

# Uniqueness of special Lagrangian submanifolds in hamiltonian isotopy orbits of Lagrangian submanifolds

*"Mirror Symmetry"*, Hori, Katz, Klemm, Pandharipande, Thomas, Vafa, Vakil, Zaslow, Clay Mathematics Institute series. 2000 (2003??)

Moment maps, monodromy and mirror manifolds, math.DG/0104196

Special Lagrangians, stable bundles and mean curvature flow, with S.-T. Yau. math.DG/0104197

Stability conditions and the braid group, math.AG/0212214

Let  $X = (X^n, \omega, \Omega)$  be an *almost Calabi-Yau manifold*:

- $X$  is a complex manifold of dimension $_{\mathbb{C}}$   $n$ ,
- $\omega$  is a Kähler form on  $X$ ,
- $\Omega$  is a nowhere-zero holomorphic  $n$ -form  
 $\Omega \in H^0(K_X)$ .      “ $\Omega = dz_1 \wedge \dots \wedge dz_n$ ”

If  $\Omega$  is compatible with the Kähler metric  $\omega$  (i.e.  $|\Omega| \equiv 1$ ; equivalently  $\Omega$  is Levi-Civita parallel) then  $X$  is Calabi-Yau. By Yau's theorem we can arrange this by altering  $\omega$  if  $X$  is compact.

Let  $L \subset X$  be a *Lagrangian* submanifold:

- $L$  has real dimension  $n$ , and
- $\omega|_L \equiv 0$ .

An easy local computation shows that  $\Omega|_L = |\Omega|e^{i\theta} \text{vol}_L$  for some phase  $\theta$ . A *graded* Lagrangian  $(L, \theta)$  is  $L$  together with a choice of a lift of  $e^{i\theta}$  to a real-valued function

$$\theta : L \rightarrow \mathbb{R}.$$

All our Lagrangians will be (implicitly) graded; in particular we consider only Lagrangians whose Maslov class

$$H^1(L; 2\pi\mathbb{Z}) \ni [d\theta] = 0.$$

$L$  is a *special Lagrangian* if  $\theta \equiv \text{constant}$ . By rotating  $\Omega$  we can assume without loss of generality that this *cohomological* constant (phase of  $\int_L \Omega$ ) is zero; that is  $L$  satisfies

- $\text{Im } \Omega|_L \equiv 0$ ; equivalently
- $\text{Re } \Omega|_L \equiv |\Omega| \text{vol}_L$ .

The special Lagrangian condition is a first order elliptic PDE describing minimal submanifolds of  $X$  (in the *almost* CY case have to define volume with respect to  $|\Omega| \text{vol}_L$ ).

We will call a Lagrangian  $L$  *close to being special Lagrangian* if  $\text{Re} \Omega|_L$  is nondegenerate on  $L$ , i.e. defines a volume form. Equivalently, the phase function  $\theta$  has bounded variation in  $(-\pi/2, \pi/2)$ . Such Lagrangians clearly have zero Maslov class and are gradeable.

## Mirror symmetry

Mirror symmetry exchanges complex geometry (governed by  $\Omega$ ) on one Calabi-Yau with symplectic geometry (governed by  $\omega$ ) on a *mirror* Calabi-Yau.

Very roughly speaking, under mirror symmetry, Lagrangians correspond to holomorphic bundles on the mirror Calabi-Yau. (The former are defined using on the symplectic structure, the latter only the complex structure.) Special Lagrangians correspond to (perturbed or MMMS) Hermitian-Yang-Mills connections on the mirror bundle.

## **Theorem 1 Donaldson-Uhlenbeck-Yau**

*A holomorphic bundle  $E$  admits a Hermitian metric such that the induced connection is Hermitian-Yang-Mills iff  $E$  is (poly)stable. The metric is then unique.*

( $E$  is stable iff for all coherent subsheaves  $F \hookrightarrow E$ , we have  $\mu(F) < \mu(E)$ . Here

$$\mu(E) = \frac{\int_X c_1(E) \cup \omega^{n-1}}{\text{rank } E}$$

is the only place that the Kähler form  $\omega$  enters, just numerically through its cohomology class. Stable bundles are “generic”; they form a Zariski open subset.)

