genetic Algorithms
  - Diffusion Methods
  - Experimental structure determination
- However, once you are near a minimum, you have to find it...



### Local Structure Optimization

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- gradient free methods (a "total energy only" approach)
- evaluate a few points
- construct new point,
   e.g. assume harmonicity
   repeat

until self-consistency



#### Local Structure Optimization

- gradient based method ("Forces")  $\mathbf{F}_i = -\frac{\delta E_{\text{tot}}}{\delta \mathbf{R}_i}$
- evaluate Force(s) on the atoms analytically
- use standard scheme to optimize the structure
- steepest descent (SD)
   *just for illustration* conjugate gradient (CG)
  - better than SD
- (Quasi-) Newton approaches (BFGS) state-of-the-art approach
   How to evaluate the forces?



# Forces (light version)

Energy 
$$E = \frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

Forces

$$\mathbf{F}_i = -\frac{dE}{d\mathbf{R}_i}$$

• factor rule:  $\frac{dE}{d\mathbf{R}_{i}} = \frac{d}{d\mathbf{R}_{i}} \frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle}$   $= \frac{\frac{d}{d\mathbf{R}_{i}} \langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} - \frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} \frac{\frac{d}{d\mathbf{R}_{i}} \langle \Psi | \Psi \rangle}{\langle \Psi | \Psi \rangle}$   $= \frac{1}{\langle \Psi | \Psi \rangle} \left( \langle \Psi | \frac{d\hat{H}}{d\mathbf{R}_{i}} | \Psi \rangle + \langle \frac{d\Psi}{d\mathbf{R}_{i}} | \hat{H} - E | \Psi \rangle + \langle \Psi | \hat{H} - E | \frac{d\Psi}{d\mathbf{R}_{i}} \rangle \right)$ 

#### Hellman-Feynman Forces

Analyze first term

$$\mathbf{F}_{i} = -\frac{1}{\langle \Psi | \Psi \rangle} \left( \langle \Psi | \frac{d\hat{H}}{d\mathbf{R}_{i}} | \Psi \rangle + \langle \frac{d\Psi}{d\mathbf{R}_{i}} | \hat{H} - E | \Psi \rangle + \langle \Psi | \hat{H} - E | \frac{d\Psi}{d\mathbf{R}_{i}} \rangle \right)$$

affects only electron-nucleon and nucleon-nucleon interaction[1]

$$\mathbf{F}_{i}^{\text{Hellman-Feynman}} = \sum_{j} \frac{\delta}{\delta \mathbf{R}_{i}} \frac{Z_{i}Z_{j}}{|\mathbf{R}_{i} - \mathbf{R}_{j}|} - \int_{\Omega} n(\mathbf{r}) \frac{\delta}{\delta \mathbf{R}_{i}} \frac{Z_{i}}{|\mathbf{R}_{i} - \mathbf{r}|} d\mathbf{r}$$

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[1] Feynman, R. P. (1939). "Forces in Molecules". Phys. Rev. 56 (4): 340

#### **Pulay Forces**

What about the second term ?

$$\mathbf{F}_{i} = -\frac{1}{\langle \Psi | \Psi \rangle} \left( \langle \Psi | \frac{d\hat{H}}{d\mathbf{R}_{i}} | \Psi \rangle + \left| \langle \frac{d\Psi}{d\mathbf{R}_{i}} | \hat{H} - E | \Psi \rangle + \langle \Psi | \hat{H} - E | \frac{d\Psi}{d\mathbf{R}_{i}} \rangle \right]$$

applying basis expansion results in
Pulay Forces = artificial force due to moving basis:  $\mathbf{F}_{i}^{\text{Pulay}} = -2 \sum_{j,j',n} f_{n}^{\text{occ}} c_{jn}^{*} \langle \frac{\delta \chi_{j}}{\delta \mathbf{R}_{i}} | \hat{H} - \epsilon_{n} | \chi_{j'} \rangle c_{j'n}$ vanishes if basis set is stationary, e.g. plane waves

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(Talk: Anderson Janotti)

