

# A hierarchy of models, from planetary to mesoscale, within a single switchable numerical framework

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Met Office

# Outline of the talk

- Overview of Met Office's Unified Model
- Spherical Geopotential Approximation
- Relaxing that approximation
- A method of implementation
- Alternative view of Shallow-Atmosphere Approximation
- Application to departure points

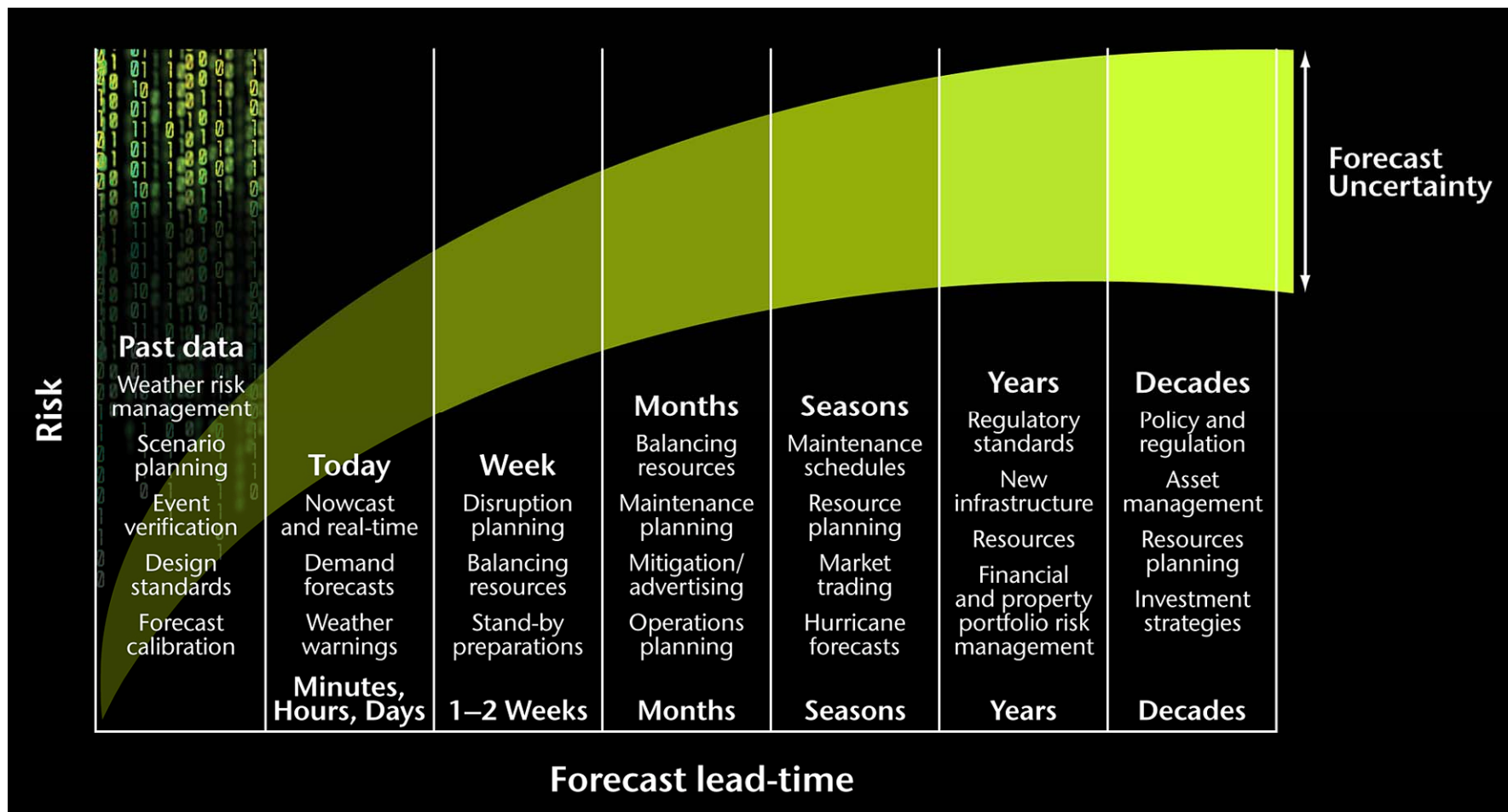


# Met Office's Unified Model

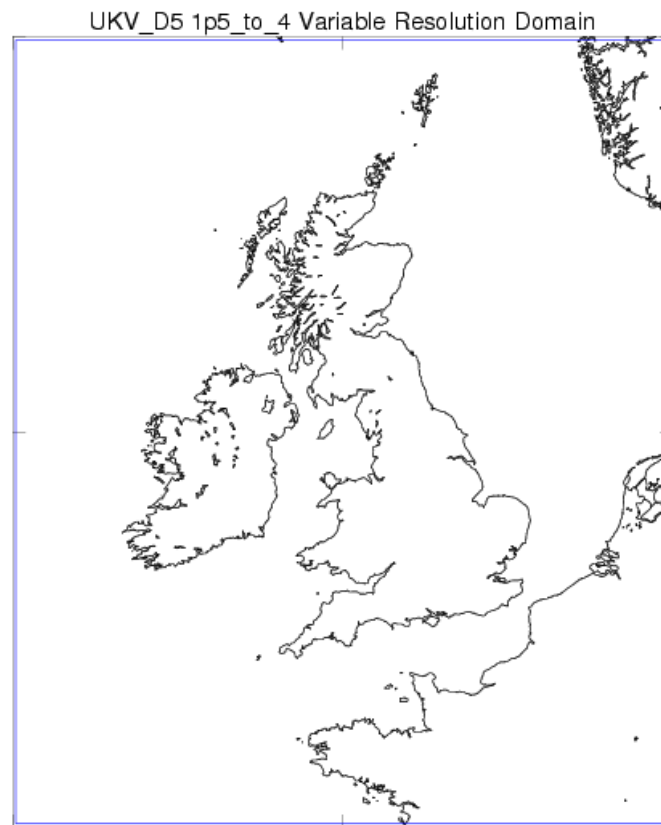
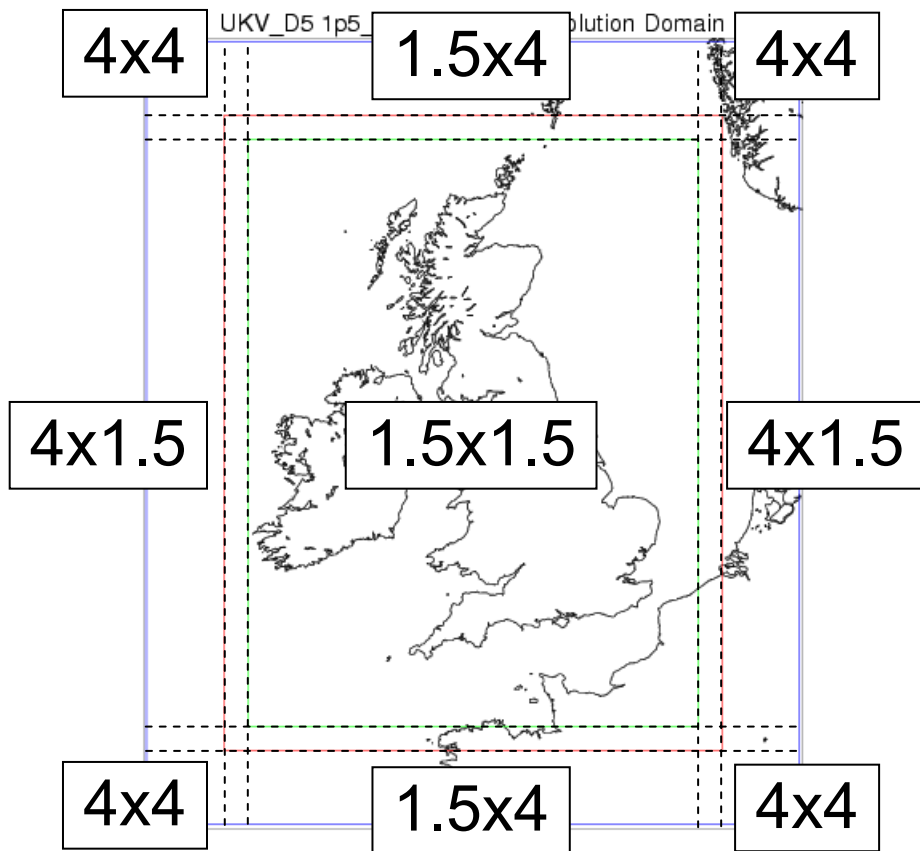
Unified Model (UM) in that *single* model for:

