

Quantum algorithms for Dynamics Simulation: Hamiltonian Simulation and Linear Differential Equations

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Outline

- 1 Time-dependent Ham. Sim.
- 2 Quantum Linear Differential Equations Solvers

Revisit: Summary of Hamiltonian Simulation

- Hamiltonian simulation: motivation; set-up
- Expected cost: No-fast-forwarding theorem and BQP-hardness
- Trotterization
- truncated Taylor series
- QSVT

Important take-home message:

QSVT + OAA \Rightarrow Optimal Hamiltonian Simulation

Time-dependent Hamiltonian Simulation

$$\left\| \mathcal{U}_{\text{app}} - \mathcal{T}e^{-i \int_0^t H(s) ds} \right\| \leq \epsilon.$$

¹[Low-Wiebe 2018]

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Besides its own applications, it can also arise from time-independent problems under the **Interaction Picture**¹

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$H = A + B$ and A has a much larger spectral norm.

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Motivation: e.g.,

$$H = -\frac{1}{2}\Delta + V(x), \quad \|\Delta_h\| \gg \|V\|$$

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$$\begin{aligned} \partial_t \psi + iA\psi &= -iB\psi \\ \xrightarrow[\text{factor}]{\text{Integrating}} e^{iAt} \partial_t \psi + i e^{iAt} A\psi &= -i e^{iAt} B\psi \\ \Rightarrow \partial_t \left(\underbrace{e^{iAt} \psi}_{\psi_I} \right) &= -i e^{iAt} B\psi = -i \underbrace{e^{iAt} B e^{-iAt}}_{H_I(t)} \psi_I \end{aligned}$$

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$$H_I(t) := e^{iAt} B e^{-iAt}, \psi_I := e^{iAt} \psi \text{ and } i\partial_t \psi_I = H_I \psi_I$$

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Time-dependent Hamiltonian Simulation

Set-up

$$H(\tau) \quad \tau \in [0, t]$$

Desire feature: log-dependence on $\|\partial_t H(t)\|$.

²continuous qDRIFT [Berry-Childs-Su-Wang-Wiebe 2020], Monte-Carlo type [Poulin-Qarry-Somma-Verstraete 2011]

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Methods (short-time integrators):

- **Trotterization**, if $H(\tau) = \sum_{k=1}^L H_k(\tau)$. ×
- **Randomized methods**², e.g., continuous qDRIFT (sample and weak convergence) ✓
- **LCU + Series truncation**³: truncated Dyson series, rescaled Dyson series, truncated Magnus series, e.g., qHOP, etc ✓

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Time-dependent Hamiltonian Simulation

- Trotterization: $H = H_1(t) + H_2(t)$

$$\mathcal{T}e^{-i \int_{t_j}^{t_{j+1}} H(s) ds} \approx e^{-ihH_2(\tau_j)} e^{-ihH_1(\tau_j)},$$

where $\tau_j \in [t_j, t_{j+1}]$ are chosen according to Suzuki construction. The number of unitaries depends on $\|\partial_t H(t)\|$.

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- **Randomized algorithms** (first-order accuracy and weak conv)
e.g., continuous qDRIFT, hybridized methods, etc ⁵

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$$\begin{aligned} \mathcal{E}(t, 0)(\rho) &= U(t, 0)\rho U^\dagger(t, 0) \\ &= \mathcal{T}_{\rightarrow} \exp\left(-i \int_0^t d\tau H(\tau)\right) \rho \mathcal{T}_{\rightarrow} \exp^\dagger\left(-i \int_0^t d\tau H(\tau)\right) \\ &\approx \mathcal{U}(t, 0)(\rho) = \int_0^t d\tau p(\tau) e^{-i \frac{H(\tau)}{p(\tau)}} \rho e^{i \frac{H(\tau)}{p(\tau)}}, \end{aligned}$$

where $p(\tau) := \frac{\|H(\tau)\|}{\int_0^t \|H(\tau)\| d\tau} := \frac{\|H(\tau)\|}{\|H\|_{\infty, 1}}$ is a probability density function defined for $0 \leq \tau \leq t$.

