

Ab initio PIC simulations and data-driven techniques for multi-fidelity or reduced plasma models

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** UCLA



Ab initio kinetic plasma simulations and the “ground truth”

Some questions

Speeding standard kinetic simulations + novel simulators

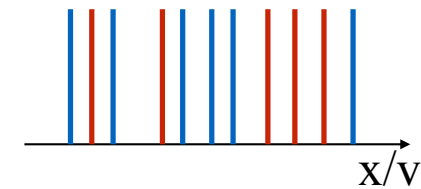
Learning advection and diffusion coefficients from PIC simulation data

Pushing PIC simulations towards capturing *ab initio* collisions: from many particles per cell to *many cells per particle*

A brief digression on plasma kinetic theory

Point like charged particles

$$N_s(\mathbf{x}, \mathbf{v}, t) = \sum_{i=1}^{N_0} \delta[\mathbf{x} - \mathbf{X}_i(t)] \delta[\mathbf{v} - \mathbf{V}_i(t)]$$



Equations of motion $\dot{\mathbf{X}}_i(t) = \mathbf{V}_i(t)$

$$m_s \dot{\mathbf{V}}_i(t) = q_s \mathbf{E}^m[\mathbf{X}_i(t), t] + \frac{q_s}{c} \mathbf{V}_i(t) \times \mathbf{B}^m[\mathbf{X}_i(t), t]$$

$$\nabla \cdot \mathbf{E}^m(\mathbf{x}, t) = 4\pi \rho^m(\mathbf{x}, t)$$

$$\nabla \cdot \mathbf{B}^m(\mathbf{x}, t) = 0$$

$$\nabla \times \mathbf{E}^m(\mathbf{x}, t) = -\frac{1}{c} \frac{\partial \mathbf{B}^m(\mathbf{x}, t)}{\partial t}$$

$$\nabla \times \mathbf{B}^m(\mathbf{x}, t) = \frac{4\pi}{c} \mathbf{J}^m(\mathbf{x}, t) + \frac{1}{c} \frac{\partial \mathbf{E}^m(\mathbf{x}, t)}{\partial t}$$

$$\rho^m(\mathbf{x}, t) = \sum_{e,i} q_s \int d\mathbf{v} N_s(\mathbf{x}, \mathbf{v}, t)$$

$$\mathbf{J}^m(\mathbf{x}, t) = \sum_{e,i} q_s \int d\mathbf{v} \mathbf{v} N_s(\mathbf{x}, \mathbf{v}, t)$$

Sources for microscopic fields

$$\frac{\partial N_s(\mathbf{x}, \mathbf{v}, t)}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} N_s + \frac{q_s}{m_s} \left(\mathbf{E}^m + \frac{\mathbf{v}}{c} \times \mathbf{B}^m \right) \cdot \nabla_{\mathbf{v}} N_s = 0$$

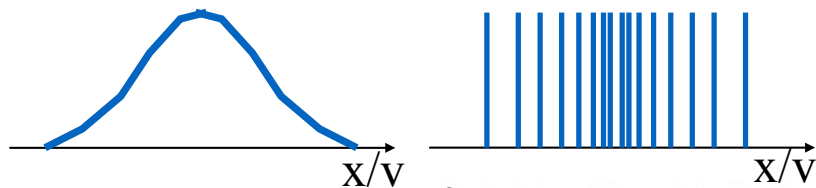
Transport equation for N_s is the Klimontovich equation

$$f_s(\mathbf{x}, \mathbf{v}, t) \equiv \langle N_s(\mathbf{x}, \mathbf{v}, t) \rangle$$

$$N_s(\mathbf{x}, \mathbf{v}, t) = f_s(\mathbf{x}, \mathbf{v}, t) + \delta N_s(\mathbf{x}, \mathbf{v}, t)$$

$$\mathbf{E}^m(\mathbf{x}, \mathbf{v}, t) = \mathbf{E}(\mathbf{x}, \mathbf{v}, t) + \delta \mathbf{E}(\mathbf{x}, \mathbf{v}, t)$$

