

# PATH INTEGRAL MOLECULAR DYNAMICS: APPLICATIONS & NEW ALGORITHMS.

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IPAM CCS Workshop  
Los Angeles, CA

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# Outline

- ▶ Introduction to Path Integrals (PI)
- ▶ Double proton transfer in DNA base pair models
- ▶ Development of free energy methods in PIMD
- ▶ Improving the convergence of PIMD
- ▶ Summary



# Outline

Introduction to Path Integrals

Double proton transfer in DNA base pair models

Development of free energy methods in PIMD

Improving the convergence of PIMD

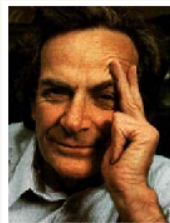
Summary

# Quantum Mechanics in Nature



## What can PI do for you?

- ▶ Delocalization
- ▶ Zero Point Energy (ZPE)
- ▶ Tunneling
- ▶ Theory development (QCD)
- ▶ quantum dynamics?



R. P. Feynman

[www.wikipedia.org](http://www.wikipedia.org) JCP **130**, 184105 (2009).

# PI Formulation of Canonical Partition Function

The canonical density matrix operator

$$\rho(\beta) = \exp(-\beta\mathcal{H}),$$

where  $\mathcal{H} = T + U$  and  $[T, U] \neq 0$ .

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where  $\mathcal{H} = T + U$  and  $[T, U] \neq 0$ .

$$Z(\beta) = \text{Tr} \left( e^{-\beta\mathcal{H}} \right) = \int dx_1 \langle x_1 | e^{-\beta(T+U)} | x_1 \rangle = \lim_{P \rightarrow \infty} \int dx_1 \langle x_1 | \Omega^P | x_1 \rangle,$$

where  $\Omega = e^{-\frac{\beta}{2P}U} e^{-\frac{\beta}{P}T} e^{-\frac{\beta}{2P}U}$ , which follows from Trotter theorem<sup>1</sup>

$$e^{-\beta(T+U)} = \lim_{P \rightarrow \infty} \left[ e^{-\frac{\beta}{2P}U} e^{-\frac{\beta}{P}T} e^{-\frac{\beta}{2P}U} \right]^P.$$

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<sup>1</sup>Tuckerman, *Stat. Mech.: Theory & Molecular Simulation*, Oxford, 2010

Inserting  $(P - 1)$  completeness relations  $I = \int dx |x\rangle \langle x|$



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$$\begin{aligned} Z(\beta) &= \lim_{P \rightarrow \infty} \int dx_1 dx_2 \cdots dx_P \langle x_1 | \Omega | x_2 \rangle \langle x_2 | \Omega | x_3 \rangle \cdots \langle x_P | \Omega | x_1 \rangle \\ &= \lim_{P \rightarrow \infty} \int \prod_{i=1}^P dx_i \langle x_i | \Omega | x_{i+1} \rangle |_{x_{P+1}=x_1}. \end{aligned}$$

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Need to compute

$$\begin{aligned} \langle x_i | \Omega | x_{i+1} \rangle &= \langle x_i | e^{-\frac{\beta}{2P} U} e^{-\frac{\beta}{P} T} e^{-\frac{\beta}{2P} U} | x_{i+1} \rangle \\ &= e^{-\frac{\beta}{2P} [U(x_i) + U(x_{i+1})]} \langle x_i | e^{-\frac{\beta}{P} T} | x_{i+1} \rangle. \end{aligned}$$

What can we do...?

Introduce complete set momentum eigenstates  $I = \int dp |p\rangle \langle p|$

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Now use  $\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar}$  in previous equality

$$\frac{1}{2\pi\hbar} \int dp e^{-\frac{\beta}{2mP} p^2} e^{ip(x_i - x_{i+1})/\hbar} = \sqrt{\frac{mP}{2\pi\beta\hbar^2}} e^{-\frac{mP}{2\beta\hbar^2} (x_{i+1} - x_i)^2},$$

where we have completed the square in the last expression.

Then, we obtain

$$\langle x_i | \Omega | x_{i+1} \rangle = \sqrt{\frac{m\mathbf{P}}{2\pi\beta\hbar^2}} e^{-\frac{m\mathbf{P}}{2\beta\hbar^2}(x_{i+1}-x_i)^2 - \frac{\beta}{2\mathbf{P}}[U(x_i)+U(x_{i+1})]}.$$

Then, we obtain

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### Path integral representation of canonical quantum PF

$$Z(\beta) = \lim_{P \rightarrow \infty} \left( \frac{mP}{2\pi\beta\hbar^2} \right)^{P/2} \int_{x_{P+1}=x_1} dx_1 \cdots dx_P e^{-\beta U_{\text{eff}}},$$

where  $U_{\text{eff}} = \sum_{i=1}^P \left[ \frac{mP}{2\beta^2\hbar^2} (x_{i+1} - x_i)^2 + \frac{1}{P} U(x_i) \right]$ .