- Operational forecasts at
  - Mesoscale (resolution approx. 12km → 4km → 1km)
  - Global scale (resolution approx. 25km)
- Global and regional climate predictions (resolution approx. 100km, run for 10-100 years)
- + Research mode (1km - 10m) and single column model
- 20 years old this year

# Timescales & applications

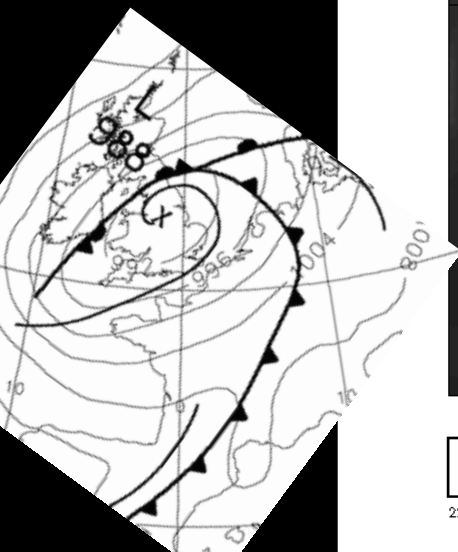


# UKV Domain

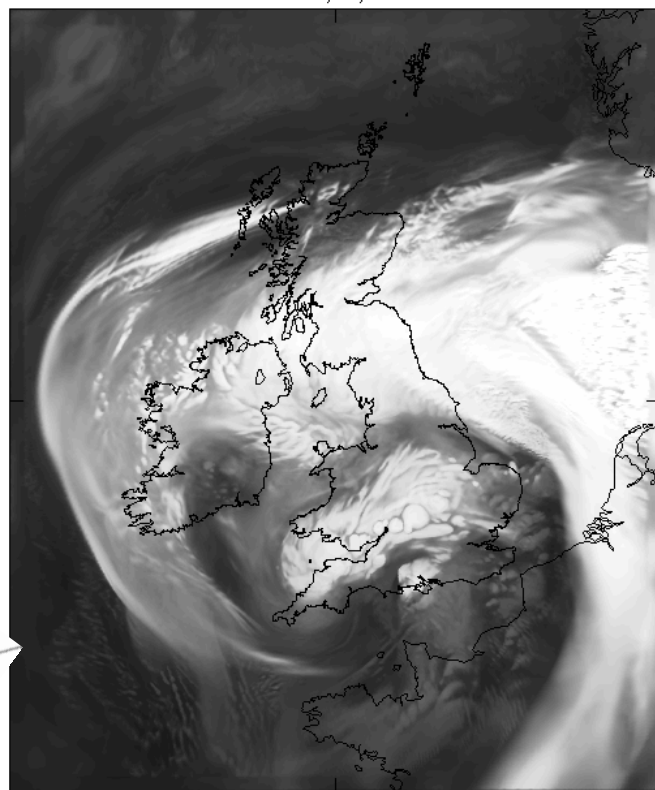


# The 'Morpeth Flood', 06/09/2008

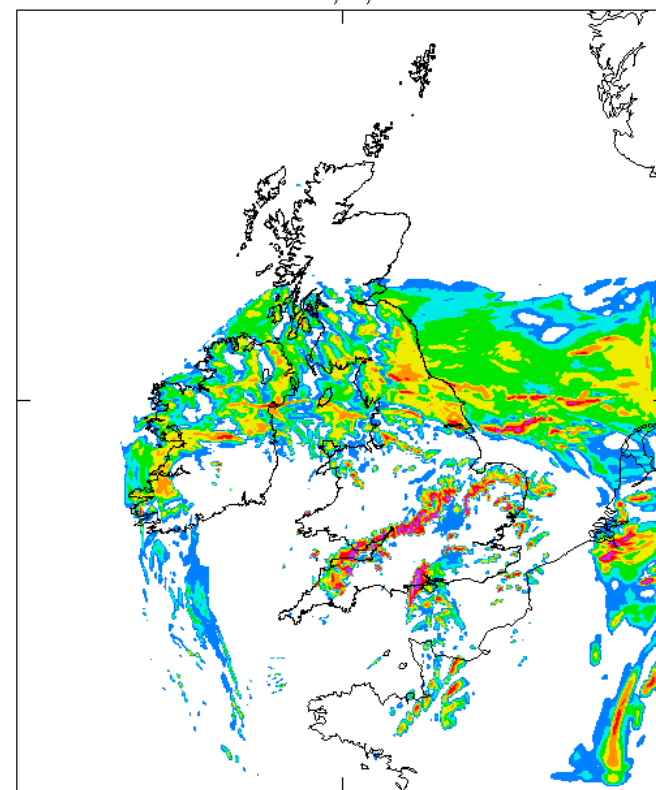
1.5 km L70  
Prototype UKV  
From  
15 UTC 05/09  
12 km