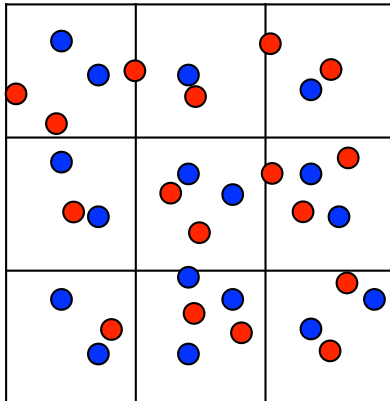
$$\mathbf{B}^m(\mathbf{x}, \mathbf{v}, t) = \mathbf{B}(\mathbf{x}, \mathbf{v}, t) + \delta \mathbf{B}(\mathbf{x}, \mathbf{v}, t)$$



$$\begin{aligned} \frac{\partial f_s(\mathbf{x}, \mathbf{v}, t)}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} f_s + \frac{q_s}{m_s} (\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}) \cdot \nabla_{\mathbf{v}} f_s \\ = - \frac{q_s}{m_s} \langle (\delta \mathbf{E} + \frac{\mathbf{v}}{c} \times \delta \mathbf{B}) \cdot \nabla_{\mathbf{v}} \delta N_s \rangle \end{aligned}$$

Particle-in-cell/particle-mesh simulations

Solving Maxwell's equations on a grid with self-consistent charges and currents due to charged particle dynamics



State-of-the-art

$\sim 10^{12}$ particles

$\sim (12000)^3$ cells

RAM \sim 1 Gbyte - 100s TByte

Run time: hours to months

Data/run \sim few MB - 100s TByte

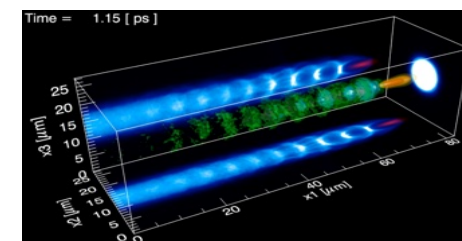
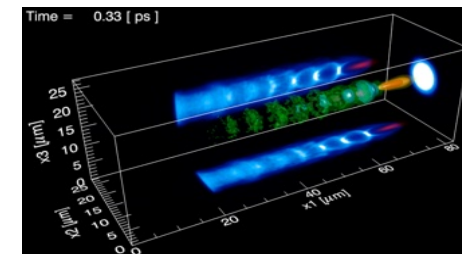
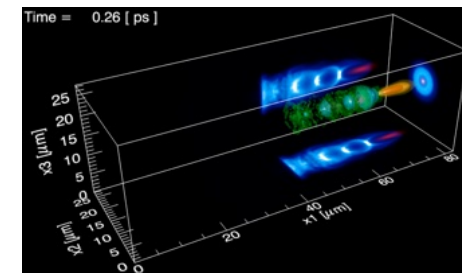
10's - 100's M cpu core hours/run

Weibel/two stream instability in astrophysics, relativistic shocks, fast igniton/inertial fusion energy, low temperature plasmas, laser-plasma interaction

