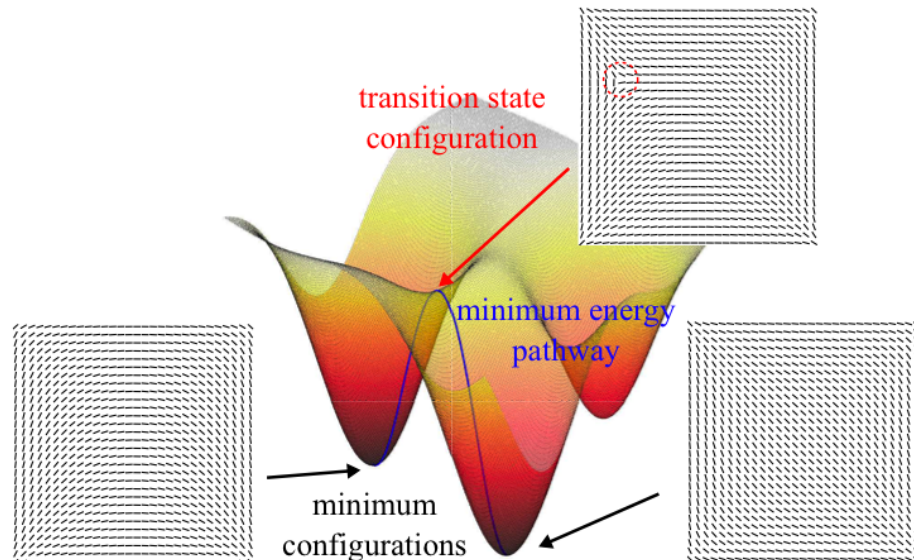


# Surveying Free Energy Landscapes: Applications to Continuum Soft Matter Systems

**Halim Kusumaatmaja**

Department of Physics, University of Durham



# Acknowledgement



**David Wales** (Cambridge)

Discussions on various numerical techniques



**Apala Majumdar** (Bath)

Collaboration on the LC example

# On the Menu Today

## The Basic Idea

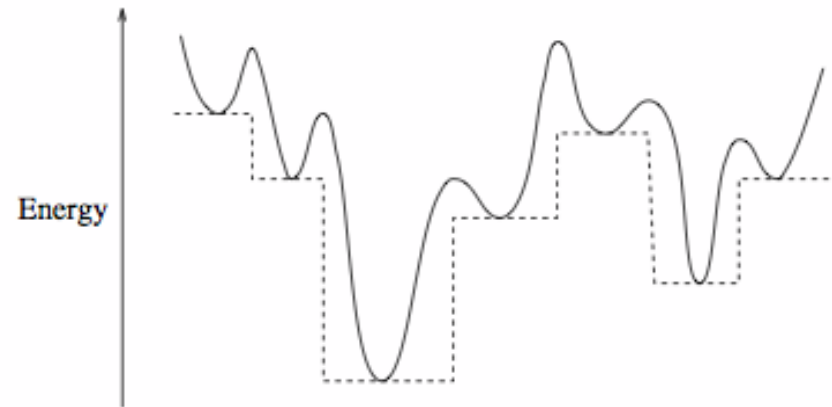
## Motivation - Typical questions

## Methods - Exploring the free energy landscapes

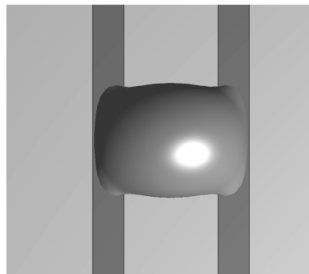
*Energy Minimization*

*Nudged Elastic Band Method*

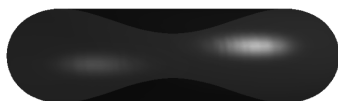
*Eigenvector Following*



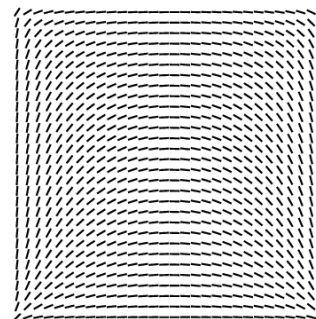
## Results - Specific Systems



*Wetting on Chemically Striped Surfaces*



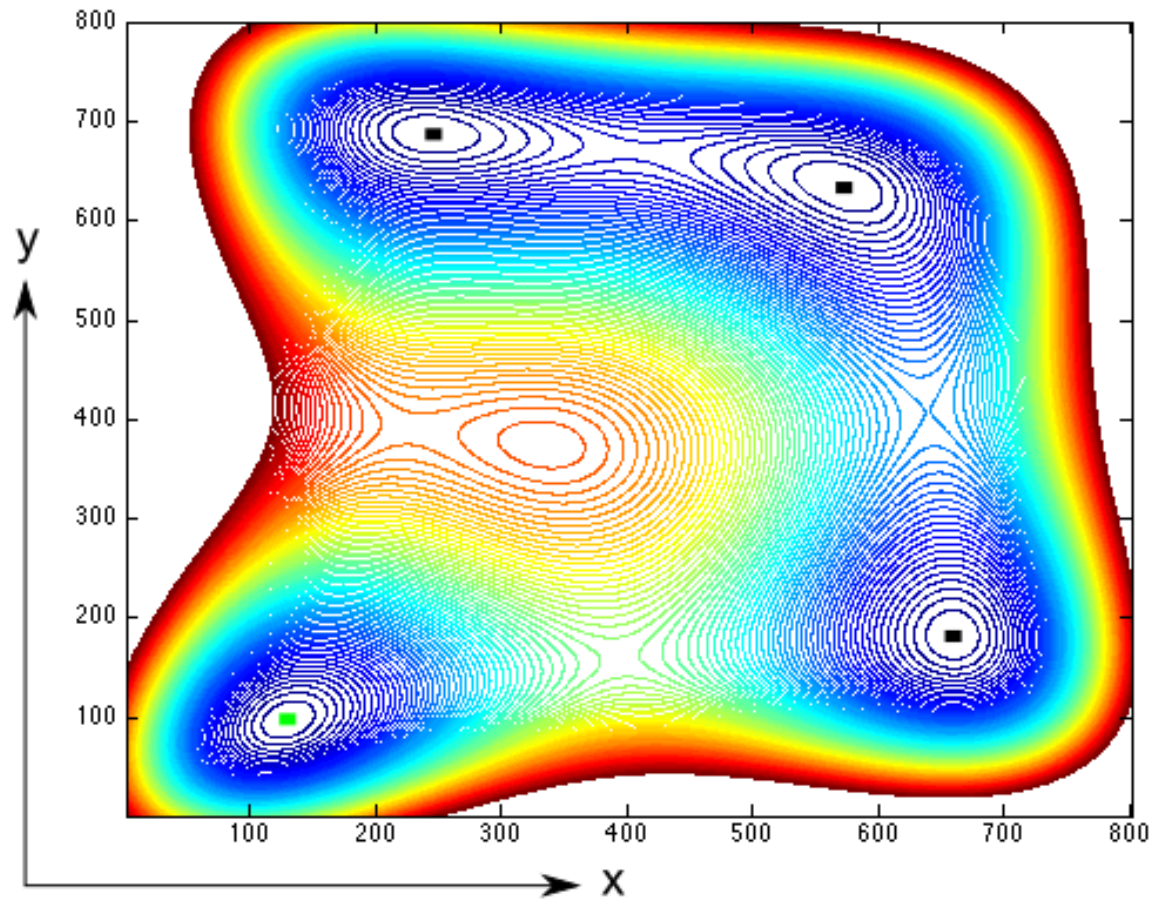
*Lipid Vesicles*



*Liquid Crystals*

# The Basic Idea

**Two-dimensional energy contour plot:  $F(x, y)$**

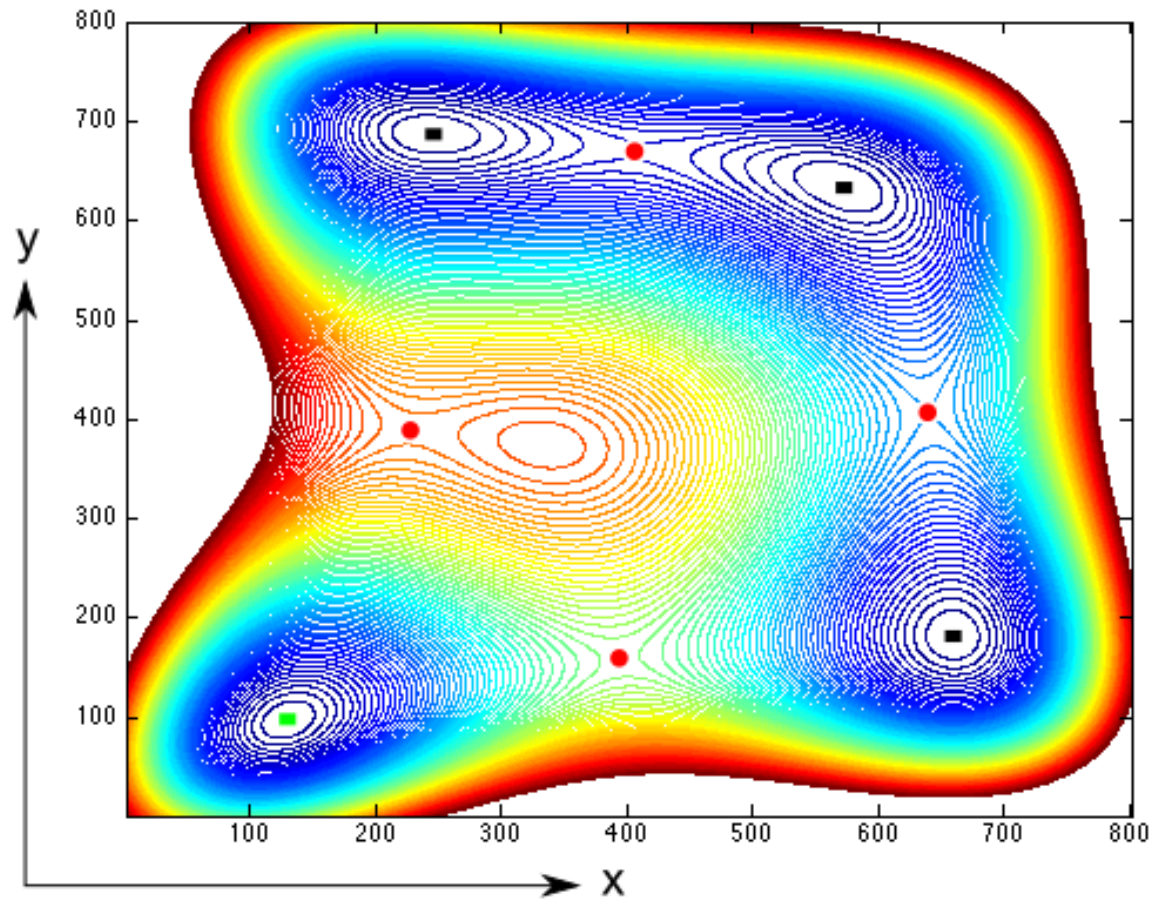


**The global minimum**

**All relevant minima**

# The Basic Idea

Two-dimensional energy contour plot:  $F(x, y)$



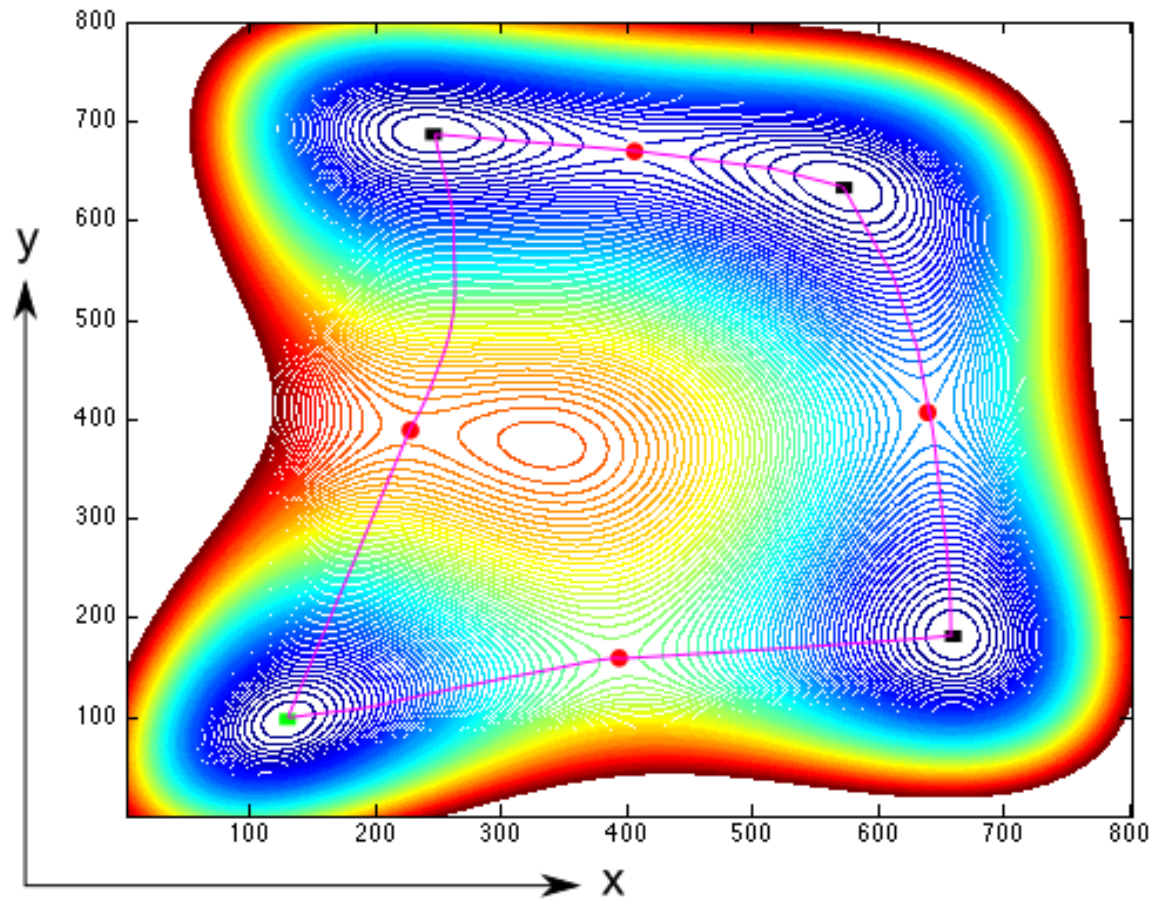
The global minimum

All relevant minima

Saddle point

# The Basic Idea

Two-dimensional energy contour plot:  $F(x, y)$



The global minimum

All relevant minima

Saddle point

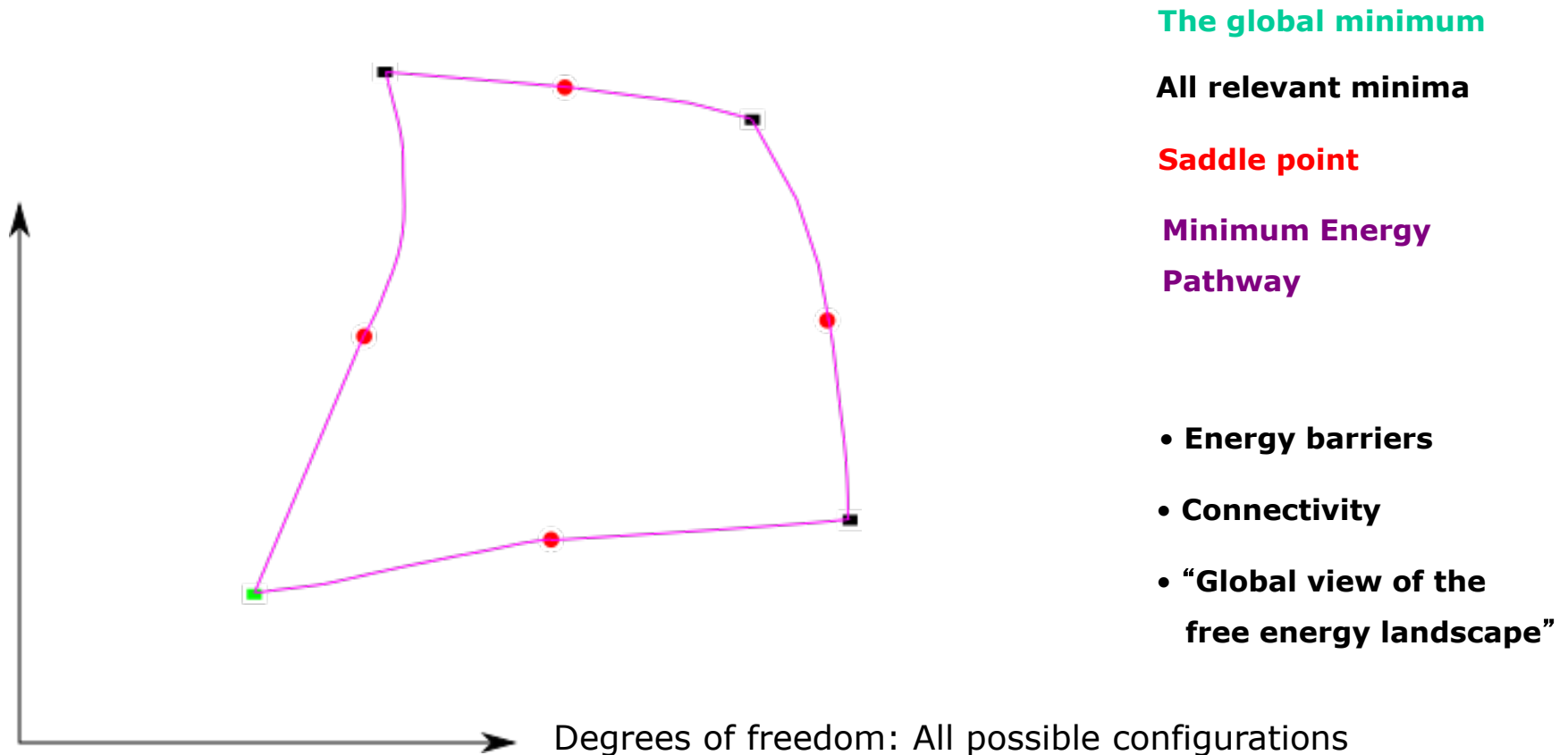
Minimum Energy  
Pathway

- Energy barriers
- Connectivity
- “Global view of the free energy landscape”