LW Radiance Temp  
1800 05/09/2008

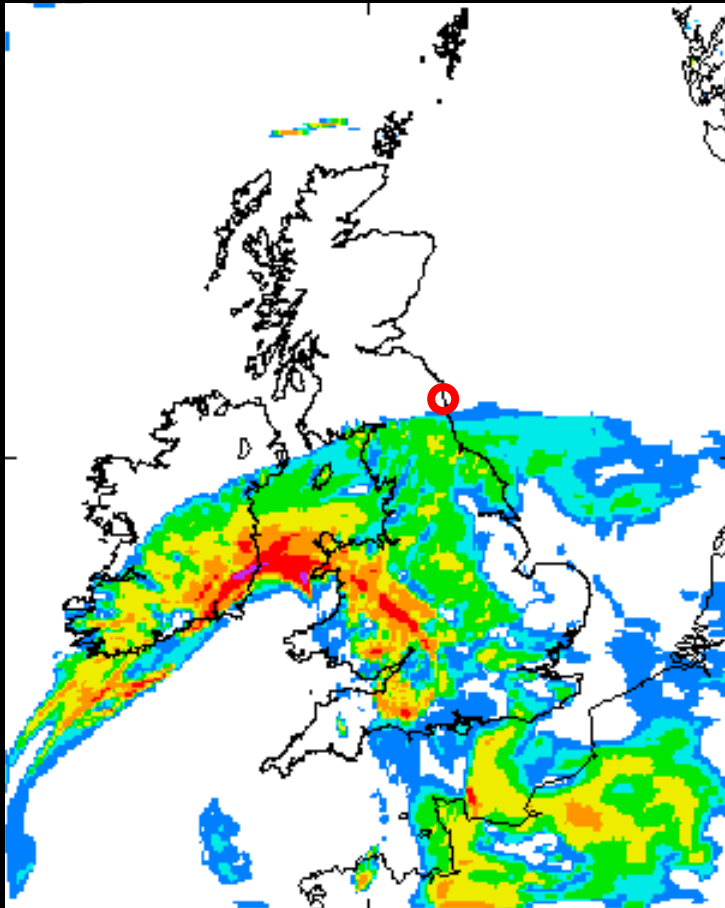


Total precipitation rate, mm/hr  
1800 05/09/2008

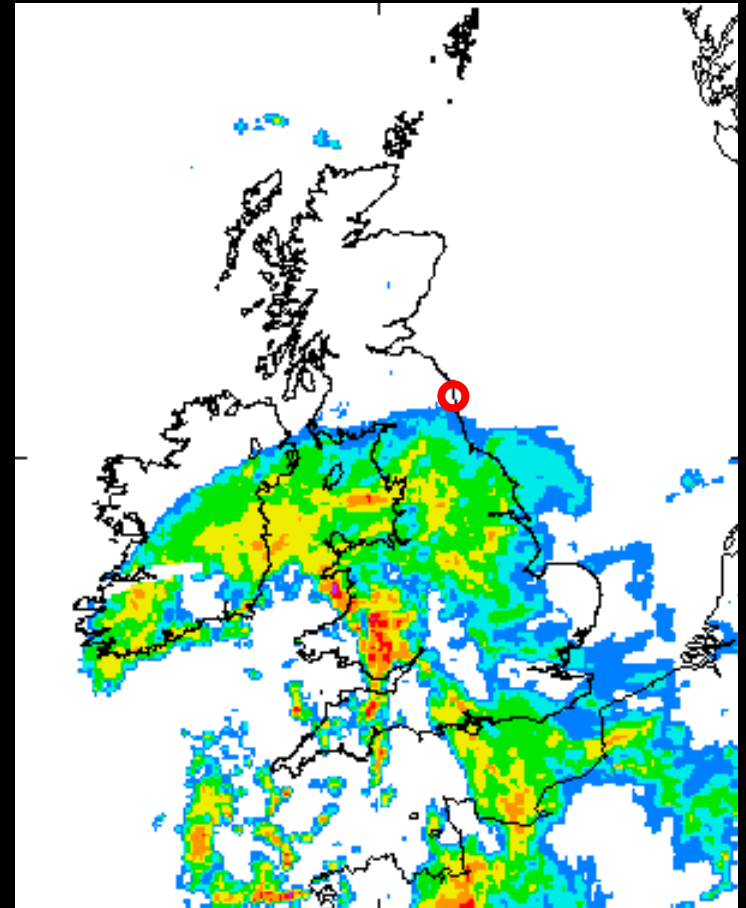


0600 UTC

# 2009: Clark et al, Morpeth Flood



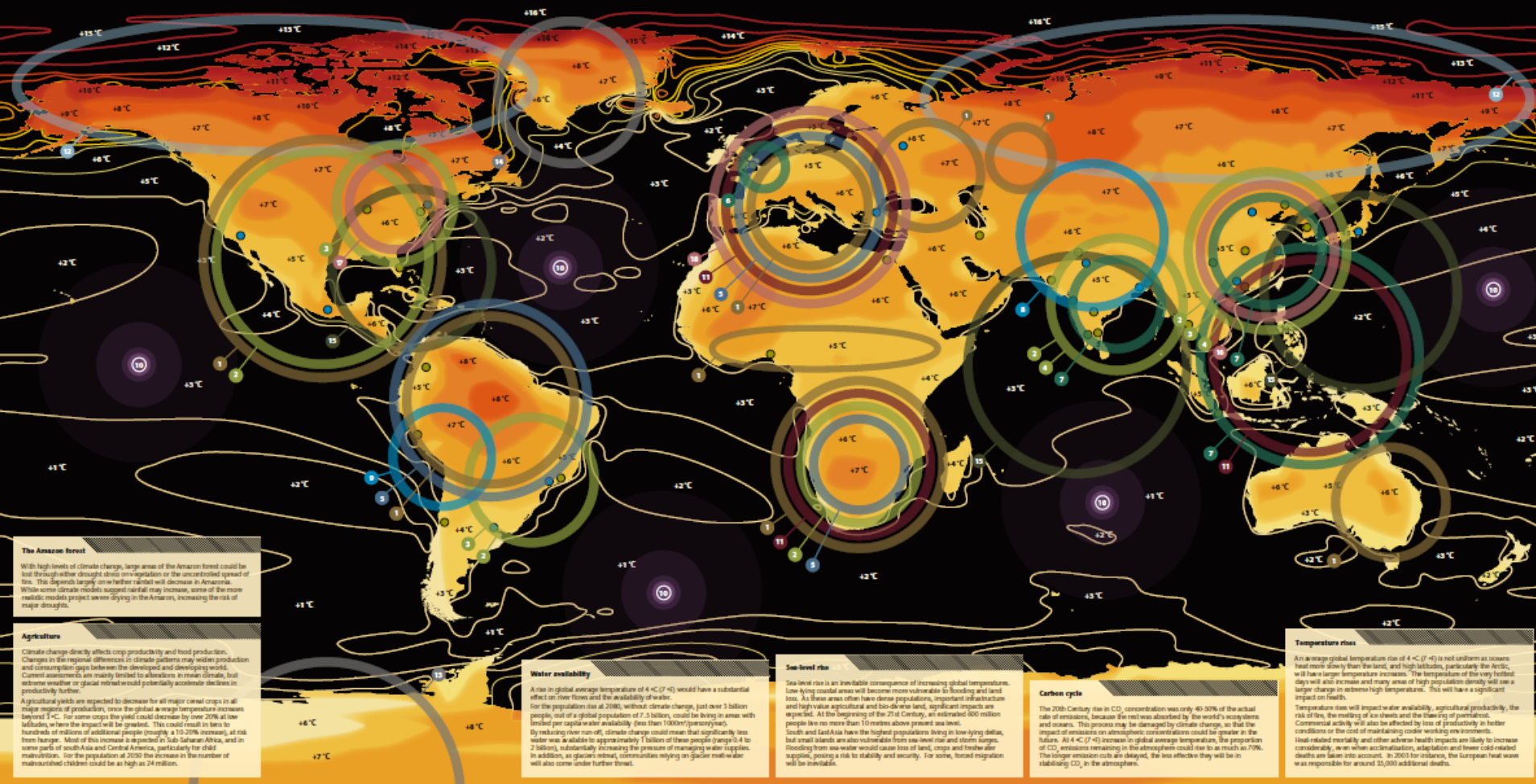
**UKV 4-20hr forecast  
1.5km gridlength  
convection permitting model**



**UK radar network**



## The impact of a global temperature rise of 4 °C (7 °F)



1 High-level fire danger projected to affect every populated continent. Regions moving into the high-danger category include: large areas of the United States, Mexico, South America, east of the Andes, southern and west Africa, the Sahel, eastern and southern Australia and southern Europe.

2 Maize and wheat yields reduced by up to 40% at low latitudes.

3 Soybean yield could decrease in all regions of production, including North and South America, southern and eastern Asia.

4 Decrease in rice yield of up to 30% in China, India, Bangladesh and Indonesia.

5 Water resources affected by up to 70% reductions in run-off around the Mediterranean, southern Africa and large areas of South America.