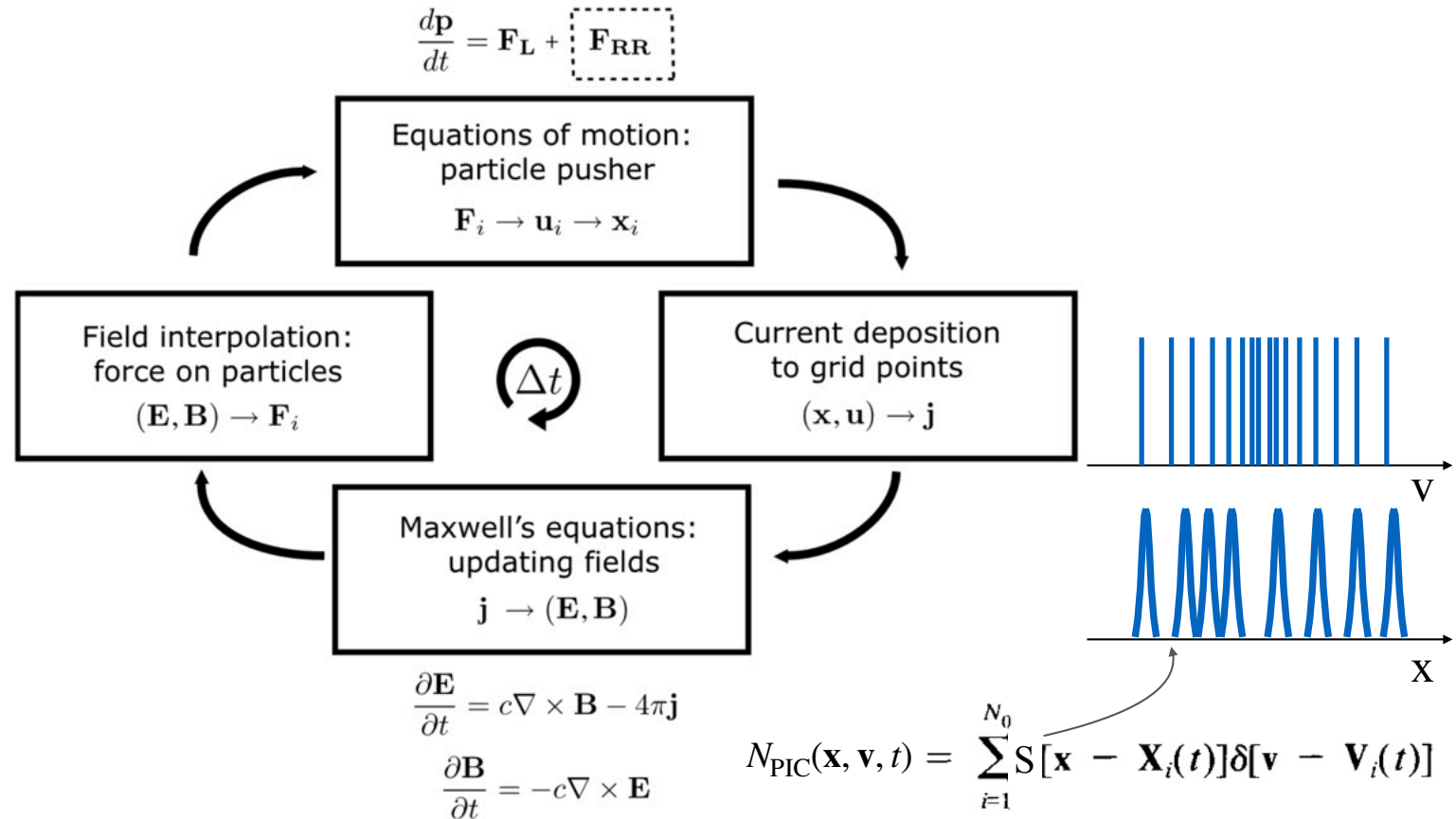
Particle-in-cell (PIC) - (Dawson, Buneman, 1960's)