Each quantum particle is mapped onto a collection of  $P$  interacting quasi-particles.<sup>1</sup>

<sup>1</sup>Chandler & Wolynes, JCP **74** 4078 (1981).

# The ring-polymer isomorphism

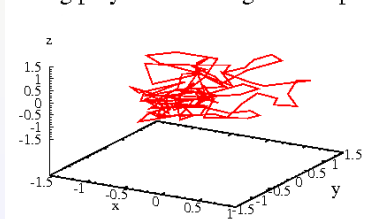
Analytical solutions to path integrals are limited:  
devise numerical techniques for complex systems.

**Discretized** version of the path integral quantum PF:

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where  $\omega_P^2 = P / (\beta\hbar)^2$  and  $x_{P+1} = x_1$ .

A ring polymer in configuration space



A quantum particle is mapped onto a collection of  $P$  “beads” coupled via harmonic springs & each under external potential  $U/P$ .



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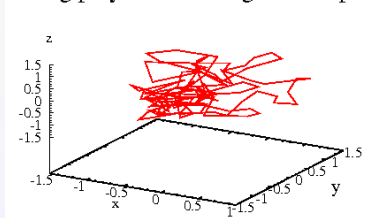
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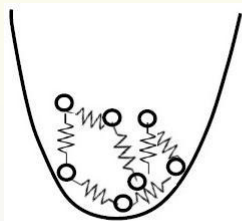


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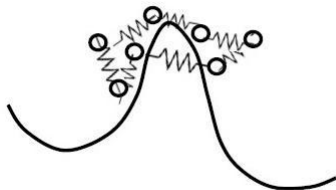
$$\text{centroid } x_c = \frac{1}{P} \sum_{i=1}^P x_i$$

# The ring-polymer isomorphism

Introducing nuclear quantum effects with the ring polymer.



zero-point energy



tunneling

[www.telluridescience.org/TullyLectures/tully.pdf](http://www.telluridescience.org/TullyLectures/tully.pdf)

# PIMD Implementation

Assuming ergodicity, sampling can be effected using MD. Introduce a set of uncoupled Gaussian integrals to enable sampling of phase space (**fictitious**  $m'_k$ )

$$Z(\beta) \approx \mathcal{N} \int dp_1 \cdots dp_P dx_1 \cdots dx_P e^{-\beta \sum_{k=1}^P \left[ \frac{p_k^2}{2m'_k} + \frac{m}{2} \omega_P^2 (x_k - x_{k+1})^2 + \frac{1}{P} U(x_k) \right]}.$$

- ▶ **Remember:** Time is only a parameter for the exploration of phase space (**unphysical dynamics**).
- ▶ *Naive* PIMD implementation suffers from **non-ergodicity** & multiple time scales issues<sup>1</sup>.

<sup>1</sup> Hall & Berne, JCP **81** 3641 (1984).

# Efficient PIMD Implementation

Solution M. Tuckerman *et al.* JCP **99** 2796 (1993).

- ▶ Transformation coordinates:  
normal modes  $u_k = \frac{1}{\sqrt{P}} \sum_{j=1}^P x_j e^{2\pi i(j-1)(k-1)/P}$  (centroid is  $u_1$ )  
or staging  $u_1 = x_1, u_k = x_k - \frac{(k-1)x_{k+1} + x_1}{k}, k = 2, \dots, P.$

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$$\tilde{H}(\mathbf{u}, \mathbf{p}) = \sum_{k=1}^P \left[ \frac{p_k^2}{2m'_k} + \frac{m_k}{2} \omega_P^2 u_k^2 + \frac{1}{P} U(x_k(\mathbf{u})) \right],$$

where  $m_1 = 0$ .  $m_k = \frac{k}{k-1}m$  for staging and  
 $m_k = 2mP [1 - \cos(2\pi(k-1)/P)]$  for normal modes.

- ▶ Fictitious masses chosen to collapse all time scales  
 $m'_1 = m, m'_k = m_k$
- ▶ Massive thermostats

Although computationally intensive, PIMD can be parallelized extremely well.

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**Double proton transfer in DNA base pair models**

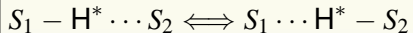
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# Proton transfer reactions

Ubiquitous and crucial process in nature

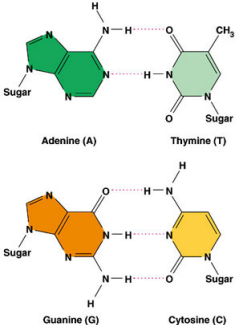
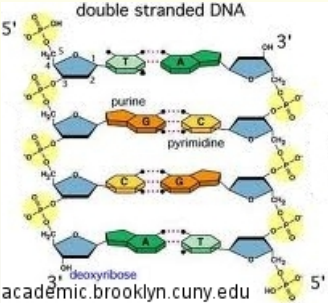


- ▶ Quantum Effects can be important even at room temperature!<sup>1</sup>
- ▶ Rearrangement of chemical bonds calls for *ab initio* description of the electronic part

<sup>1</sup> Tuckerman & Marx, PRL **86** 4946 (2001).