# The Basic Idea

## **Our Goal:**

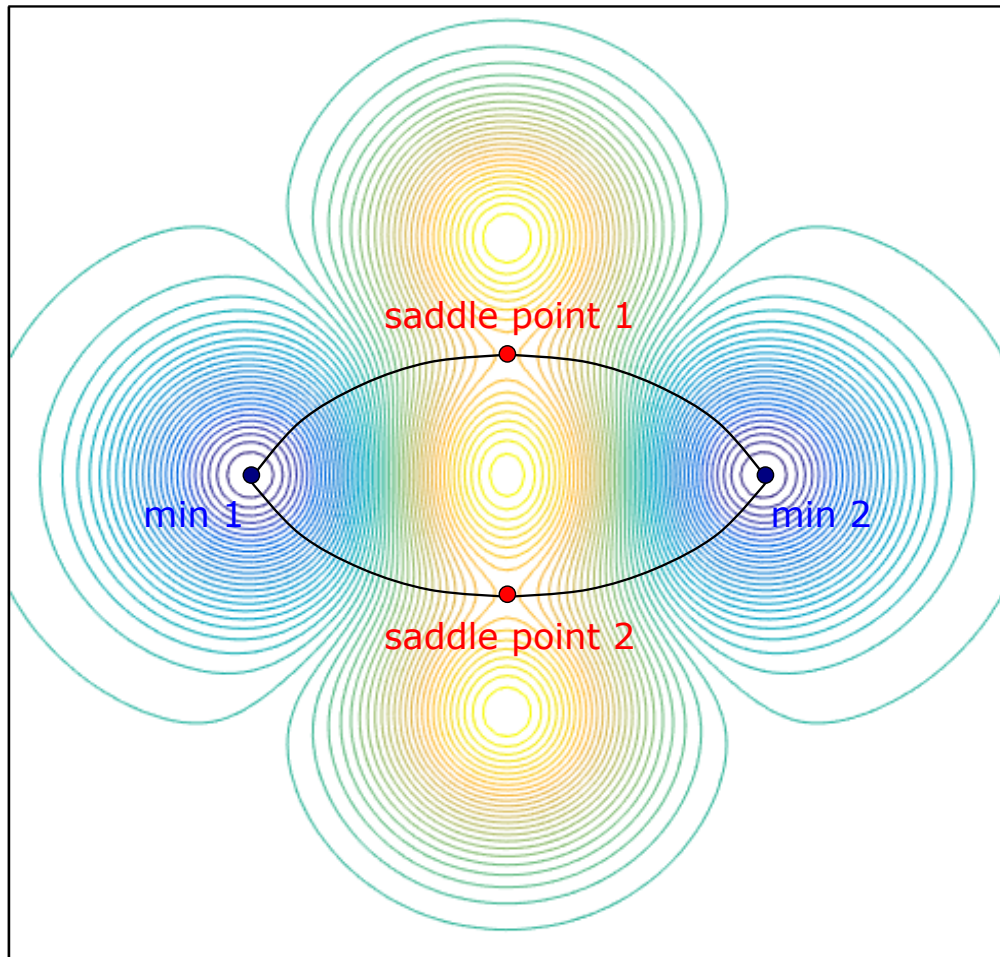
Study the evolution of the **free energy landscapes** as function of the system parameters.  
The skeleton represents perhaps the most interesting points on the landscape.



# Competing Pathways

***More than one possible minimum energy pathways***

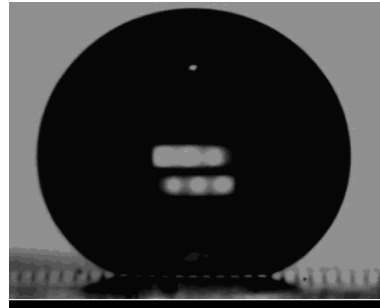
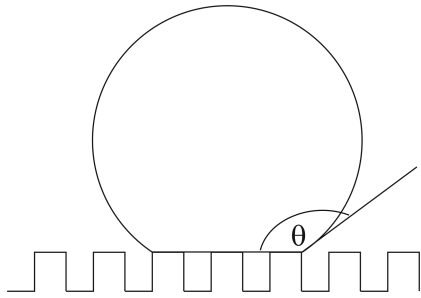
Systematic algorithm to enumerate these pathways?



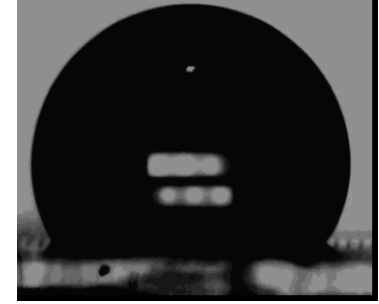
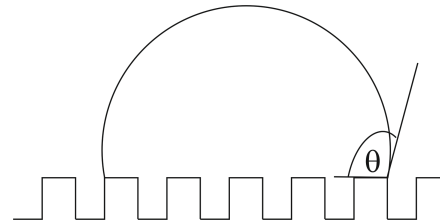
# Motivation

# Wetting States

- **Superhydrophobic Surfaces**

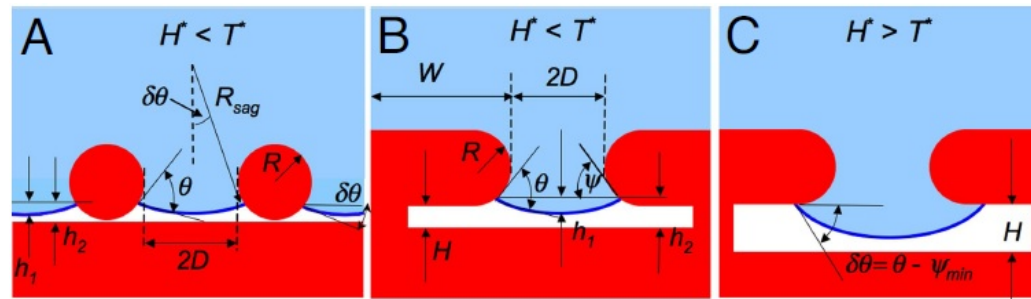


**Suspended - 'Good'**



**Collapsed - 'bad'**

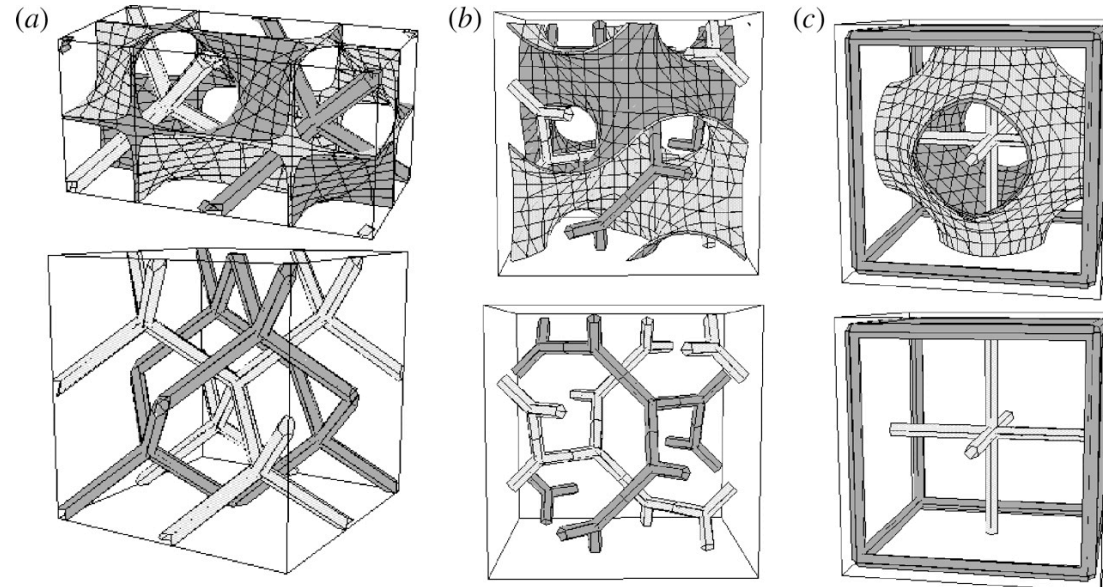
**Reentrant  
Post Geometry:**



Tuteja et al.

# Membranes

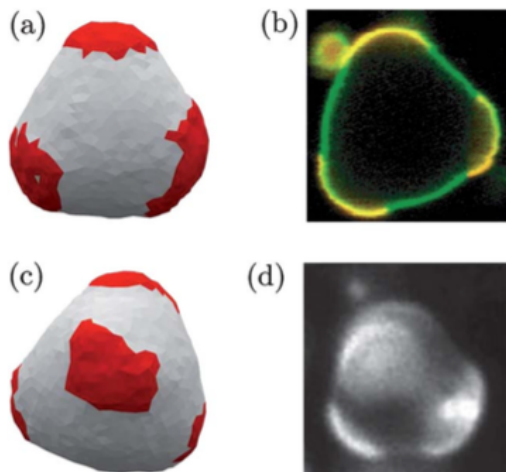
**D Surface** ↔ **G Surface** ↔ **P Surface**



Seddon et al.

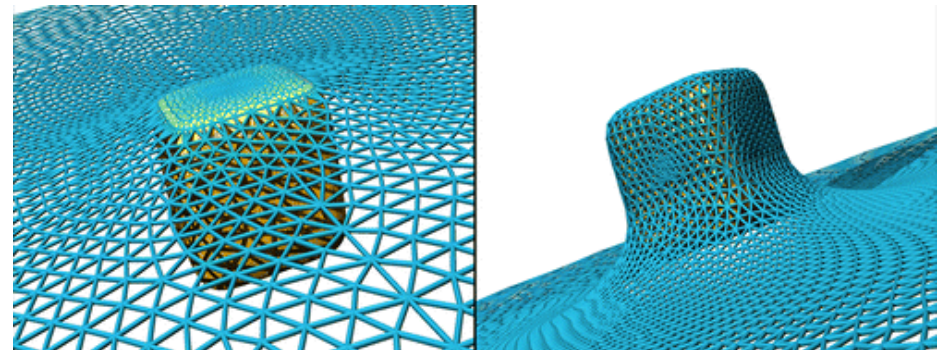
## • Cubic Membranes

## • Multicomponent Membranes



Hu et al.

## • Encapsulation of Nanoparticles

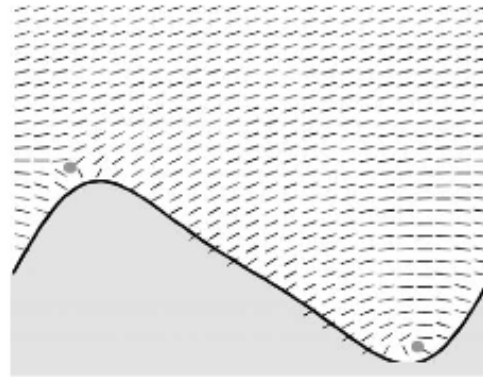


Dasgupta et al.

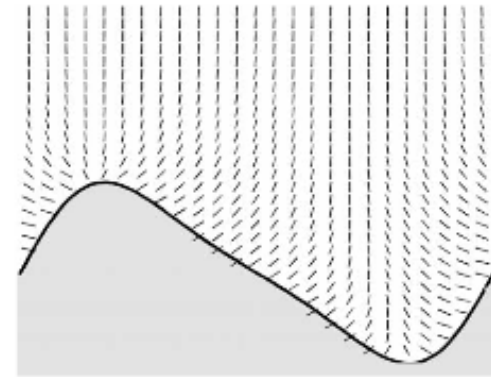
# Multistable Liquid Crystal Devices

## Zenithal Bistable Display

Spencer et al.

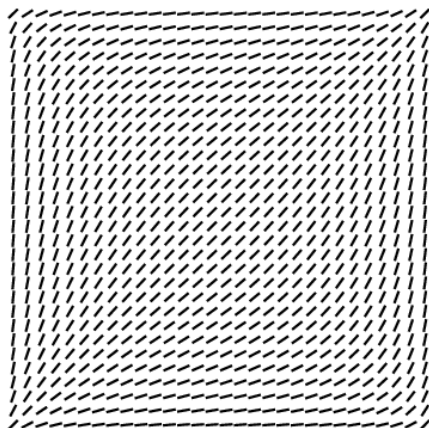


Defect State - bright

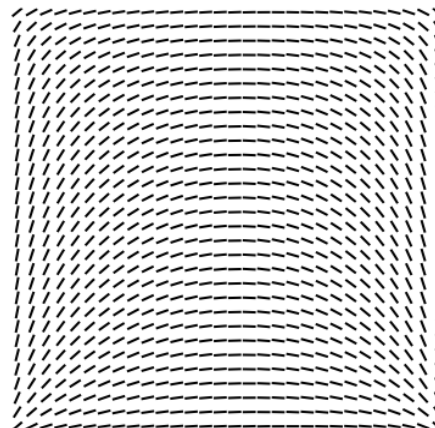


Continuous State - dark

## Nematic Liquid Crystal Wells

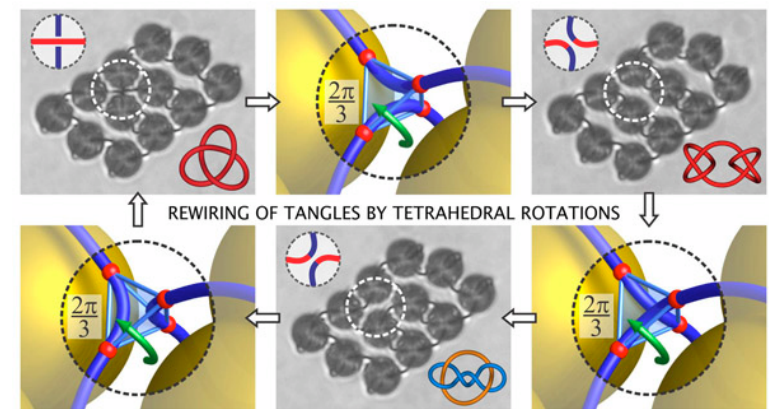


Diagonal State



Rotated State

## Liquid Crystal Colloids



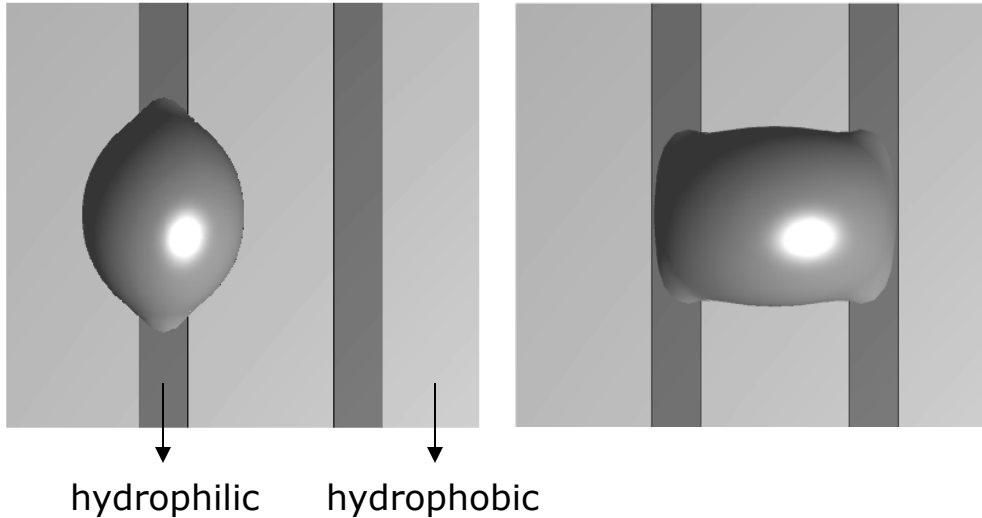
Tkalec et al.

# Methods

***H. Kusumaatmaja***, J. Chem. Phys. **142**, 124112 (2015)

# Wetting on Chemically Striped Surfaces

## Wetting on Striped Surfaces



## Contact Angle

$$\cos \theta = \frac{\gamma_{SV} - \gamma_{SL}}{\gamma_{LV}}$$

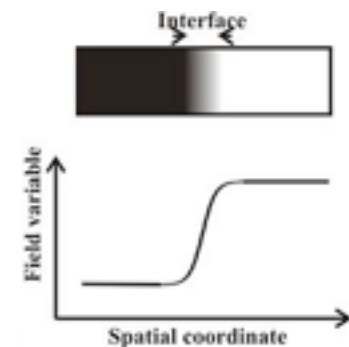
## Free Energy Function

Surface Energies       $\Psi = \gamma_{LV} A_{LV} + \gamma_{SV} A_{SV} + \gamma_{SL} A_{SL}$

Landau Free Energy       $\Psi = \int_V \left( \psi_b + \frac{\varepsilon}{2} |\nabla \phi|^2 \right) dV + \int_A \psi_A dA$

Double well potential       $\psi_b = \frac{1}{4\varepsilon} (\phi^2 - 1)^2$

Surface Energy       $\psi_A = -h\phi_A$



# Continuum vs Discrete

## Continuum

$$\Psi = \int_V \left( \psi_b + \frac{\epsilon}{2} |\nabla \phi|^2 \right) dV + \int_A \psi_A dA$$

## Minimum/Saddle Points

$$\text{1st Derivative} \quad \frac{\delta \Psi}{\delta \phi(\vec{x}_1)} = 0$$

$$\text{2nd Derivative} \quad \frac{\delta^2 \Psi}{\delta \phi(\vec{x}_1) \delta \phi(\vec{x}_2)}$$