6 Sea-level rise combined with storm surges could pose a serious threat to people and assets in the Netherlands and south-western parts of the UK.

7 Sea levels could rise as much as 80 cm by the end of the century. Larger rises, 4°C (7°F) could result in a much higher rise in sea level. Sea level increases are likely to be even greater at low latitudes, disproportionately affecting tropical islands and low-lying regions such as Bangladesh.

8 For the population at 2050, a mean sea level rise of 15 cm means that up to an additional 150 million people per year would be flooded due to extreme sea levels. Three-quarters of these people live in Asia. Up to 54 million people could be flooded along the Indian Ocean coast. 25 million along the south-east Asian coast and 31 million people would be flooded along the South-East Asian coast.

9 Other vulnerable regions include Africa, Caribbean Islands, Indian Ocean Islands and Pacific small islands.

10 Half of all glaciers on glaciers significantly reduced by 2050, even at a global average temperature rise below 4°C. The ice in the Swiss Alps contains 70% of its summer flow from glacial melt. In China, 27% of the population lives in the western regions where glacial melt provides the principal dry season water source.

11 Complete disappearance of glaciers from many regions in South America. In Peru's Cordillera Blanca summer rain-off from glaciers reduced by up to 60% as the glacier area falls by 70%.

12 Marine ecosystems could be fundamentally altered by ocean acidification which would have a significant impact on fisheries. This could cause substantial loss in marine and jobs. The loss of coral reef habitats due to acidification may severely affect many commercial fish species and could pose dangers for coastal communities relying on industries fishing of reef species.

13 Drought events occur twice as frequently across southern Africa, South-East Asia and the Mediterranean Basin.

14 Almost complete disappearance of near surface permafrost from Northern Siberia. Reduction of permafrost in Canada and Alaska. Infrastructure built on the permafrost becomes at risk.

15 It is not known how stable the West Antarctic Ice Sheet is, or whether a 4°C (7°F) global temperature rise will send it into irreversible decline. If this is the case, it would contribute a further 3.2 metres to long-term sea-level rise globally.

16 Greenland Ice Sheet has a 60% likelihood of irreversible decline. This would result in a very long-term sea-level rise of up to 7 metres globally.

17 Tropical systems could be more intense and destructive. Global population increases, particularly in coastal areas, and sea-level rise mean greater cyclones and hurricane related losses, disruptions to infrastructure and loss of life as a result of storm surges. For major systems, flooding from storm surges has been the primary cause of death.

18 Hottest days of the year could be as much as 6°C (11°F) warmer on or highly populated areas of eastern China.

19 Hottest days of the year could become as much as 10-12°C (18-22°F) warmer over eastern North America, affecting Toronto, Chicago, Ottawa, New York and Washington DC.

20 Hottest days of the year across Europe could be as much as 6°C (14°F) warmer.

Climate patterns have changed with an overall increase in droughts, water stress, droughts such as in India and droughts, hurricanes and the health impacts of weather events such as flooding and drought.

+Tabelle		Change in temperature from pre-industrial climate															
1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
2	4	5	7	9	11	13	14	16	18	20	22	23	25	27	28	29	30
+Tabelle																	
City populations																	
5 - 10 Million																	
10 - 20 Million																	



*The underlying equations...*

## Traditional Spherically Based Equations

$$\frac{D_{\textcolor{red}{r}}u}{Dt} - \frac{uv \tan \phi}{\textcolor{red}{r}} - (2\Omega \sin \phi) v + \frac{c_{pd}\theta_v}{\textcolor{red}{r} \cos \phi} \frac{\partial \Pi}{\partial \lambda} = \boxed{-\frac{uw}{\textcolor{red}{r}} - (2\Omega \cos \phi) w} + S^u$$

$$\frac{D_{\textcolor{red}{r}}v}{Dt} + \frac{u^2 \tan \phi}{\textcolor{red}{r}} + (2\Omega \sin \phi) u + \frac{c_{pd}\theta_v}{\textcolor{red}{r}} \frac{\partial \Pi}{\partial \phi} = \boxed{-\frac{vw}{\textcolor{red}{r}}} + S^v$$

$$\boxed{\frac{D_{\textcolor{blue}{r}}w}{Dt}} + c_{pd}\theta_v \frac{\partial \Pi}{\partial r} + \frac{\textcolor{green}{\partial \Phi}}{\textcolor{green}{\partial r}} = \boxed{\frac{u^2 + v^2}{\textcolor{red}{r}} + (2\Omega \cos \phi) u + S^w}$$

**Red** = Shallow-Atmosphere Approx (eliminate boxed terms)

**Blue** = Hydrostatic Approx (eliminate boxed term)

**Green** = Spherical Geopotential Approx ( $\partial \Phi / \partial r = g(r)$  only)

## The Spherical Geopotential Approximation

- All geopotentials - including the Earth's surface - are represented by concentric *spheres*
- Apparent gravity acts towards the centre of the Earth
- To prevent spurious vorticity sources/sinks  $g$  cannot vary with latitude (White et al., 2005)
- Nearly all numerical models of the global atmosphere are based on the SGA

## Why relax the SGA?

- Good to get Earth's shape right(ish) in models!
- Good to include the small (  $\approx 0.5\%$  ) increase of  $g$  from Equator to Poles without spurious vorticity sources/sinks
- Most likely to be evident in long term climate simulations
- Good to quantitatively test the accuracy of the spherical geopotential approximation
- Important though that signal is not influenced by different numerics

# The Shape of the Earth

- The Earth is approximately an **oblate spheroid**:

- Equatorial radius

$$a = 6378 \text{ km}$$

- Polar radius

$$c = 6357 \text{ km}$$

- **Ellipticity small**

$$\varepsilon \equiv \frac{a - c}{a} \approx \frac{1}{298}$$

- **But  $a - c = 21 \text{ km}$ , more than twice the height of Everest**



## Geopotential near rotating fluid spheroid

- Classical problem (**Clairaut** in 18th century) to account for the observed increase of  $g$  with latitude
- **Re-examined by White et al. (2008) to find geopotentials**
- **Two small parameters:**

$$\varepsilon \equiv \frac{a - c}{a} \approx \frac{1}{298}$$

$$m \equiv \frac{\Omega^2 a^3}{\gamma M_E} \sim \left| \frac{\text{centrifugal force}}{\text{Newtonian gravity}} \right| \approx \frac{1}{289},$$

- **To  $O(\varepsilon, m)$  geopotentials are spheroids**

$$x^2 + z^2 \left[ 1 + (2\varepsilon - m) \frac{a^2}{R^2} + m \frac{R^3}{a^3} \right] = R^2$$

## Variation of Gravity

- Near to the surface  $R = a$

$$\frac{g_P}{g_{Eq}} = 1 + \frac{5}{2}m - \epsilon$$

- **Earth's actual mass distribution**  $m = \epsilon$  (to within 3%):

$$\boxed{\frac{g_P}{g_{Eq}} = 1 + \frac{3}{2}\epsilon} \Rightarrow 0.5\% \text{ increase}$$

- **For uniform mass distribution**  $m = 4\epsilon/5$  (Newton)

$$\boxed{\frac{g_P}{g_{Eq}} = 1 + \epsilon}$$

**captures 2/3rds of the variation**

## *Spheroidal Coordinates...*

## Confocal Oblate Spheroids

- “Oblate spheroidal coordinates” of e.g. Gill (1982) and Gates (2004) are in fact *confocal* oblate spheroidal coordinates:

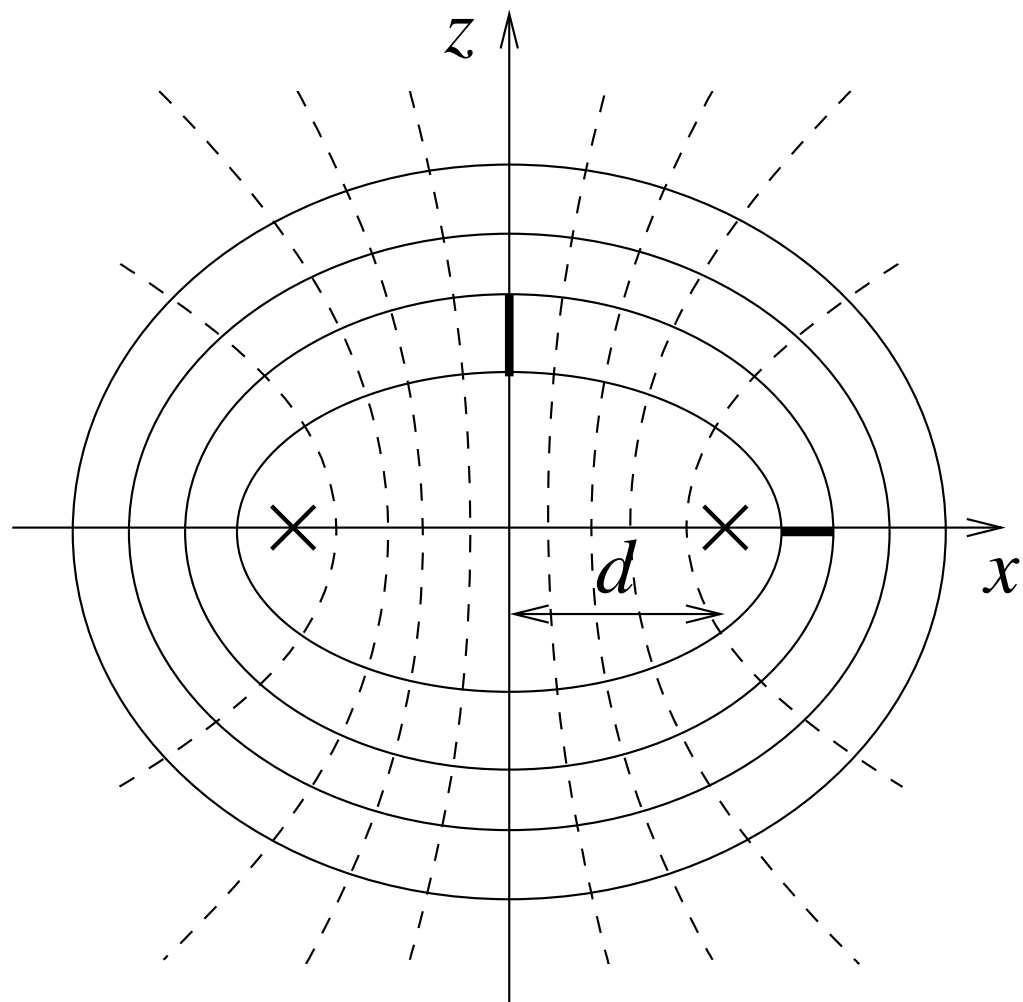
$$\frac{x^2}{\cosh^2 \eta} + \frac{z^2}{\sinh^2 \eta} = d^2$$

with  $d$  fixed,  $\eta$  variable

- Match to Earth's surface  $\Rightarrow$

$$\frac{g_P}{g_{Eq}} = \tanh \eta_0 \approx 1 - \epsilon \neq 1 + \frac{3}{2}\epsilon!$$

- Separation along minor axis *greater* than along major axis



**CONFOCAL ELLIPSES and HYPERBOLAS**



## Similar ellipses

Equation is:

$$x^2 + (1 + \mu) z^2 = r^2$$

Now  $\mu$  fixed,  $r$  variable

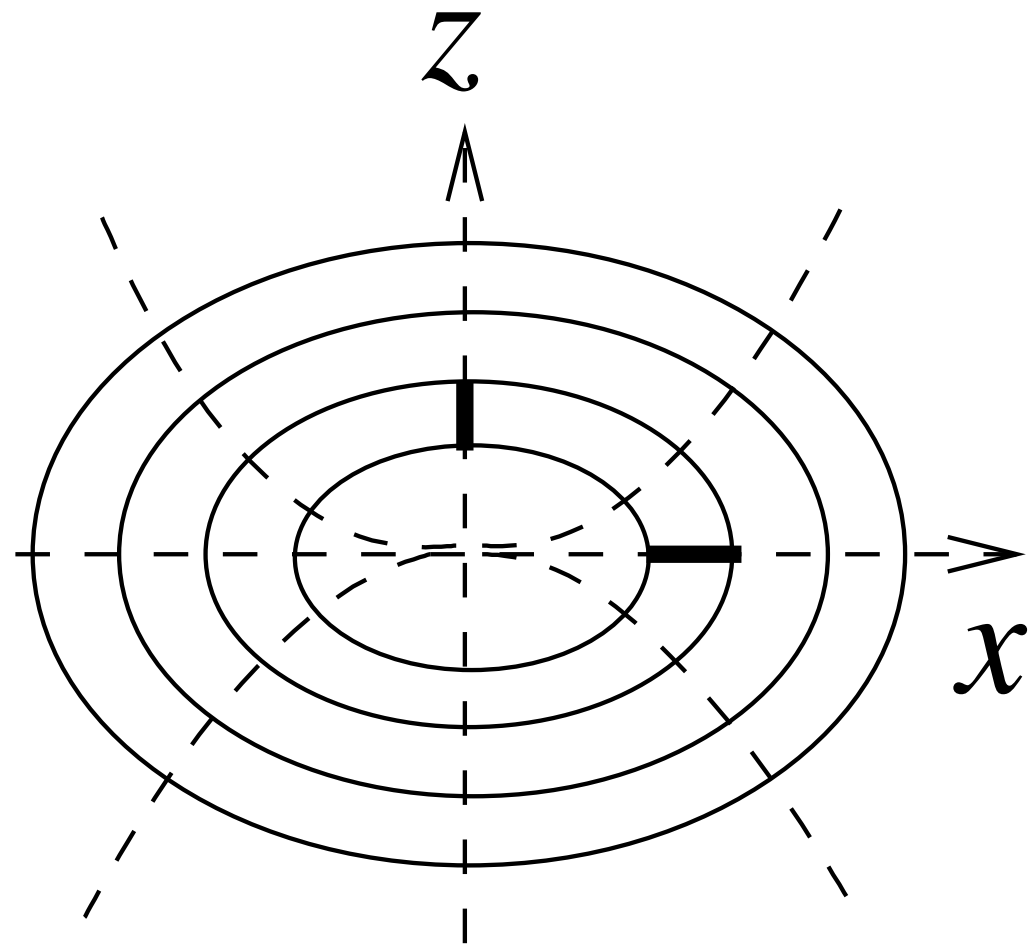
Match to Earth's surface  $\Rightarrow$

$$\mu = 2\epsilon$$

and

$$\frac{g_P}{g_{Eq}} = (1 + \mu)^{1/2} \approx 1 + \epsilon$$

- Sign right!
- Captures 2/3rds actual variation for actual Earth
- Exact for an Earth with uniform mass distribution



**SIMILAR ELLIPSES**



***Implementation...***

## Semi-Lagrangian scalar advection

**Consider**

$$\frac{DF}{Dt} = R$$

**Integrate *along trajectory*:**

$$\int_t^{t+\Delta t} \frac{DF}{Dt} dt = \frac{F^{t+\Delta t}(x) - F^t(x - U\Delta t)}{\Delta t} = \int_t^{t+\Delta t} R dt \approx \frac{R^{t+\Delta t}(x) + R^t(x - U\Delta t)}{2}$$

**i.e.**

$$\left( F - \frac{\Delta t}{2} R \right)_A^{t+\Delta t} = \left( F + \frac{\Delta t}{2} R \right)_D^t$$

## Semi-Lagrangian vector advection

For the vector equation

$$\frac{D\mathbf{v}}{Dt} = \mathbf{S}$$

Integrate *along trajectory* to obtain, as before:

$$\left( \mathbf{v} - \frac{\Delta t}{2} \mathbf{S} \right)_A^{t+\Delta t} = \left( \mathbf{v} + \frac{\Delta t}{2} \mathbf{S} \right)_D^t$$

All well and good...