# Intermolecular proton transfer

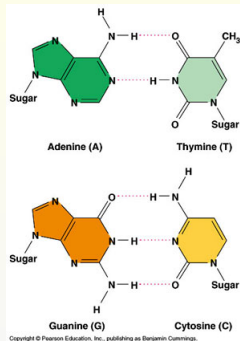
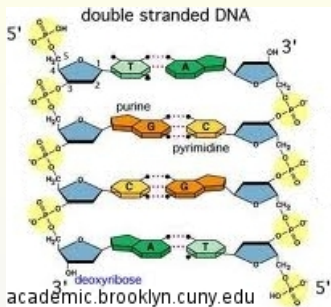
## DNA base pairs





# Intermolecular proton transfer

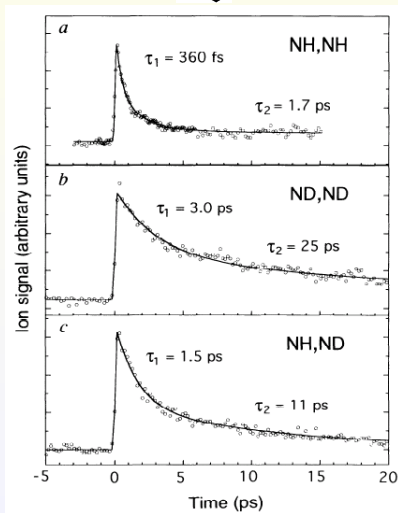
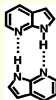
## DNA base pairs



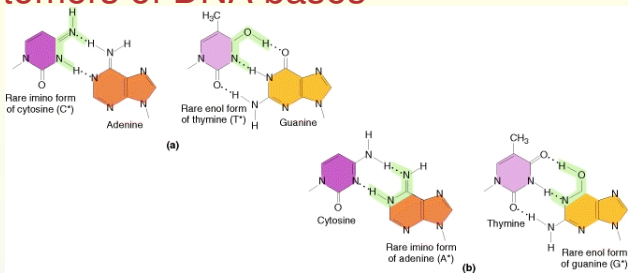
- ▶ “Rare” tautomers implicated in **DNA mutations**
- ▶ Proton **tunneling** was hypothesized to favor “rare” tautomers<sup>1</sup>
- ▶ Experimental evidence only for analogues (large **KIE**)<sup>2</sup>

<sup>1</sup> Löwdin, RMP **35** 724 (1963). <sup>2</sup> Zewail *et al.*, Nature **378** 260 (1995).

# Experiment on 7-azaindole dimer(g)

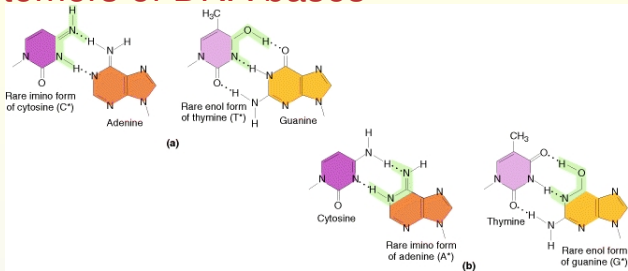


# Rare tautomers of DNA bases



<http://www.ncbi.nlm.nih.gov/>

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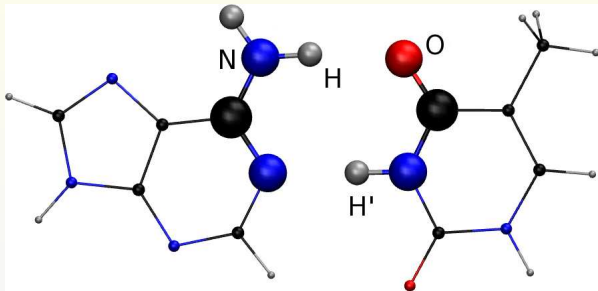


<http://www.ncbi.nlm.nih.gov/>



# Intermolecular double proton transfer

Formamidinium-Formamide (FIFA) tautomer complex<sup>1</sup>

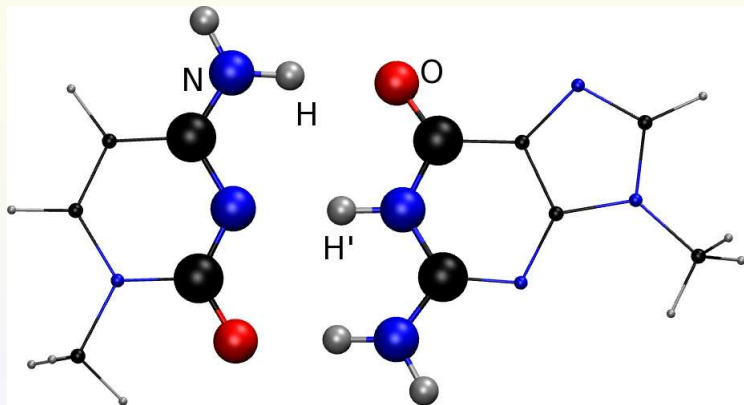


- ▶ Model system for A-T DNA base pair
- ▶ Concerted mechanism.