## Discrete

$$\int_V \psi_b dV = \sum_{ijk} \frac{1}{\epsilon} \left( -\frac{1}{2} \phi_{ijk}^2 + \frac{1}{4} \phi_{ijk}^4 \right) \Delta x \Delta y \Delta z$$

$$\int_V \frac{\epsilon}{2} |\nabla \phi|^2 dV = \sum_{ijk} \frac{\epsilon}{2} |\nabla \phi|_{ijk}^2 \Delta x \Delta y \Delta z$$

$$\int_A -h \phi_A dA = \sum_{ijk \in A} -h \phi_{ijk} \Delta x \Delta y$$

## Minimum/Saddle Points

$$\text{1st Derivative} \quad \frac{\partial \Psi}{\partial \phi_{ijk}} = 0$$

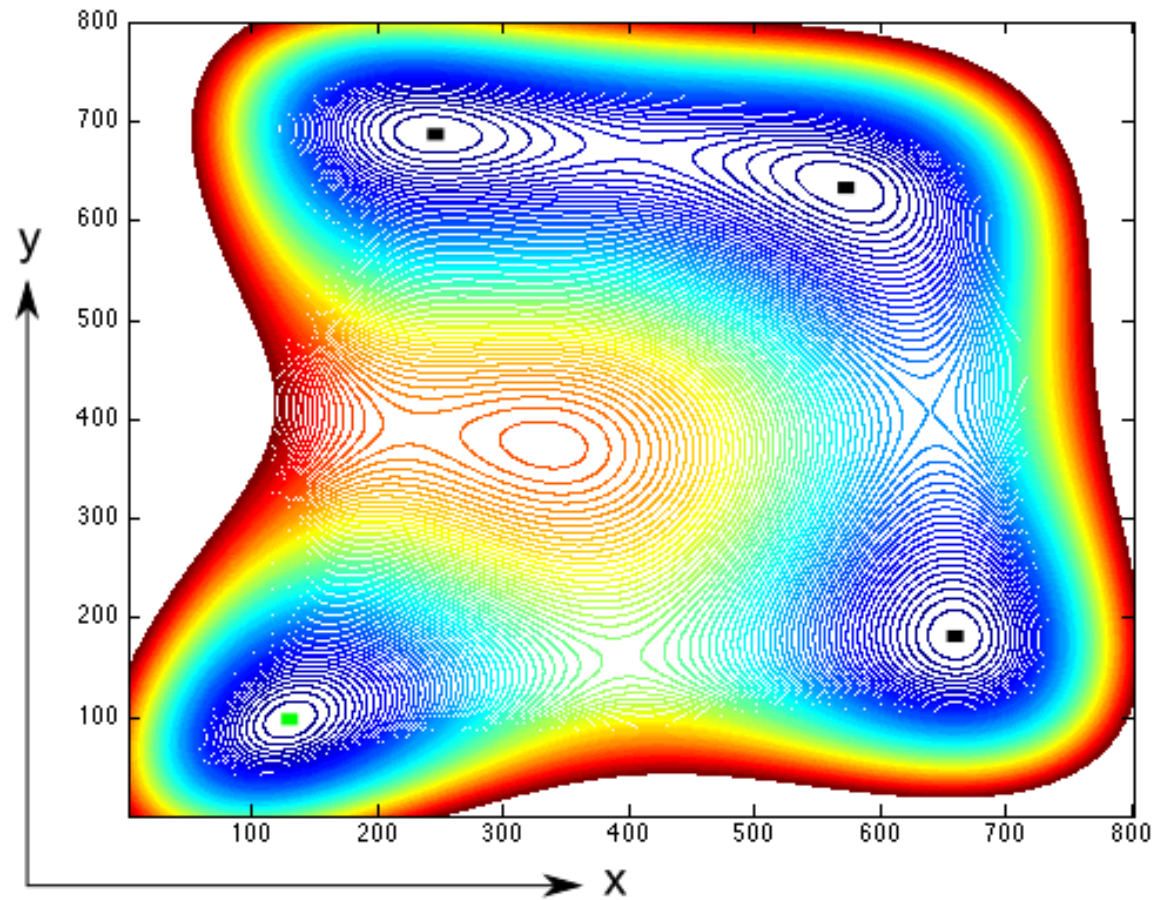
$$\text{2nd Derivative} \quad \frac{\partial^2 \Psi}{\partial \phi_{ijk} \partial \phi_{i'j'k'}}$$

**All positive eigenvalues** for a **minimum**.

**One negative eigenvalue**, the others are positive for a **transition state**.

Typically  $10^6$  degrees of freedom.

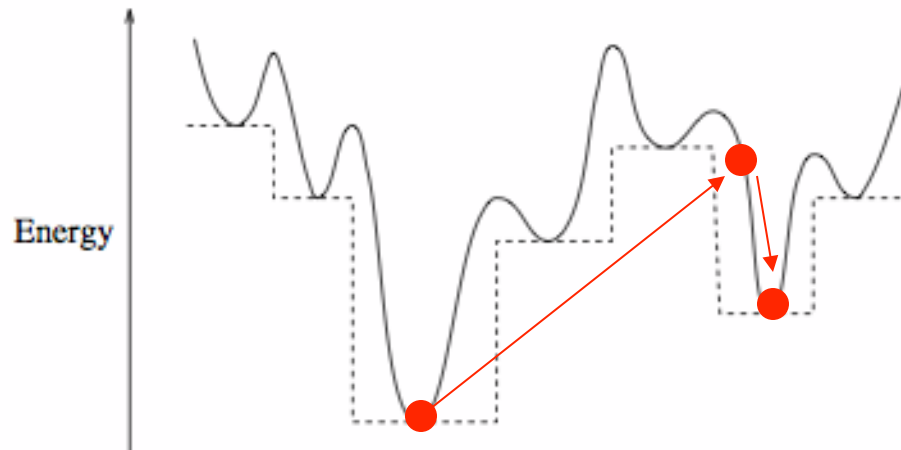
# Finding Minimum Free Energy States



The global minimum

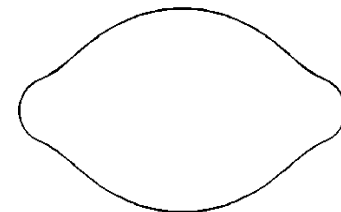
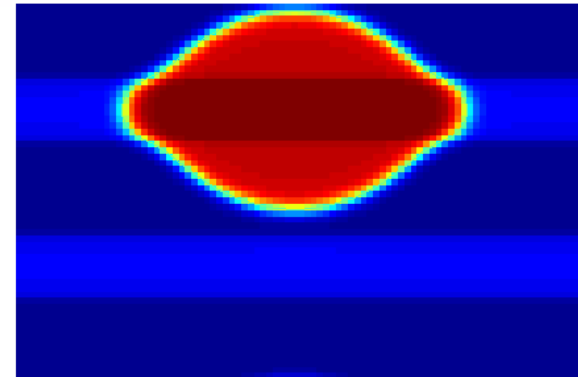
All relevant minima

# Basin-Hopping



## Three sub-steps:

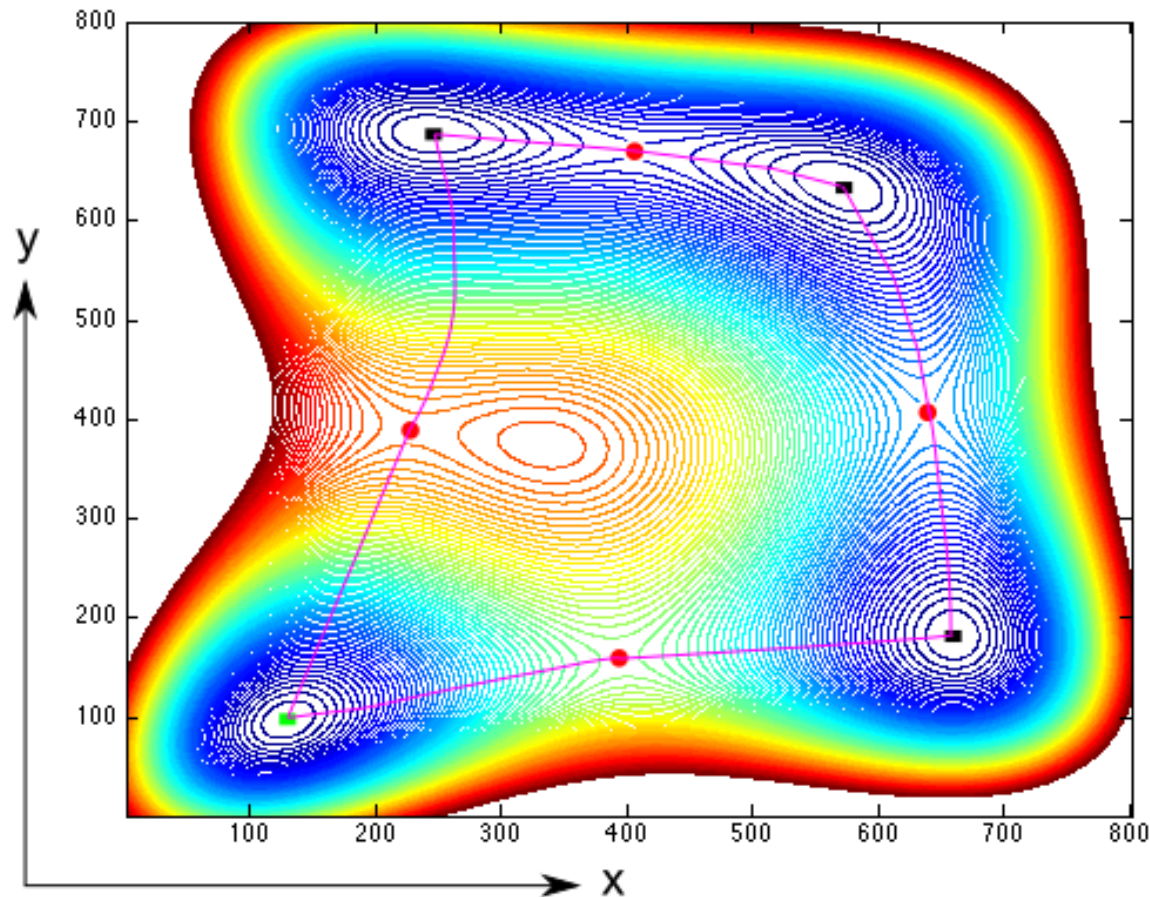
1. Random perturbations  
Option 1:  $\phi_{ijk} \rightarrow \phi_{ijk} + d \xi$   
Option 2: Deform the contour of the shape
2. Minimizations – use LBFGS algorithm (or others)
3. Acceptance criterion  
Option 1: use Metropolis algorithm  
Option 2: accept all minima (random search)



# Finding barriers and pathways

## Generally two classes of method:

1. Single ended methods – start from a minimum
2. Double ended methods – start with two minima of interests



The global minimum

All relevant minima

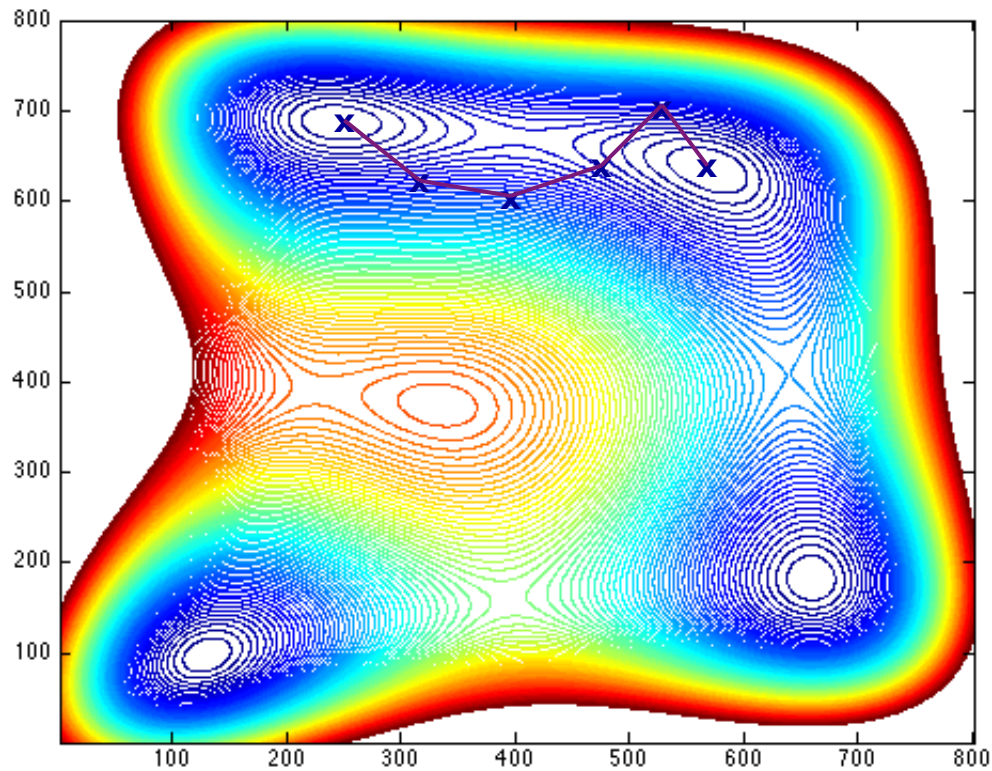
Saddle point

Minimum Energy  
Pathway

# (Doubly-) Nudged Elastic Band

## Basic Algorithm:

1. Make an initial guess for a set of images between the two minima

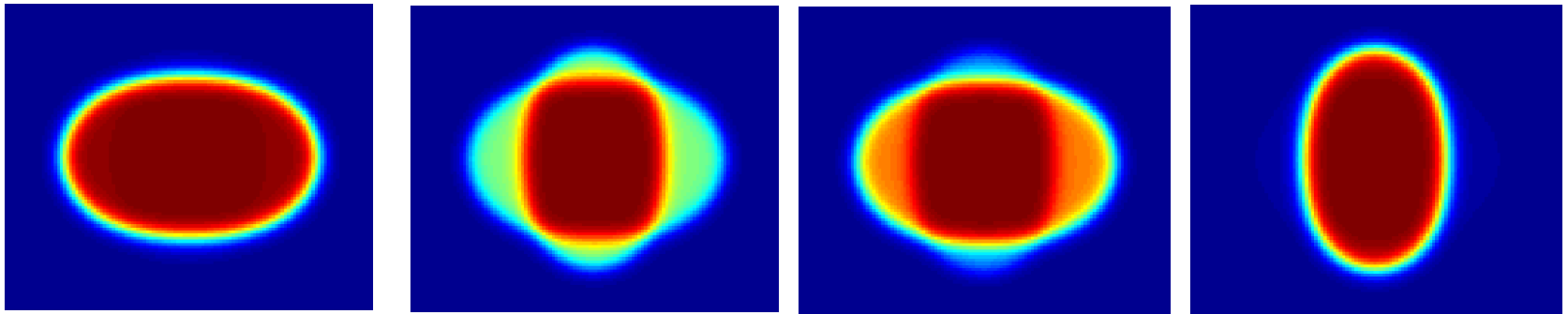


# (Doubly-) Nudged Elastic Band

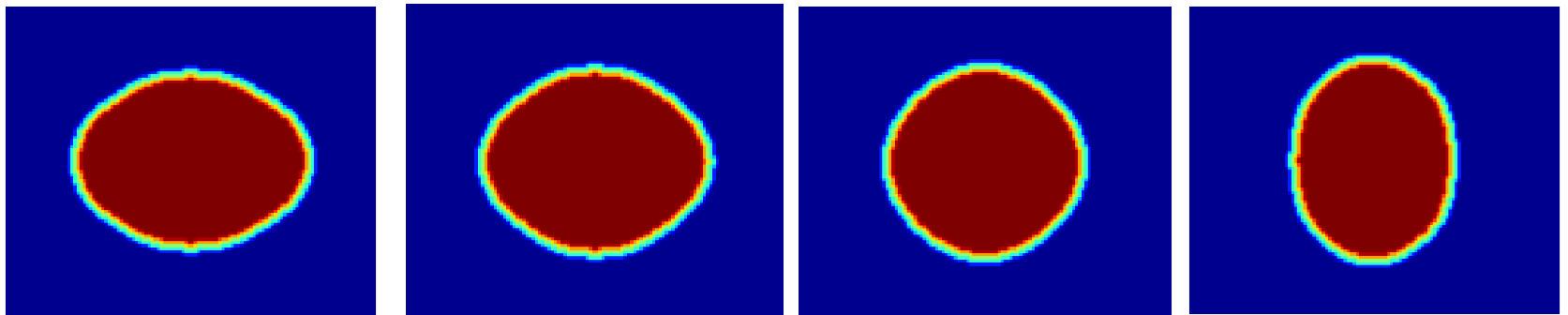
## Basic Algorithm:

1. Make an initial guess for a set of images between the two minima

## Linear Interpolations - Inefficient



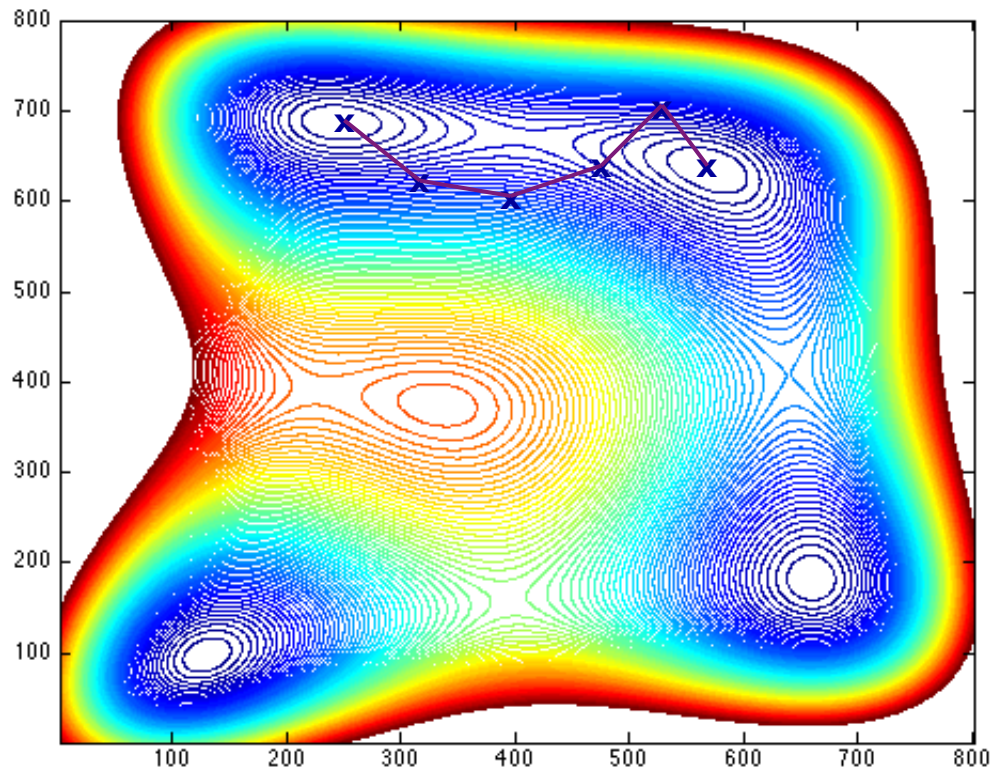
## Need to “Morph” the shape



# (Doubly-) Nudged Elastic Band

## Basic Algorithm:

1. Make an initial guess for a set of images between the two minima
2. Connect these images by springs



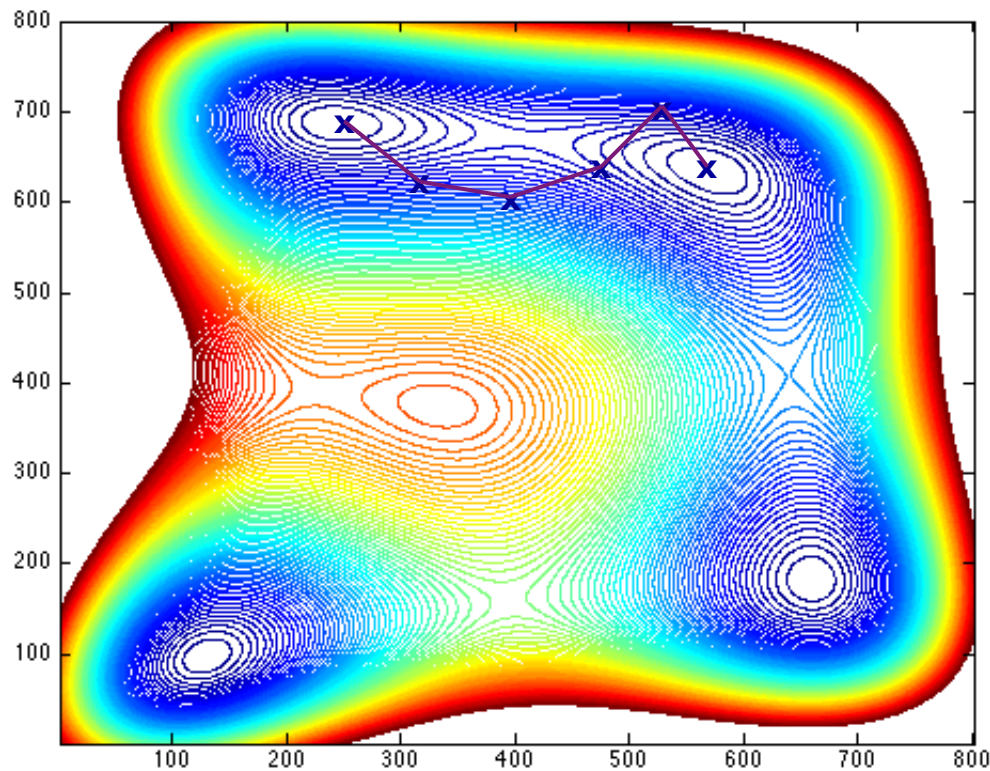
## Spring potential energy:

$$V_{spring}^{\alpha} = \frac{k}{2} \left( (s^{\alpha,-})^2 - (s^{\alpha,+})^2 \right)$$

