

Bounds on mean vertical heat transport in convection driven by internal heating

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Joint work with:

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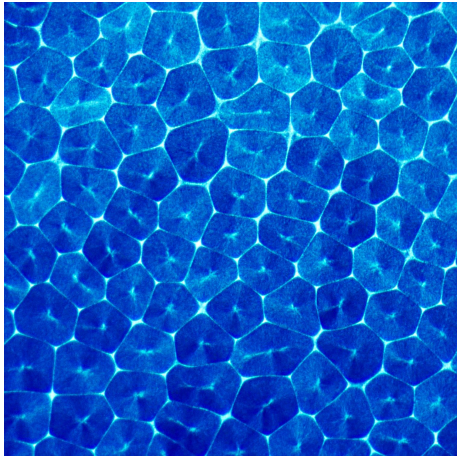
IPAM Worskshop

“Transport and Mixing in Complex and Turbulent Flows”

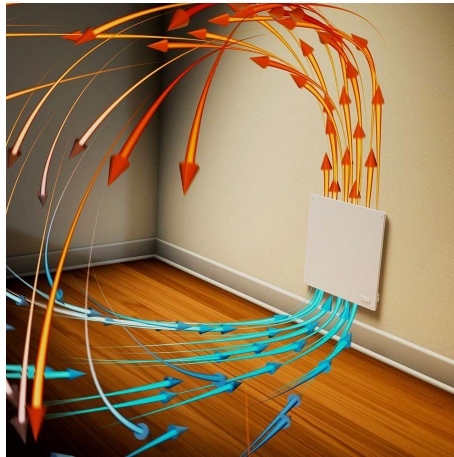
11 January 2021

Convection is important!

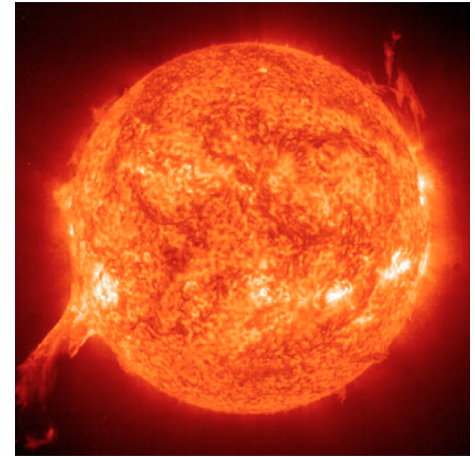
Drying paint



Central heating

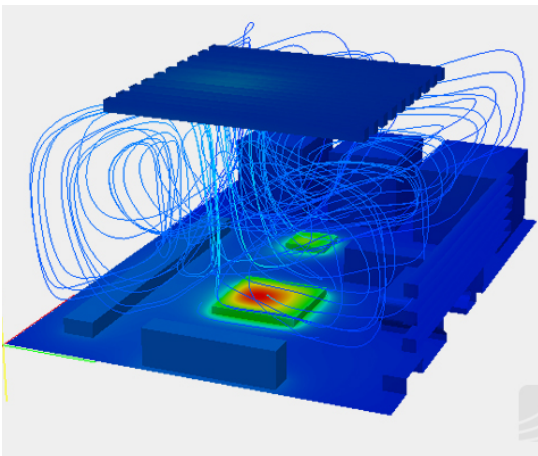


Stars



scale →

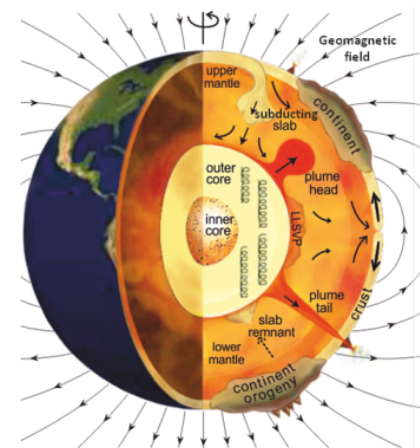
Chip cooling



Soup

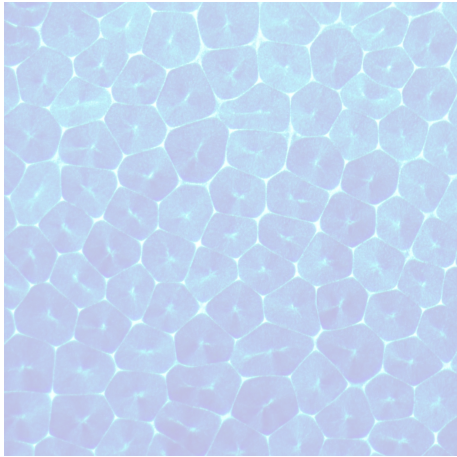


Mantle dynamics

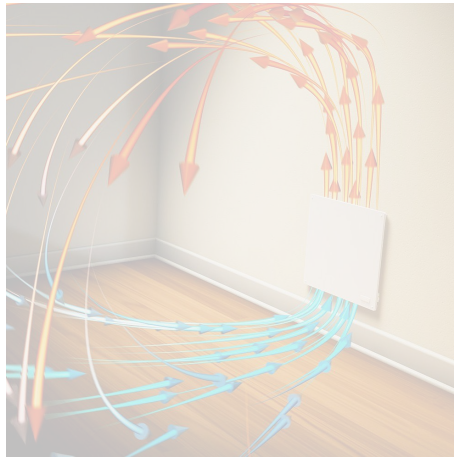


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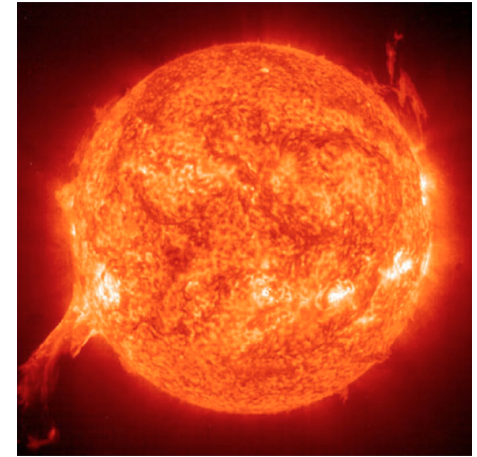
Drying paint



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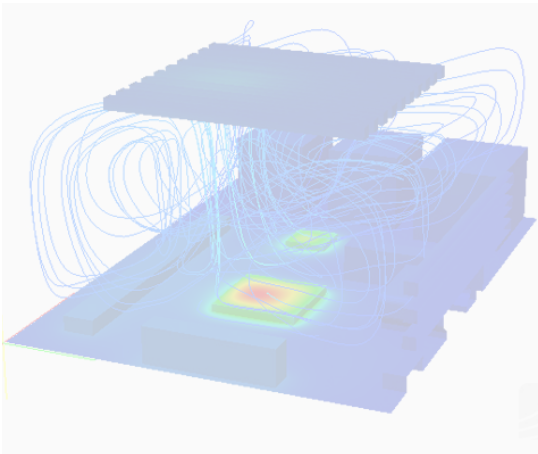


Stars



scale →

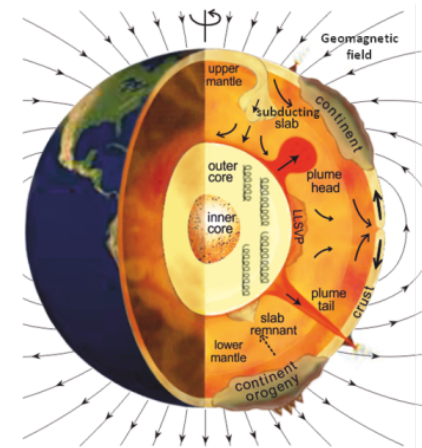
Chip cooling



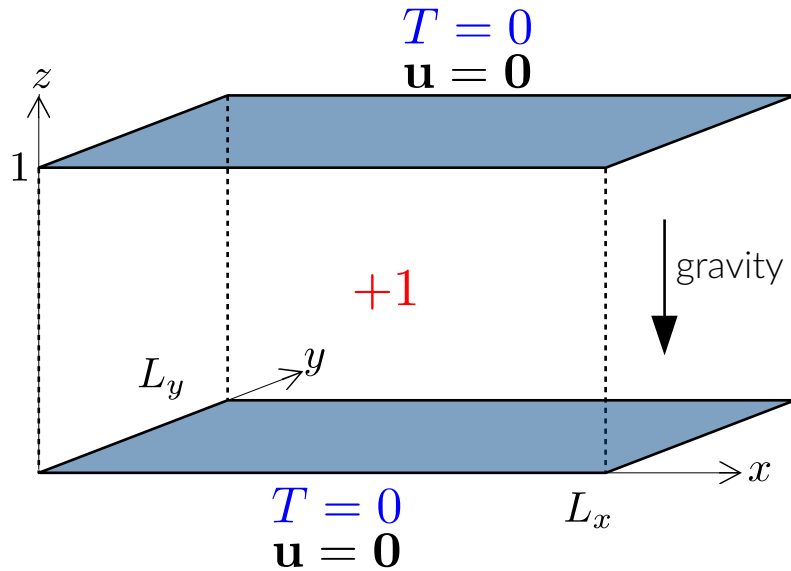
Soup



Mantle dynamics



An idealized nondimensional model

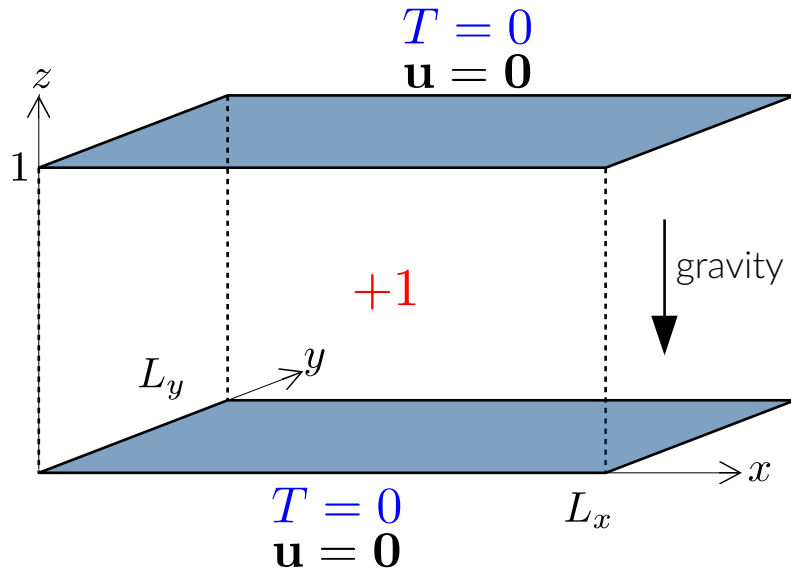


$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p = Pr(\nabla^2 \mathbf{u} + RT\hat{z})$$