The theorem reflects the infinite dimensional Kähler geometry of the action of the gauge group on the space of connections on  $E$ . Moment maps, Yang-Mills flow, convexity...

To get similar structures on the space of Lagrangians, we need to complexify.

$$\mathcal{A} := \{ (L, A) : L \subset X \text{ Lagrangian,} \\ A \text{ a flat } U(1) \text{ connection on } L \}$$

We restrict attention to Lagrangians close to being special Lagrangian, and connections on the trivial line bundle  $L \times \mathbb{C}$ .

The tangent space to  $\mathcal{A}$  at  $(L, A)$

$$Z^1(L, \mathbb{R}) \oplus iZ^1(L, \mathbb{R}) = \text{Defs of } L \oplus \text{Defs of } A$$

is naturally complex, and we can induce a non-degenerate 2-form on it by defining a metric on  $Z_L^1 \subset \Omega_L^1$  by

$$\begin{aligned} \langle a, b \rangle &= \int_L a \wedge (b \lrcorner \omega^{-1}) \lrcorner \text{Im } \Omega \\ &= \int_L (a \wedge *b) \text{Re } \Omega|_L. \end{aligned}$$

$C^\infty(X, U(1)) \ni \phi$  acts on  $\mathcal{A}$  as the gauge transformation

$$e^{i\phi}|_L : A \mapsto A + id\phi.$$

This clearly complexifies via the above complex structure to give an action of  $C^\infty(L, \mathbb{R})$  by exact deformations  $d\phi \in \Omega_L^1$  of  $L$  – *hamiltonian deformations*.

Combining these two actions gives an action of  $C^\infty(X, \mathbb{C}^\times) \ni f = he^{i\phi}$ : as gauge transformations on  $A$  by  $e^{i\phi}|_L$ , and as hamiltonian isotopies of  $(L, A)$  induced from hamiltonian diffeomorphisms of  $X$  with hamiltonian  $h$ .

(I.e. by the time-one flow of its hamiltonian vector field  $X_h = dh \lrcorner \omega^{-1}$ .)



$C^\infty(X, \mathbb{C}^\times)$  preserves the complex structure on  $\mathcal{A}$ ;  $C^\infty(X, U(1))$  also preserves the symplectic structure and has a *moment map*

$$\mu : \mathcal{A} \rightarrow (C^\infty(X, \mathbb{R}))^*, \quad \mu(L, A) = \text{Im } \Omega|_L.$$

(Here the current  $\text{Im } \Omega|_L$  is the linear functional on  $C^\infty(X, \mathbb{R})$  taking  $f$  to  $\int_L (f \text{Im } \Omega)|_L$ .)

The upshot of the general theory is that we expect the generic (“*stable*”) Lagrangian (close to being special Lagrangian) to have in its hamiltonian isotopy class a zero of the moment map,

$$\text{Im } \Omega|_L \equiv 0,$$

i.e. a *special Lagrangian*; this should then be *unique*.

We also expect to reach this special Lagrangian (zero of the moment map) via the gradient flow of  $-||\mu||^2$ ; this can be shown to be *mean curvature flow*, which is indeed hamiltonian for Lagrangians of Maslov class zero.

Other considerations (angle criterion, an example of Harvey-Lawson and Joyce, Kontsevich's mirror conjecture) lead to the following Hitchin-Kobayashi-type conjecture.

**Conjecture 1** *Take graded Lagrangians  $(L_1, \theta_1)$  and  $(L_2, \theta_2)$ , hamiltonian isotoped to intersect cleanly, and such that the graded (relative) Lagrangian connect sums  $(L_1 \# L_2, \theta_1 \# \theta_2)$  exist. Then a Lagrangian  $L$  of Maslov class zero is said to be destabilised by the  $L_i$  if it is hamiltonian isotopic to such an  $L_1 \# L_2$ , and the phases (real numbers, induced by the gradings) satisfy*

$$\phi(L_1) \geq \phi(L_2).$$

*If  $L$  is close to being a special Lagrangian and is not destabilised by any such  $L_i$  then it should have a special Lagrangian in its hamiltonian deformation class, and this should be unique.*

*If  $L$ 's phase variation is sufficiently small then this special Lagrangian should be the infinite time limit of MCF of  $L$ .*

Infinitesimally, the linearised formal picture is clear; it becomes simple Hodge theory.