# (Doubly-) Nudged Elastic Band

## Basic Algorithm:

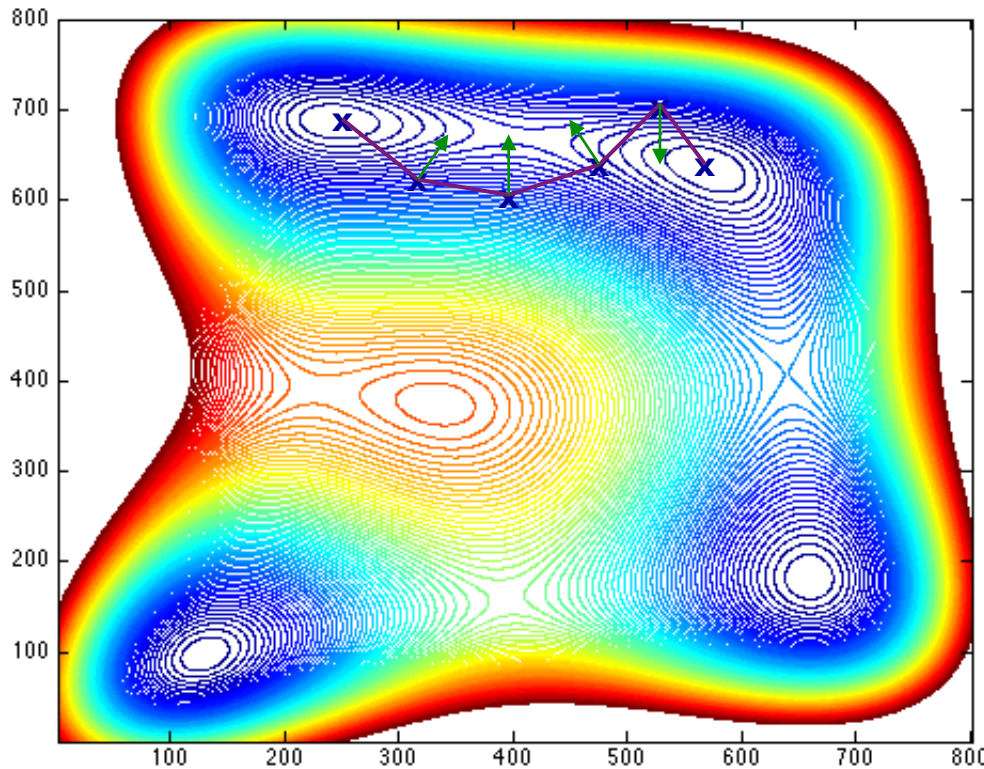
1. Make an initial guess for a set of images between the two minima
2. Connect these images by springs
3. Relax using projections from the “true gradient” and from the “spring force” component



# (Doubly-) Nudged Elastic Band

## Basic Algorithm:

1. Make an initial guess for a set of images between the two minima
2. Connect these images by springs
3. Relax using projections from the “true gradient” and from the “spring force” component



Perpendicular component of  
the true gradient:

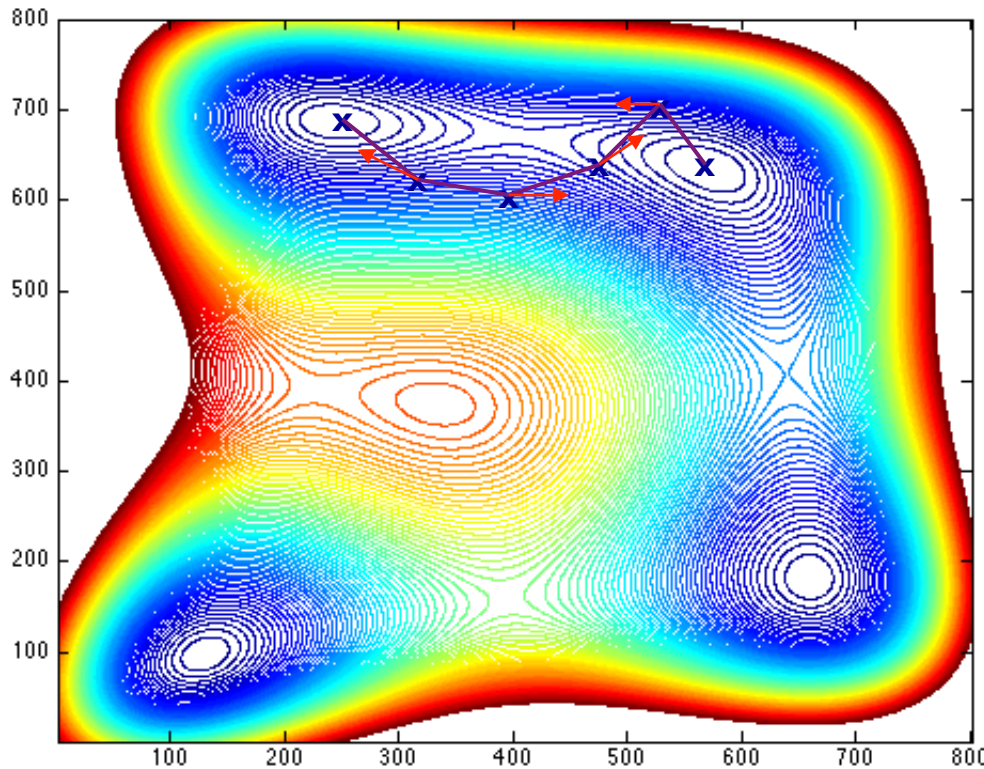
$$g_{\perp}^{\alpha} = g^{\alpha} - (g^{\alpha} \cdot \hat{t}^{\alpha}) \hat{t}^{\alpha}$$

Avoid all images collapsing to minima

# (Doubly-) Nudged Elastic Band

## Basic Algorithm:

1. Make an initial guess for a set of images between the two minima
2. Connect these images by springs
3. Relax using projections from the “true gradient” and from the “spring force” component



Parallel component of the spring force:

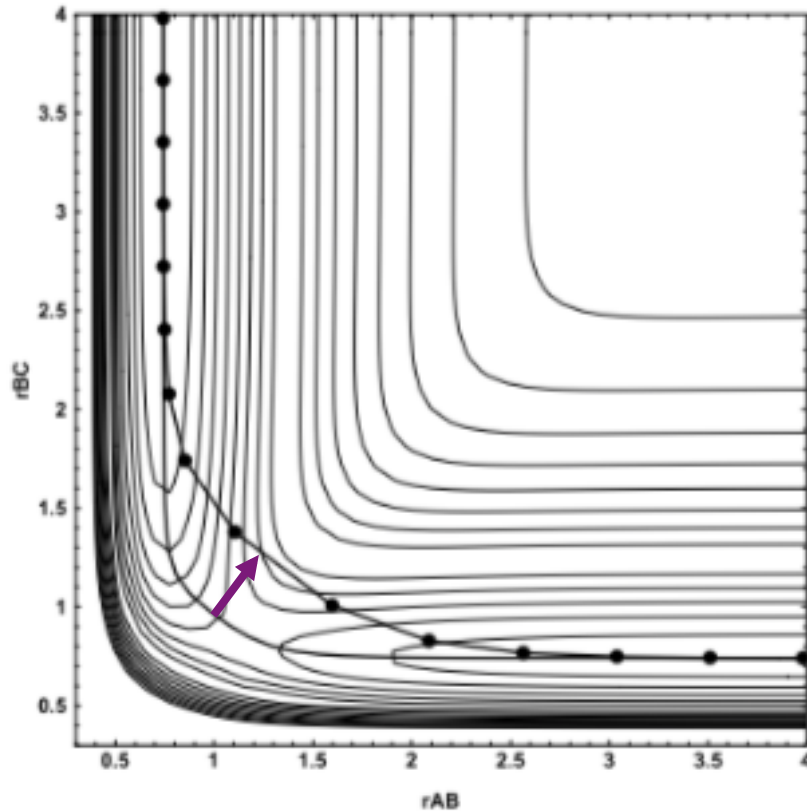
$$\tilde{g}_{\perp}^{\alpha} = k(s^{\alpha,+} - s^{\alpha,-})\hat{t}^{\alpha}$$

Maintain equal distance between images  
Avoid “corner-cutting”

# (Doubly-) Nudged Elastic Band

## Basic Algorithm:

1. Make an initial guess for a set of images between the two minima
2. Connect these images by springs
3. Relax using projections from the “true gradient” and from the “spring force” component



Parallel component of the spring force:

$$\tilde{g}_{\perp}^{\alpha} = k(s^{\alpha,+} - s^{\alpha,-})\hat{t}^{\alpha}$$

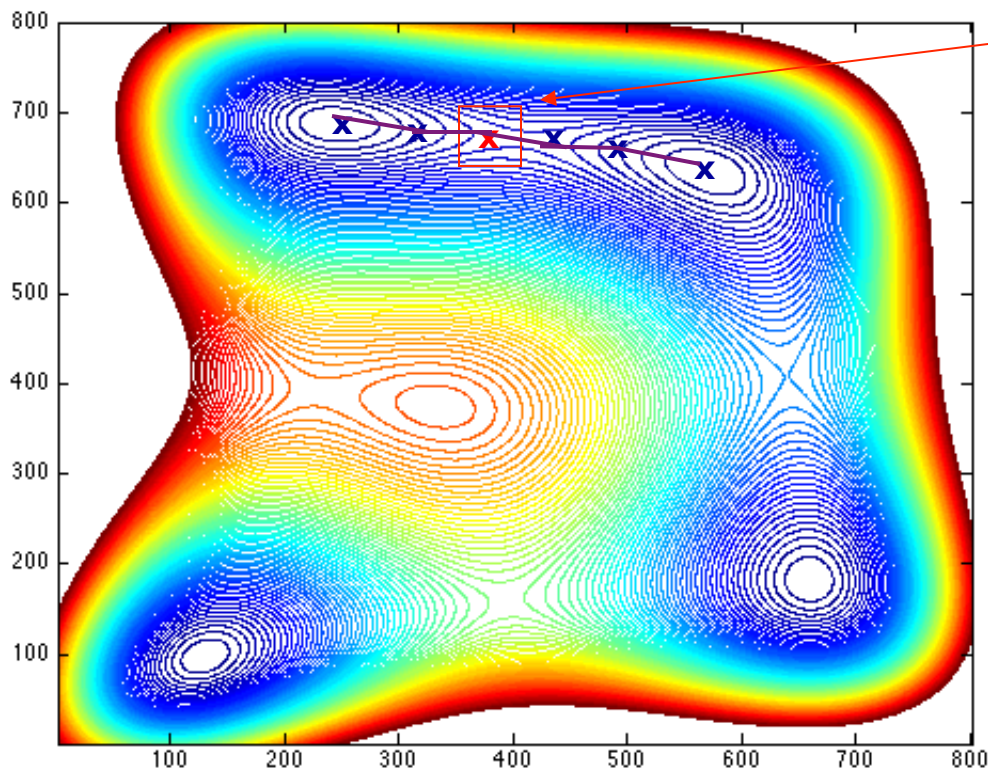
Maintain equal distance between images

**Avoid “corner-cutting”**

# Hybrid Eigenvector-Following

## Further optimise using hybrid eigenvector-following:

1. First, remember that we want to find saddle point of index 1
2. Uphill step in one eigendirection
3. Minimization in the tangent space



Candidate for a transition state  
Optimize using eigenvector-following

Obtain a “**minimum energy path**”

Keep iterating until all minima  
and transition states are found!

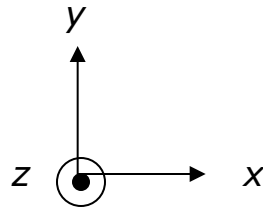
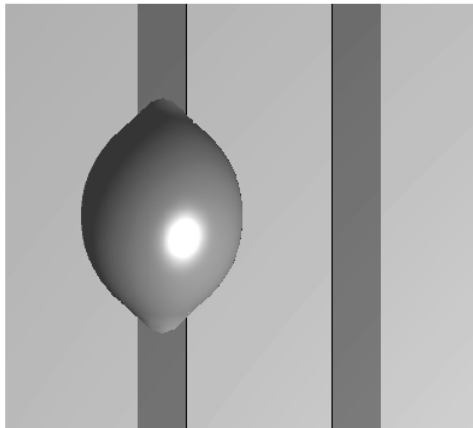
# Hybrid Eigenvector-Following

## Further optimise using hybrid eigenvector-following:

1. First, remember that we want to find saddle point of index 1
2. Uphill step in one eigendirection
3. Minimization in the tangent space
4. Remove zero eigenvectors

## Zero Eigenvectors

Arise when periodic boundary conditions are used

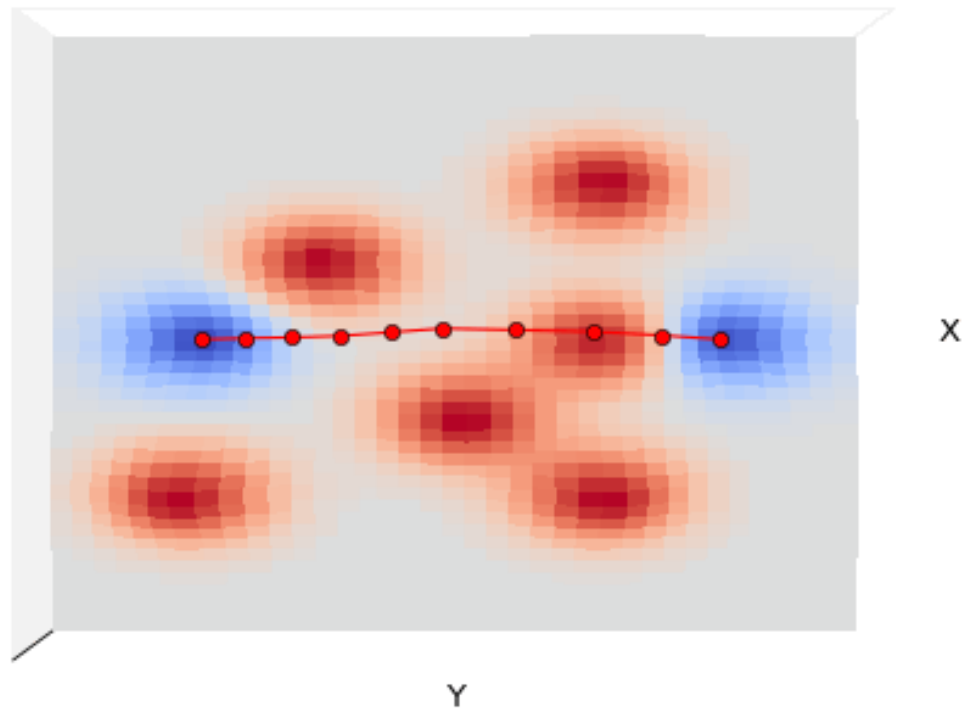


In this example, for the y-direction:

$$\hat{e} = \frac{1}{N} \left( \frac{\partial \phi_{111}}{\partial y} \dots \frac{\partial \phi_{ijk}}{\partial y} \dots \frac{\partial \phi_{NxNyNz}}{\partial y} \right)$$