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Why? How to derive?

$$\partial_t U(t, 0) = -iH(t)U(t, 0), \quad U(0, 0) = I$$

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Part 2: General Linear Differential Equations (non-unitary dynamics)

Outline of Quantum Linear Differential Equation Solvers

- Definition of the task
- Challenges
- Ways to address
 - QSLA + Padding
 - Time-marching
 - LCHS

Linear differential equations

$$\frac{d}{dt} |\psi(t)\rangle = A(t) |\psi(t)\rangle, \quad |\psi(0)\rangle = |\psi_0\rangle,$$

Task: To prepare a quantum state that is proportional to the final solution $|\psi(T)\rangle$ with certain precision ϵ .

To prepare a quantum state $|\tilde{\psi}(T)\rangle$ that satisfies

$$\left\| \frac{|\psi(T)\rangle}{\| |\psi(T)\rangle \|} - \frac{|\tilde{\psi}(T)\rangle}{\| |\tilde{\psi}(T)\rangle \|} \right\| = \mathcal{O}(\epsilon).$$

Note that

$$|\psi(T)\rangle = \mathcal{T}e^{\int_0^T A(s)ds} |\psi_0\rangle,$$

and hence it is reasonable to construct $\mathcal{T}e^{\int_0^T A(s)ds}$ and then apply it to the quantum state.

Challenge

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Question: Why it works for Hamiltonian Simulation?

For general A , V_{num} are NOT close to unitaries! Cannot use **OAA**!

Challenge

Major Challenge:
Short time \Rightarrow Long time?

Sidenote for simple cases

Consider the simple time-independent case $A(t) \equiv A$. We seek to implement e^{tA} .

- When A is **anti-Hermitian**, the task becomes Hamiltonian simulation and QSVT gives the best asymptotic scaling.

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Issue: matrix exponential is defined via eigenvalue decomposition that does not agree with singular value decomposition unless the matrix is normal.

Remedy: **Contour integral formulation**⁷.

$$e^A = \frac{1}{2\pi i} \oint_{\Gamma} e^z (z - A)^{-1} dz \approx \frac{1}{K} \sum_{k=0}^{K-1} e^{z_k} z_k (z_k - A)^{-1}.$$

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Challenge

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Ideas:

- Postpone A.A. as much as possible
- Device new ways to boost success probability at each time step

Short time to Long time: Approach 1

Approach 1: Apply **QLSA + padding** to history state ⁸

$$\mathcal{L}\mathbf{x} = \mathbf{b}, \quad \mathbf{x} = ((\psi^1)^T (\psi^2)^T \dots (\psi^L)^T)^T$$

⁸[Berry 2014], [Berry-Childs-Ostrander-Wang 2017], [Childs-Liu 2019], [Krovi 2023], etc

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$$\begin{pmatrix} I & 0 & 0 & \dots & 0 & 0 \\ -(I + \Delta t A_1) & I & 0 & \dots & 0 & 0 \\ 0 & -(I + \Delta t A_2) & I & \dots & 0 & 0 \\ \vdots & & & & & \vdots \\ 0 & 0 & 0 & \dots & -(I + \Delta t A_{L-1}) & I \end{pmatrix} \begin{pmatrix} \psi^1 \\ \psi^2 \\ \psi^3 \\ \vdots \\ \psi^L \end{pmatrix} = \begin{pmatrix} (I \\ \vdots \\ I) \end{pmatrix}$$

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Cost of applying QLSA:

$$\mathcal{O}(s\kappa \text{ polylog}(Ns\kappa/\epsilon))$$

queries to oracles for \mathcal{L} and \mathbf{b}

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$$\begin{pmatrix} I & 0 & 0 & \dots & 0 & 0 \\ -(I + \Delta t A_1) & I & 0 & \dots & 0 & 0 \\ 0 & -(I + \Delta t A_2) & I & \dots & 0 & 0 \\ \vdots & & & & & \vdots \\ 0 & 0 & 0 & \dots & -(I + \Delta t A_{L-1}) & I \end{pmatrix} \begin{pmatrix} \psi^1 \\ \psi^2 \\ \psi^3 \\ \vdots \\ \psi^L \end{pmatrix} = \begin{pmatrix} I \\ 0 \\ 0 \\ \vdots \\ 0 \end{pmatrix}$$

Cost of applying QLSA:

$$\mathcal{O}(s\kappa \text{ polylog}(Ns\kappa/\epsilon))$$

queries to oracles for \mathcal{L} and \mathbf{b} (which in turn gives query complexity to oracles for A and for preparing ψ^0).