Maxwell's equation solved on simulation grid

Particles pushed with Lorentz force



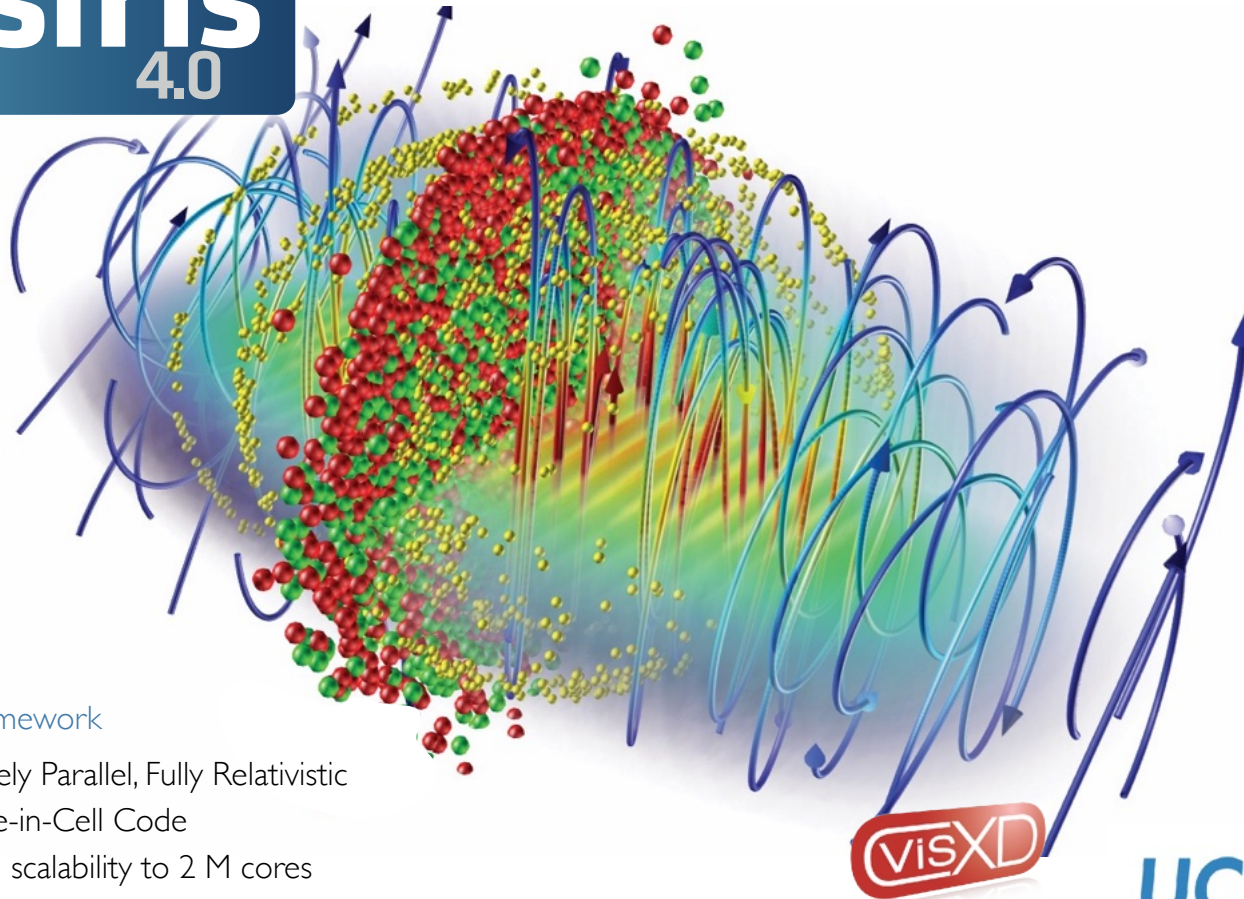
The most fundamental (classical) model in plasma physics



Particle-in-cell simulations underpin many scientific advances in plasma physics



Osiris 4.0



OSIRIS framework

- Massively Parallel, Fully Relativistic Particle-in-Cell Code
- Parallel scalability to 2 M cores
- Explicit SSE / AVX / QPX / Xeon Phi / CUDA support
- Extended physics/simulation models

Open-source version available

Open-access model

- 40+ research groups worldwide are using OSIRIS
- 300+ publications in leading scientific journals
- Large developer and user community
- Detailed documentation and sample inputs files available

Using OSIRIS 4.0

- The code can be used freely by research institutions after signing an MoU
- Find out more at:
<http://epp.tecnico.ulisboa.pt/osiris>

UCLA

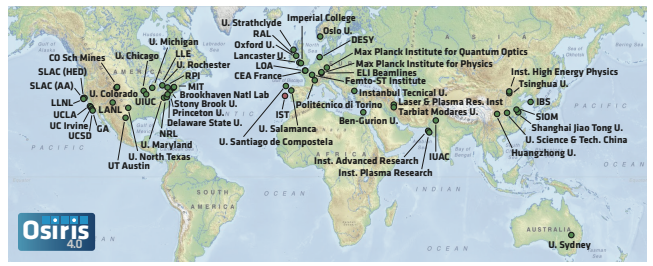
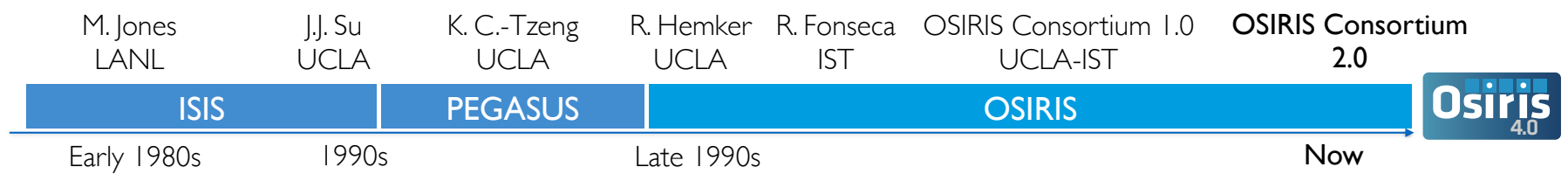