$$\nabla \cdot \mathbf{u} = 0$$

$$\partial_t T + \mathbf{u} \cdot \nabla T = \nabla^2 T + 1$$

An idealized nondimensional model

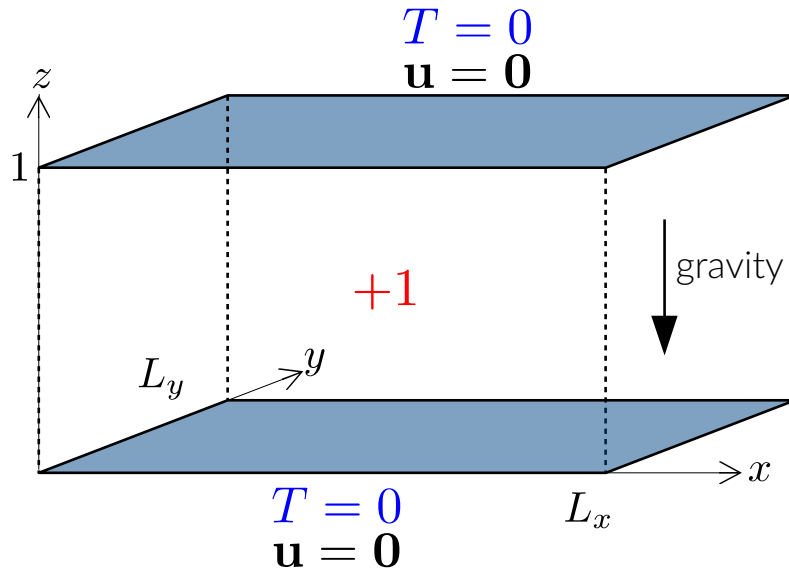


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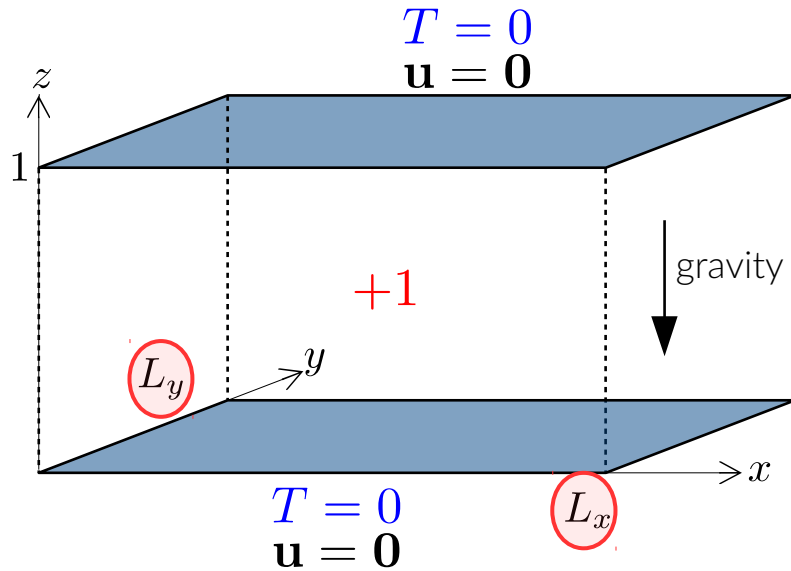


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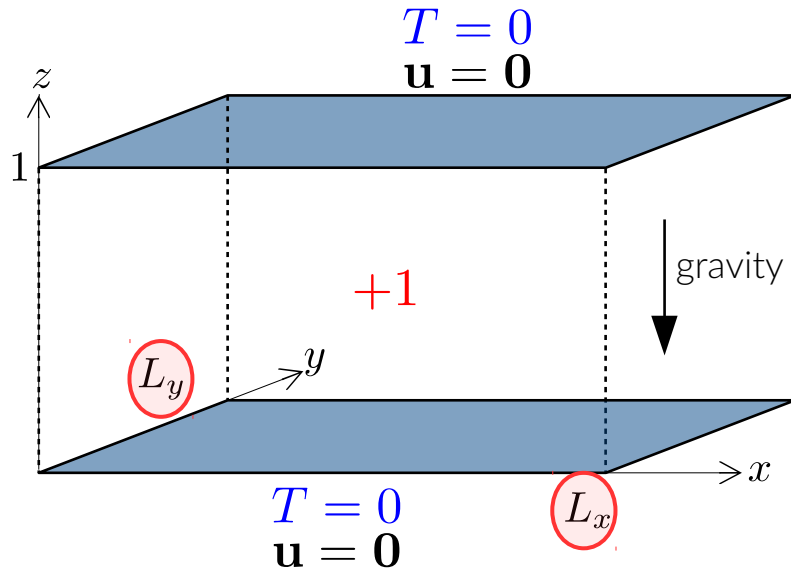


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An idealized nondimensional model



$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p = \textcircled{Pr} (\nabla^2 \mathbf{u} + \textcircled{R} T \hat{z})$$

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$$\partial_t T + \mathbf{u} \cdot \nabla T = \nabla^2 T + 1$$

Mean vertical heat flux:

$$\langle wT \rangle = \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_{\Omega} wT \, d\mathbf{x} \, dt$$

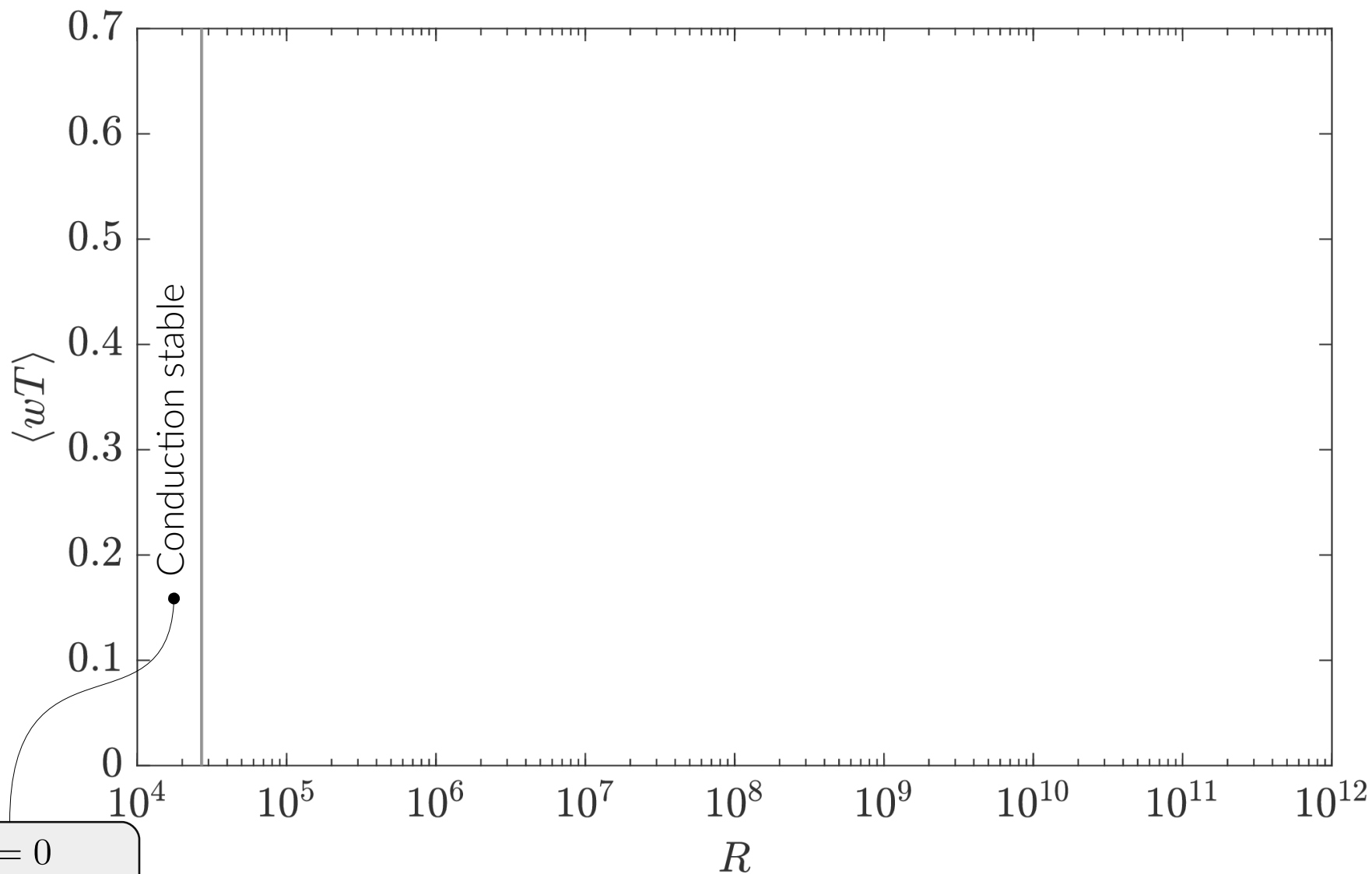
Heat flux through top and bottom boundaries:

$$\mathcal{F}_T = -\overline{\partial_z T}|_{z=1} = \frac{1}{2} + \langle wT \rangle$$

$$\mathcal{F}_B = \overline{\partial_z T}|_{z=0} = \frac{1}{2} - \langle wT \rangle$$

Q: Variation with heating rate (R)?

What do we know?

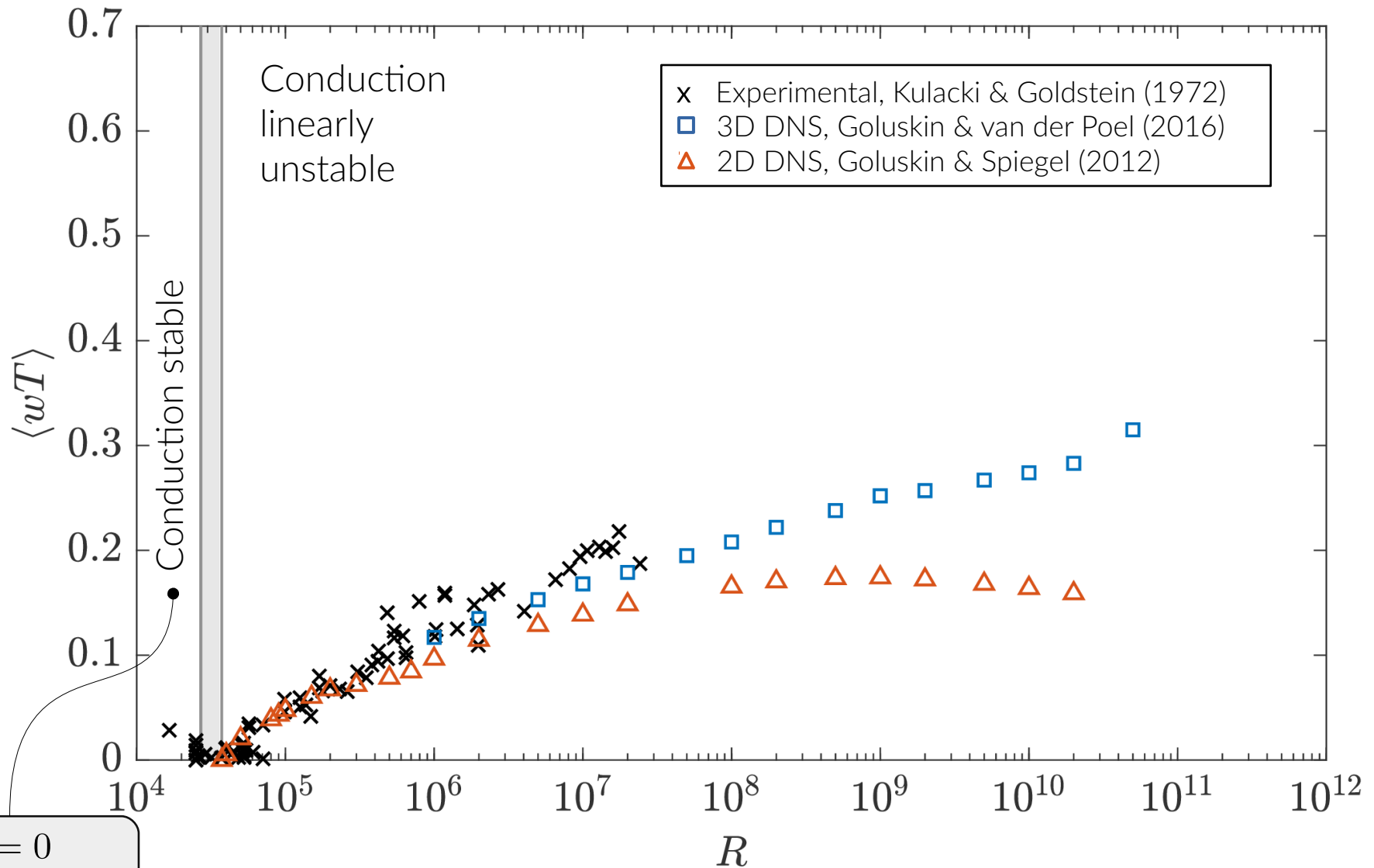


$$\mathbf{u} = 0$$

$$T = \frac{1}{2}z(1 - z)$$

$$\langle wT \rangle = 0$$

What do we know?