<sup>1</sup> J. Leszczynski *et al.*, JPCA **106** 12103 (2002)

# Intermolecular double proton transfer

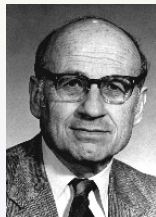
Model for G-C DNA base pair



# Forces from “First Principles”

Combining the best of two worlds...

- ▶ The quantum nature of nuclei is described using the **PI** formalism
- ▶ The electronic part (governing interaction between nuclei) is approximated using **DFT**



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Car & Parrinello Molecular Dynamics (CPMD), PRL **55** 2471 (1985).



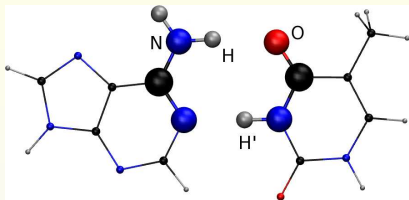
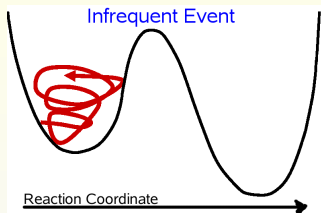
# Choice of DFT XC functional

Energetics (kcal/mol) from TURBOMOLE code/**TZVPP** basis set.  
Structures correspond  $\approx$  to reactants (R), the transition state (TS), and products (P)

method	GC model			AT model		
	$\Delta E_{P-R}$	$\Delta E_{TS-P}$	$\Delta E_{TS-R}$	$\Delta E_{P-R}$	$\Delta E_{TS-P}$	$\Delta E_{TS-R}$
RHF	13.50	14.65	28.15	11.80	9.47	21.27
MP2	9.37	4.04	13.41	8.62	4.37	12.99
CC2	<b>10.24</b>	<b>2.67</b>	<b>12.91</b>	<b>9.24</b>	<b>3.08</b>	<b>12.32</b>
PBE0	10.95	3.00	13.95	8.93	2.81	11.74
PBE	10.63	0.55	11.18	8.56	1.19	9.75
B3LYP	11.39	4.78	16.17	9.47	3.52	12.99
BLYP	<b>11.16</b>	<b>3.15</b>	<b>14.31</b>	<b>9.30</b>	<b>2.44</b>	<b>11.74</b>

**BLYP** represents best compromise between accuracy and cost!

# Free energy methods & rare events



- ▶ Reaction coordinate:  $\xi(\mathbf{r}) = d_{\text{HN}} - d_{\text{HO}}$
- ▶ Umbrella sampling<sup>1</sup> and WHAM technique<sup>2</sup>
- ▶ Drive system along reaction coordinate using harmonic bias potential (**centroid**)<sup>3</sup>  $\xi_c = \frac{1}{P} \sum_{i=1}^P \xi_i$

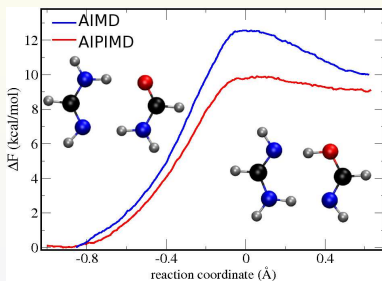
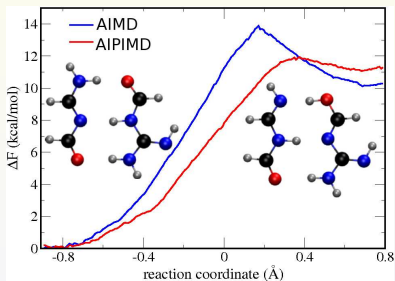
<sup>1</sup> Torrie & Valleau, Chem. Phys. Lett. **28** 578 (1974)

<sup>2</sup> Ferrenberg & Swendsen, Phys. Rev. Lett. **61** 2635 (1988)

<sup>3</sup> Gillan, PRL **58** 563 (1987). Voth, JCP **97** 8365 (1993).

# Free Energy Profiles Double Proton Transfer

Nuclear quantum effects have a dramatic impact on free energy profiles.



**Rare tautomers predicted dynamically unstable:  
therefore, not involved in mutations!**

# Free Energy Profiles Double Proton Transfer

**J | A | C | S**  
A R T I C L E S

## **Enol Tautomers of Watson–Crick Base Pair Models Are Metastable Because of Nuclear Quantum Effects**

Alejandro Pérez,<sup>†</sup> Mark E. Tuckerman,<sup>‡</sup> Harold P. Hjalmarson,<sup>¶</sup> and  
O. Anatole von Lilienfeld<sup>\*,¶</sup>

J. AM. CHEM. SOC. **132**, 11510 (2010).

[www.youtube.com](http://www.youtube.com)



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# Computing Thermodynamic Quantities: Estimators

An **estimator**: function whose average approximates a physical observable.

