

# Quasi-relative entropy: the closest separable state and reversed Pinsker inequality

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# Outline of the talk

- Quasi-relative entropy (*f*-divergence)
- Pinsker and reversed Pinsker inequalities
- Reversed Pinsker for Tsallis entropy
- Reversed Pinsker for quasi-relative entropy
- The 'closest' separable state
- Future Research

The quantum relative entropy between two states  $\rho$  and  $\sigma$  is as follows

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Additivity of relative entropy

 $S(\rho_1 \otimes \rho_2 || \sigma_1 \otimes \sigma_2) = S(\rho_1 || \sigma_1) + S(\rho_2 || \sigma_2)$ 

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Monotonicity of Quantum Relative Entropy (Data Processing Inequality)  $S(\rho \| \sigma) \geq S(\mathcal{N}(\rho) \| \mathcal{N}(\sigma))$ 

Completely Positive Trace Preserving (CPTP) map

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In particular,

$$S(\rho^{AB} \| \sigma^{AB}) \ge S(\rho^A \| \sigma^A)$$

Strong Sub-additivity of quantum entropy

For a tri-partite state  $\rho^{ABC}$ 

 $S(\rho^{ABC}) + S(\rho^B) \le S(\rho^{AB}) + S(\rho^{BC})$ 

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Joint Convexity of Quantum Relative Entropy

For 
$$\rho = \sum_{j} p_{j} \rho_{j}$$
 and  $\sigma = \sum_{j} p_{j} \sigma_{j}$   
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Strong Sub-additivity of quantum entropy For a tri-partite state  $\rho^{ABC}$  $S(\rho^{ABC}) + S(\rho^B) \leq S(\rho^{AB}) + S(\rho^{BC})$ Joint Convexity of Quantum Relative Entropy For  $\rho = \sum_{j} p_{j} \rho_{j}$  and  $\sigma = \sum_{j} p_{j} \sigma_{j}$  $S(\rho \| \sigma) \leq \sum_{j} p_j S(\rho_j \| \sigma_j)$ all "equivalent" Monotonicity of Quantum Relative Entropy  $S(\rho \| \sigma) > S(\mathcal{N}(\rho) \| \mathcal{N}(\sigma))$ Monotonicity under partial traces  $S(\rho^{AB} \| \sigma^{AB}) \ge S(\rho^A \| \sigma^A)$ e.g. Monotonicity under partial traces holds for all states iff Strong sub-additivity relation holds for all states

#### Stinespring factorization theorem

 $\mathcal{N}(\rho) = \mathrm{Tr}_E(U^*(\rho \otimes 1_E)U)$ 

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see "Strong Subadditivity of Quantum Entropy" on Wikipedia

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Monotonicity (Data Processing inequality)

$$\|\rho - \sigma\|_1 \ge \|\mathcal{N}(\rho) - \mathcal{N}(\sigma)\|_1$$

$$\smile \quad \mathsf{CPTP map}$$

In particular,

$$\|\rho^{AB} - \sigma^{AB}\|_1 \ge \|\rho^A - \sigma^A\|_1$$

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Let  $f: (0, \infty) \to \mathbb{R}$  be operator monotone decreasing, and f(1) = 0, *quasi-relative entropy*, or *f-divergence* is

 $S_f(\rho \| \sigma) = \text{Tr}\{f(\Delta_{\sigma,\rho})\rho\}$ 

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For  $f(x) = -\log x$ , the quasi-relative entropy becomes the relative entropy

$$S_{-\log}(\rho \| \sigma) = S(\rho \| \sigma) = \operatorname{Tr}(\rho \log \rho - \rho \log \sigma)$$

For  $q \in (0,2)$ , the function  $f_q(x) = \frac{1}{1-q}(1-x^{1-q})$  gives *Tsallis q-entropy* 

$$S_q(\rho \| \sigma) = \frac{1}{1-q} \left( 1 - \operatorname{Tr}(\rho^q \sigma^{1-q}) \right)$$

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$$S_q(\rho \| \sigma) = \frac{1}{1-q} \left( 1 - \operatorname{Tr}(\rho^q \sigma^{1-q}) \right)$$

is used in entanglement and thermodynamics, nonextensive statistics, optical lattice theory, particle charging, statistical mechanics, and others - see "Quantum Entropies" on Scholarpedia