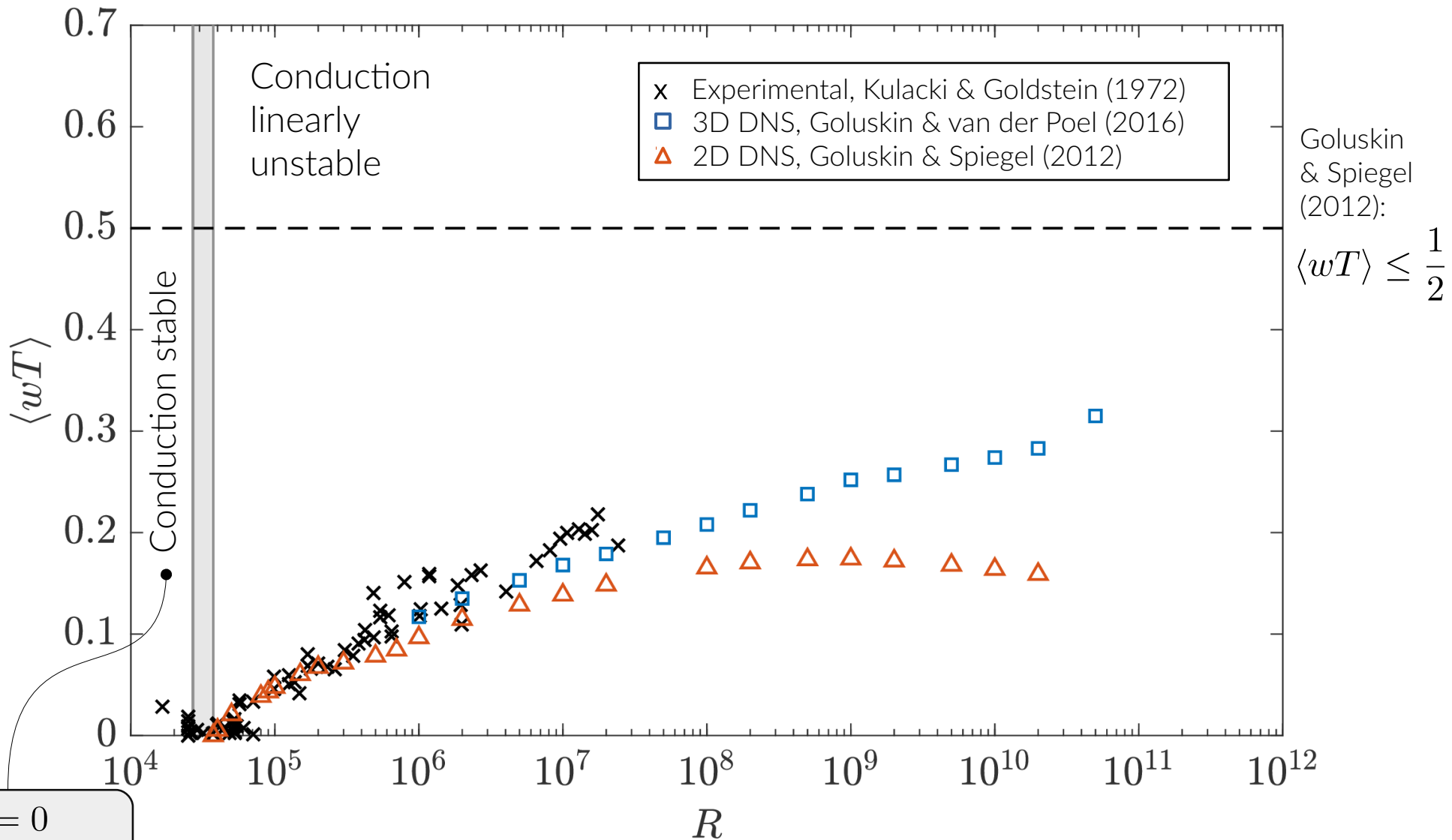


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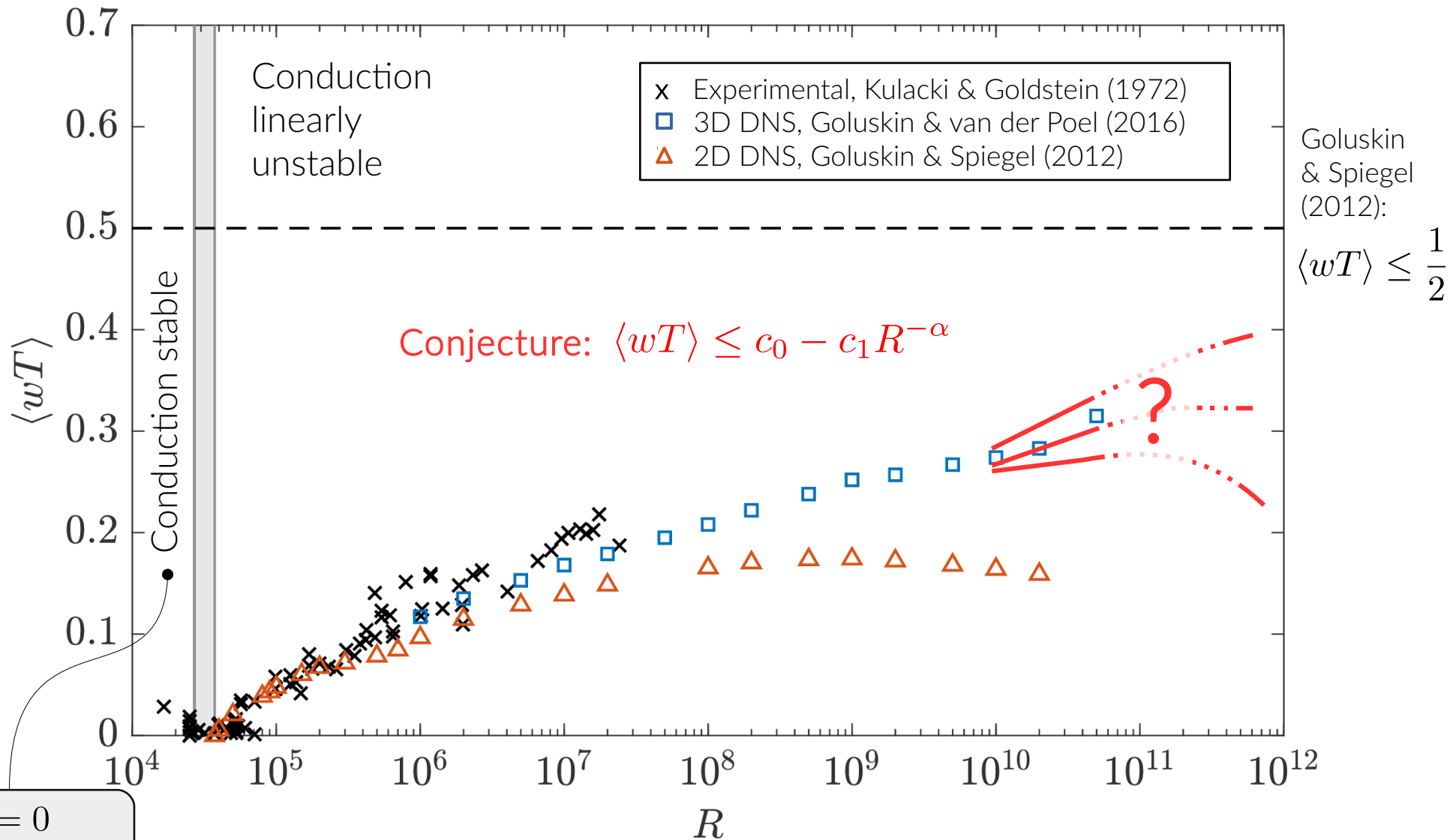


$$\mathbf{u} = 0$$

$$T = \frac{1}{2}z(1 - z)$$

$$\langle wT \rangle = 0$$

What do we know?



$$\mathbf{u} = 0$$

$$T = \frac{1}{2}z(1 - z)$$

$$\langle wT \rangle = 0$$

Upper bounds on heat transfer

- Infinite-time averages of time derivatives of bounded functionals vanish!

$$\begin{aligned} \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \frac{d}{dt} \mathcal{V}\{\mathbf{u}(\cdot, t), T(\cdot, t)\} dt &= 0 \\ &= \frac{\mathcal{V}\{\mathbf{u}(\cdot, \tau), T(\cdot, \tau)\} - \mathcal{V}\{\mathbf{u}(\cdot, 0), T(\cdot, 0)\}}{\tau} \end{aligned}$$

Upper bounds on heat transfer

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 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_{\Omega} wT + \frac{\delta \mathcal{V}}{\delta \mathbf{u}} \cdot \partial_t \mathbf{u} + \frac{\delta \mathcal{V}}{\delta T} \cdot \partial_t T \, d\mathbf{x} \, dt
 \end{aligned}$$

Upper bounds on heat transfer

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 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_\Omega wT + \frac{\delta \mathcal{V}}{\delta \mathbf{u}} \cdot \partial_t \mathbf{u} + \frac{\delta \mathcal{V}}{\delta T} \cdot \partial_t T \, d\mathbf{x} \, dt \\
 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_\Omega wT + \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x} \, dt
 \end{aligned}$$

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 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_\Omega wT + \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x} \, dt + U - U
 \end{aligned}$$

Upper bounds on heat transfer

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 \langle wT \rangle &= \langle wT \rangle + \overbrace{\limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \frac{d}{dt} \mathcal{V}\{\mathbf{u}(\cdot, t), T(\cdot, t)\} dt}^{=0} \\
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 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \oint_{\Omega} wT + \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x} \, dt + \mathcal{U} - \mathcal{U} \\
 &= \mathcal{U} - \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \oint_{\Omega} \mathcal{U} - wT - \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x} \, dt
 \end{aligned}$$

Upper bounds on heat transfer

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 &= \mathcal{U} - \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \underbrace{\int_\Omega \mathcal{U} - wT - \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x}}_{\mathcal{S}\{\mathbf{u}, T\}} dt
 \end{aligned}$$

Upper bounds on heat transfer

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 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_{\Omega} wT + \frac{\delta \mathcal{V}}{\delta \mathbf{u}} \cdot \partial_t \mathbf{u} + \frac{\delta \mathcal{V}}{\delta T} \cdot \partial_t T \, d\mathbf{x} \, dt \\
 &= \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \int_{\Omega} wT + \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x} \, dt + U - U \\
 &= U - \limsup_{\tau \rightarrow \infty} \frac{1}{\tau} \int_0^\tau \underbrace{\int_{\Omega} U - wT - \mathcal{D}(\mathbf{u}, T) \, d\mathbf{x}}_{\mathcal{S}\{\mathbf{u}, T\}} dt
 \end{aligned}$$

If $\mathcal{S}\{\mathbf{u}, T\} \geq 0$ for all $\mathbf{u}(\mathbf{x})$ and $T(\mathbf{x})$ consistent with physical constraints, then $\langle wT \rangle \leq U$

Reinterpreting the background method

$$\mathcal{V}\{\mathbf{u}, T\} = \int_{\Omega} \frac{a}{2PrR} |\mathbf{u}|^2 + \frac{b}{2} |T|^2 - [1 - z + \psi(z)] T \, d\mathbf{x}$$

- After some algebra:

$$\begin{aligned} \mathcal{S}\{\mathbf{u}, T\} = & U - \frac{1}{2} + \psi(1) \int_{z=1} \partial_z T \, dx dy - [\psi(0) - 1] \int_{z=0} \partial_z T \, dx dy \\ & + \int_{\Omega} \frac{a}{R} |\nabla \mathbf{u}|^2 + b |\nabla T|^2 - (a - \psi') w T + (bz - \psi' - 1) \partial_z T + \psi \, d\mathbf{x} \end{aligned}$$

- A **linear** optimization problem:

$$\begin{aligned} \langle wT \rangle \leq & \inf_{U, a, b, \psi(z)} U \\ \text{s.t.} \quad & \mathcal{S}\{\mathbf{u}, T\} \geq 0 \quad \forall \mathbf{u}, T : \begin{cases} \text{BCs} \\ \nabla \cdot \mathbf{u} = 0 \end{cases} \end{aligned}$$