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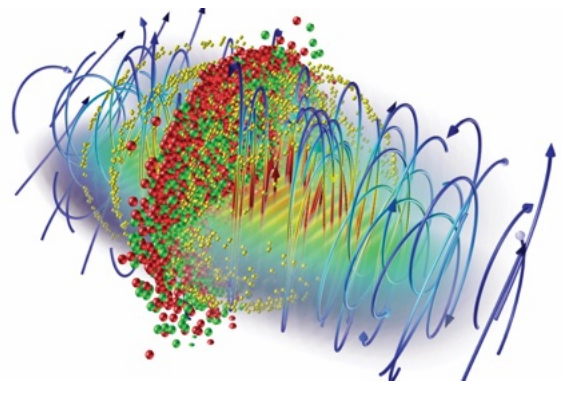


Ricardo Fonseca: ricardo.fonseca@tecnico.ulisboa.pt
Luis O. Silva | IPAM, UCLA | April 16, 2026 | 8

A brief history of OSIRIS



+50 groups worldwide
Fully developed by universities in the US and in the EU



Running on all top 10 machines + **open source**

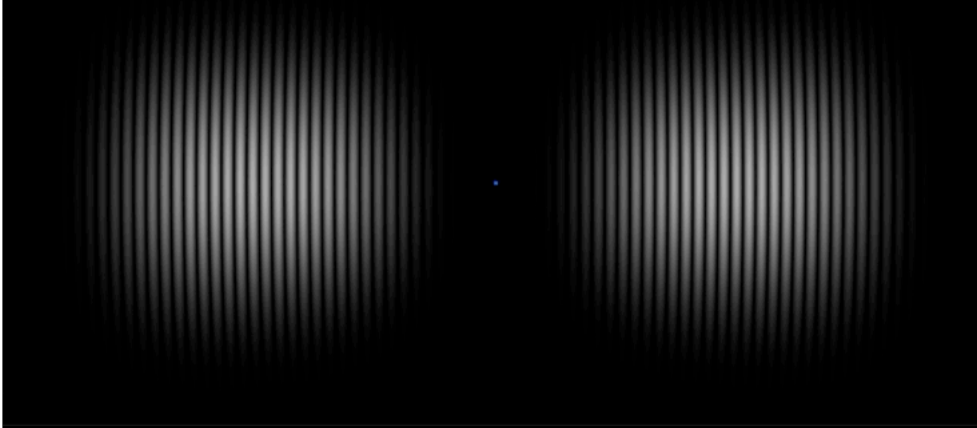
Massively parallel on modern architectures, including GPUs, dynamic load balancing, large suite of field solvers + particle pushers + boundary conditions + AI

Extended physics models including QED, general relativity, spherical grids, quasi-3D, etc.