#### Steepest Descent

- Naive approach: Follow gradient in negative direction
    $\mathbf{R}^n = \mathbf{R}^{n-1} \alpha \mathbf{F}[\mathbf{R}^{n-1}]$ 
  - Step length as parameter
     Convergence depends on the step length
     Slow, oscillates around path
- Improvements:
   Line minimization for optimal step length
   Conjugate Gradient method



# (Quasi-) Newton Method

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• Approximate energy surface as quadratic form  $E(\mathbf{R} + \Delta \mathbf{R}) \approx E(\mathbf{R}) + \mathbf{F}^T(\mathbf{R})\Delta \mathbf{R} + \frac{1}{2}(\Delta \mathbf{R})^T \mathbf{H}(\Delta \mathbf{R})$ 

with Hesse matrix

$$H_{ij} = \frac{dE}{d\mathbf{R}_i d\mathbf{R}_j}$$

Minimum:

 $\Delta \mathbf{R} = \mathbf{H}^{-1} \mathbf{F}(\mathbf{R})$ 

 But H is unknown / expensive:
 cheaper construction of Hesse matrix needed



#### Hessian construction "on the fly"

Make a guess for the Hessematrix

scaled unity

 $\tilde{\mathbf{H}}^0 = \beta \mathbf{1}$ 

force field model based on internal coordinates [1] (e.g. Lindh initialization)

H stays positive definite
[1] Lindh, Bernhardsson, Karlström, Malmqvist, Chem. Phys. Lett. 241, 423 (1995)
[2] Broyden, C. G. (1970), Journal of the Institute of Mathematics and Its Applications 6: 76–90
[3] Fletcher, R. (1970), Computer Journal 13 (3): 317–322
[4] Goldfarb, D. (1970), Mathematics of Computation 24 (109): 23–26
[5] Shanno, David F. (1970), Math. Comput. 24 (111): 647–656

#### Stability and Improvements

Hesse Matrix is an approximation

Step control: Allow only finite step sizes to avoid overshooting

$$\Delta \mathbf{R} = \alpha H^{-1} \mathbf{F}(\mathbf{R})$$

Trust radius method:

• Evaluate quality of quadratic model  $q = \frac{E_{\text{true}}}{E_{\text{model}}}$ 

Adjust maximal step size based on q
 Introduce internal coordinates like rotations rather the cartesian steps

### Performance

α-helical peptide Ac-LysH<sup>+</sup>-Ala<sub>19</sub>, 220 atoms

> initial guess: MD snapshot





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[1] Lindh, Bernhardsson, Karlström, Malmqvist, Chem. Phys. Lett. 241, 423 (1995)

#### Vibrations

- Based on harmonic approximation (Harmonic oscillator)
- Perform Taylor expansion of energy up to second order

$$k = \mu \omega^{2}$$
  

$$E(r) = E_{0} - F \cdot r + \frac{1}{2} \omega^{2} r^{2} + O(r^{3})$$
  

$$Forces vanish for optimized geometries:$$
  

$$E(r) \approx E_{0} + \frac{1}{2} \omega^{2} r^{2}$$

#### Vibrations

 Assume harmonic approximation in 3N space:

 $E(\mathbf{R}) \approx E_0 + \frac{1}{2} \Delta \mathbf{R}^T \mathbf{H} \Delta \mathbf{R}$ 

Hesse matrix H needed



 Numerical evaluation of Hessian via finite differences (BFGS Hessian not accurate enough)
  $H_{\alpha\beta} = \frac{d^2E}{dR_{\alpha}dR_{\beta}} \approx -\frac{\mathbf{F}_{\alpha}(\mathbf{R} + h\mathbf{e}_{\beta}) - \mathbf{F}_{\alpha}(\mathbf{R} - h\mathbf{e}_{\beta})}{2h}$ 

### Vibrations

n

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Exponential Ansatz to solve Newtons equations

 $\mathbf{M}\frac{d^2}{dt^2}\Delta\mathbf{R} = -\mathbf{H}\Delta\mathbf{R}$ 

 $\Delta \mathbf{R} = \mathbf{u}e^{-\mathrm{i}\omega t}$ 

• Generalized eigenvalue problem:  $\mathbf{H}\mathbf{u}_n = \omega_n^2 \mathbf{M}\mathbf{u}_n$ 