Reinterpreting the background method

$$\mathcal{V}\{\mathbf{u}, T\} = \int_{\Omega} \frac{\overset{\text{balance parameters}}{\overset{a}{2PrR}} |\mathbf{u}|^2 + \frac{\overset{b}{2}}{2} |T|^2 - \underbrace{[1 - z + \psi(z)]}_{\text{(rescaled) background temperature field}} T \, d\mathbf{x}$$

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Reinterpreting the background method

$$\mathcal{V}\{\mathbf{u}, T\} = \int_{\Omega} \frac{\overset{\text{balance parameters}}{\color{red}a}}{2PrR} |\mathbf{u}|^2 + \frac{\color{red}b}{2} |T|^2 - \underbrace{[1 - z + \color{red}\psi(z)]}_{\text{(rescaled) background temperature field}} T \, d\mathbf{x}$$

- After some algebra:

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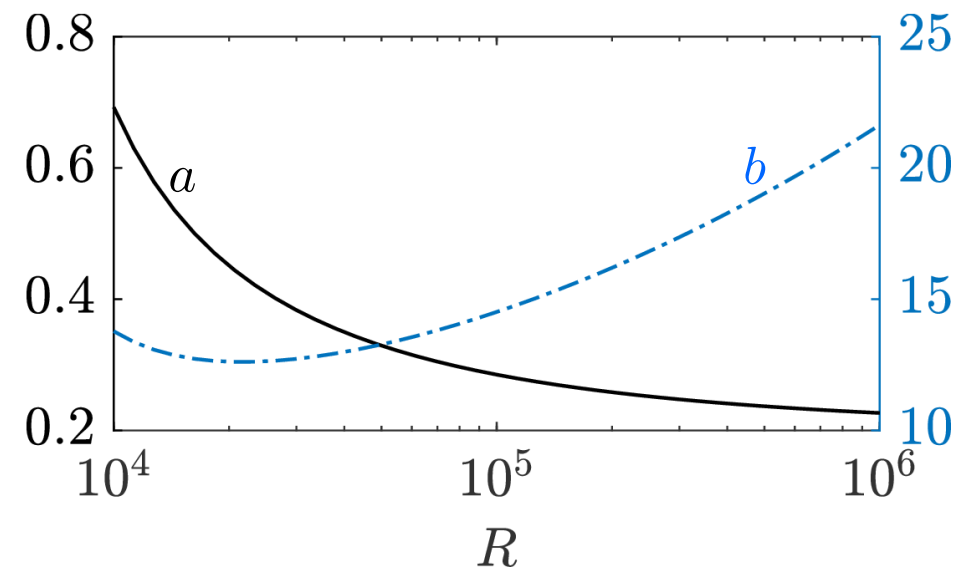
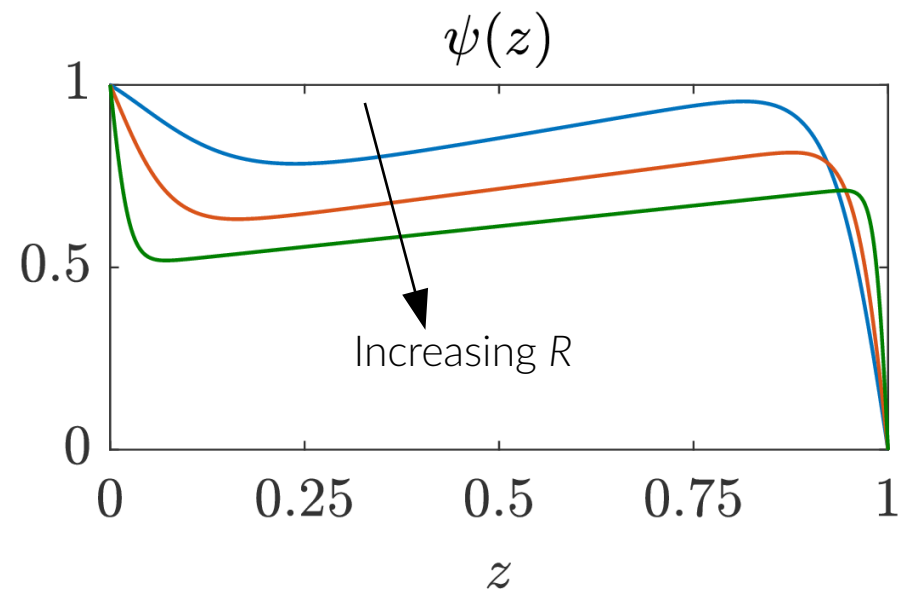
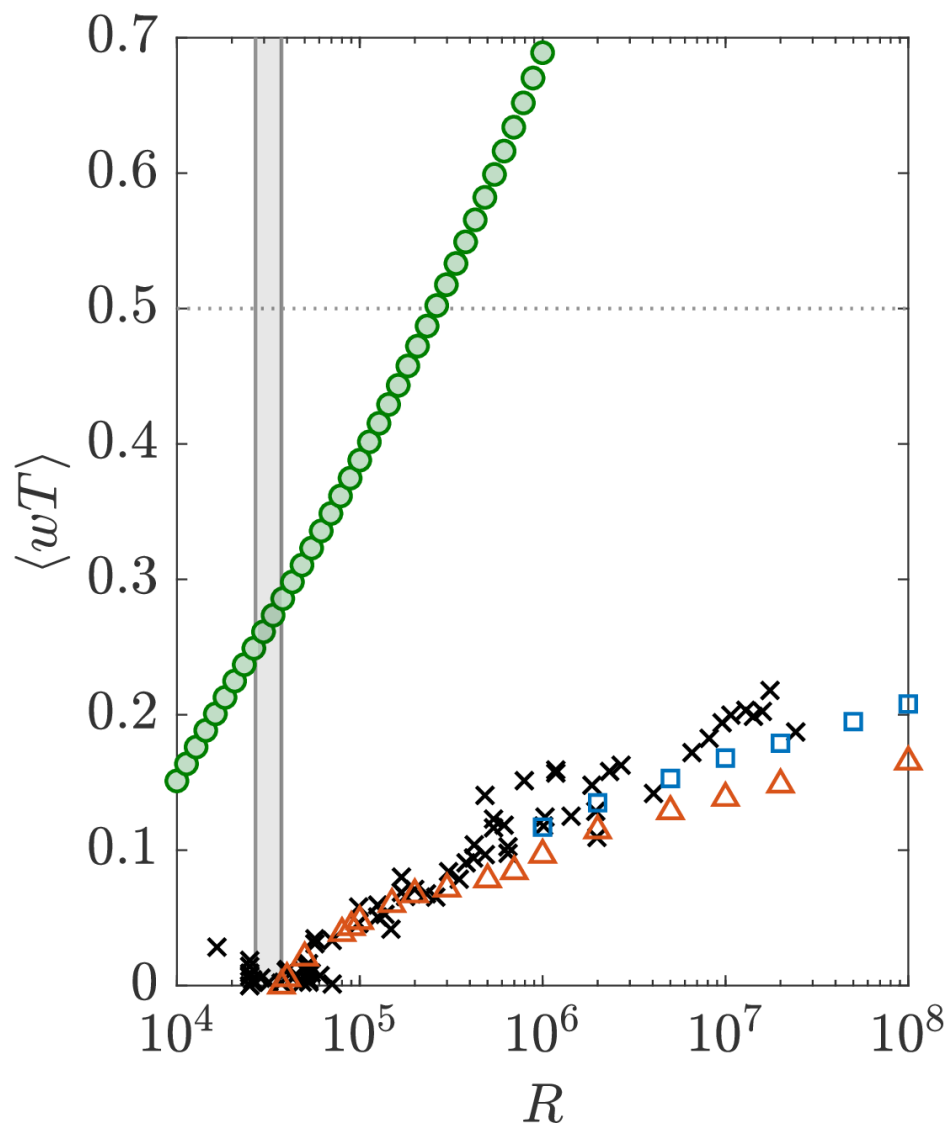
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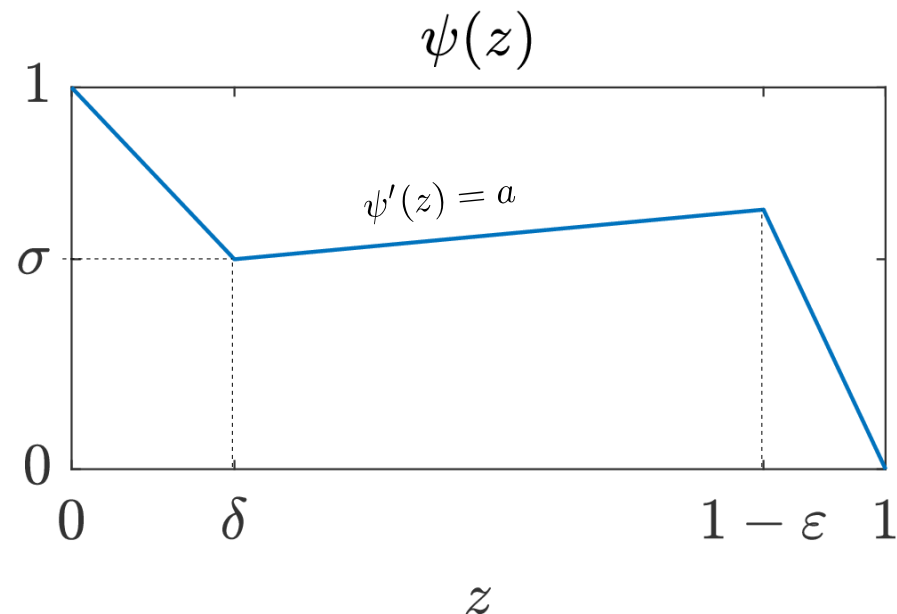
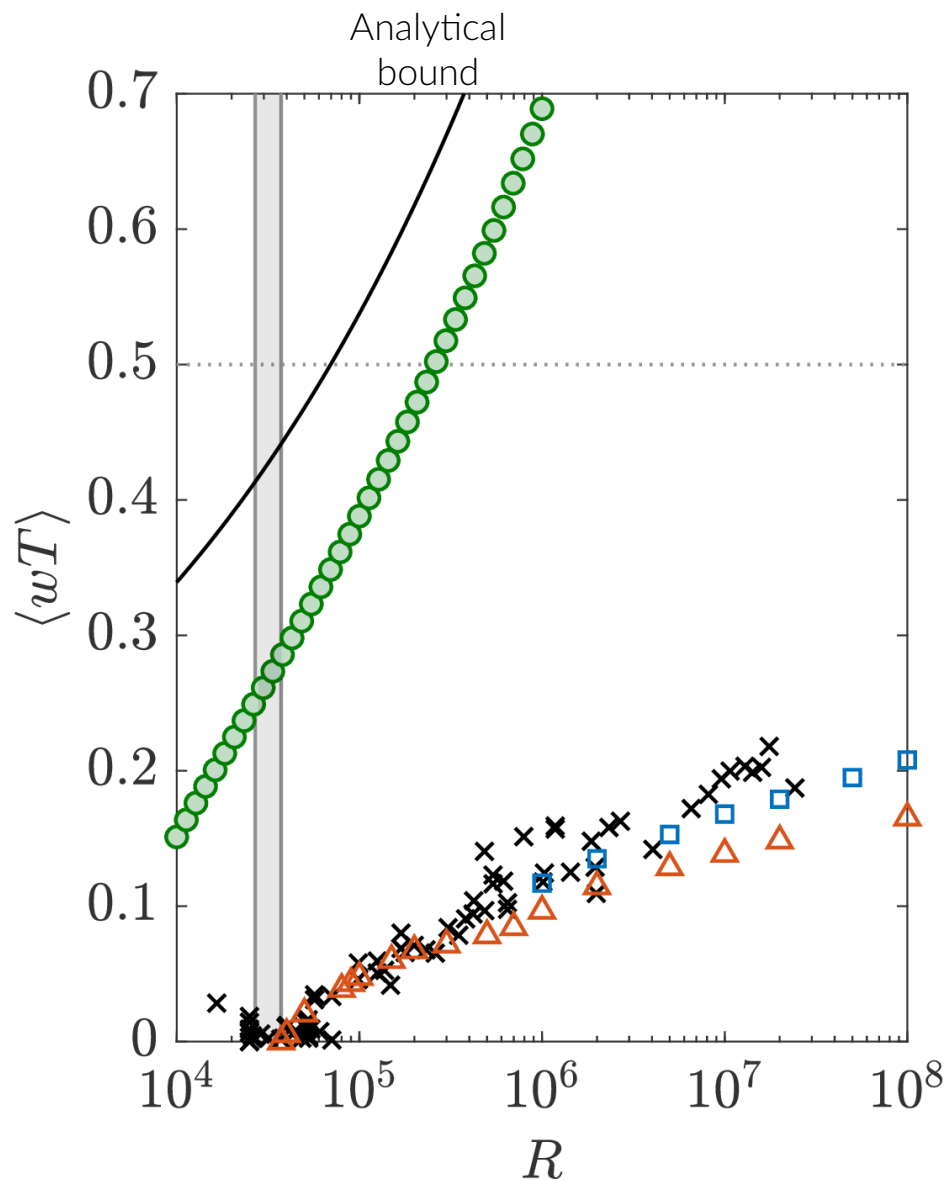
- A **linear** optimization problem:

$$\begin{aligned} \langle wT \rangle &\leq \inf_{U, a, b, \psi(z)} U \\ \text{s.t.} \quad &\inf_{\substack{\mathbf{u}, T: \\ \text{BCs} \\ \nabla \cdot \mathbf{u} = 0}} \mathcal{S}\{\mathbf{u}, T\} \geq 0 \end{aligned}$$