\Rightarrow estimate the condition number of $\mathcal{L} \sim 1/\Delta t^\alpha$ for $\alpha > 0$.

⁸[Berry 2014], [Berry-Childs-Ostrander-Wang 2017], [Childs-Liu 2019], [Krovi 2023], etc

Short time to Long time: Approach 1

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So far we only get the history state $\mathbf{x} \Rightarrow$ still needs to extract information of ψ^L out of \mathbf{x} .

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Success probability issue: padding + A.A.

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Generic Result: Under certain assumptions (sparse **smooth/analytic dissipative** $A(t) = V(t)\Lambda(t)V(t)^{-1}$ is diagonalizable for all time and the condition number of $V(t)$ has a uniform(-in- t) upper bound; all derivatives of the solution have a uniform upper bound),

$$\mathcal{O}(qT \text{ polylog}(T/\epsilon)) \quad (q := \|\psi(0)\|/\|\psi(t)\|)$$

queries to both Oracle for A and Oracle preparing the initial state.

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Lower bound: linear in q is needed, but the optimal queries to the oracle preparing the initial state is $\mathcal{O}(q)$. ¹⁰

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Natural Question: Can we achieve optimal state preparation cost?

⁹[Berry 2014], [Berry-Childs-Ostrander-Wang 2017], [Childs-Liu 2019], [Krovi 2023], etc

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Short time to Long time: Approach 2

Approach 2: Remedy **time-marching** strategy by Uniform Singular Value Amplification (USVA) ¹¹

Motivation:

$$\begin{pmatrix} \frac{I+\Delta t A}{1+\Delta t \|A\|} & * \\ * & * \end{pmatrix} \Rightarrow \begin{pmatrix} \frac{(I+\Delta t A)^L}{(1+\Delta t \|A\|)^L} & * \\ * & * \end{pmatrix}$$

Success probability: $\Omega \left(\frac{\|\psi(t)\|^2}{\|\psi(0)\|^2 (1+\Delta t \|A\|)^{2L}} \right)$

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Observation: $\|I + \Delta t A\|^L \approx \|e^{\Delta t A}\|^L \sim \|e^{tA}\|$. For dissipative system and unitary dynamics, it is ok!

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Needs:

$$\begin{pmatrix} \overline{\alpha} & * \\ * & * \end{pmatrix} \xrightarrow{\text{oblivious to the state}} \begin{pmatrix} \frac{\overline{\alpha}}{\|\overline{\alpha}\|} & * \\ * & * \end{pmatrix}$$

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Short time to Long time: Approach 2

Needs:

$$\begin{pmatrix} \alpha \|u\| & * \\ * & * \end{pmatrix} \xrightarrow[\text{QSVT}]{\text{oblivious to the state}} \begin{pmatrix} \frac{\|u\|}{\|v\|} & * \\ * & * \end{pmatrix}$$

Short time to Long time: Approach 2

Needs:

$$\begin{pmatrix} \alpha \| \Xi \| & * \\ * & * \end{pmatrix} \xrightarrow[\text{QSVT}]{\text{oblivious to the state}} \begin{pmatrix} \Xi & * \\ * & * \end{pmatrix}$$

Remark: For unitary dynamics, **OAA** works the magic! OAA only needs to boost a singular value, say $1/2$ to 1 using $3x - 4x^3$.)