But what of components?



## The Problem of Curvilinear Coordinates...

**The familiar problem that**

$$\mathbf{i} \cdot \frac{D\mathbf{v}}{Dt} \equiv \mathbf{i} \cdot \frac{D}{Dt} (u\mathbf{i} + v\mathbf{j} + w\mathbf{k}) \neq \frac{D(\mathbf{i} \cdot \mathbf{v})}{Dt} \quad \left[ = \frac{Du}{Dt} \right]$$

**translates into the SL equivalent that**

$$\boxed{\mathbf{i}_A \cdot (\mathbf{v}_D) \neq (\mathbf{i}_A \cdot \mathbf{v}_A)_D \quad [= u_D]}$$

**In fact**

$$\mathbf{v}_D = u_D \mathbf{i}_D + v_D \mathbf{j}_D + w_D \mathbf{k}_D$$

**so that, e.g.,**

$$\mathbf{i}_A \cdot (\mathbf{v}_D) = u_D (\mathbf{i}_A \cdot \mathbf{i}_D) + v_D (\mathbf{i}_A \cdot \mathbf{j}_D) + w_D (\mathbf{i}_A \cdot \mathbf{k}_D)$$

## The Matrix Formulation

- Extending this to all three directions:

$$\begin{pmatrix} u \\ v \\ w \end{pmatrix}_D = \mathbf{M} \begin{pmatrix} u_{\mathcal{D}} \\ v_{\mathcal{D}} \\ w_{\mathcal{D}} \end{pmatrix}$$

where

$$\mathbf{M} \equiv \begin{pmatrix} \mathbf{i}_A \cdot \mathbf{i}_D & \mathbf{i}_A \cdot \mathbf{j}_D & \mathbf{i}_A \cdot \mathbf{k}_D \\ \mathbf{j}_A \cdot \mathbf{i}_D & \mathbf{j}_A \cdot \mathbf{j}_D & \mathbf{j}_A \cdot \mathbf{k}_D \\ \mathbf{k}_A \cdot \mathbf{i}_D & \mathbf{k}_A \cdot \mathbf{j}_D & \mathbf{k}_A \cdot \mathbf{k}_D \end{pmatrix}$$

- $\mathbf{M}$  transforms
  - **from:** vector components in the *departure-point* frame
  - **to:** vectors in the *arrival-point* frame

## The Consequence...

**Component form of momentum equations can be written as:**

$$\left( \mathbf{v} - \frac{\Delta t}{2} \mathbf{S} \right)_{\mathcal{A}}^{t+\Delta t} = \mathbf{M} \left( \mathbf{v} + \frac{\Delta t}{2} \mathbf{S} \right)_{\mathcal{D}}^t$$

**where**

$$\mathbf{X}_{\mathcal{A}} \equiv (X_{\mathcal{A}}, Y_{\mathcal{A}}, Z_{\mathcal{A}})^T$$

$$\mathbf{X}_{\mathcal{D}} \equiv (X_{\mathcal{D}}, Y_{\mathcal{D}}, Z_{\mathcal{D}})^T$$

- **No explicit metric terms**
- **No singularity at the pole**

## M for Spherical Polar Coordinates

$$M_{11} = \cos \Delta \lambda, \quad M_{12} = \sin \phi_D \sin \Delta \lambda, \quad M_{13} = -\cos \phi_D \sin \Delta \lambda$$

$$M_{21} = -\sin \phi_A \sin \Delta \lambda$$

$$M_{22} = \cos \phi_A \cos \phi_D + \sin \phi_A \sin \phi_D \cos \Delta \lambda$$

$$M_{23} = \cos \phi_A \sin \phi_D - \sin \phi_A \cos \phi_D \cos \Delta \lambda$$

$$M_{31} = \cos \phi_A \sin \Delta \lambda$$

$$M_{32} = \sin \phi_A \cos \phi_D - \cos \phi_A \sin \phi_D \cos \Delta \lambda$$

$$M_{33} = \sin \phi_A \sin \phi_D + \cos \phi_A \cos \phi_D \cos \Delta \lambda$$

## Where have the Metric Terms Gone?

- SL form holds for finite displacements
- Reduces to Eulerian form as displacement ( $\Delta\lambda \equiv \lambda_A - \lambda_D$ ) and  $\Delta t \rightarrow 0$ :
- In this limit:

$$\sin \Delta\lambda \rightarrow \Delta\lambda, \quad \frac{\Delta\lambda}{\Delta t} \rightarrow \frac{u_A}{r_A \cos \phi_A}, \quad u_D \rightarrow u_A, \quad v_D \rightarrow v_A$$

and so:

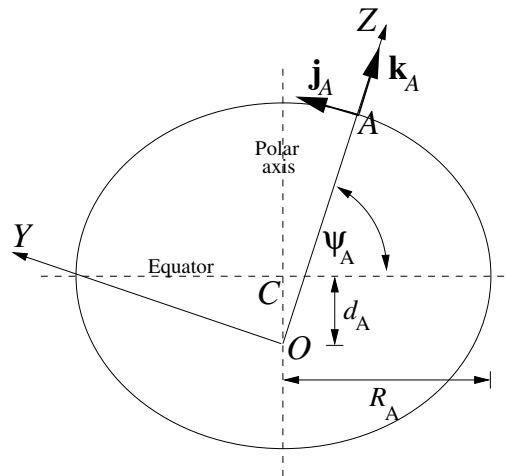
$$M_{12}v_D = v_D \sin \phi_D \sin \Delta\lambda \rightarrow \frac{u_A v_A \tan \phi_A}{r_A} \Delta t$$

- Each off-diagonal element of **M** generates one metric term



## Similar Oblate Spheroidal Form

- Basic derivation is the same
- “Trick”: use secondary meridional coordinate,  $\psi$ , *geographic latitude* (not orthogonal to vertical coordinate)



- **Significantly  $M$  has exactly the same form but with  $\phi \rightarrow \psi$**
- **Contrast Eulerian formulation where (for COS at least) two extra metric terms appear**

***Shallow-atmosphere...***

## Shallow-Atmosphere Approximation

$$\frac{D_{\textcolor{red}{r}}u}{Dt} - \frac{uv \tan \phi}{\textcolor{red}{r}} - (2\Omega \sin \phi) v + \frac{c_{pd}\theta_v}{\textcolor{red}{r} \cos \phi} \frac{\partial \Pi}{\partial \lambda} = \boxed{-\frac{uw}{\textcolor{red}{r}} - (2\Omega \cos \phi) w} + S^u$$

$$\frac{D_{\textcolor{red}{r}}v}{Dt} - \frac{u^2 \tan \phi}{\textcolor{red}{r}} + (2\Omega \sin \phi) u + \frac{c_{pd}\theta_v}{\textcolor{red}{r}} \frac{\partial \Pi}{\partial \phi} = \boxed{-\frac{vw}{\textcolor{red}{r}}} + S^v$$