# Computing Thermodynamic Quantities: Estimators

An **estimator**: function whose average approximates a physical observable. **Example**:  $E(\beta) = -\frac{\partial}{\partial\beta} \ln Z(\beta)$

$$\langle E \rangle \approx \frac{1}{Z_P} \left( \frac{mP}{2\pi\beta\hbar^2} \right)^{P/2} \int dx_1 \cdots dx_P \epsilon_{prim} e^{-\beta \sum_{k=1}^P \left[ \frac{m}{2} \omega_P^2 (x_k - x_{k+1})^2 + \frac{U(x_k)}{P} \right]},$$

where the *primitive* estimator for energy is

$$\epsilon_{prim} = \frac{P}{2\beta} + \sum_{k=1}^P \left[ \frac{U(x_k)}{P} - \frac{m}{2} \omega_P^2 (x_k - x_{k+1})^2 \right]_{x_{P+1}=x_1}.$$

Use **MC**, **MD** to sample this configuration integral.

A more efficient estimator (**virial**) was proposed:

Herman *et al* JCP **76** 5150 (1982).

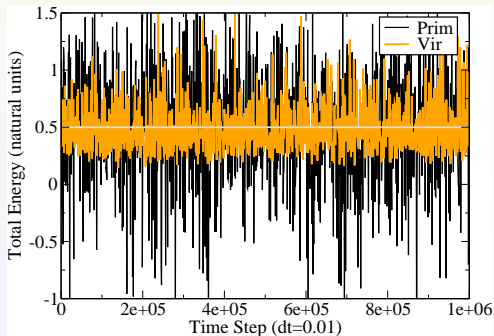


## Internal Energy: Virial vs Primitive Estimators

There is a more efficient estimator for the internal energy called the “**virial**”. Standard Deviation: 0.522(prim); 0.220(vir).

$$\epsilon_{vir} = \frac{1}{2P} \sum_{k=1}^P x_k \partial V / \partial x_k + \frac{1}{P} \sum_{k=1}^P V(x_k),$$

Instantaneous values of virial (orange) and primitive (black) estimators for  $V(x) = x^2/2$  at  $\beta = 10, P = 64$ .



# Isotope Effect via Thermodynamic Integration (TI)

**Motivation:** No PI method to study Isotope Effect (IE).  
Enzyme catalysis: proton transfer is the rate limiting step in many enzymatic reactions (**dehydrogenases**): large IE.  
Acid dissociation, etc.  
Typically, Site-H reacts faster than Site-D.

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$$\Delta F = F(m_f) - F(m_i) = \int_0^1 d\lambda \left( \frac{dF}{d\lambda} \right),$$

where  $F = -\frac{1}{\beta} \ln Z$  is the (quantum) free energy.  
For a linear path  $m(\lambda) = \lambda m_f + (1 - \lambda) m_i$ .

# Isotope Effect via Thermodynamic Integration (TI)

$$Z(\beta) \approx \prod_{i=1}^n \left( \frac{m_i^{\text{P}}}{2\pi\beta\hbar^2} \right)^{3\text{P}/2} \int d^{n\text{P}} \mathbf{r} e^{-\beta \sum_{i=1}^n \sum_{s=1}^{\text{P}} \left[ \frac{m_i}{2} \omega_{\text{P}}^2 (\mathbf{r}_{i,s} - \mathbf{r}_{i,s+1})^2 + \frac{U(\{\mathbf{r}_s\})}{\text{P}} \right]}.$$

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$$Z(\beta) \approx \prod_{i=1}^n \left( \frac{m_i^{\text{P}}}{2\pi\beta\hbar^2} \right)^{3P/2} \int d^{nP} \mathbf{r} e^{-\beta \sum_{i=1}^n \sum_{s=1}^P \left[ \frac{m_i}{2} \omega_{\text{P}}^2 (\mathbf{r}_{i,s} - \mathbf{r}_{i,s+1})^2 + \frac{U(\{\mathbf{r}_s\})}{\text{P}} \right]}.$$

Two PI estimators for  $\frac{dF}{d\lambda}$ :  
**primitive**

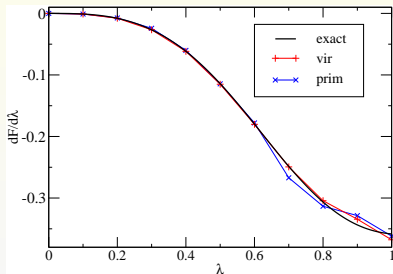
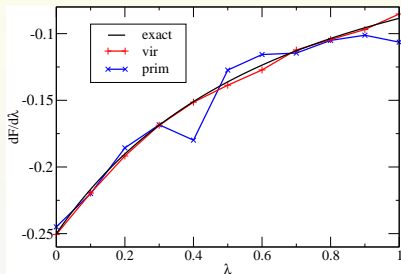
$$\left( \frac{dF}{d\lambda} \right)^{\text{prim}} = - \sum_{I=1}^N \left( \frac{m'_I}{m_I} \right) \left[ \frac{3P}{2\beta} - \left\langle \frac{m_I}{2} \omega_{\text{P}}^2 \sum_{s=1}^P (\mathbf{r}_{I,s} - \mathbf{r}_{I,s+1})^2 \right\rangle_{\lambda} \right],$$