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$$\begin{pmatrix} \alpha \|\Xi\| & * \\ * & * \end{pmatrix} \xrightarrow[\text{QSVT}]{\text{oblivious to the state}} \begin{pmatrix} \frac{\Xi}{\|\Xi\|} & * \\ * & * \end{pmatrix}$$

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Denote $\gamma := \alpha / \|\Xi\|$.

We seek for a polynomial approximation of γx on $[-\gamma^{-1}, \gamma^{-1}]$.

Short time to Long time: Approach 2

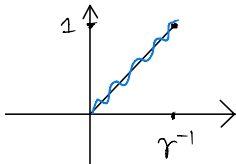
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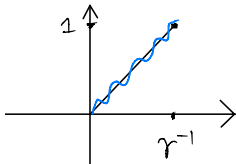
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Gibbs phenomena

Short time to Long time: Approach 2

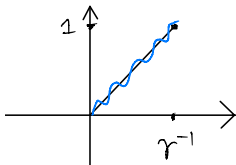
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Gibbs phenomena

Lemma (**Uniform Singular Value Amplification**)

$$d = \mathcal{O}(\delta^{-1} \gamma \log(\gamma/\epsilon)).$$

Short time to Long time: Approach 2

Algorithm per time step:

- 1 Numerical integrator (Dyson, Magnus, Euler, etc)
- 2 Uniform Singular Value Amplification (QSVT)

[Fang-Lin-Tong 2023 / arXiv 2022]

Short time to Long time: Approach 2

Algorithm per time step:

- 1 Numerical integrator (Dyson, Magnus, Euler, etc)
- 2 Uniform Singular Value Amplification (QSVT)

Construct a block-encoding of long-time evolution by compression gadget.

Generic Result for dissipative or near-unitary dynamics: The quantum algorithm makes

$$\mathcal{O}(qT^2 \text{polylog}(T/\epsilon))$$

queries to the oracle for A and

$$\mathcal{O}(q)$$

queries to the oracle preparing the initial state. (smoothness is not required; bounded variation is sufficient)

[Fang-Lin-Tong 2023 / arXiv 2022]

Short time to Long time: Approach 3

Approach 3: Complex analysis (Linear combination of Hamiltonian simulation)¹²

$$A(t) = L(t) + iH(t), \quad L(t) = \frac{A(t) + A^\dagger(t)}{2}, \quad H(t) = \frac{A(t) - A^\dagger(t)}{2i}.$$

Assume $L(t) \preceq 0$ for all $t \in \mathcal{I}$.

$$\mathcal{T}e^{-\int_0^t A(s)ds} = \int_{\mathbb{R}} \frac{1}{\pi(1+k^2)} \mathcal{T}e^{-i\int_0^t (H(s)+kL(s))ds} dk, \quad t \in \mathcal{I}.$$

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Generic Result for dissipative dynamics: The quantum algorithm makes

$$\mathcal{O}\left(q^{2+2/p} T^{1+1/p} / \epsilon^{1+2/p}\right)$$

queries to the oracle for H and L and

$$\mathcal{O}(q)$$

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¹²[An-Lin-Liu 2023]

Improved solver for special cases

Everything we discussed is on **general quantum solvers** (for general $A(t)$) and attaining lower-bound in the worst-case scenario.

This doesn't mean that the scaling for a **specific case** can not be further improved.

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a specific linear system with conservation laws.

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For simplicity, consider $E(0) \equiv E(t) = \sum_{j=1}^N \frac{\dot{x}_j(t)^2}{2} + \frac{x_j(t)^2}{2}$.
Hamiltonian ODE: $\dot{y} = Ay$, $y = (x_1, \dots, x_N, \dot{x}_1, \dots, \dot{x}_N)$.

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Results (informal):

- It can be mapped to Hamiltonian simulation, and the quantum algorithm makes

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- still BQP-hard.

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Summary of Quantum Linear Differential Equation Solvers

- Definition of the task
- Challenges
- Ways to address (general solvers)
 - QSLA + Padding
 - Time-marching
 - LCHS
- Improved solver for specific cases, e.g., simulating coupled classical oscillators

Thank you for your attention!