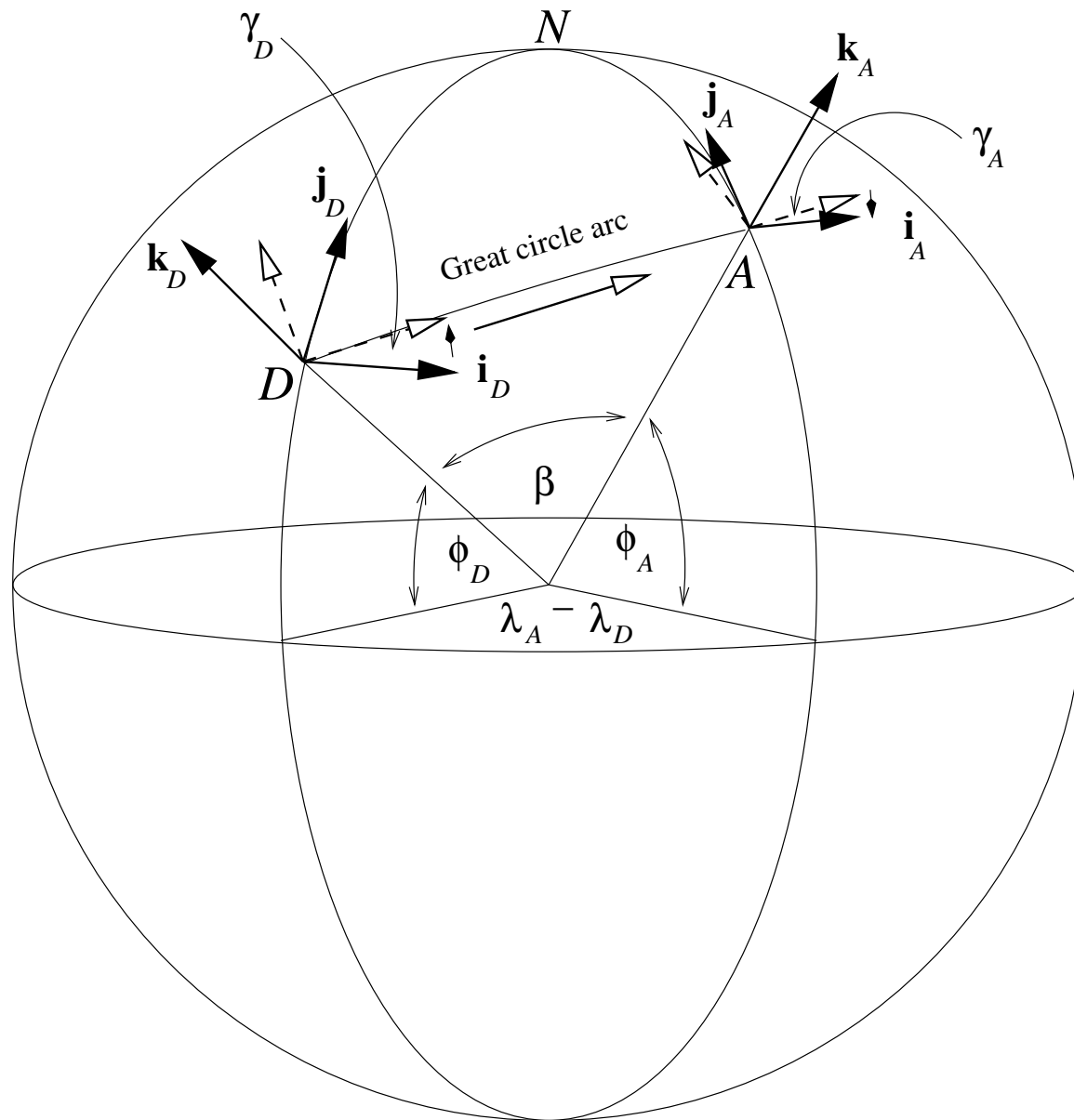
$$\boxed{\frac{D_{\textcolor{blue}{r}}w}{Dt}} + c_{pd}\theta_v \frac{\partial \Pi}{\partial r} + \frac{\partial \Phi}{\partial r} = \boxed{\frac{u^2 + v^2}{\textcolor{red}{r}} + (2\Omega \cos \phi) u} + S^w$$

### • Traditional approach:

- Eliminate all boxed red terms
- Replace  $\textcolor{red}{r}$  with  $a$  in all algebraic expressions

## Shallow-atmosphere - a different perspective

- **M is a rotation matrix**
- **Can be factored in infinite number of ways**
- **One is of particular interest:**
  - 1. Rotate departure-point UVT about radial direction to line up i with Great Circle connecting  $r_D$  with  $r_A$**
  - 2. Rotate UVT along Great Circle arc about new j direction**
  - 3. Rotate UVT about new radial direction to line up i direction**



**A PARTICULAR DECOMPOSITION OF  $M$**

## Shallow-atmosphere - a different perspective

- Result is  $M = ACD$
- **Shallow-atmosphere approximation is obtained by  $C \rightarrow I$**
- $M \rightarrow AD \equiv Q$ :

$$Q = \begin{pmatrix} p & q & 0 \\ -q & p & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

with

$$p = \frac{M_{11} + M_{22}}{1 + M_{33}}$$

and

$$q = \frac{M_{12} - M_{21}}{1 + M_{33}}$$

**giving the same result as in Temperton et al. (2001)**



*Departure points...*

## Rotation Matrix for Departure Points?

- Departure points now take on greater significance
- And they are also governed by a vector equation

$$\frac{D\mathbf{x}}{Dt} = \mathbf{v}$$

- So can we apply consistent approach (and avoid polar singularity issues)?
- Discrete vector form is

$$\mathbf{x}_D = \mathbf{x}_A - \frac{\Delta t}{2} (\mathbf{v}_A + \mathbf{v}_D)$$



## Rotation Matrix for Departure Points II

**As before component form can be written as**

$$\mathbf{M}\mathbf{x}_{\mathcal{D}} = \mathbf{x}_{\mathcal{A}} - \frac{\Delta t}{2} (\mathbf{v}_{\mathcal{A}} + \mathbf{M}\mathbf{v}_{\mathcal{D}})$$

**Or, since  $\mathbf{x}_{\mathcal{A}} = (0, 0, r_A)^T$  and  $\mathbf{x}_{\mathcal{D}} = (0, 0, r_D)^T$ :**

$$M_{13}r_{\mathcal{D}} = -\frac{\Delta t}{2} [u_{\mathcal{A}} + (M_{11}u_{\mathcal{D}} + M_{12}v_{\mathcal{D}} + M_{13}w_{\mathcal{D}})]$$

$$M_{23}r_{\mathcal{D}} = -\frac{\Delta t}{2} [v_{\mathcal{A}} + (M_{21}u_{\mathcal{D}} + M_{22}v_{\mathcal{D}} + M_{23}w_{\mathcal{D}})]$$

$$M_{33}r_{\mathcal{D}} = r_A - \frac{\Delta t}{2} [w_{\mathcal{A}} + (M_{31}u_{\mathcal{D}} + M_{32}v_{\mathcal{D}} + M_{33}w_{\mathcal{D}})]$$