and the **virial**

$$\left( \frac{dF}{d\lambda} \right)^{\text{vir}} = - \sum_{I=1}^N \left( \frac{m'_I}{m_I} \right) \left[ \frac{3}{2\beta} + \left\langle \frac{1}{2P} \sum_{s=1}^P \mathbf{r}_{I,s} \cdot \partial_{\mathbf{r}_I} V(\mathbf{r}_s) \right\rangle_{\lambda} \right].$$

# Isotope Effect via PIMD

Results for an analytically solvable system (harmonic potential)



First derivative of quantum free energy with respect to mass.

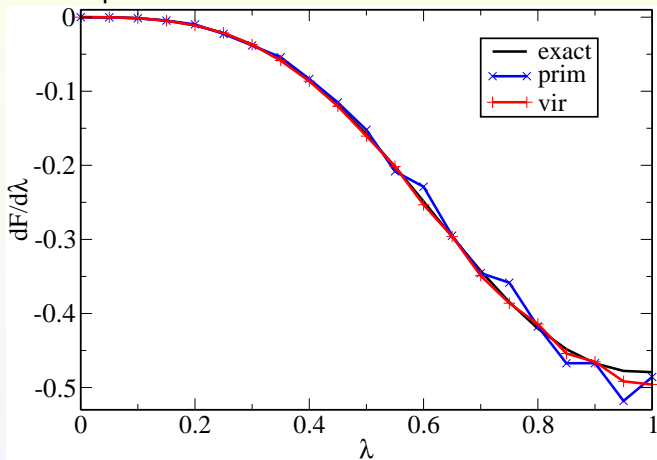
Left:  $m(\lambda) = \lambda m_f + (1 - \lambda) m_i$ .

Right:  $m(\lambda) = \lambda^4 m_f + (1 - \lambda^4) m_i$

**Black:** exact result. **Red:** virial and **Blue:** primitive estimators.

# Isotope Effect via TI-PIMD

Two-oscillator model for  $O - H(D) \cdots O$ . (Antisymmetric Mode)  
Useful to compute IE & understand **Ubbelohde** effect.

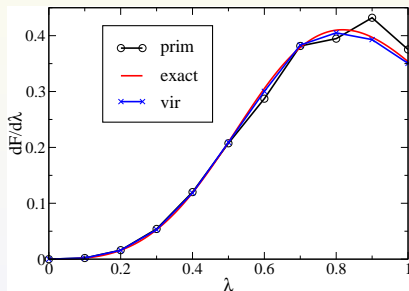
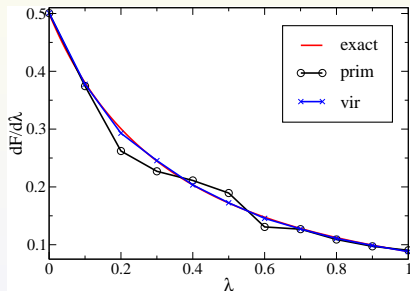


100,000 PIMD steps/window: prim=-0.2035, vir=-0.2052,  
exact=-0.2042.

# Simultaneous change in mass and potential

Change of identity/interaction potential: Useful to study H binding to active sites.

$$\left(\frac{dF}{d\lambda}\right) = \frac{\hbar}{4\sqrt{(\kappa/m)}} \coth\left(\frac{\beta\hbar}{2}\sqrt{\kappa/m}\right) \left[\frac{m\kappa' - m'\kappa}{m^2}\right].$$



First derivative of the quantum free energy with respect to mass & force constant for harmonic potential.

Left: Linear path. Right: quartic path.

**Black:** primitive result. **Red:** exact and **Blue:** virial estimators.



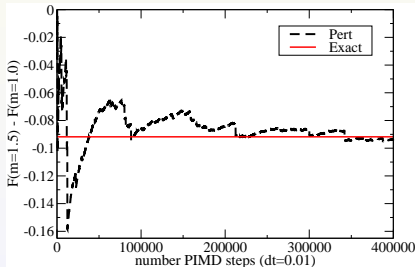
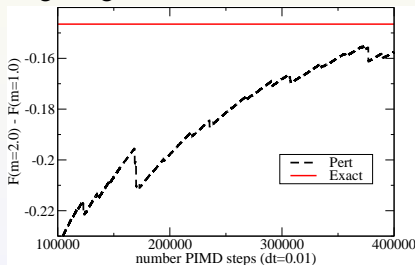
# Perturbation Theory

**Motivation:** PI calculations are expensive!

$$\Delta F = -\frac{1}{\beta} \left[ \frac{3NP}{2} \ln \left( \frac{m_f}{m_i} \right) + \ln \left\langle e^{-\beta(m_f - m_i) \frac{\omega_P^2}{2} A} \right\rangle_i \right],$$

where  $A = \sum_{I=1}^N \sum_{s=1}^P (\mathbf{r}_{I,s} - \mathbf{r}_{I,s+1})^2$ .