• Solution gives eigenmodes and eigenfrequencies 3N-6 $\mathbf{R} = \mathbf{R}_0 + \sum_{n=1}^{3N-6} c_n \mathbf{u}_n \cos(\omega_n t + \varphi_n)$ 

*◆ 3N - 6:* 

3 modes describe translations

3 modes describe rotations



## Interpreting Vibrations

hard modes
 high ω
 e.g. stretching



 imaginary modes indicates sattle point, e.g. flat NH3 Reaction coordinate

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Reaction coordinate

## Beyond classical nuclei

• Free energy for T> o with quantum vibrational effects [1]  $F(T) = E_0 + \sum_{i} \frac{\hbar \omega_i}{2} + k_b T \sum_{i} \ln \left(1 - e^{-\hbar \omega_i / k_b T}\right)$ 

Zero-point vibrational energy correction

$$F(0) = E_0 + \sum_i \frac{\hbar\omega_i}{2}$$

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- Connect to experimental data
  - Infrared spectroscopy
  - Raman spectroscopy
- More details (Tutorial: Christian Carbogno)

[1] F. Jensen, Introduction to Computational Chemistry

## Summary Geometry Optimization

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- Forces in electronic structure theory
  - Hellman-Feynman and Pulay Forces
- Minimization techniques
  - Steepest descent
  - Quasi-Newton (BFGS)
- Atomic Vibrations
  - structure stability
  - zero-point energy
  - one first step to go beyond T=oK



 $rac{\delta E_{
m tot}}{\delta {f R}_{i}}$ 

 $\mathbf{F}_i = -$ 

## Beyond conventional CPUs



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http://www.futuretimeline.net/21stcentury/2019.htm#.U8RgGagmAZs

Every 24 months ...

#### Moore's observation



## Towards the ExaFLOP scale

- Moore's observation for FLoating-point OPerations per second
- Prediction: Overcome ExaFLOP barrier in 2018



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## Are we ready for the ExaScale ?

- Current ExaScale model predicts a machine with 1 billion cores
- If 1 billion cores would be typical CPUs then the predicted energy consumption would be more than 1 GW



## Modern Compute Architectures

- New architectures: GPU computing
- Power consumption:
   10.5 Watt / CPU core
   0.08 Watt / GPU core
   0.77 % Watt / core
- GPUs are a possibility towards exaFLOP scaling
- ... and are available

Intel Xeon E7-4830 V2 up to 1536 GB RAM 10 cores (20 threads) 105 W

NVIDIA K 6 GB GPU RAM 2880 CUDA cores 235 W





## The TOP 10 Supercomputers

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|--|

	Sile	Manufacturer	Computer	Country	Cores	Rmax [Pflops]	Power [MW]
1	National University of Defense Technology	NUDT	Tianhe-2 NUDT TH-IVB-FEP, Xeon 12C 2.20Hz, IntelXeon Phi	China	3,120,000	33.9	17.8
2	Oak Ridge National Laboratory	Cray	Titon Cray XK7, Opteron 16C 2.2GHz, Gemini, NVIDIA K20x	USA	560,640	17.6	8.21
3	Lawrence Livermore National Laboratory	IBM	Sequeia BlueGene/Q, Power BQC 16C 1.5 GHz, Custom	USA	1,572,864	17.2	7.89
4	RIKEN Advanced Institute for Computational Science	Fujitsu	K Computer SPARC64 VIIItx 2.0GHz, Tolu Interconnect	Japan	795,024	10.5	12.7
6	Argonne National Laboratory	IBM	Blue Gene IQ, Power BQC 16C 1.6 GHz, Custom	USA	786,432	8.59	3.95
6	Swiss National Supercomputing Centre (CSCS)	Crwy	Piz Daint Cray XC30, Xeon E5 8C 2.6GHz, Aries, NVIDIA K20x	Switzer-land	115,984	6.27	2.33
7	Texas Advanced Computing Center UT	Dell	Stampede PowerEdge C8220, Xeon ES BC 2.70Hz, Intel Xeon Phi	USA	462,462	5.17	4.51
	Forschungszentrum Juelich (FZJ)	ВМ	JuQUEEN BlueGene/0, Power BQC 16C 1.6GHz, Custom	Germany	458,752	5.01	2.30
9	Lawrence Livermore National Laboratory	BM	Vulcan BlueGeneiQ, Power BOC 16C 1.8GHz, Custom	USA	393,216	4.29	1.97
10	Government	Cray	Cray XC30, Xeon ES 12C 2.7GHz, Arles	USA	225,984	3.14	