Deformations of Lagrangians are given by closed 1-forms  $\ker d : \Omega^1(L; \mathbb{R}) \rightarrow \Omega^2(L; \mathbb{R})$ , so that dividing by hamiltonian deformations we get

$$\begin{aligned} T_L(\{\text{Lagrangians/hamiltonian isotopies}\}) \\ = \ker d / \text{im } d = H^1(L; \mathbb{R}). \end{aligned}$$

If instead of dividing by hamiltonian deformations we impose the *special* condition  $\text{Im } \Omega|_L = 0$ , we get a  $\ker d^*$  slice

$$\begin{aligned} T_L(\{\text{Special Lagrangians}\}) = \\ \ker d \cap \ker d^* = H^1(L; \mathbb{R}), \end{aligned}$$

to the hamiltonian deformations.

The Kähler quotient and moment map picture is a vast nonlinear generalisation of this linear fact that the quotient of a vector space by a subspace is isomorphic to the orthogonal complement of the subspace.

## Uniqueness

Even in infinite dimensional Kähler quotient problems, convexity properties of the moment map usually prove uniqueness of its zeros (modulo the action of the real group) in a complexified group orbit. The proof only relies on the existence of 1-parameter subgroups connecting any two points of the orbit.

Unfortunately this need *not* be the case here as most hamiltonian isotopies are *not* time independent. They also need not preserve the “close to special Lagrangian” condition.

But for those time-independent isotopies which *do* preserve this, we can demonstrate the general picture in our context as follows.

**Theorem 2** *If two SLags  $L_0, L_1$  are time-independent hamiltonian deformations of each other through Lagrangians  $L$  with  $\text{Re } \Omega|_L > 0$ , then  $L_0 = L_1$ .*

*Proof* We are supposing that  $L_0 = \Phi(L_1)$ , where  $\Phi$  is the time-one hamiltonian isotopy of a hamiltonian  $h \in C^\infty(X, \mathbb{R})$ ; that is  $\Phi = \Phi_1$ , where

$$\Phi_0 = \text{id}, \quad \frac{d}{dt} \Phi_t = X_h = dh \lrcorner \omega^{-1}.$$

(Or, we can allow  $h$  to be time independent as it follows  $L$ , in the sense that we have  $h = h_t$  with  $\Phi_t^* h_t|_L = \text{constant}$ . Then there is no  $X_h(h)$  term below.)

We compute, down the flow,

$$\begin{aligned} \frac{d}{dt} \int_L h \text{Im } \Omega|_L &= \int_L \left\{ X_h(h) \text{Im } \Omega + h \mathcal{L}_{X_h}(\text{Im } \Omega) \right\} \\ &= \int_L h d(X_h \lrcorner \text{Im } \Omega) \\ &= \int_L (dh \wedge *dh) \text{Re } \Omega|_L. \end{aligned}$$

So for phase  $\theta$  lying in  $(-\pi/2, \pi/2)$ , i.e.  $\operatorname{Re} \Omega|_L$  strictly positive,

$$\frac{d}{dt} \int_L h \operatorname{Im} \Omega|_L \geq 0,$$

with equality iff  $h|_L = \text{constant}$ .

But  $\int_L h \operatorname{Im} \Omega$  is zero at  $t = 0, 1$ . Thus  $h|_L \equiv \text{constant}$  and the two special Lagrangians coincide.  $\square$

For the general case we mirror the algebro-geometric argument that a non-zero map between stable bundles of the same slope is an isomorphism. So we need to use the Floer homology of Fukaya-Oh-Ohta-Ono.