Full PT for free energy difference due to isotope transformation using single PIMD simulation.



Left:  $m_i = 1 \rightarrow m_f = 2$ . Right:  $m_i = 1 \rightarrow m_f = 1.5$ .

Exact: red horizontal line.

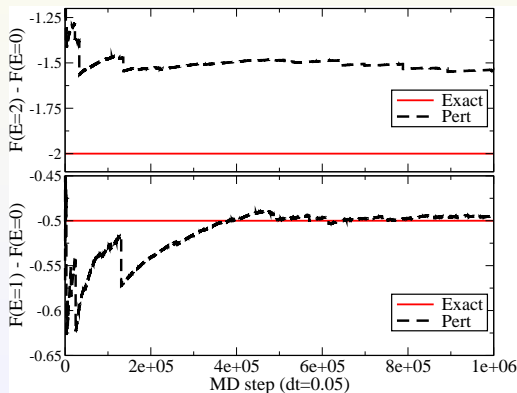
PT works provided change is small ... **Multistage-PT.**

# Perturbation Theory (II)

Effect of a static electric field:

$$\Delta F = -\frac{1}{\beta} \ln \langle \exp [-\beta (q\mathbf{E} \cdot \bar{\mathbf{r}})] \rangle_{E=0}.$$

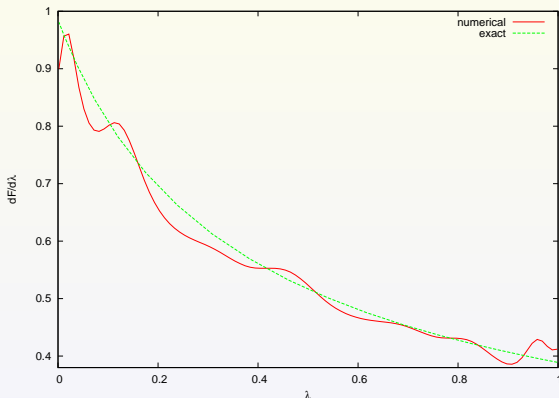
Useful to understand ferro/paraelectric hydrogen-bonded solids (KDP), Stark Effect.



# $\lambda$ -dynamics PIMD

Allow  $\lambda$  to vary in a continuous way using single PIMD simulation.

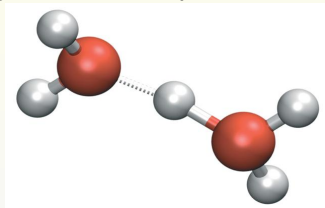
Example: transmutation of diatomics  $\text{H}_2 \rightarrow \text{HCl}$  at 300 K.



$\Delta F = F(\text{HCl}) - F(\text{H}_2) = -0.00332 \text{ a.u. (exact)}, -0.00332 \text{ a.u. (numerical)}$ .

# Isotope effects in the Zundel cation

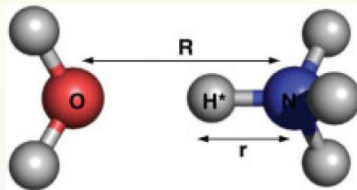
Important in hydrogen-bonded liquids & in enzymatic reactions.



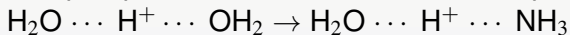
CPMD Details: *Ab initio* PIMD ( $\text{H}_2 \rightarrow \text{D}_2$ ).  
BLYP/100 Ry, Cubic Box  $12^3 \text{ \AA}^3$ ,  $P = 32$ ,  $T=300 \text{ K}$ .  
Numerical (PT)  $\approx 4 \text{ kcal/mol}$ .  
CPL **329**, 36 (2000).

# Goal: change of identity in PIMD simulations

Feasible in force fields-PIMD, harder in *ab initio* PIMD (pseudos).



Example: protonated water/ammonia complex



Towards predicting changes in free energy barriers heights?

Predict kinetic isotope effect  $k_H/k_D \approx e^{-\beta(\Delta F_H^\ddagger - \Delta F_D^\ddagger)}$ .