Computational results



Analytical results

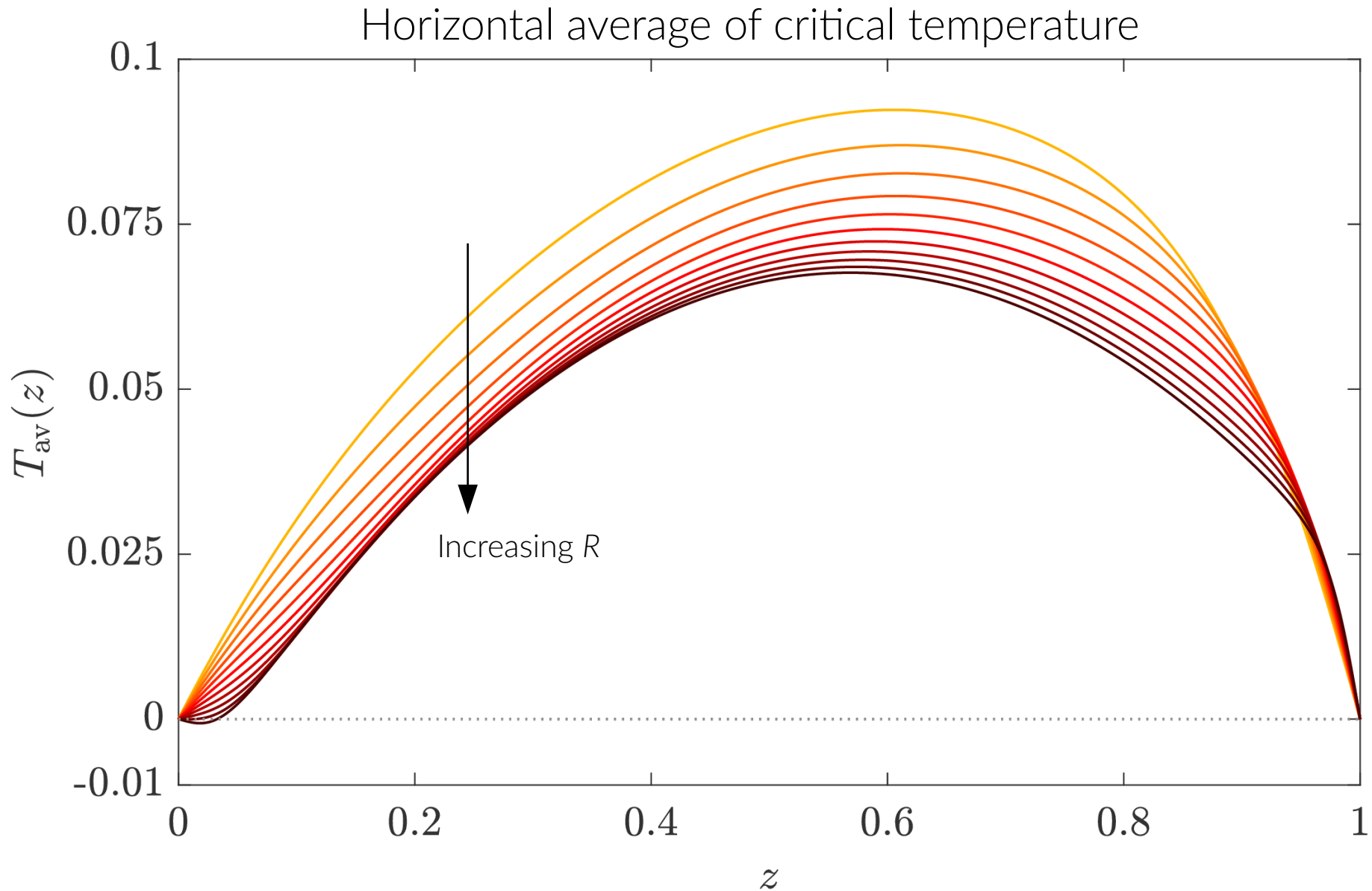


$$\langle wT \rangle \leq \frac{1}{16} [8(3\sqrt{2} - 4)^2]^{\frac{1}{5}} R^{\frac{1}{5}}$$

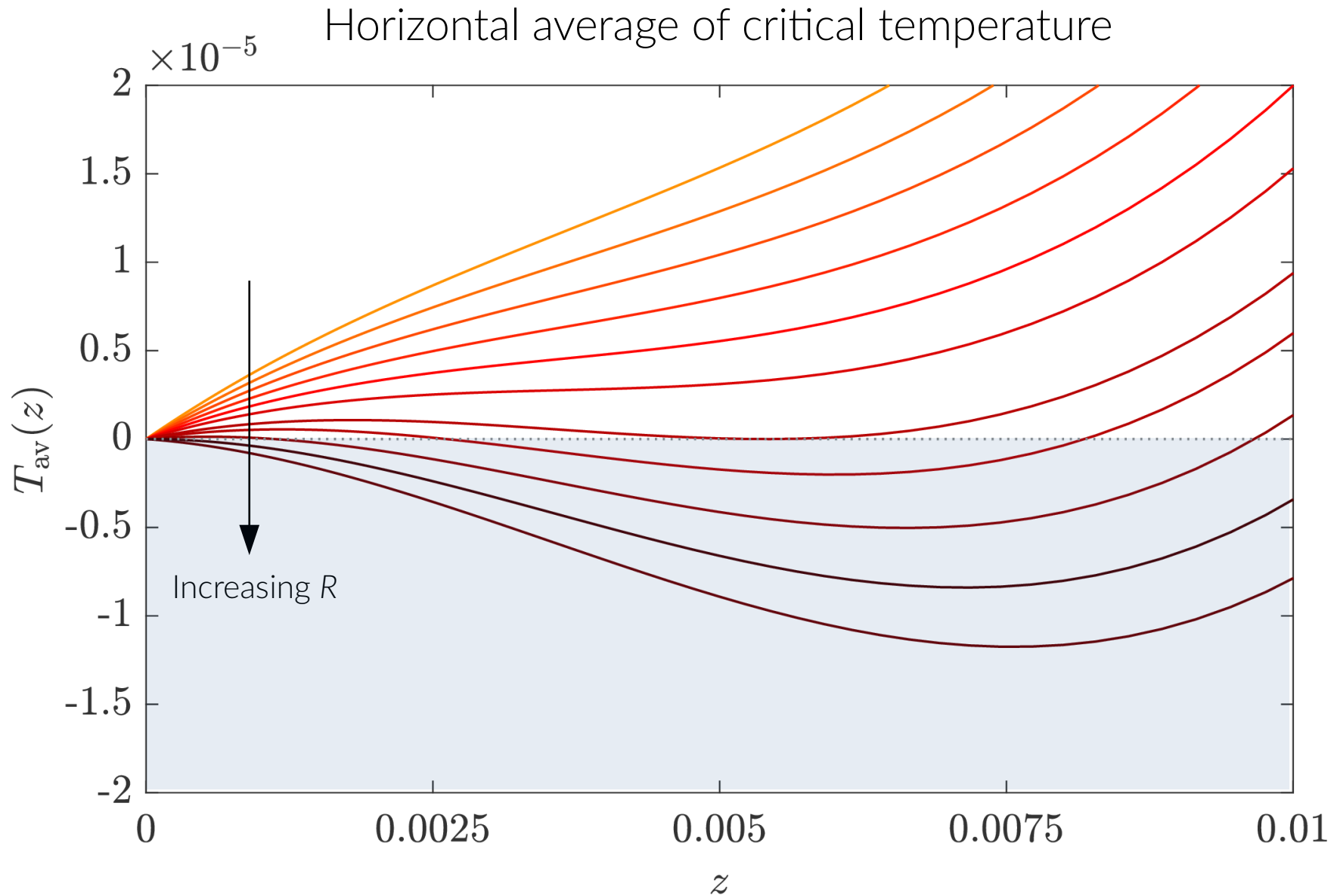
$$\leq \frac{1}{2} \text{ if } R \leq 69572$$

What is missing?

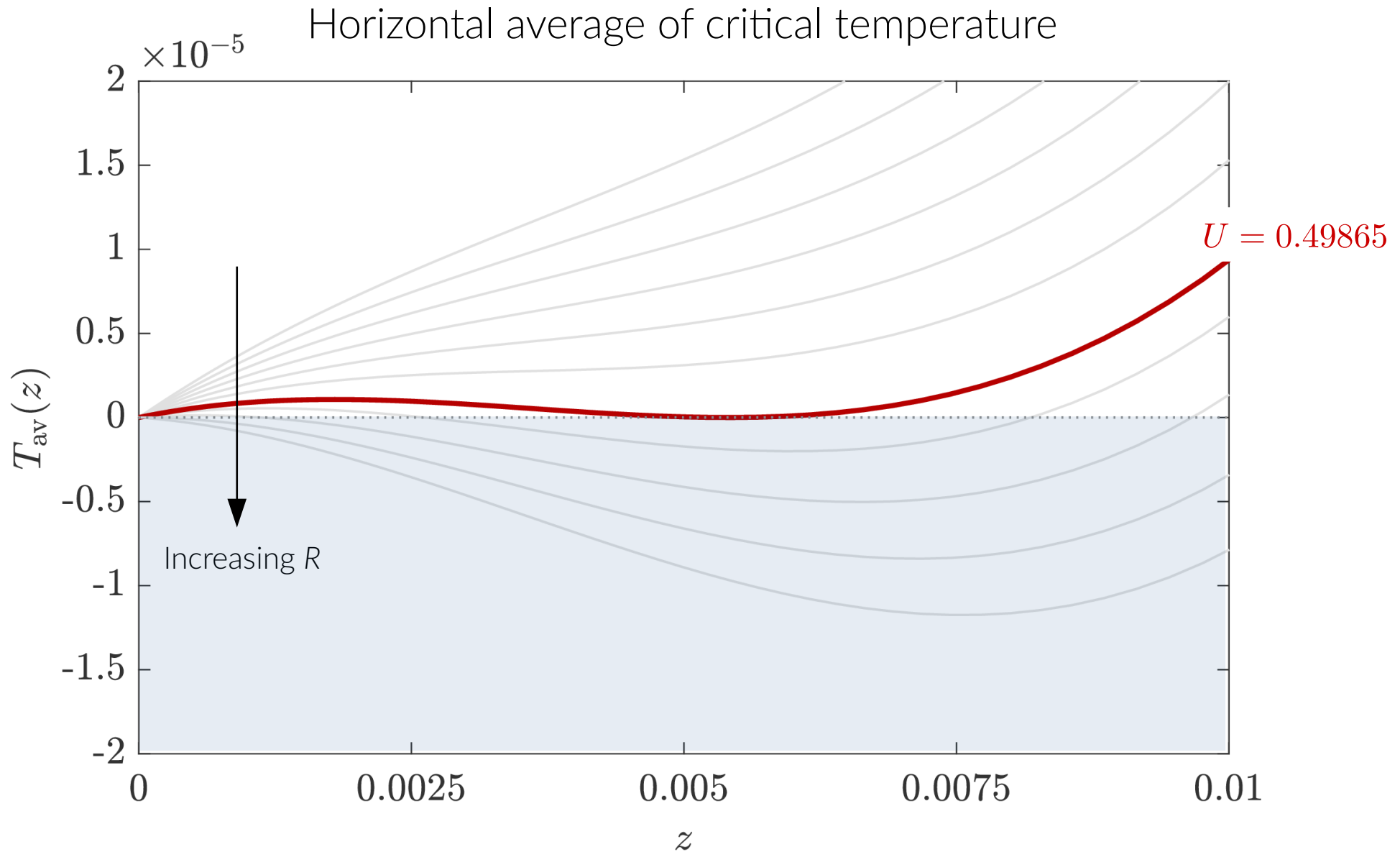
Physically realistic T should be positive!