Call a connected graded Lagrangian  $L$  *unobstructed* if  $w_2(L)$  is the restriction of a class  $\in H^2(X; \mathbb{Z}/2)$  (e.g. if  $L$  is spin), and the [FOOO] obstructions to the existence of its Floer cohomology vanish. E.g. Lagrangian homology spheres are unobstructed in dimension 3 and above.

Graded Lagrangians have graded Floer homology groups. Suppose  $L_2$  has phase 0 at an intersection point  $p$  with  $L_1$ , and let its tangent space be the  $x$ -axes:

$$L_2 = \{y_i = 0\}.$$

Then

$$L_1 = \{z_i = re^{i\alpha_i}\}, \quad \alpha_i \in (0, \pi) \forall i.$$

Then  $\sum \alpha_i = \theta_p(L_1) \bmod \pi$  and we define the Floer index at  $p$  to be

$$\text{ind}_p(L_2, L_1) := \frac{1}{\pi} \left( \sum \alpha_i - \theta_p(L_1) \right).$$

Notice that  $\text{ind}_p(L_2, L_1) + \text{ind}_p(L_1, L_2) = n$ .



It is not clear if the following theorem can be extended to obstructed Lagrangians. The role of Floer homology is analogous to that in the proofs of the Arnol'd conjecture.

**Theorem 3** *There is at most one smooth special Lagrangian in the hamiltonian deformation class of an unobstructed Lagrangian  $L$ .*

*In particular, special Lagrangian homology spheres are unique in their hamiltonian deformation class in dimension 3 and above.*

*Proof* The key points are that Floer cohomology is independent of hamiltonian deformations [FOOO], and the zeroth order piece of  $H^*(L)$  survives in  $HF^*(L, L)$  for  $L$  with Maslov class zero [FOOO].

Therefore two graded special Lagrangians  $L_1, L_2$  in the same hamiltonian deformation class satisfy

$$HF^0(L_1, L_2) = H^0(L_1; \mathbb{C}) = \mathbb{C},$$

and there must be at least one intersection point  $p$  of  $L_1$  and  $L_2$ .

The (constant) phases of the  $L_i$  differ by  $r\pi$  for some  $r \in \mathbb{Z}$ ; we first show that  $r = 0$ .

Hamiltonian perturbation  $\implies$  transverse intersection point of  $L_i$ s of Floer index 0, with the phases of the  $L_i$  at this point differing by  $r\pi + \epsilon$ .

Locally  $L_2$  is graph of  $df$  in  $T^*L_1$ ;  $f$  Morse index  $r$ , so  $r \geq 0$ .

Similarly there is a point of Floer index  $n$  (index 0 on swapping  $L_1, L_2$ ); locally  $dg$  with  $g$  of Morse index  $n + r$ ; so  $r \leq 0$  also. Therefore  $r = 0$  as required.

If the intersection points are isolated, one must have Floer index zero:

$$\sum \alpha_i = 0, \quad \alpha_i \in (0, \pi) \quad \forall i.$$

So the relative angles  $\alpha_i = 0 \quad \forall i$  and the  $L_i$  are tangent at  $p$ ; there is no isolated transverse intersection point of Maslov class zero.

Near  $p$  write  $L_1$  as the graph in  $T^*L_2$  of  $df$  with  $d^*df = 0$ :  $f$  harmonic. By the maximum principle,  $f$  has no local maxima or minima. This, roughly, is the proof, but we must also deal with non-Morse  $f$ , i.e. non-transverse intersections.

We want to show that  $f$  is constant, so assume for a contradiction that the critical set of  $f$  is not all of  $L_2$ .

Perturb  $f$  inside a connected component of a small neighbourhood of its critical set such

that its value is unchanged on the boundary, *where it attains its global maximum and minimum*, and is Morse in the interior.

By Morse theory we can then perturb  $f$  further to arrange its index 1 critical points to be lower (with respect to  $f$ ) than all higher index points and then cancel any local minima with them. (There must be index 1 critical points if there are interior minima, by connectivity of our neighbourhood.)

So we get a hamiltonian perturbation of  $L_1$  with no Floer index zero intersection points with  $L_2$ . Thus  $HF^0(L_1, L_2) = 0$ , a contradiction. □