For  $\alpha \in (0,1)$ , the function  $f_{\alpha}(x) = 1 - x^{1-\alpha}$  gives

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which defines Renyi relative entropy

$$S_{\alpha}^{R}(\rho \| \sigma) = \frac{1}{\alpha - 1} \log \operatorname{Tr}(\rho^{\alpha} \sigma^{1 - \alpha}) = \frac{1}{\alpha - 1} \log(1 - S_{\alpha}(\rho \| \sigma)))$$

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is used in hypothesis testing (Csiszar '95); entanglement-assisted LOCC conversion; strong converse problem in quantum hypothesis testing (Mosonyi, Ogawa, '15); strong converse problem for the classical capacity of a quantum channel (Wilde et. al., '14)

Unitary invariance

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Unitary invariance

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#### Positivity

 $S_f(\rho \| \sigma) \ge 0$  and  $S_f(\rho \| \sigma) = 0$  if and only if  $\rho = \sigma$ 

Data Processing Inequality

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#### Trace-distance

$$\|\rho - \sigma\|_1 \ge \|\mathcal{N}(\rho) - \mathcal{N}(\sigma)\|_1$$

# Relative entropy vs Trace distance

Pinsker inequality

$$\frac{1}{2} \|\rho - \sigma\|_1^2 \le S(\rho \|\sigma)$$
### Relative entropy vs Trace distance

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## Pinsker inequality for quasi-relative entropy

By Hiai and Mosonyi '16

$$\frac{f''(1)}{2} \|\rho - \sigma\|_1^2 \le S_f(\rho \|\sigma)$$

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### Reversed Pinsker inequality

By Audenaert, Eisert '11

$$S(\rho \| \sigma) \le (\alpha_{\sigma} + T) \log(1 + T/\alpha_{\sigma}) - \alpha_{\rho} \log(1 + T/\alpha_{\rho})$$

 $T = \frac{1}{2} \| \rho - \sigma \|_1^2$   $\alpha_{\omega}$  is the minimal non-zero eigenvalue of the state  $\omega$ 

For  $q \in (0, 2)$ , the *Tsallis* q-entropy is

$$S_q(\rho \| \sigma) = \frac{1}{1-q} \left( 1 - \operatorname{Tr}(\rho^q \sigma^{1-q}) \right)$$

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Rastegin '11

 $1 < q \leq 2$ 

$$S_q(\rho \| \sigma) \le \frac{1}{q-1} \frac{\lambda_{\rho}^q}{\alpha^q} \| \rho - \sigma \|_1$$

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Rastegin '11

 $1 < q \le 2 \qquad \qquad 0 < q < 1$ 

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### Reversed Pinsker for quasi-relative entropy

### Theorem (V. '19)

Let f any operator monotone decreasing function. Let states  $\rho$  and  $\sigma$  be either 2-dimensional qubit states or classical states. Assume one of two conditions: 1)  $\rho$  is full rank; 2)  $a_f = 0$ . Then the following holds

$$S_f(\rho \| \sigma) \le \|\rho - \sigma\|_1 \left[ \frac{\lambda_\rho}{\lambda_\rho - \alpha_\sigma} f(\lambda_\rho^{-1} \alpha_\sigma) - a_f \right]$$

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 $\lambda_{\rho} \in (0,1]$  is the largest eigenvalue of  $\rho$ 

 $\alpha_{\sigma} \in (0,1]$  is the smallest eigenvalue of  $\sigma$ 

$$a_f = -\lim_{y \uparrow \infty} \frac{f(iy)}{iy}$$

Theorem

In any finite dimensions,

$$S(\rho \| \sigma) \le \| \rho - \sigma \|_1 \lambda_\rho \frac{\log(\alpha_\rho) - \log(\alpha_\sigma)}{\alpha_\rho - \alpha_\sigma} \le \frac{\lambda_\rho}{\alpha} \| \rho - \sigma \|_1$$

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For qubits, taking  $f(x) = -\log(x)$ , we have a slightly improved bound

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Before,

$$S(\rho \| \sigma) \le (\alpha_{\sigma} + T) \log(1 + T/\alpha_{\sigma}) - \alpha_{\rho} \log(1 + T/\alpha_{\rho})$$

 $T = \frac{1}{2} \|\rho - \sigma\|_1^2$ 

In any dimension larger than four there are states for which our bound is better.