00500

1) Tianhe-2 (Heavenrider-2) China 32,000 Intel Xeon E5 12 Cores **48,000 Intel Xeon Phi** 33.86 PFLOPS

2) Titan USA 18,688 AMD Opteron 6274 16 Cores 18,688 Nvidia Tesla K20X GPUs 17.59 PFLOPS (27 theoretical)

6) Piz Daint Switzerland 5272 Intel Xeon E5 8Core 5272 Nvidia Tesla K20X GPUs 7.78 PFLOPS

# Modern Computer Architectures

"classical" CPU Part

Data Communication

Accelerator Part



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distinct memory for CPUs and GPUs
 data has to be copied using (slow) PCIe bus
 GPU memory (4-12 GB / GPU) < CPU RAM (2-4 GB / CPU)</li>

#### Program structure

typically suggested way:



What about the the CPU cores, while the GPU is working
 CPUs should be working as well
 One GPU "beating" 10 - 20 CPU cores is a much harder task

#### Example: Construct Hamiltonian matrix

The Hamilton matrix construction is a compute intense operation

$$H_{ij} = \int_{\Omega} \chi_i^*(\mathbf{r}) \hat{H} \chi_j(\mathbf{r}) d\mathbf{r}$$

- Operation is performed point-wise in real space
- Operation is organised in batches = collection of points in real space



# Parallelization + Loadbalancing



#### e.g. space integration



# Load Balancing Hamiltonian

Ongoing work: Example C (diamond) - Gamma point only 8000 basis functions, 512 batches a 70 points



We have to aim for a Speedup of 21 to achieve a node speedup of 2.

### Summary

CPUs cannot be the way for ExaFLOP scale (Energy consumption)

 Accelerator cards such as GPUs are a promising architecture

- Accelerator cards are available on compute facilities
- GPUs and CPUs have to work together (Load balancing)

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Ongoing work





						Rmax [Pflops]	Power [MN]
1	National University of Defense Technology	NUDT	Tianha-2 NUDT TH-IVB-FEP, Xeon 12C 2.20Hz, IntelXeon Phi	China	3,120,000	33.9	17
2	Oak Ridge National Laboratory	Cray	Titan Gray XK7, Opteron 16C 2.2GHz, Gemini, NVIDIA K20x	USA	560,640	17.6	8.2
3	Lawrence Uvermore National Laboratory	BM	BlueGenelQ, Power BOC 16C 1.5 CHz, Custom	USA	1,572,864	17.2	7.8
4	RIKEN Advanced Institute for Computational Science	Fujitsu	K Computer SPARC64 VIIItx 2.0GHz, Tolu Interconnect	Japan	795,024	10.5	12
6	Argonne National Laboratory	8M	Blue Gene 10, Power BGC 16C 1.6 GHz, Custom	USA	786,432	8.59	3.9
6	Swias National Supercomputing Centre (CSCS)	Cray	Pir Daint Crey XC30, Xeon E5.8C 2.6GHz, Aries, NVIDIA K20x	Switzer-land	115,984	6.27	2.1
,	Texas Advanced Computing Center UT	Dell	Stampede PowerEdge C8220, Xeon ES 8C 2.70Hz, Intel Xeon Phi	USA	462,462	6.17	43
•	Forschungszentrum Juelich (FZJ)	84	JuGUEEN Blue Gene 0, Power BQC 16C 1.6 GHz, Custom	Germany	458,752	5.01	2.1
9	Lawrence Livermore National Laboratory	BM	Vulcan BlueGenelQ, Power BOC 16C 1.6 CHz, Custom	USA	383,216	4.29	1.5
0	Government	Cray	Cray XC30, Xeon ES 12C 2.7GHz, Arles	USA	225,984	3.14	



## Thank you !

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