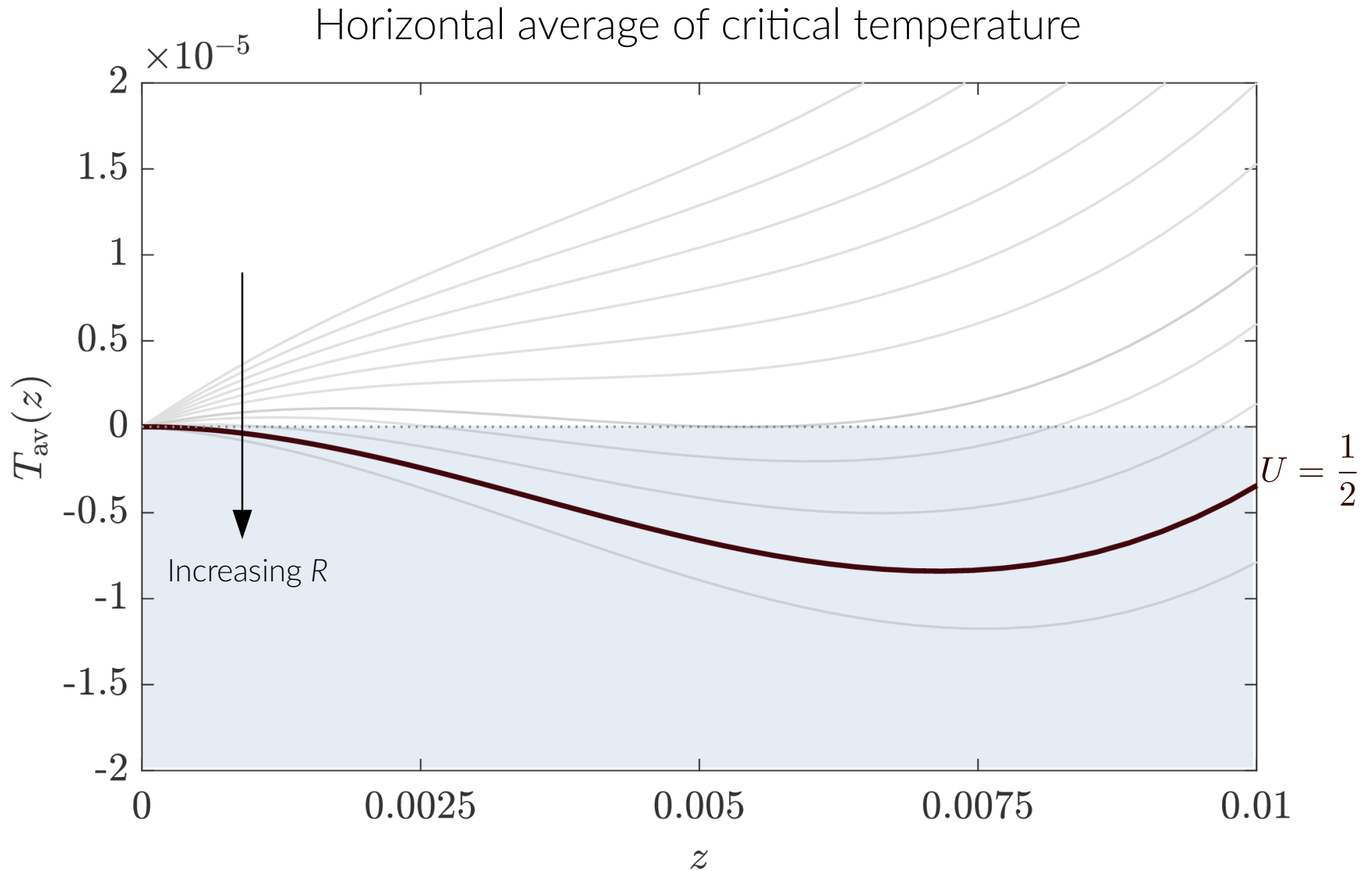
Physically realistic T should be positive!



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Physically realistic T should be positive!



Revised bounding framework

$$\langle wT \rangle \leq \inf_{U, a, b, \psi(z)} U$$

$$\text{s.t.} \quad \mathcal{S}\{\mathbf{u}, T\} \geq 0 \quad \forall \mathbf{u}, T : \begin{cases} \text{BCs} \\ \nabla \cdot \mathbf{u} = 0 \\ T(\mathbf{x}) \geq 0 \text{ on } \Omega \end{cases}$$

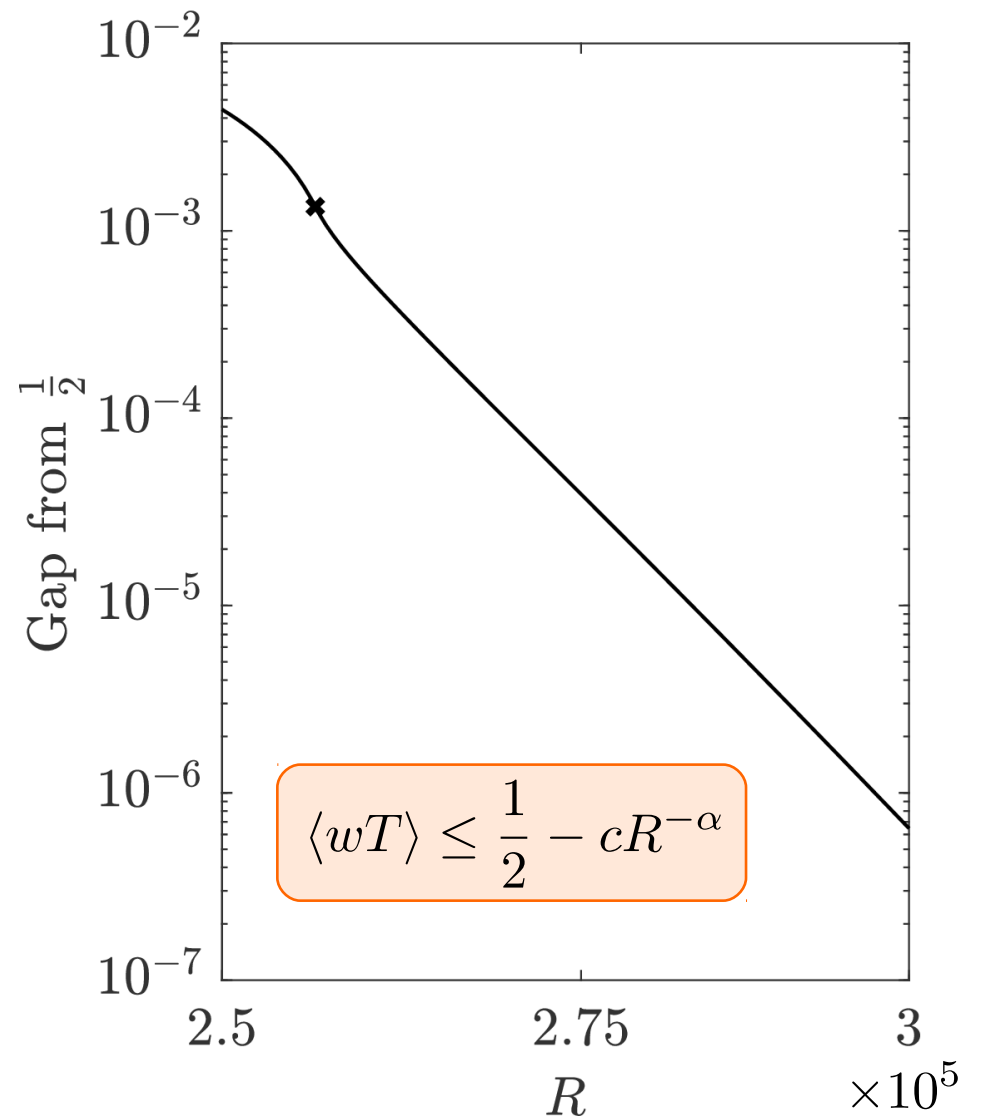
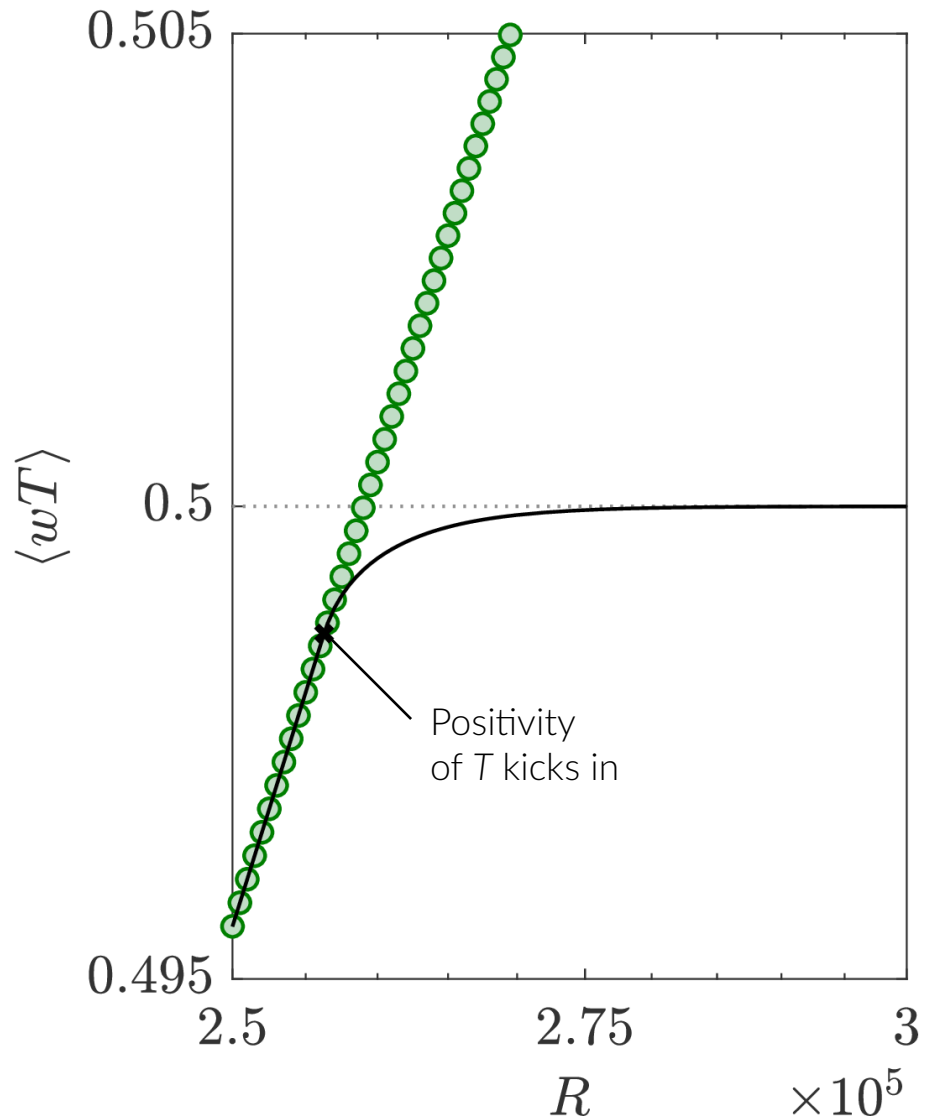
Enforce using a nonnegative Lagrange multiplier:

$$\langle wT \rangle \leq \inf_{U, a, b, \psi(z), q(z)} U$$

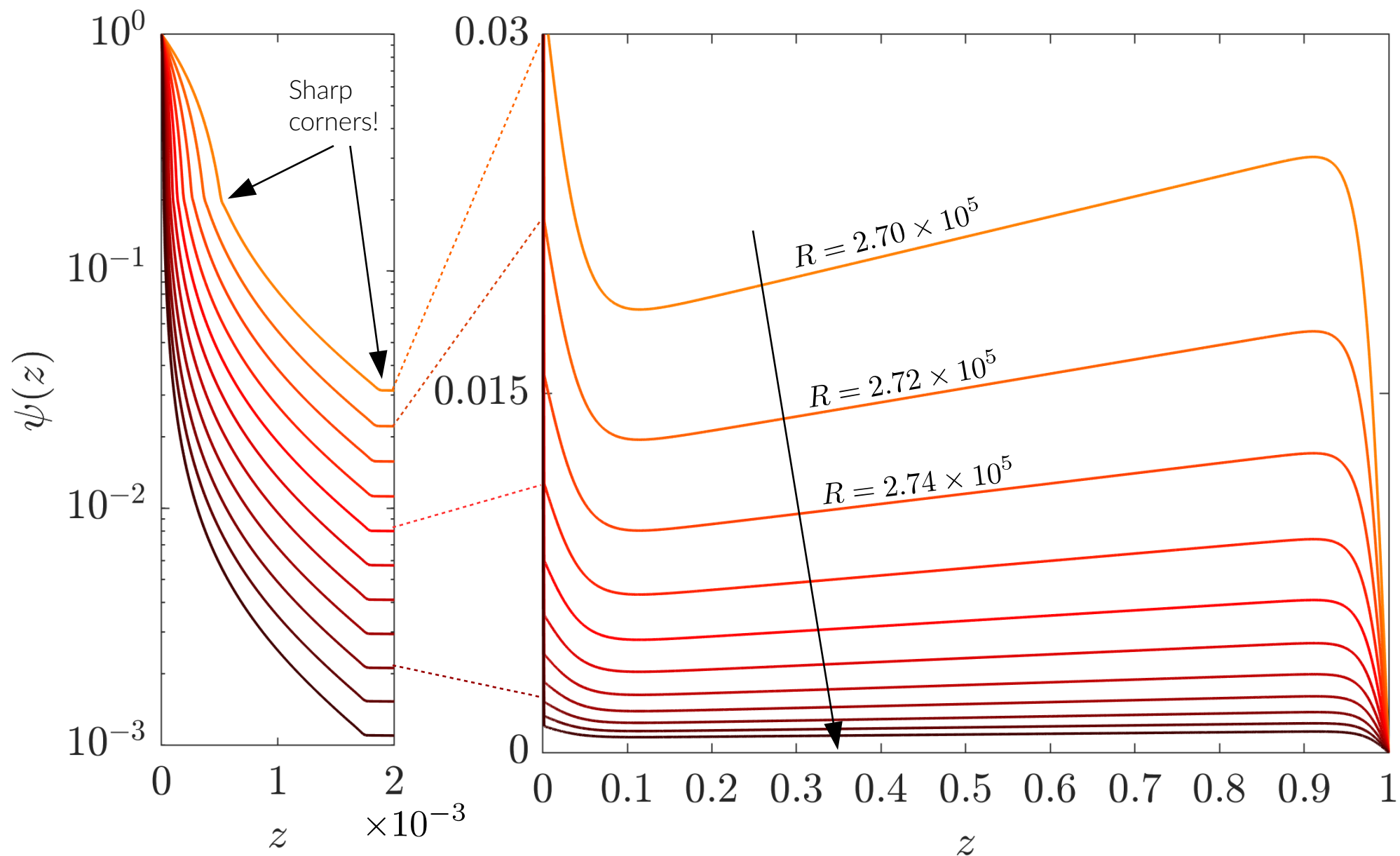
$$\text{s.t.} \quad \mathcal{S}\{\mathbf{u}, T\} \geq \int_{\Omega} q'(z) T \, d\mathbf{x} \quad \forall \mathbf{u}, T : \begin{cases} \text{BCs} \\ \nabla \cdot \mathbf{u} = 0 \end{cases}$$