# Outline

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Double proton transfer in DNA base pair models

Development of free energy methods in PIMD

**Improving the convergence of PIMD**

Summary

# Higher-order methods to improve PIMD averages





# Higher-order methods to improve PIMD averages



**IDEA** Why not to use the classical forces to improve averages?

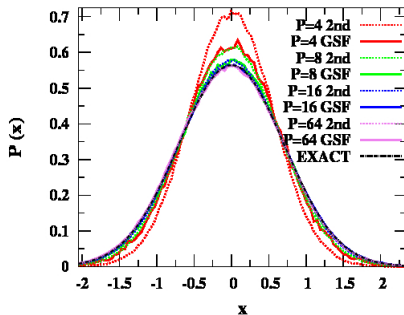
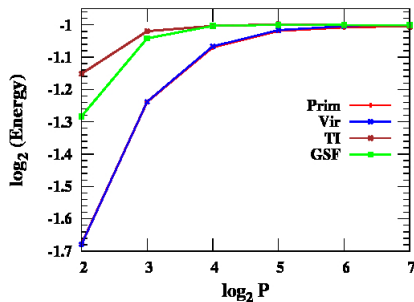
$$\langle \hat{O} \rangle \approx \frac{\langle O(\mathbf{r}_1) w(\mathbf{r}_1, \dots, \mathbf{r}_P) \rangle}{\langle w(\mathbf{r}_1, \dots, \mathbf{r}_P) \rangle}.$$

$w(\{\mathbf{r}\}; \beta)$  weighting function constructed using classical forces.  
In some cases, fourth-order convergence can be attained  
**without** additional overhead!

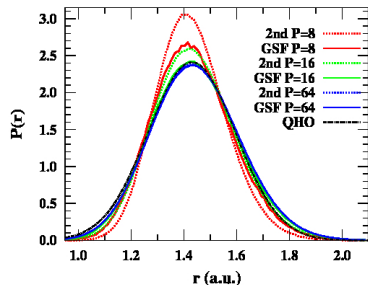
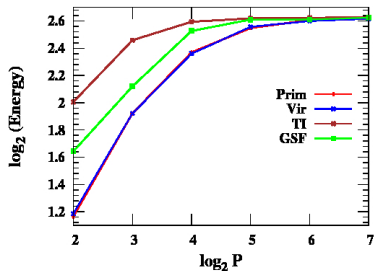
G. Voth et al, JCP 115, 7832 (2001).

A. Perez & M. Tuckerman, JCP (in review).

# Harmonic potential at $\beta\hbar\omega = 10$

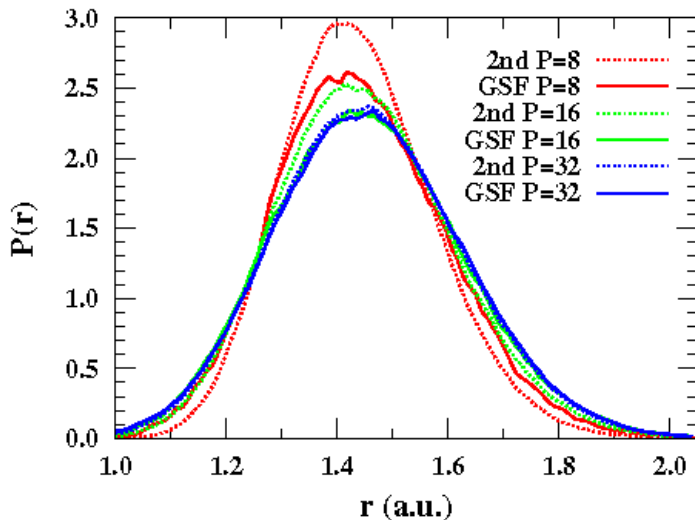


# Hydrogen molecule at 300K: $D [1 - \exp(-a(x - x_{eq}))]$



# Hydrogen molecule: *Ab initio* PIMD

Details:  $T=300\text{K}$ , CPMD BLYP/75 Ry, Cubic Box  $9^3 \text{ \AA}^3$ .



# High order PIMD methods

Liquid water at ambient temperature (q-SPC/Fw force field<sup>1</sup>)

P	2nd	4th
4	-0.2778	0.5989
8	0.0400	0.3911
16	0.2699	0.3892
32	0.3688	0.4055
64	0.3968	0.4067

Energy values in a.u.

- ▶ Useful to converge faster averages in quantum simulations of bulk phases.
- ▶ Very easy to implement in existing *ab initio* PIMD codes.
- ▶ No overhead associated.

<sup>1</sup> JCP **125**, 184507 (2006).

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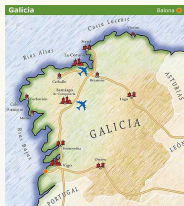
- ▶ The Feynman's path integral formalism was reviewed & applied to several systems
- ▶ Nuclear quantum effects are important even at room temperature
- ▶ “Rare” tautomers dynamically unstable & not implicated in mutations
- ▶ New estimators were derived & implemented to investigate alchemical transformations
- ▶ Higher-order algorithm to improve convergence of PIMD estimators

# Acknowledgments



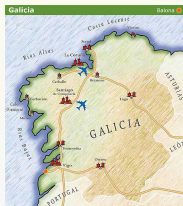
# Acknowledgments

The beautiful town of Baiona



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Thanks for your attention! Questions?