## A Local Cartesian Transform Approach

The result can be cast into a local Cartesian form:

$$X_{DA} = -\frac{\Delta t}{2} (U_A + U_{DA})$$

$$Y_{DA} = -\frac{\Delta t}{2} (V_A + V_{DA})$$

$$Z_{DA} = r_A - \frac{\Delta t}{2} (W_A + W_{DA})$$

where

$$\mathbf{X}_{DA} \equiv \mathbf{M}\mathbf{x}_D,$$

and

$$\mathbf{V}_{DA} \equiv \mathbf{M}\mathbf{v}_D$$

are Departure point *coordinates* and *velocities* as seen in the Arrival-point Cartesian system

## The Inverse Transformation

**From  $(X_{DA}, Y_{DA}, Z_{DA})$  obtain departure point as:**

$$\tan(\lambda_A - \lambda_D) = \frac{-X_{DA}}{Z_{DA} \cos \phi_A - Y_{DA} \sin \phi_A}$$

$$r_D^2 = X_{DA}^2 + Y_{DA}^2 + Z_{DA}^2$$

$$\sin \phi_D = \frac{Y_{DA} \cos \phi_A + Z_{DA} \sin \phi_A}{r_D}$$

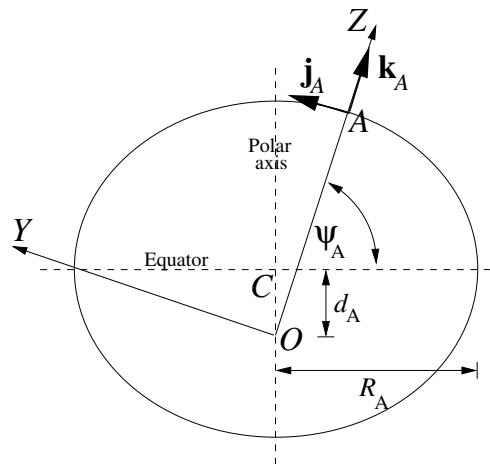
## Spheroidal Case

**Spheroidal case follows spherical case closely except:**

$$\mathbf{x}_A = (0, 0, r_A)^T - d_A (0, \cos \phi_A, \sin \phi_A)^T$$

**and**

$$\mathbf{x}_D = (0, 0, r_D)^T - d_D (0, \cos \phi_D, \sin \phi_D)^T$$



## Spheroidal Case

- Applying  $M$  as before, but to modified position vectors, results in

$$X_{DA} = M_{13}r_{\mathcal{D}}$$

$$Y_{DA} = M_{23}r_{\mathcal{D}} + (d_A - d_D) \cos \phi_A$$

$$X_{DA} = M_{33}r_{\mathcal{D}} + (d_A - d_D) \sin \phi_A$$

- Inversion is a little more complicated

## Shallow-atmosphere case

**Form obtained by constraining departure point to remain on arrival sphere:**

$$\mathbf{x}_D = \mathbf{x}_A - \frac{\Delta t}{2} (\mathbf{v}_A + \mathbf{v}_D) + \mathcal{B} (\mathbf{x}_A + \mathbf{x}_D)$$

**1. Cartesian transformation as before**

$$\mathbf{X}_{DA} \equiv \mathbf{M}\mathbf{x}_D,$$

**2. but velocities transform using momentum shallow-atmosphere rotation matrix**

$$\mathbf{V}_{DA} \equiv \mathbf{Q}\mathbf{v}_D,$$

**3. and**

$$\mathbf{X}_D^H = - \left( \frac{1 + M_{33}}{2} \right) \frac{\Delta t}{2} \left( \mathbf{v}_A^H + \mathbf{v}_D^H \right)$$

## What does it all mean?

**Single set of equations:**

$$(u_i - \alpha \Delta t \Psi_i)^{n+1} = M_{ij} (u_j + \beta \Delta t \Psi_j)_D^n$$

**with similar ones for departure point equations**

**Need only to specify:**

- 1.  $\nabla$ -operator in terms of the coordinate metric factors**
- 2. Components of gravity vector (only vertical is non-zero)**
- 3. Matrix elements  $M_{ij}$**
- 4. Components of Earth's rotation vector in chosen coordinate system**

# Similar Oblate Spheroidal and Spherical Polar Coordinates

	Similar oblate spheroids (SOS) ( $\varepsilon \equiv$ common eccentricity)	Spherical polars ( $\varepsilon = 0$ )
Coords	$(\lambda, \varphi, \mathbf{r})$ $\lambda =$ longitude $\varphi =$ meridional coordinate $\mathbf{r} =$ semi – major axis of spheroid	$(\lambda, \phi, r)$ $\lambda =$ longitude $\phi =$ latitude $r =$ radius of sphere
Metric factors	$h_\lambda = \left(\frac{\mathbf{r}}{h_r}\right) \cos \phi$ $h_\varphi = \left(\frac{\mathbf{r}}{h_r}\right) \left(\frac{\sin \phi \cos \phi}{\sin \varphi \cos \varphi}\right)$ $h_r = (1 - \varepsilon^2 \sin^2 \phi)^{1/2}$	$h_\lambda = r \cos \phi$ $h_\phi = r$ $h_r = 1$
Gravity	$g = \frac{g_a}{(1 - \varepsilon^2 \sin^2 \phi)^{1/2}} \left(\frac{a}{\mathbf{r}}\right)^2$	$g = g_a \left(\frac{a}{r}\right)^2$
Notes	$\phi =$ geographic latitude $\varphi \rightarrow \phi$ & $\mathbf{r} \rightarrow r$ as $\varepsilon \rightarrow 0$	Geographic latitude = spherical polar latitude when $\varepsilon = 0$



## Result is...

**Identical numerical system encompassing hierarchy of model approximations:**

- 1. Similar Oblate Spheroidal geopotential approximation**
- 2. Deep-atmosphere Spherical geopotential approximation  
(no latitudinal variation of  $g$ )**
- 3. Shallow-atmosphere approximation  
(non-Euclidean geometry)**
- 4. Cartesian**

**Hydrostatic versions apply to (2)-(4)**

## Where are we?

- Elements of approach form basis of deep-atmosphere, non-hydrostatic, spherical New Dynamics in MetUM
- **Next version of dynamical core (ENDGame) will incorporate full functionality**
- **Allow clean assessment of impact of various approximations**

## Summary

- **Can relax Spherical Geopotential Approximation**
  - Retain 2/3rds latitudinal variation of gravity
- **Semi-Lagrangian method can be straightforwardly applied to non-spherical coordinate systems**
  - Especially when use geographic latitude
- **Eulerian component equations can be derived from the rotation matrix forms**

## Summary

- **Matrix view gives insight into the non-Euclidean nature of the shallow-atmosphere approximation**
  - **neglect of great circle curvature**
- **Therefore desirable to avoid shallow-atmosphere approximation:**
  - **avoid loss of some (significant) Coriolis terms**
  - **and avoid non-Euclidean, space-distorting character**

## Questions?

### Further details can be found in:

- Spheroidal coordinate systems for modelling global atmospheres  
White, Staniforth and Wood 2008 *Q.J.R.Meteorol.Soc.* 134 pp 261-270
- Rotation matrix treatment of vector equations in  
semi-Lagrangian models of the atmosphere I: Momentum equation  
Staniforth, White and Wood. 2010 *Q.J.R.Meteorol.Soc.* 136 pp 497-506
- Rotation matrix treatment of vector equations in  
semi-Lagrangian models of the atmosphere II: Kinematic equation  
Staniforth, White and Wood. 2010 *Q.J.R.Meteorol.Soc.* 136 pp 507-516