# Many Ways to Rome

*With Joao Louis Carabetta*

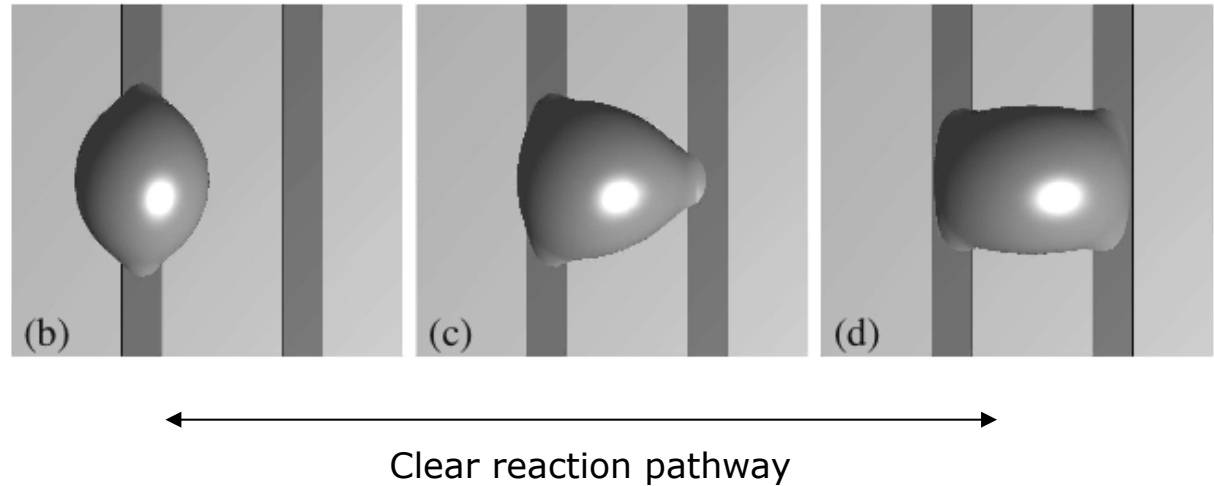
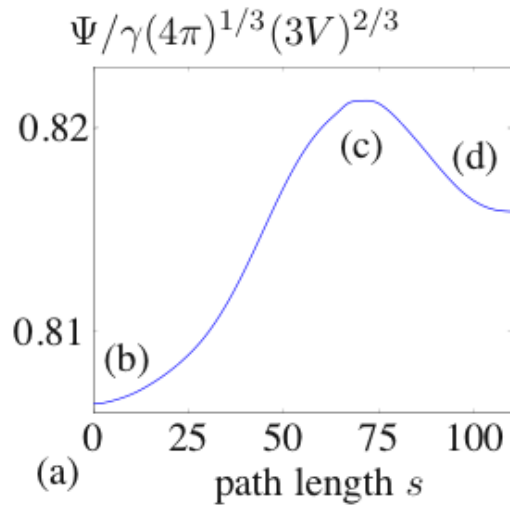


# Results

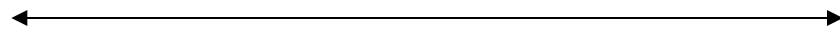
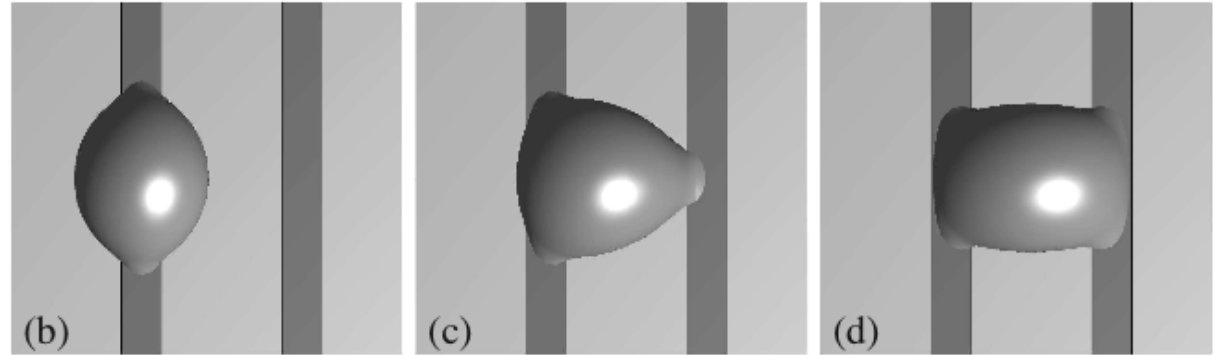
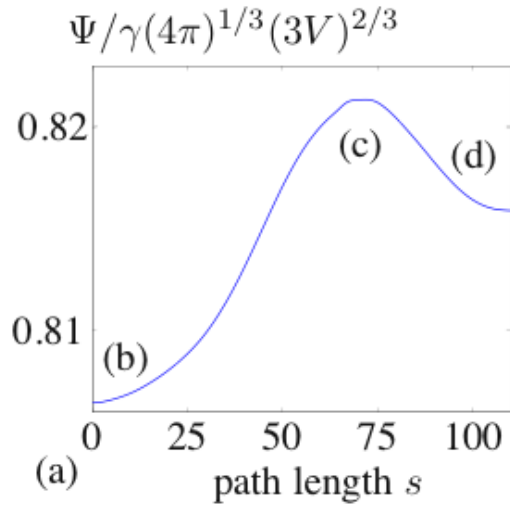
# Wetting on Chemically Striped Surfaces

**Surface Energies**  $\Psi = \gamma_{LV}A_{LV} + \gamma_{SV}A_{SV} + \gamma_{SL}A_{SL}$

**Contact Angle**  $\cos\theta = \frac{\gamma_{SV} - \gamma_{SL}}{\gamma_{LV}}$

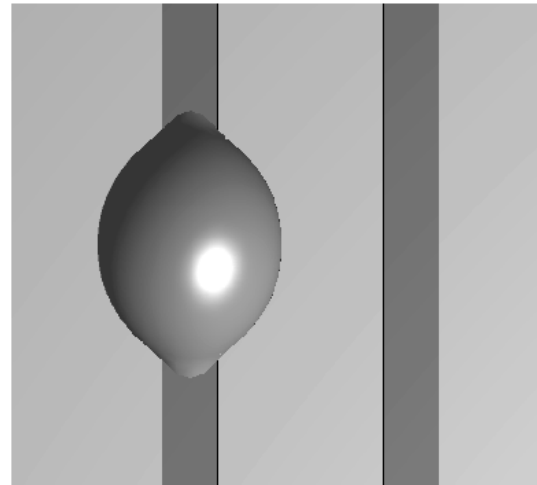


# Wetting on Chemically Striped Surfaces



Clear reaction pathway

**Minimum Energy Path**



# Vesicle Shapes I

Helfrich Free Energy

$$\Psi = \int_A \frac{\kappa}{2} \left( \frac{1}{R_1} + \frac{1}{R_2} \right)^2 dA + PV + \gamma A$$

Curvature Energy      Volume Conservation      Area Conservation

## Energy scales for bending:

e.g. for a vesicle of size 1 micron.

$$F \sim \kappa_b \left( \frac{1}{R^2} \right) A \sim \kappa_b \approx 10 - 20 k_B T$$

## Energy scales for shearing: $\sim 0$ .

This is a consequence of membrane fluidity. For example, when where there is a spectrin (cytoskeleton) network (e.g. red blood cells), this contribution is non-zero and can be important.

## Energy scales for stretching:

e.g. for a vesicle of size 1 micron to be stretched by 1%.

$$F \sim \kappa_s A \left( \frac{\Delta A}{A} \right)^2 \approx 60 \frac{k_B T}{\text{nm}^2} \times (1000 \text{nm})^2 \left( \frac{1}{100} \right)^2 \approx 6000 k_B T$$

For this reason, the membrane area is also assumed to be constant, when considering vesicle shapes that arise due to bending energy.

# Vesicle Shapes I

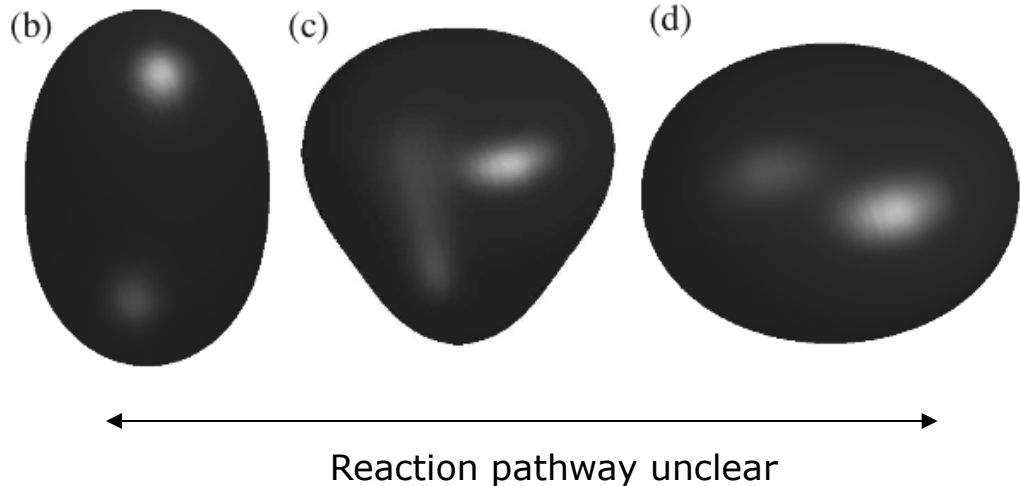
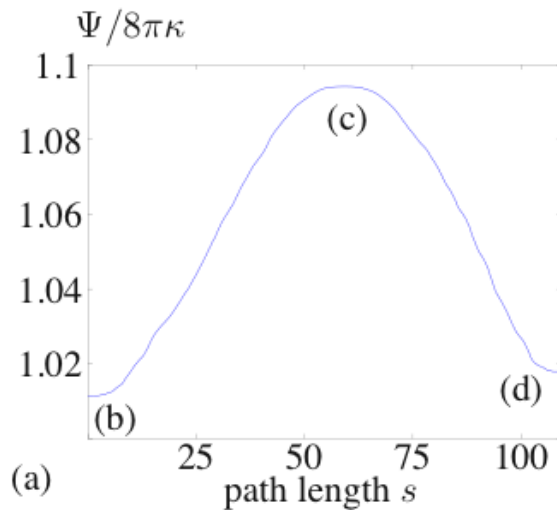
Helfrich Free Energy  $\Psi = \int_A \frac{\kappa}{2} \left( \frac{1}{R_1} + \frac{1}{R_2} \right)^2 dA + PV + \gamma A$

Curvature Energy
Volume Conservation
Area Conservation

Landau Free Energy  $\Psi = \int_V \frac{\kappa \varepsilon}{2} \left( \Delta \phi - \frac{1}{\varepsilon^2} (\phi^2 - 1) \phi \right)^2 dV + \frac{1}{2} k_V (V - V_0)^2 + \frac{1}{2} k_A (A - A_0)^2$

[Du et al. (2004)]

Reduced volume  $\nu = \frac{V_0 / (4\pi/3)}{(A_0 / 4\pi)^{3/2}} = 0.987$



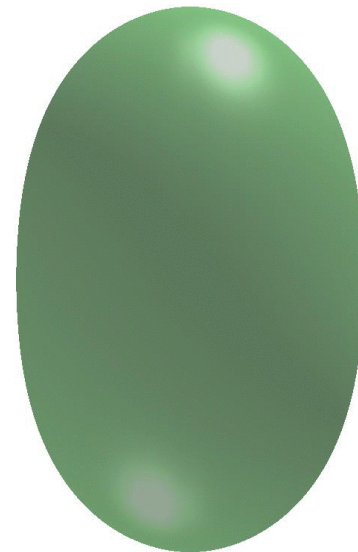
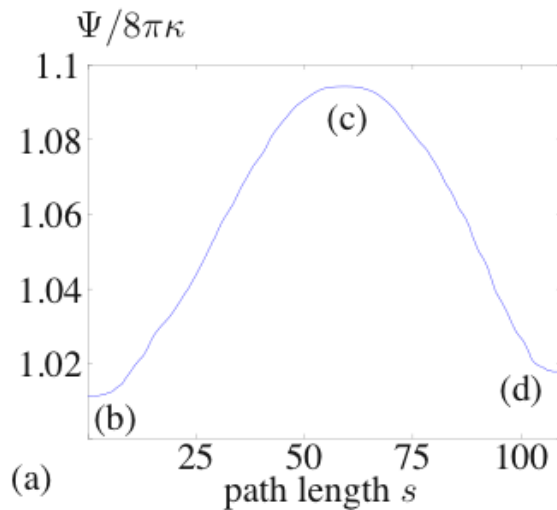
# Vesicle Shapes I

Helfrich Free Energy  $\Psi = \int_A \frac{\kappa}{2} \left( \frac{1}{R_1} + \frac{1}{R_2} \right)^2 dA + PV + \gamma A$

Curvature Energy
Volume Conservation
Area Conservation

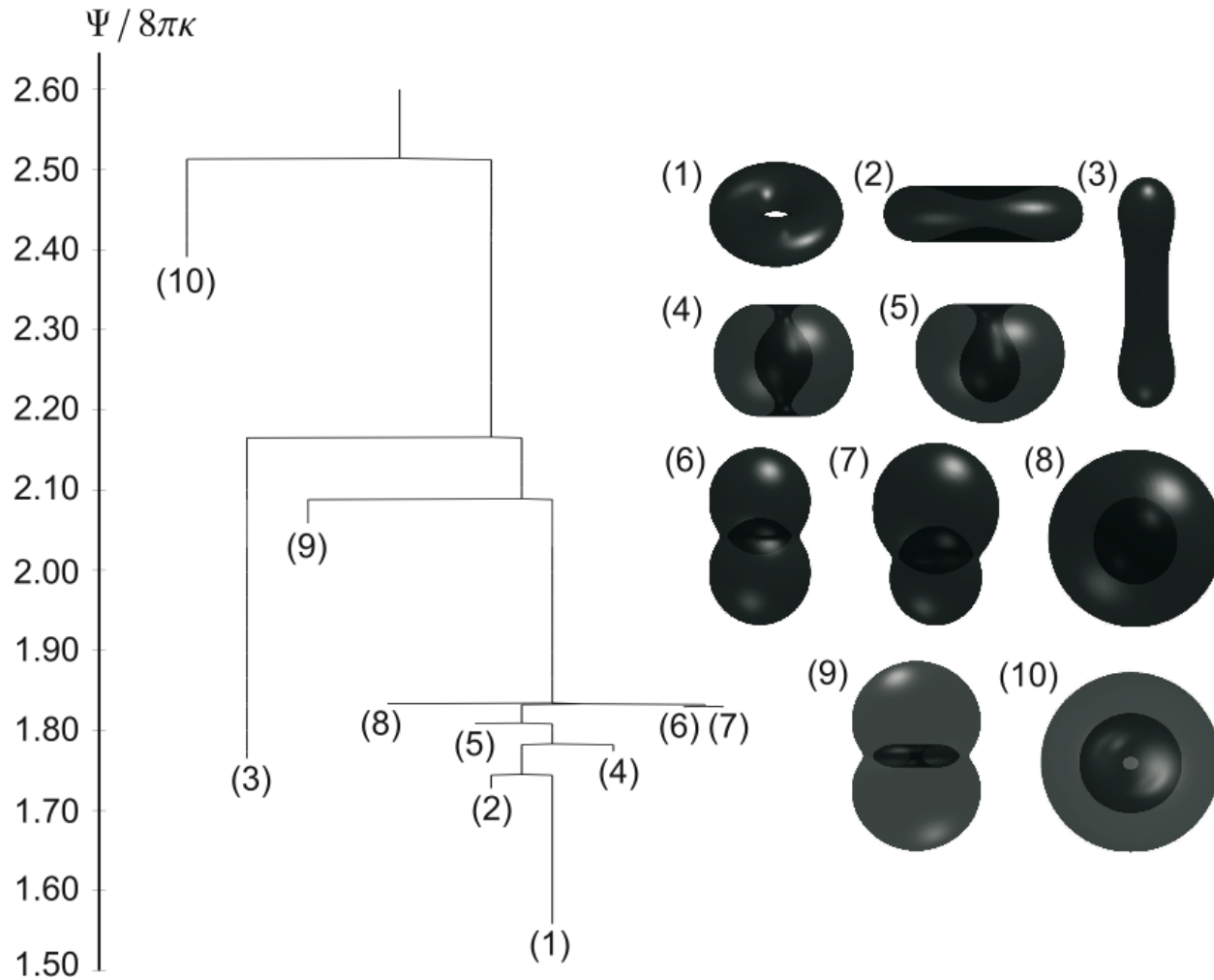
Landau Free Energy  $\Psi = \int_V \frac{\kappa \varepsilon}{2} \left( \Delta \phi - \frac{1}{\varepsilon^2} (\phi^2 - 1) \phi \right)^2 dV + \frac{1}{2} k_V (V - V_0)^2 + \frac{1}{2} k_A (A - A_0)^2$