$$q'(z) \geq 0 \quad \forall z \in [0, 1]$$

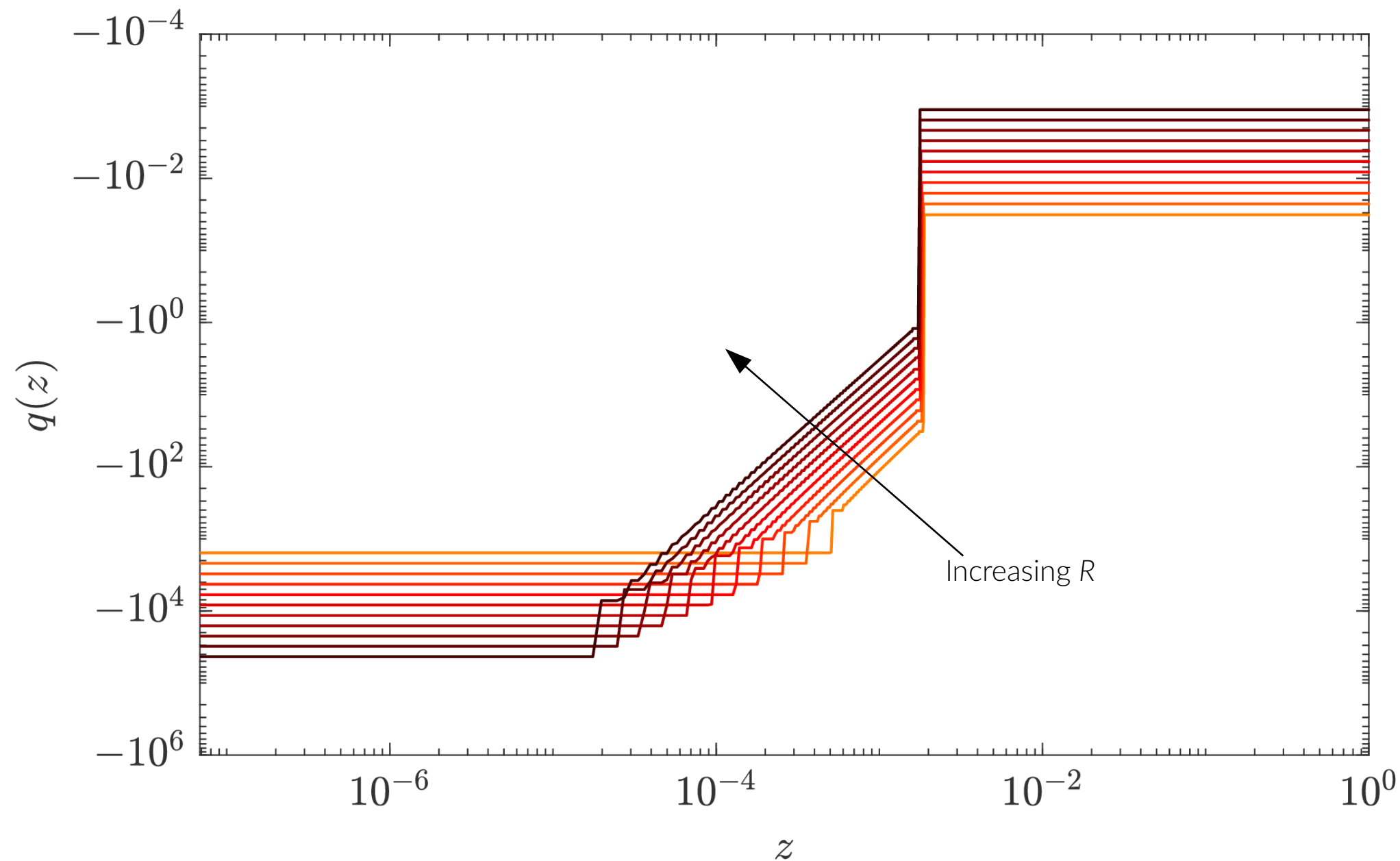
Computational upper bounds



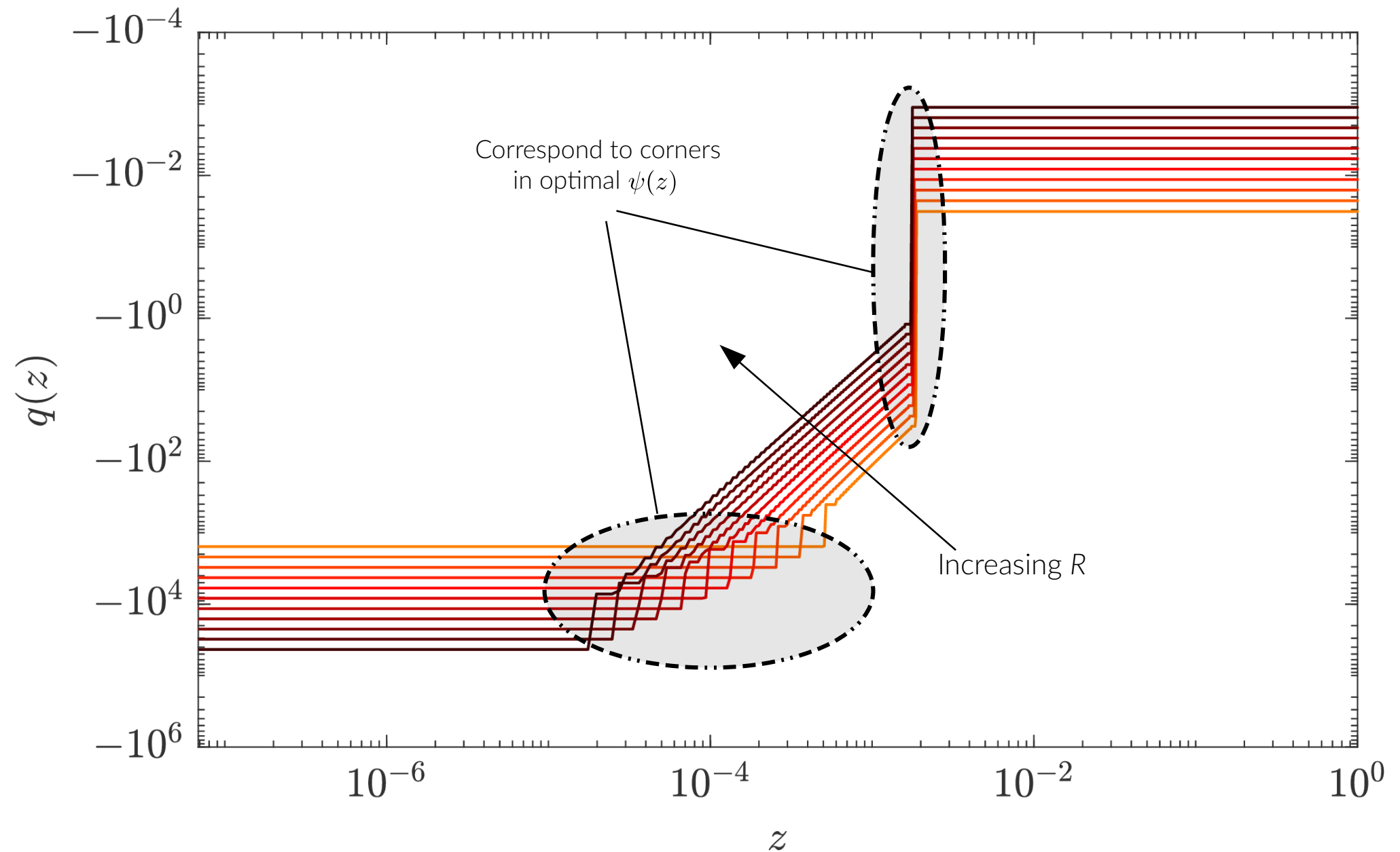
Optimal $\psi(z)$



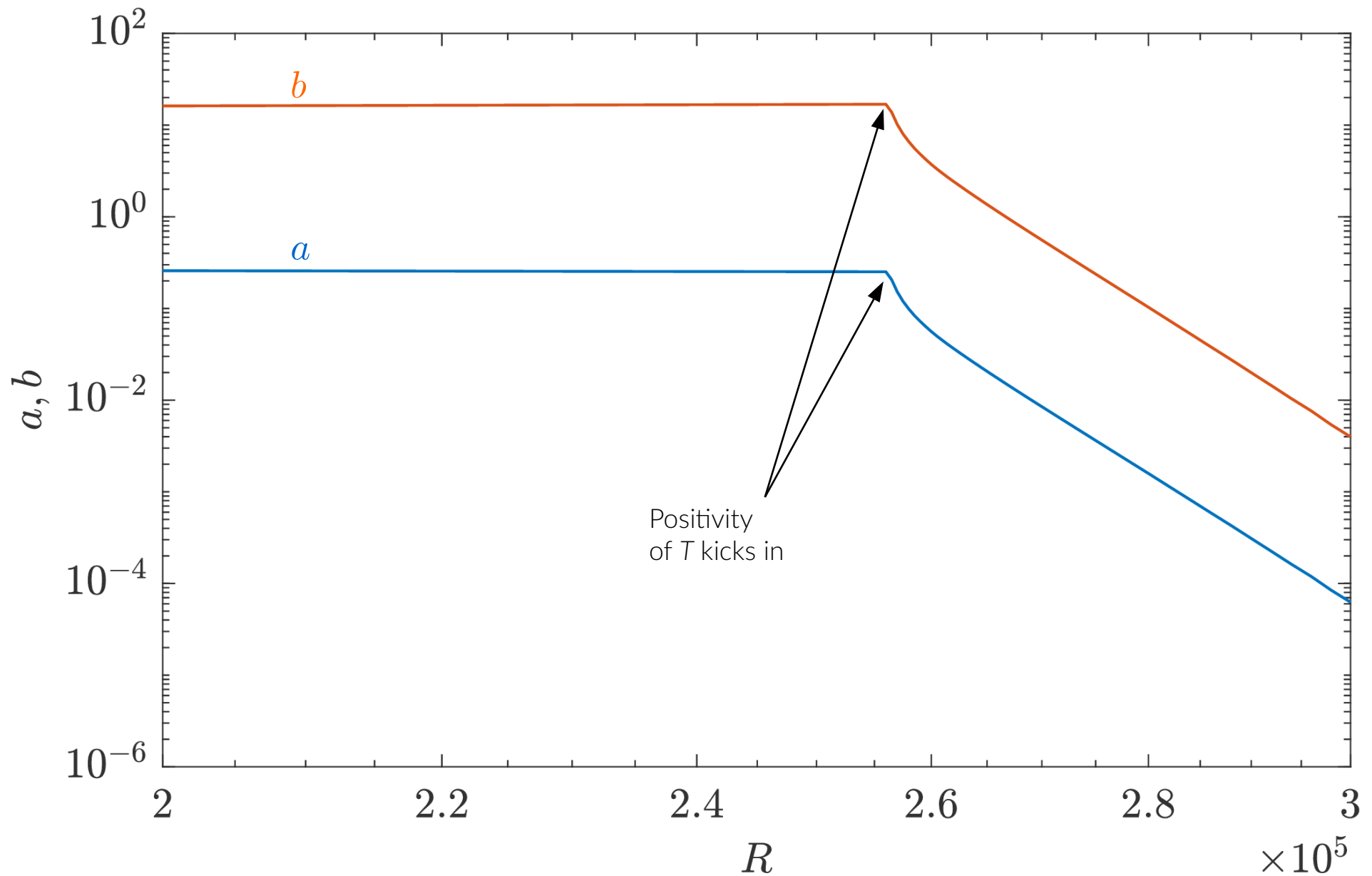
Optimal Lagrange multipliers



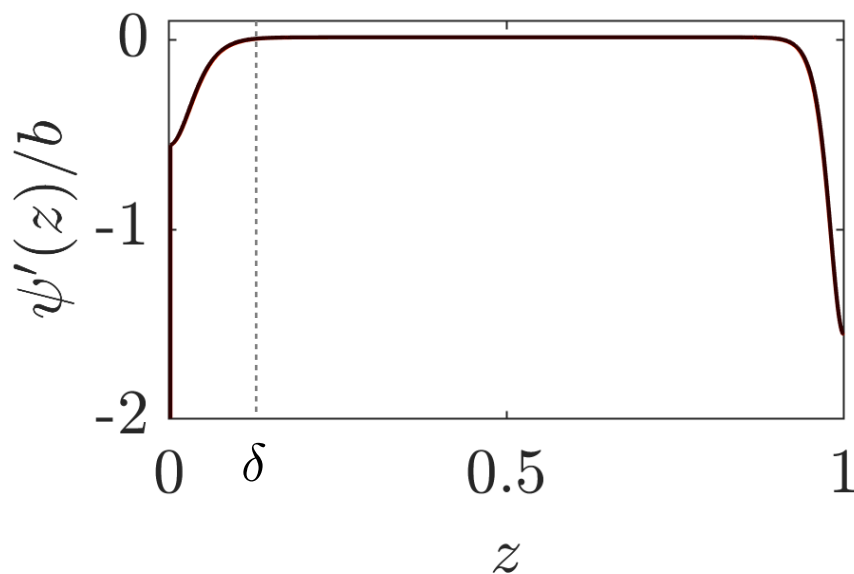
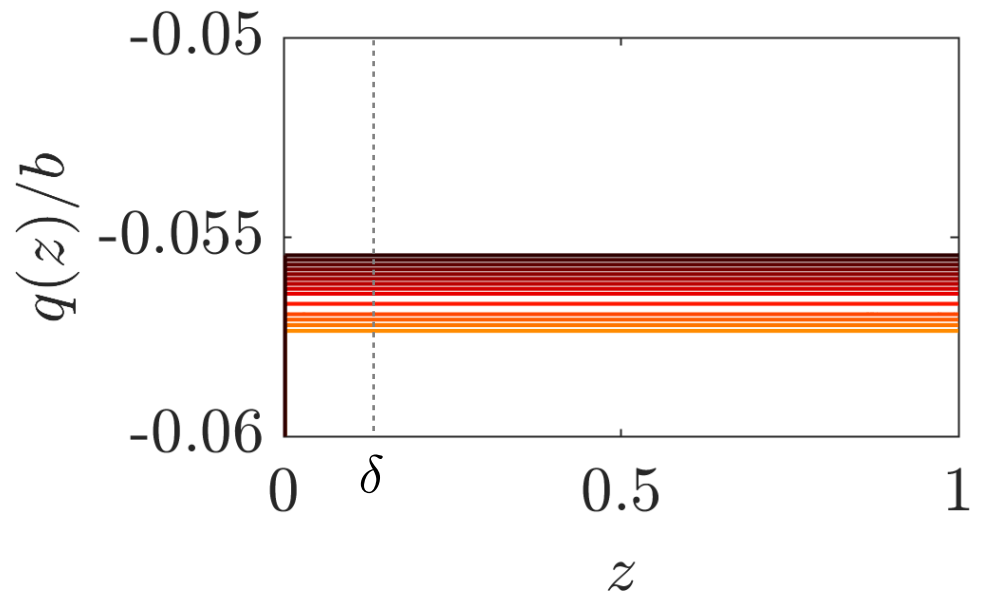
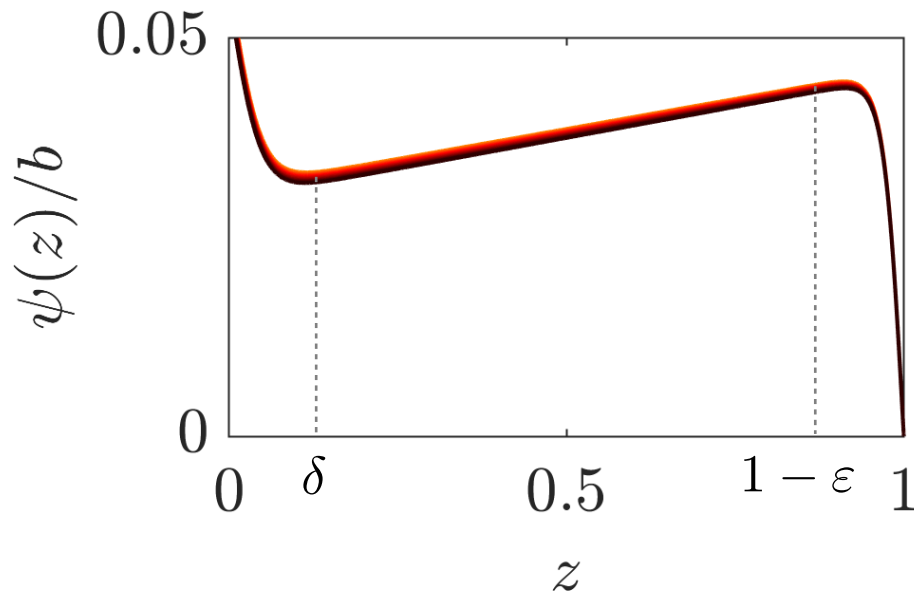
Optimal Lagrange multipliers



Optimal balance parameters



Approx. scaling with b away from $z=0$



$$\|\psi\|_{L^\infty(\delta,1)} \leq \mathcal{O}(b)$$

$$\|\psi'\|_{L^\infty(\delta,1)} \leq \mathcal{O}(b)$$

$$q(z) = \text{const} \quad \text{if } z > \delta$$

$$\psi'(z) \approx \text{const} \quad \text{in the bulk}$$

δ = width of bottom boundary layer
 ε = width of top boundary layer

Analysis?

$$\langle wT \rangle \leq \frac{1}{2} + \frac{1}{4b} \left\| bz - \frac{b}{2} - \psi' + q \right\|_2^2 - \int_0^1 \psi dz$$

Provided that

$$\psi(0) \leq 1,$$

$$\psi(1) \leq 0,$$

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Analysis?

$$\langle wT \rangle \leq \frac{1}{2} + \frac{1}{4b} \left\| bz - \frac{b}{2} - \psi' + q \right\|_2^2 - \int_0^1 \psi dz$$

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Suppose that:

- Agree well with numerics
- 1) $\psi(0) = 1, \psi(1) = 0$
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 - 3) $\psi'(z) \geq 0$ on an interval $(\delta, 1 - \varepsilon)$
 - 4) $\psi'(z)$ and $q(z)$ are constant on some interval $(\frac{1}{2} - c, \frac{1}{2} + c)$ with $c = \mathcal{O}(1)$
 - 5) $a \leq b$
 - 6) We use the “classical” estimates, e. g.

$$\left| \int_0^\delta (a - \psi') w T dz \right| \leq \delta^2 \|a - \psi'\|_{L^\infty(0,\delta)} \|\partial_z w\|_2 \|\partial_z T\|_2$$
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Then, the best possible bound U^* that one can prove satisfies

$$U^* \geq \frac{1}{2} + b \left[\frac{c^3}{6} - \mathcal{O}(R^{-\frac{1}{4}}) \right] > \frac{1}{2} \text{ if } R \gg 1$$

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Conclusions

- Bounds on vertical heat transfer for convection with internal heating are elusive!
- “Modern background method” leads to a sensible and computationally tractable problem
- New analytical bound

$$\langle wT \rangle \leq \frac{1}{16} [8(3\sqrt{2} - 4)^2]^{\frac{1}{5}} R^{\frac{1}{5}} \quad \left(\leq \frac{1}{2} \text{ if } R \leq 69\,572 \right)$$

- Positivity of temperature is necessary to obtain

$$\langle wT \rangle \leq \frac{1}{2} - cR^{-\alpha} < \frac{1}{2} \quad \forall R$$

- Simplest type of analytical constructions fail
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Thank you!

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