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$$\geq 1 \text{ for } q \in (1, 2]$$

# Reversed Pinsker for Tsallis relative entropy For $q \in (0, 1)$

$$S_q(\rho \| \sigma) \le \frac{1}{1-q} \| \rho - \sigma \|_1 \lambda_\rho^q \frac{\alpha_\rho^{1-q} - \alpha_\sigma^{1-q}}{\alpha_\rho - \alpha_\sigma} \le \| \rho - \sigma \|_1 \frac{\lambda_\rho^q}{\alpha^q}$$

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original 
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Every operator monotone decreasing function f has the following integral representation (Donoghue '74)

$$f(x) = -a_f x - b_f + \int_0^\infty \left(\frac{1}{t+x} - \frac{t}{t^2+1}\right) d\mu_f(t)$$

here  $a_f := -\lim_{y \uparrow \infty} \frac{f(iy)}{iy} \ge 0$ ,  $b_f := -\operatorname{Re} f(i) \in \mathbb{R}$  and  $\mu$  is a positive measure on  $(0, \infty)$  s.t.

$$\int_0^\infty \frac{1}{t^2 + 1} \mathrm{d}\mu_f(t) < \infty$$

and  $\mu_f(x_1) - \mu_f(x_0) = -\lim_{y \downarrow 0} \frac{1}{\pi} \int_{x_0}^{x_1} \operatorname{Im} f(-x + iy) \, \mathrm{d}x$ 

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and  $\mu_f(x_1) - \mu_f(x_0) = -\lim_{y \downarrow 0} \frac{1}{\pi} \int_{x_0}^{x_1} \operatorname{Im} f(-x + iy) \, \mathrm{d} x$ If f(1) = 0, then

$$f(x) = a_f(1-x) + \int_0^\infty \left(\frac{1}{t+x} - \frac{1}{t+1}\right) d\mu_f(t)$$

## $S_f(\rho \| \sigma) = \operatorname{Tr}\{(f(\Delta_{\sigma,\rho}) - f(\Delta_{\rho,\rho}))\rho\}$

 $S_f(\rho \| \sigma) = \operatorname{Tr}\{(f(\Delta_{\sigma,\rho}) - f(\Delta_{\rho,\rho}))\rho\}$ = 0

$$\begin{split} S_f(\rho \| \sigma) &= \operatorname{Tr}\{(f(\Delta_{\sigma,\rho}) - f(\Delta_{\rho,\rho}))\rho\} & \text{using integral representation} \\ &= \int_0^\infty d\mu_f(t) \,\operatorname{Tr}\{\left((tI + \Delta_{\sigma,\rho})^{-1} - (tI + \Delta_{\rho,\rho})^{-1}\right)\rho\} \end{split}$$

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using  $A^{-1} - B^{-1} = A^{-1}(B - A)B^{-1}$  and  $\Delta_{A,B} = L_{A}R_{B^{-1}}$   
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$$= \int_0^\infty d\mu_f(t) \operatorname{Tr}\left\{ \left( (tI + \Delta_{\sigma,\rho})^{-1} (L_\rho - L_\sigma) (tI + \Delta_{\rho,\rho})^{-1} \right) (I) \right\}$$

$$= \int_0^\infty d\mu_f(t) \ (t+1)^{-1} \text{Tr}\{D_t(\rho-\sigma)\}$$

where  $D_t = \sum_{jk} \left( t + \frac{\mu_k}{\lambda_j} \right)^{-1} \langle \psi_j || \phi_k \rangle |\psi_j \rangle \langle \phi_k |$ , using spectral decomposition of  $\rho$ ,  $\sigma$ 

$$\rho = \sum_{j} \lambda_{j} |\psi_{j}\rangle \langle \psi_{j}|, \quad \sigma = \sum_{k} \mu_{k} |\phi_{k}\rangle \langle \phi_{k}|$$

Lemma For orthogonal bases  $\{|\psi_j\rangle\}$  and  $\{|\phi_k\rangle\}$ , let

 $D = \sum_{kj} C_{kj} \langle \psi_j || \phi_k \rangle |\psi_j \rangle \langle \phi_k |$ 

such that  $0 \le C_{kj} \le C$  for all k, j and some C. Consider two cases:

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$$S_f(\rho \| \sigma) \le \| \rho - \sigma \|_1 \int_0^\infty \frac{1}{t + \lambda_\rho^{-1} \alpha_\sigma} \cdot \frac{1}{t+1} d\mu_f(t) ,$$

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How much entanglement does a state have?