Reduced volume  $v = \frac{V_0 / (4\pi/3)}{(A_0 / 4\pi)^{3/2}} = 0.987$



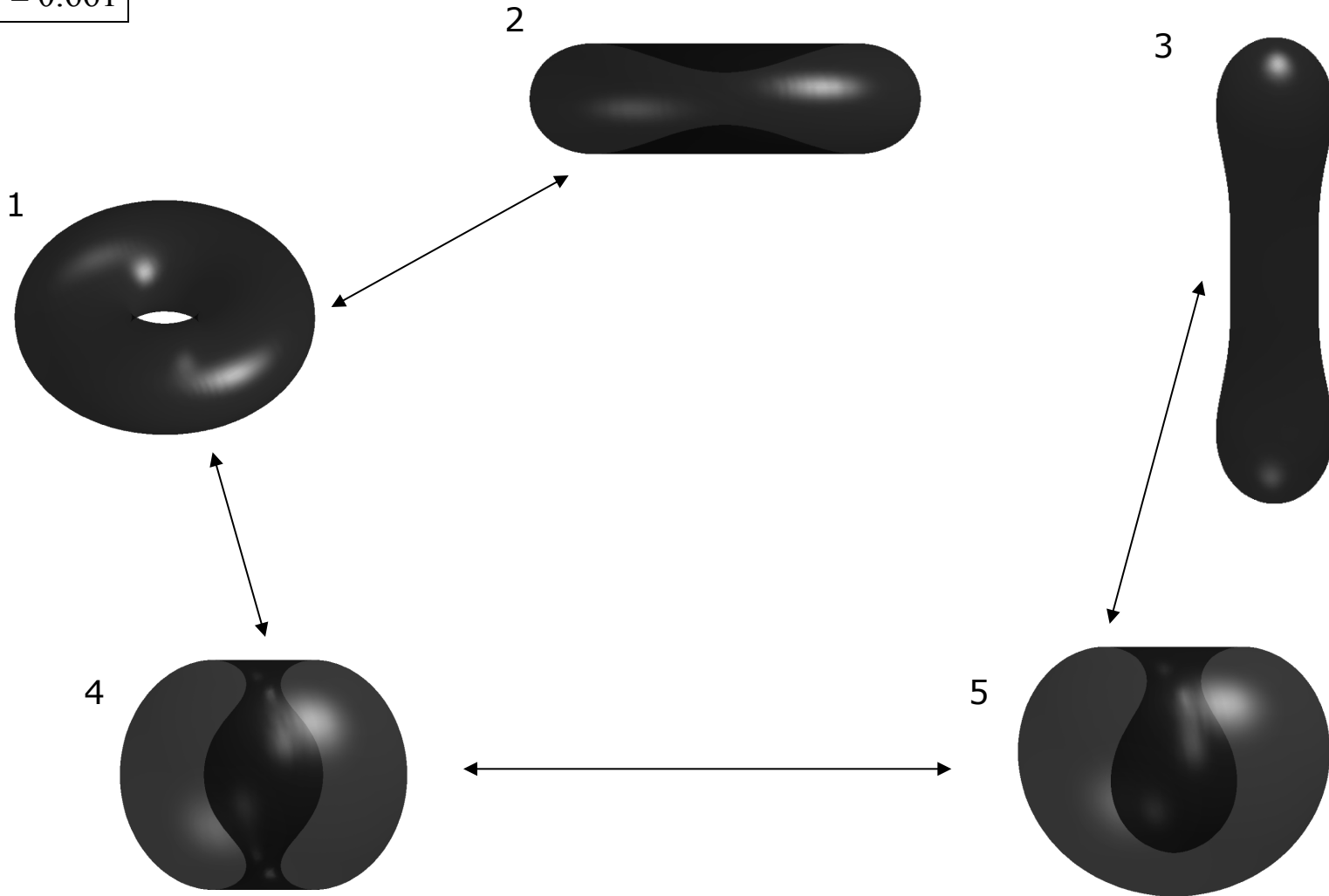
# Vesicle Shapes II

$\nu = 0.661$



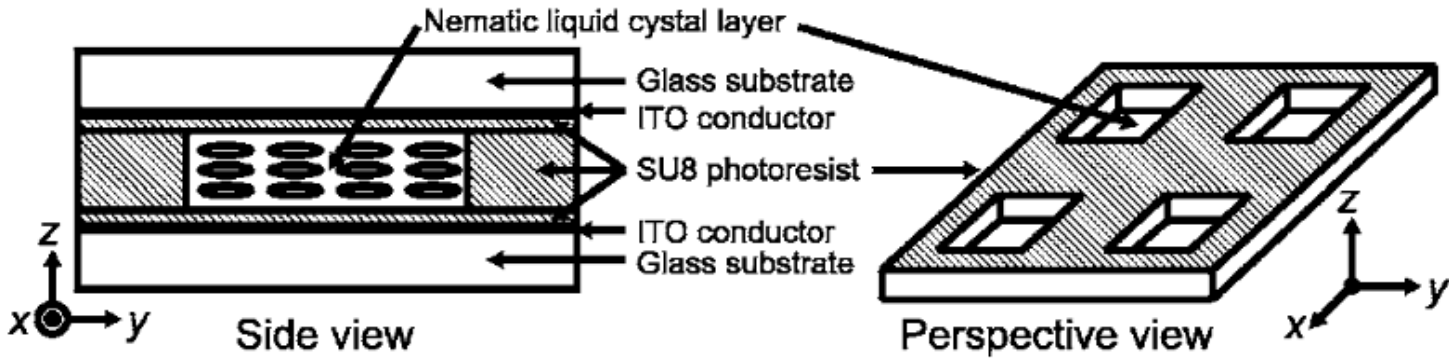
# Vesicle Shapes II

$\nu = 0.661$

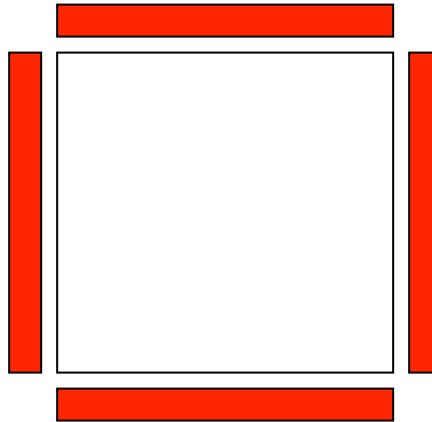


# Liquid Crystals

**Experimental Setup** – Tsakonas et al., APL 90, 111913 2007



+ surface anchoring



# Liquid Crystals

**H. Kusumaatmaja** and A. Majumdar, Soft Matter **11**, 4809 (2015)

Landau-de Gennes  
Free Energy

$$\Psi = \int_A \frac{\kappa_{el}}{2} |\nabla \mathbf{Q}|^2 dA + \int_A \left[ -\alpha \text{Tr} \mathbf{Q}^2 - \frac{b^2}{3} \text{Tr} \mathbf{Q}^3 + \frac{c^2}{4} (\text{Tr} \mathbf{Q}^2)^2 \right] dA$$

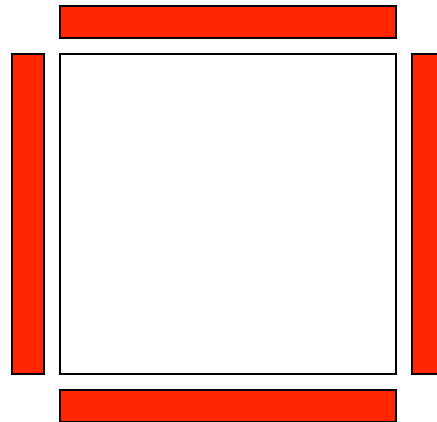
Bending Energy Bulk Energy

Where the  $\mathbf{Q}$ -tensor is defined as  $\mathbf{Q} = \begin{bmatrix} Q_{11} & Q_{12} \\ Q_{12} & -Q_{11} \end{bmatrix} = s(2\mathbf{n} \otimes \mathbf{n} - \mathbf{I})$  in two dimensions.

Dimensionless form  $\tilde{\Psi} = \int_{\tilde{A}} \tilde{\kappa}_{el} \left[ |\tilde{\nabla} \tilde{Q}_{11}|^2 + |\tilde{\nabla} \tilde{Q}_{12}|^2 \right] d\tilde{A} + \int_{\tilde{A}} (\tilde{Q}_{11}^2 + \tilde{Q}_{12}^2 - 1) d\tilde{A}$

As for a scalar order parameter  $\Psi = f(\{Q_{11}, Q_{12}\}_{ijk})$

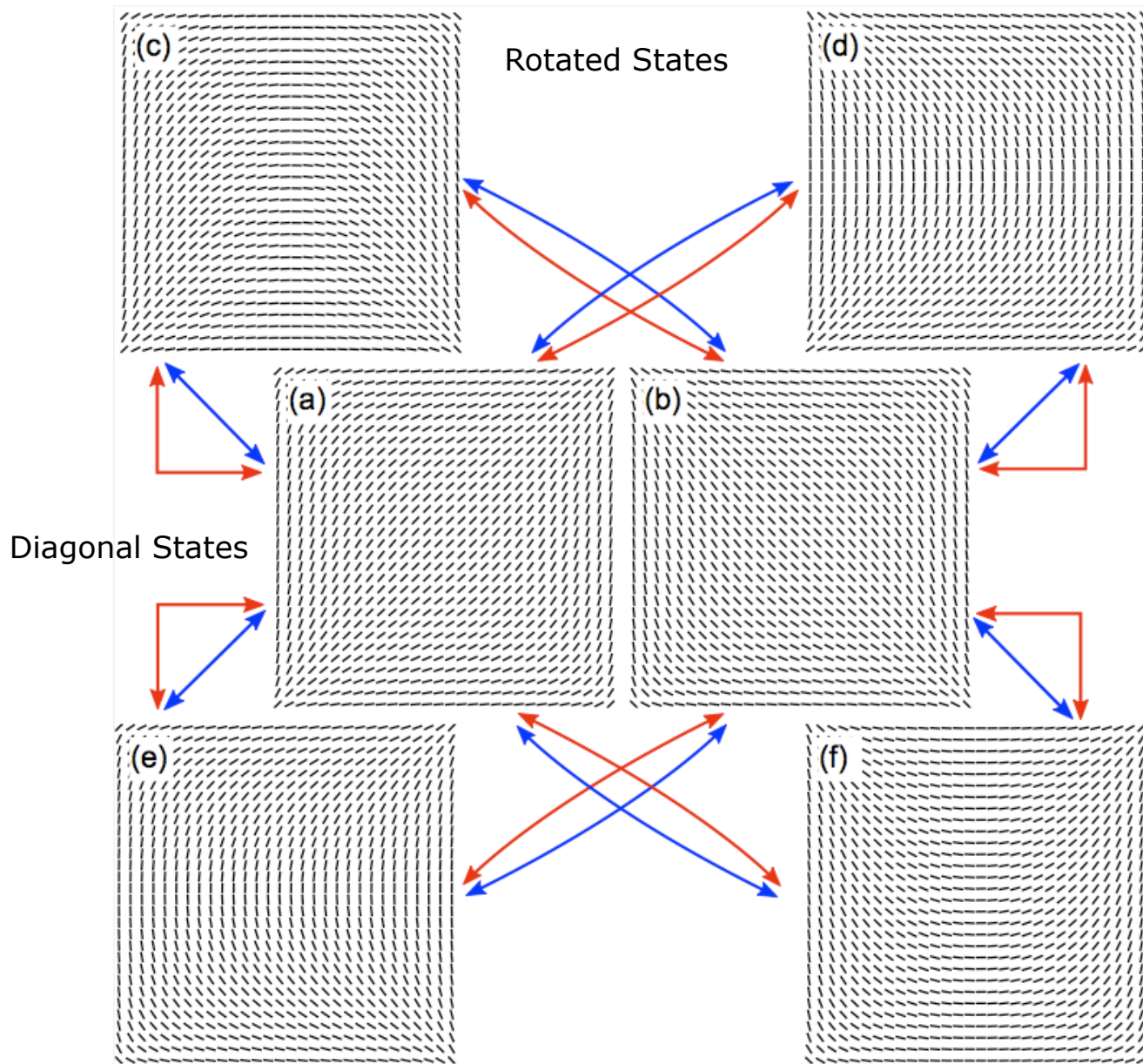
+ surface anchoring



$$\int_{\partial A} W (\tilde{Q} - \tilde{Q}_{wall})^2 dl$$

Anchoring Strength  $W$

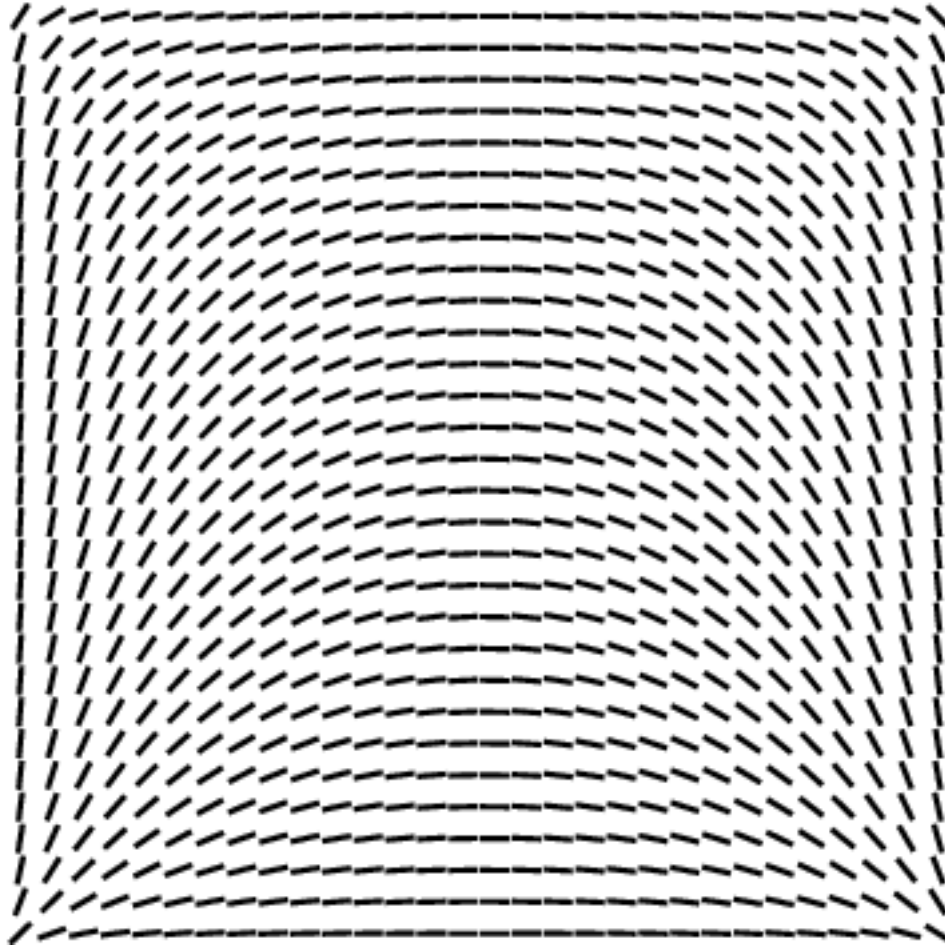
# Rotated and Diagonal States



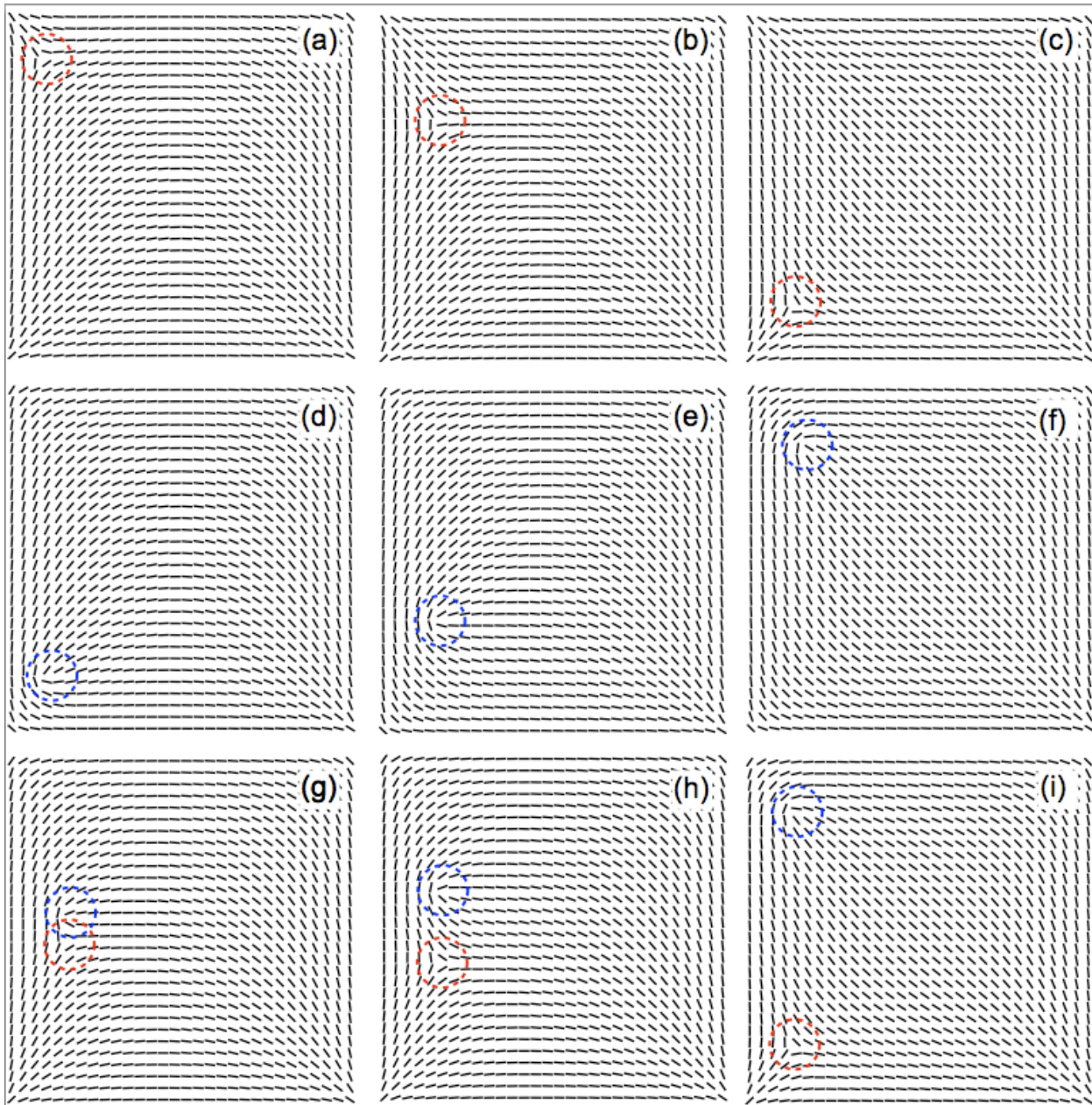
## Beyond the minimizers:

1. Pathways and barriers
2. Competing pathways?
3. Connectivity
4. Stability of Minima

# One Possible Pathway



# Strong Anchoring Pathways

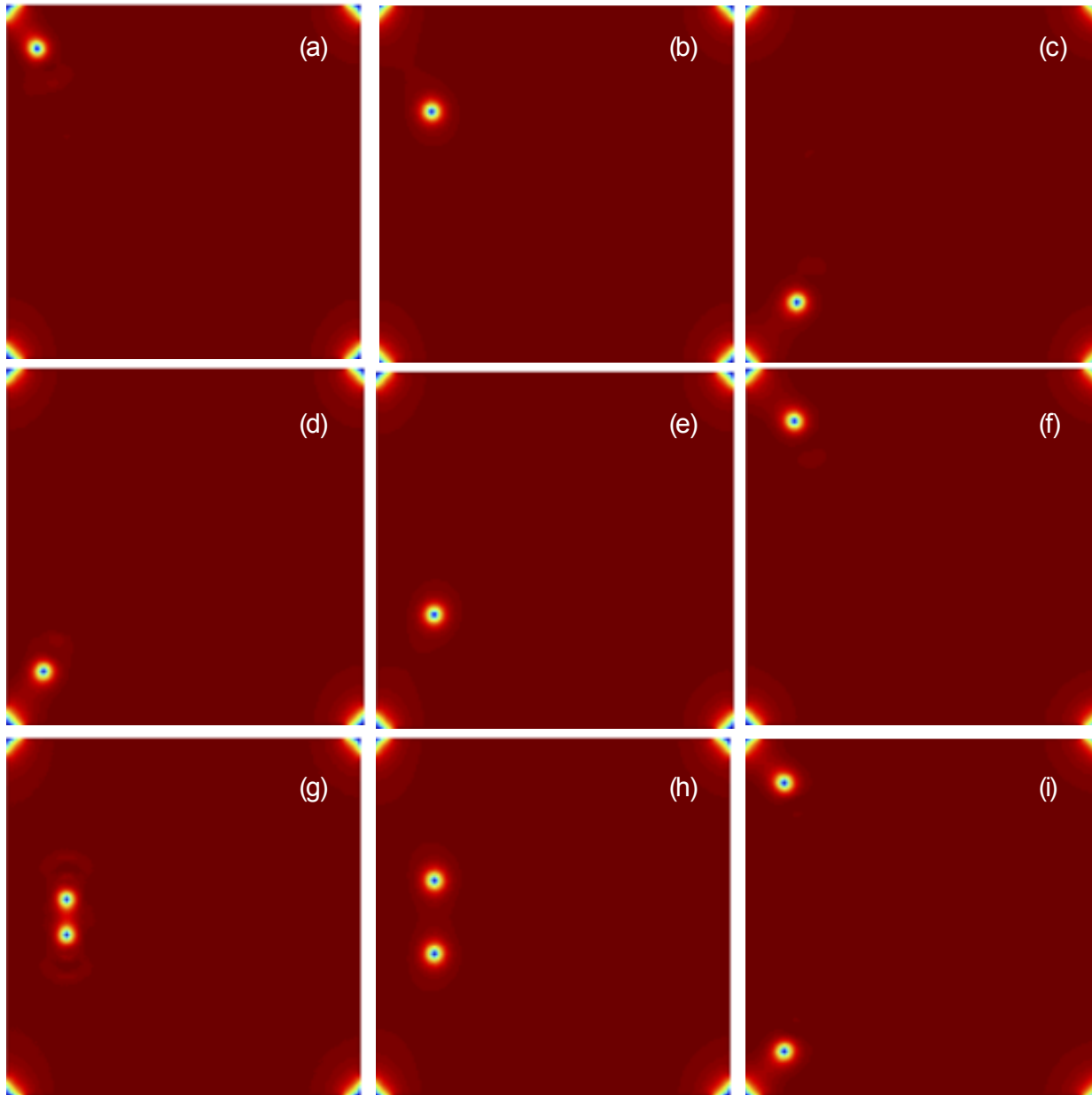


-  $1/2$  defect

+  $1/2$  defect

Pair of defects

# Order Parameter $s$



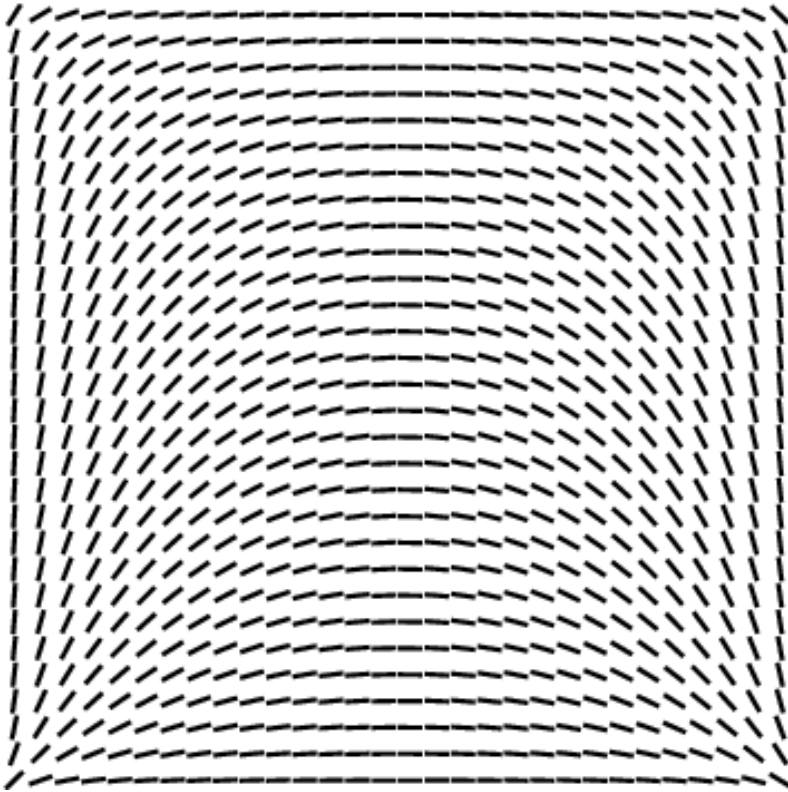
- 1/2 defect

Remember that:  
 $\mathbf{Q} = s(2\mathbf{n} \otimes \mathbf{n} - \mathbf{I})$

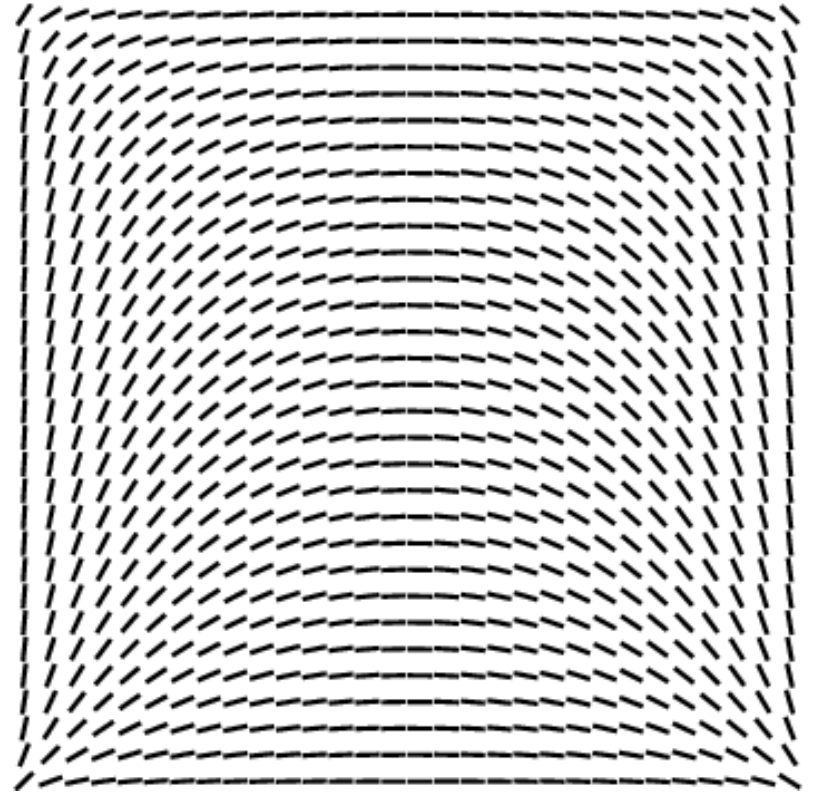
+ 1/2 defect

Pair of defects

# Strong Anchoring Pathways



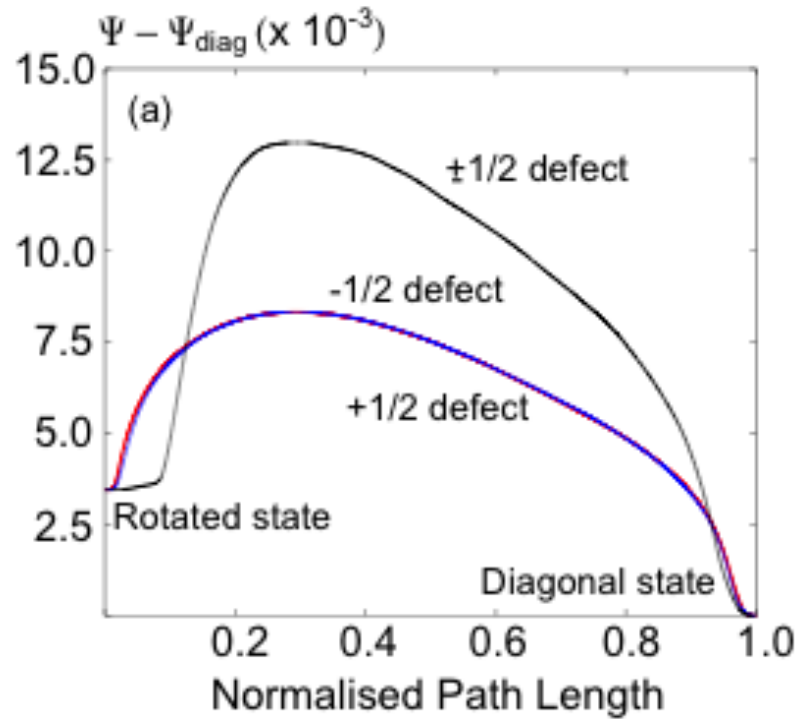
+  $\frac{1}{2}$  defect



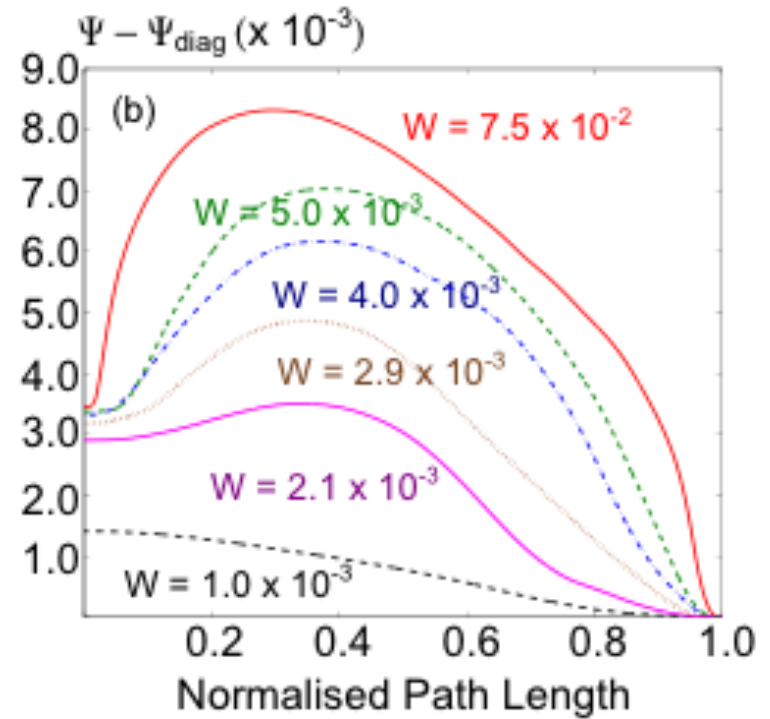
Pair of defects

# Energy Profiles

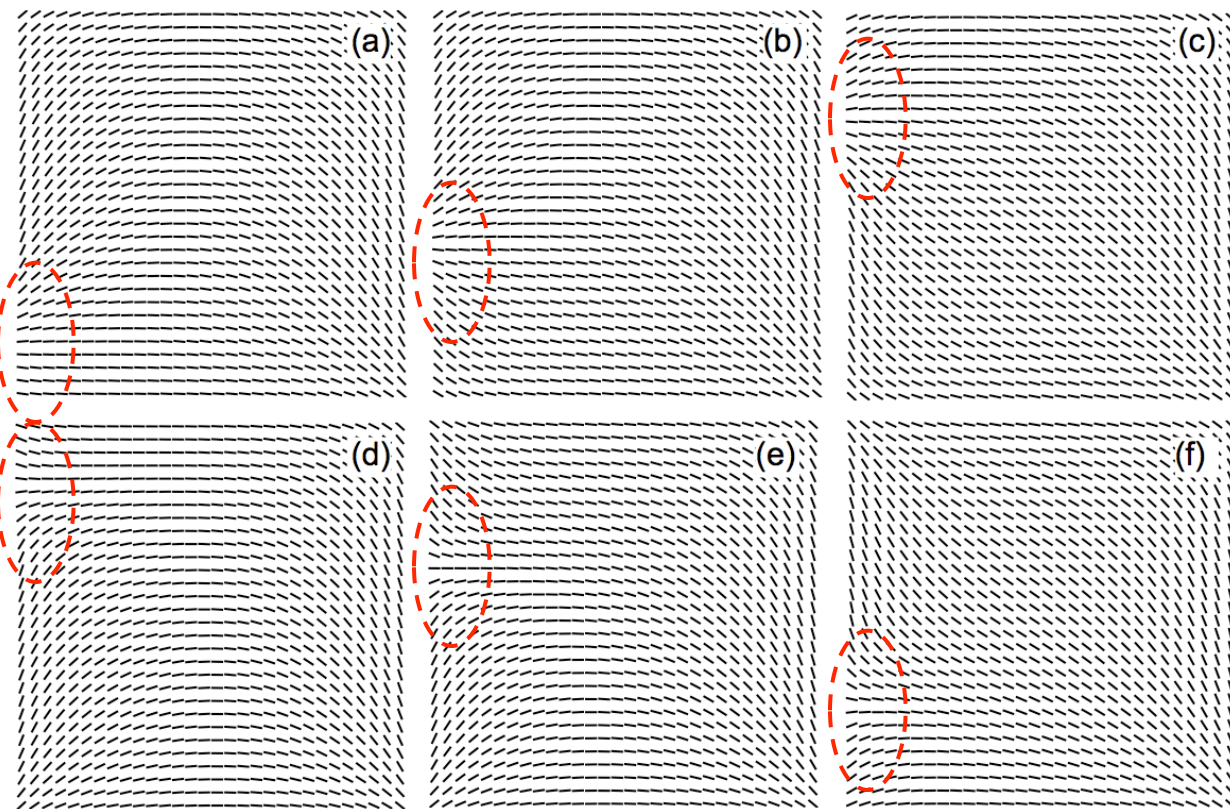
## Strong Anchoring



## Varying Anchoring Strengths

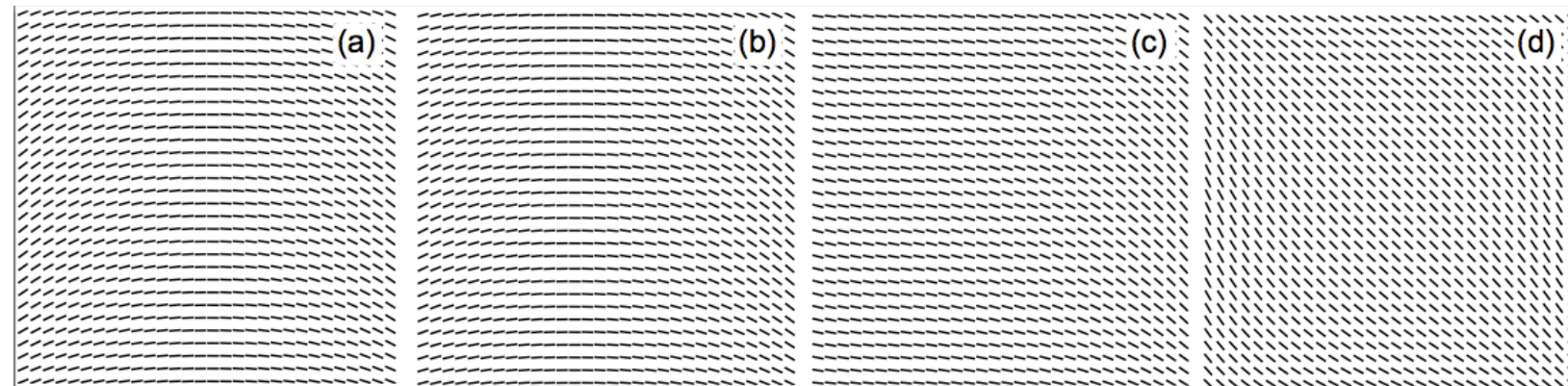


# Pathways Depend on Anchoring Strengths



**Medium Anchoring Strength**

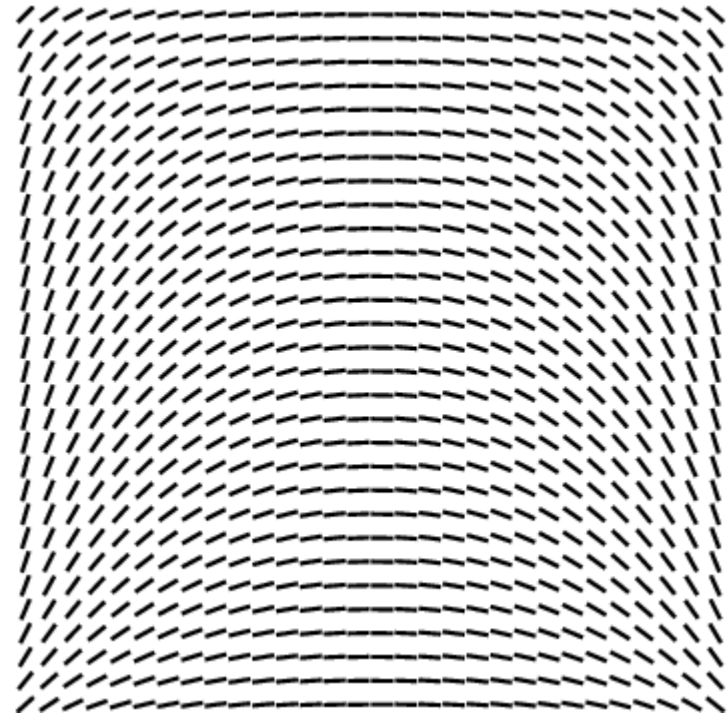
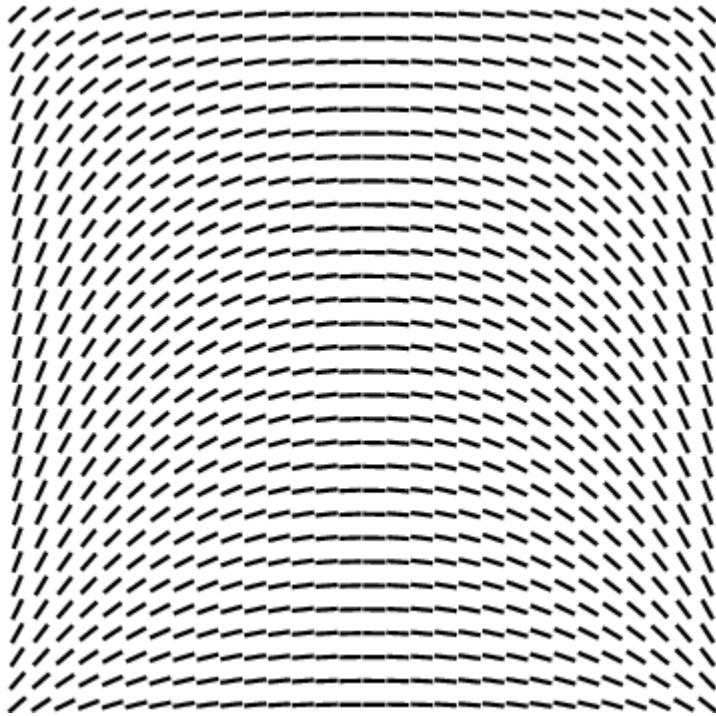
Two Degenerate Pathways



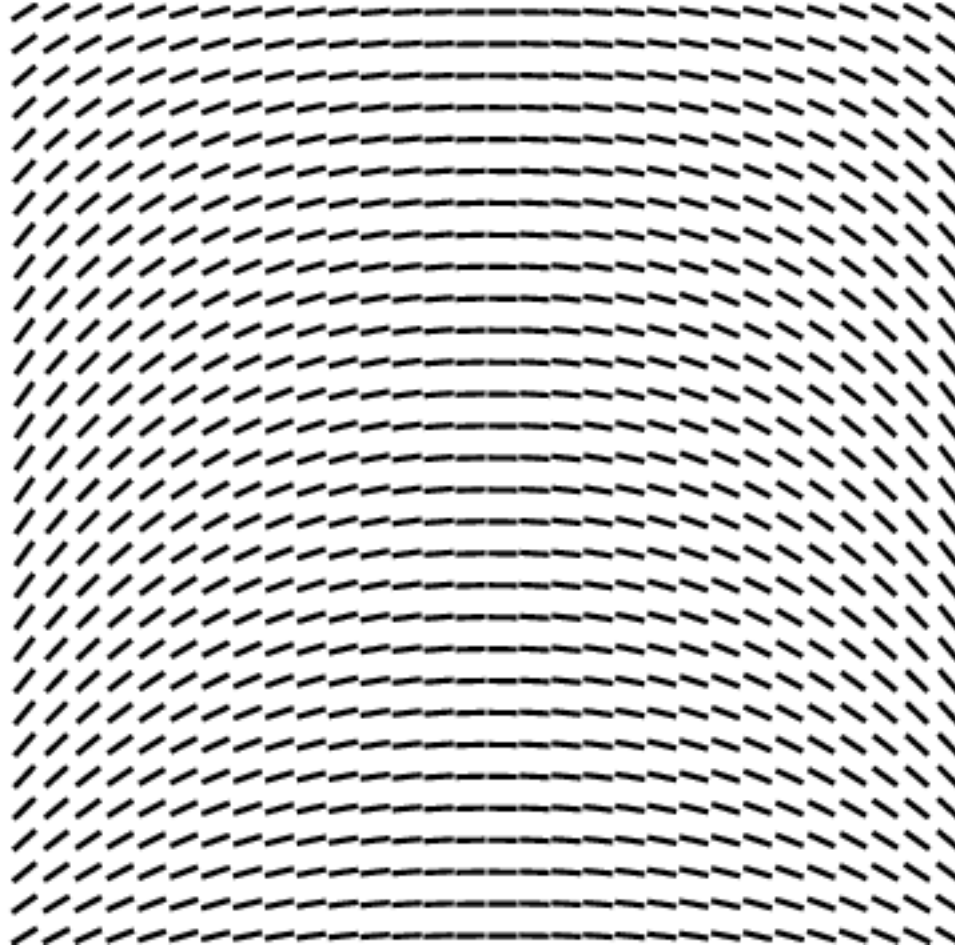
**Weak  
Anchoring  
Strength**

# Medium Anchoring

## *Two Degenerate Pathways*

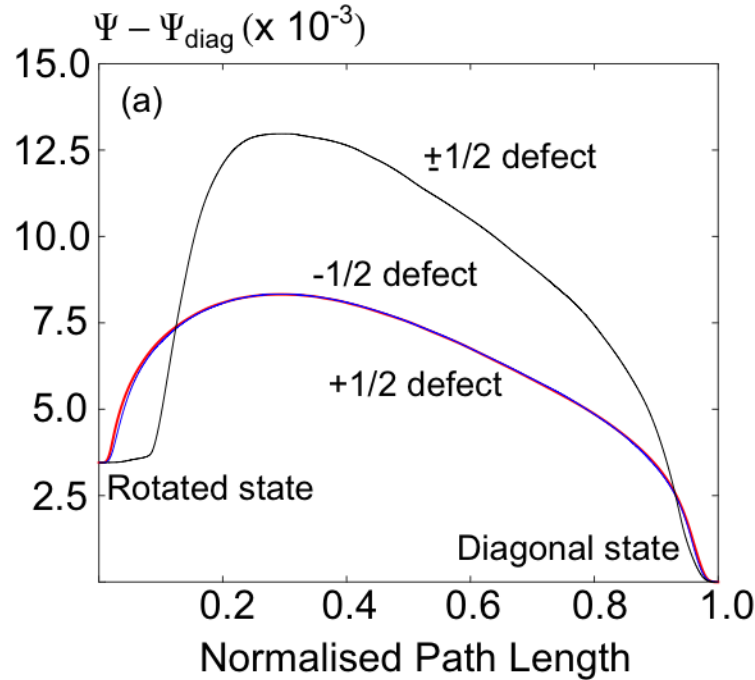


# Weak Anchoring

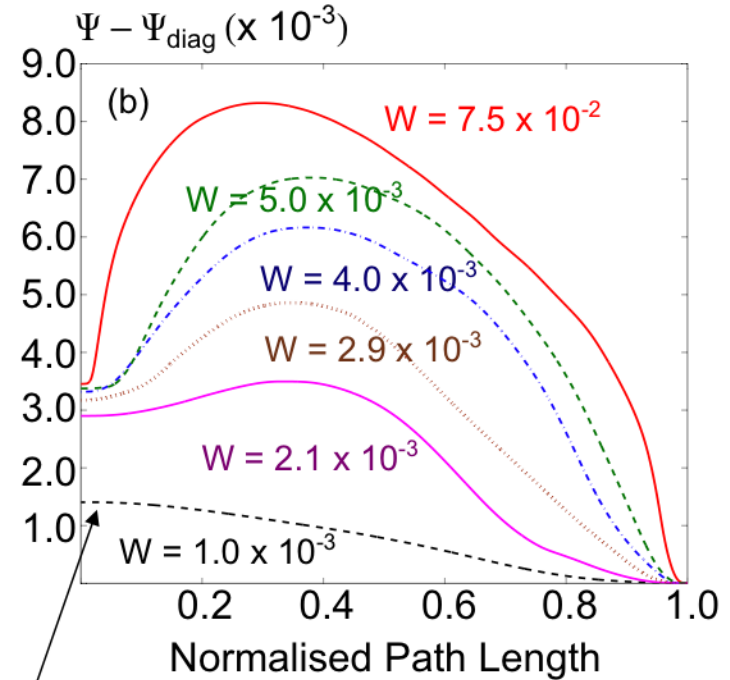


# Energy Profiles

## Strong Anchoring



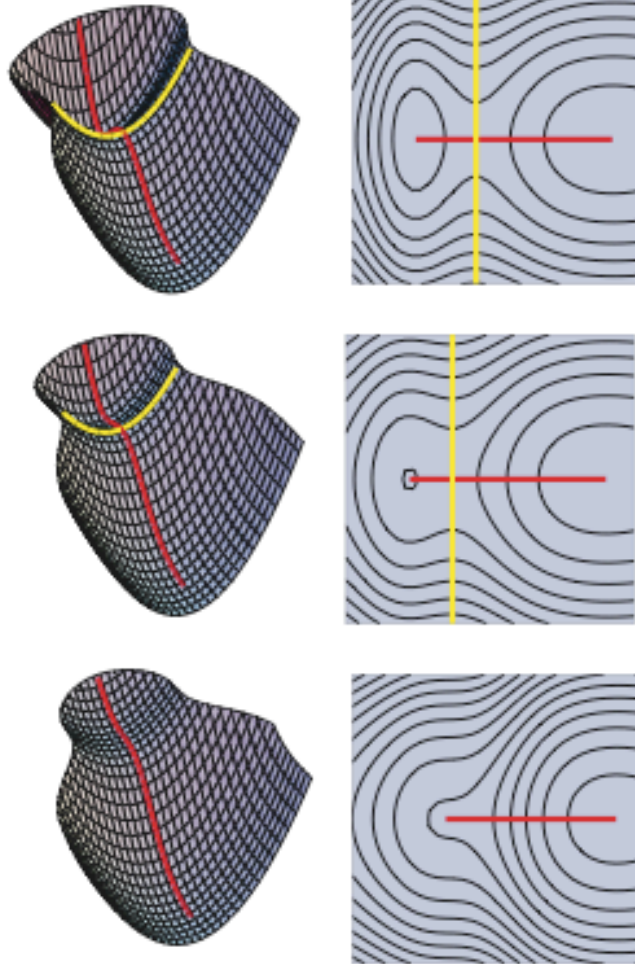
## Varying Anchoring Strengths



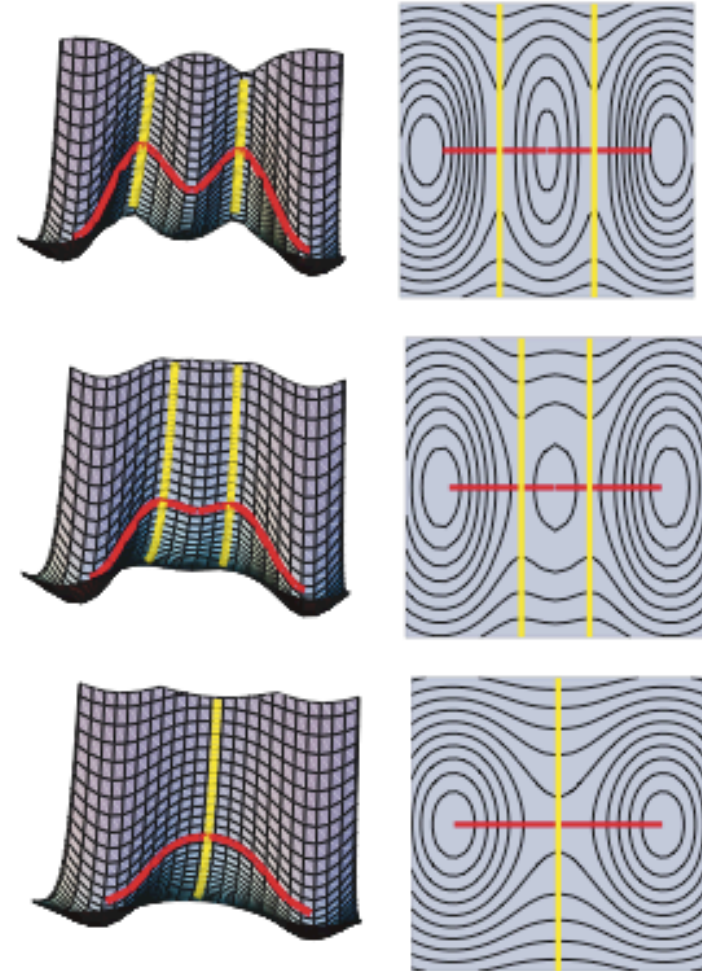
Disappearance of minima  
(Catastrophe event)

# Catastrophe Events

**Fold Catastrophe**



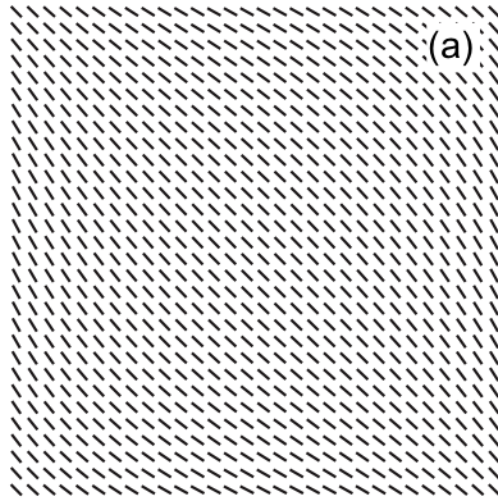
**Cusp Catastrophe**



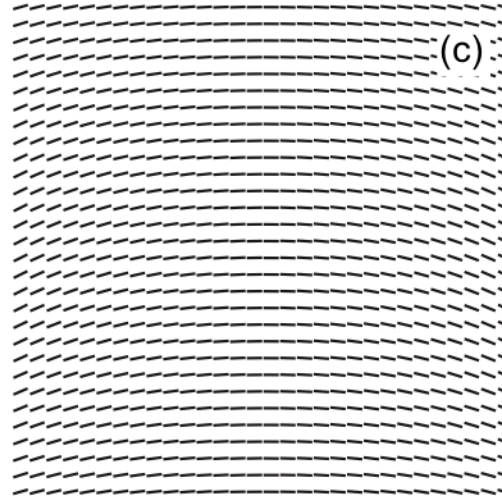
# Only Diagonal States Are Stable

Reminiscent of Rotated State

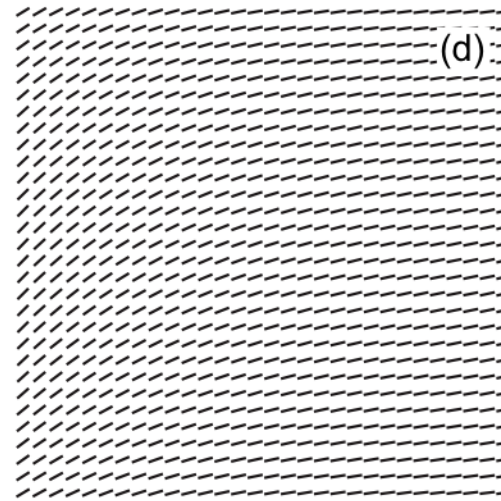
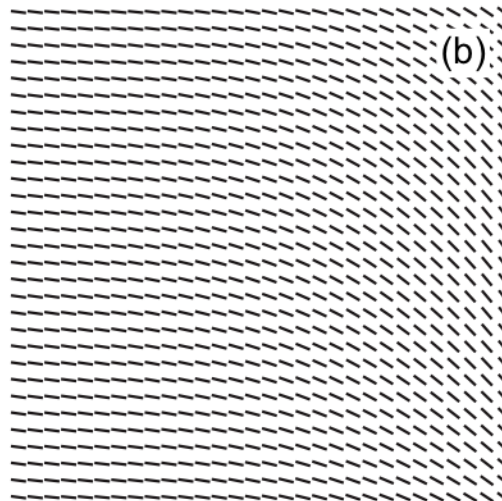
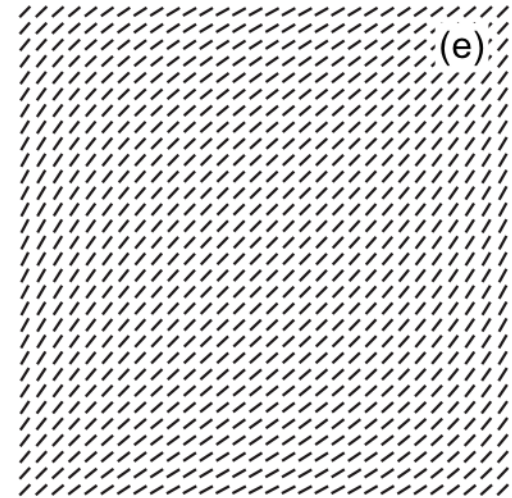
Minimum



Transition State

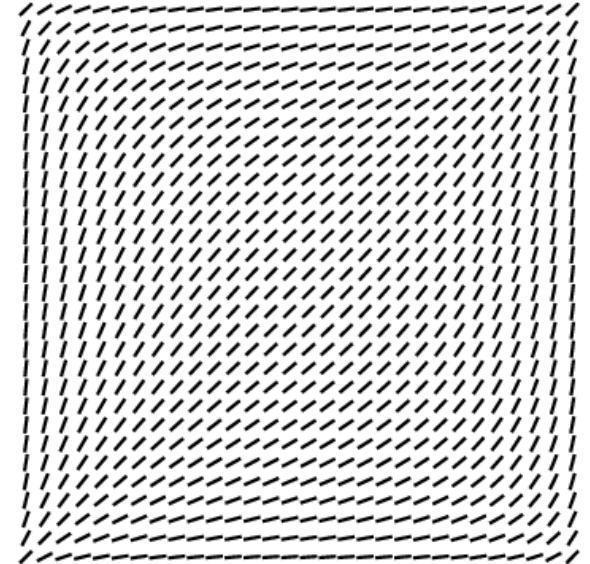
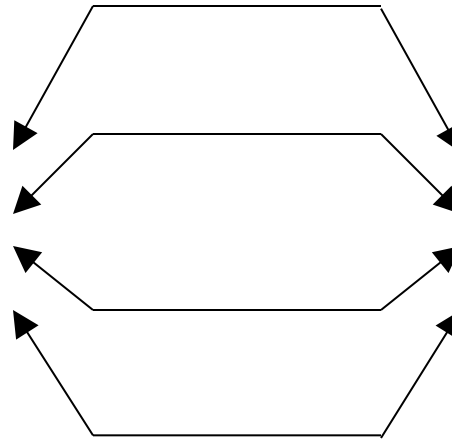
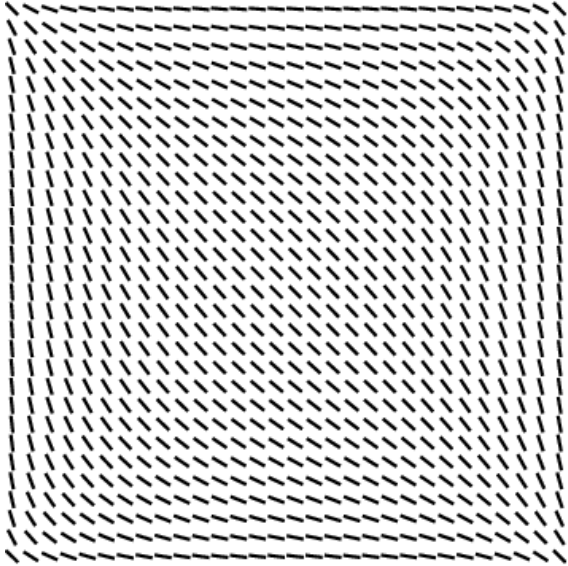


Minimum



# Only Diagonal States Are Stable

**Four equivalent pathways**



# Take Home Messages

1. Tools to **survey the free energy landscapes** of Landau models in details, including how they **vary with system parameters, boundary conditions, and external perturbations**.
  - Most, if not all, relevant minima
  - Most, if not all, transition states and competing minimum energy paths
  - Connectivity of the landscapes
  - “Catastrophe” events/phase transitions
2. **Cannot do dynamics**. However, the descriptions and discretizations of the functionals are fully **compatible with other methods**.
3. Can take advantage of schemes already developed for handling **complex boundaries** (corrugations, curved boundaries, etc) from other methods.
4. **Scalar, vector, and tensor** fields can be treated on **equal footing**.

Landau free energies are very popular in physics, chemistry, and materials science.  
We are interested in testing the limits of the methods.
5. Relatively **inexpensive** (especially compared to atomistic calculations).
  - Typical minimization runs take  $O(\text{minutes})$ .
  - Typical pathway runs take  $O(1 \text{ hour})$ .

} Run in parallel

Thank you for listening!