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A relative entropy of entanglement of a state  $\rho_{AB}$  is given by

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[Vedral, Plenio '98, +Rippin, Knight '97] A distance should satisfy

- $D(\rho, \sigma) = 0$  if and only if  $\rho = \sigma$
- D is invariant under unitary operations:  $D(U\rho U^*, U\sigma U^*) = D(\rho, \sigma)$
- Data Processing Inequality:  $D(\mathcal{N}(\rho), \mathcal{N}(\sigma)) \leq D(\rho, \sigma)$  for any CPTP map  $\mathcal{N}$

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Theorem (V. '20)

For a maximally entangled state  $|\Psi^+\rangle = \sum_j \frac{1}{\sqrt{d}} |jj\rangle$ , the quasi-relative entropy of entanglement is reached for a state  $\sigma_+ = \sum_j \frac{1}{d} |jj\rangle\langle jj|$ , and becomes

$$E_f(|\Psi^+\rangle\langle\Psi^+|) = f(1/d)$$

Theorem

For a pure state  $|\Psi\rangle = \sum_{j} \sqrt{p_j} |jj\rangle$ , the quasi-relative entropy of entanglement for a class of functions, including  $-\log(x)$ , is reached for a state  $\sigma = \sum_{j} p_j |jj\rangle\langle jj|$ , and becomes

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#### Theorem

For a pure 2-qubit state  $|\Psi\rangle = \sqrt{p}|00\rangle + \sqrt{1-p}|11\rangle$ , the quasi-relative entropy of entanglement for a class of functions, including  $1 - x^{1-\alpha}$ , is reached for a state in the form  $\sigma = q|00\rangle\langle00| + (1-q)|11\rangle\langle11|$ , and becomes

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 $f(x) = -\log x$  gives relative entropy  $S_{-\log}(\rho \| \sigma) = \operatorname{Tr}(\rho \log \rho - \rho \log \sigma)$ 

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Then, for  $0 < x \leq 1$ , since  $S_f$  is jointly convex

$$F(x,\sigma) = S_f(\rho || (1-x)\sigma^* + x\sigma) \le (1-x)S_f(\rho || \sigma^*) + xS_f(\rho || \sigma) = (1-x)F(0,\sigma) + xF(1,\sigma)$$

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This implies that

$$\frac{1}{x}[F(x,\sigma) - F(0,\sigma)] \le F(1,\sigma) - F(0,\sigma) < 0$$

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The goal is to prove that for any separable state  $\frac{\partial F}{\partial x}(0,\sigma) \ge 0$ 

To show that  $\sigma^*$  is the closest, suppose that there is a separable state  $\sigma$  s.t.

$$S_f(\rho \| \sigma) < S_f(\rho \| \sigma^*)$$

Then, for  $0 < x \leq 1$ , since  $S_f$  is jointly convex

$$F(x,\sigma) = S_f(\rho \| (1-x)\sigma^* + x\sigma) \le (1-x)S_f(\rho \| \sigma^*) + xS_f(\rho \| \sigma) = (1-x)F(0,\sigma) \neq xF(1,\sigma)$$

This implies that

impossible due to

$$\frac{1}{x}[F(x,\sigma) - F(0,\sigma)] \le F(1,\sigma) - F(0,\sigma) < 0$$

Constrain on the function

$$H_f(p) = \int_0^\infty p(t+p)^{-2} d\mu_f(t)$$

is either 1) constant; or 2) monotonically increasing or decreasing

Is this true?

For a pure state  $|\Psi\rangle = \sum_{j} \sqrt{p_j} |jj\rangle$ , the quasi-relative entropy of entanglement for ANY function is reached for a state in the form  $\sigma = \sum_{j} q_j |jj\rangle \langle jj|$ , and becomes

$$E_f(|\Psi\rangle\langle\Psi|) = \sum_j p_j f(q_j)$$

Thank you!