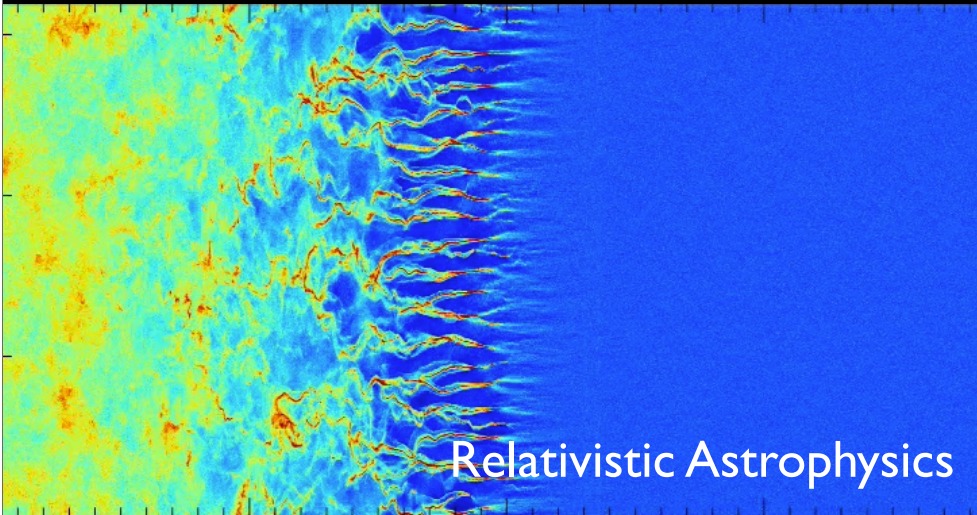
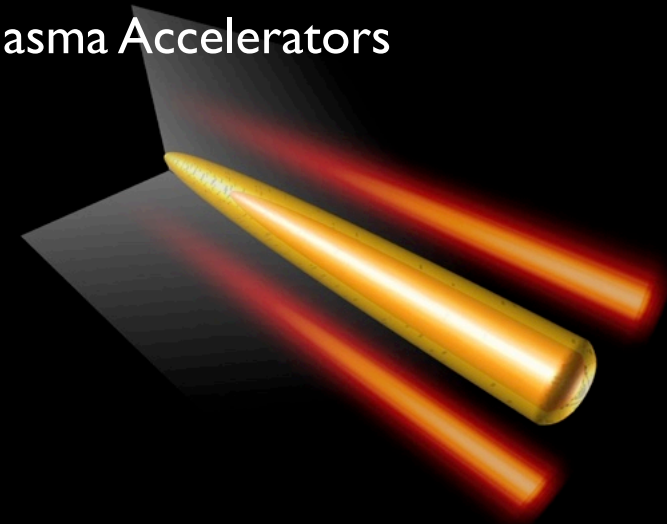
Physics results covering plasma accelerators, laser fusion, relativistic astrophysics, SF QED, beam physics, fundamental plasma physics, space physics, laboratory astrophysics, etc.



Strong field QED

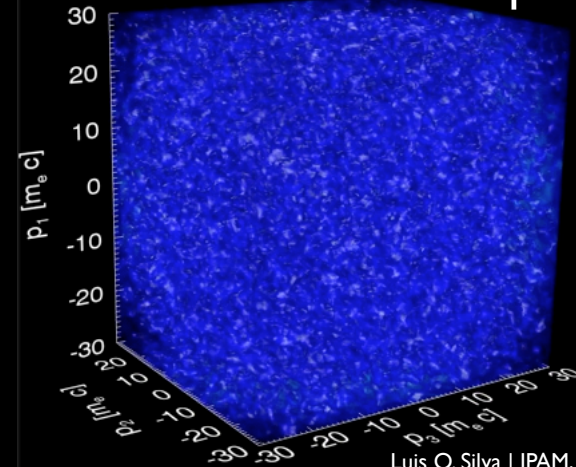


Plasma Accelerators



Relativistic Astrophysics

Fundamental plasma physics



Why PIC simulations?

Very few approximations - unique source of high fidelity data

At the basis of the hierarchy of plasma physics descriptions - Klimontovich equation - can provide high resolution information to inform other scales

Well tested and benchmarked - against experiments, theory, other models

Can take advantage of large scale computational resources - across many different architectures/scales

The most fundamental (classical) physics model for plasmas- very well understood statistical mechanics/statistical properties - used across many sub-fields

Some questions

Can ML help us speed up standard plasma simulators? Early attempts: ML replaces Monte Carlo modules in PIC - Badiali et al., JPP 2022; Amaro et al., JCP 2026

Can we build faster ML based simulators? Rethinking architecture of simulators to match ML uniqueness: 1D collisional plasma model - Carvalho et al., MLST 2024

What can we learn from data-driven approaches + ML? Learning physics (following pioneering work in plasma physics by Alves et al.) e.g. collision operators: Carvalho et al., to appear PPCF (2026); *idem*, submitted JPP (2026)

Can standard plasma simulators provide “high quality data” for data-driven discovery? Capturing self-consistent collisions in PIC codes: D. Carvalho et al., in preparation

Can we understand qualitative modifications of plasma behavior from “Learning what we already know”? e.g. Waterbag vs Maxwellian; nonlinear waves vs unstable (and then turbulent) scenarios; nonrelativistic to relativistic, from conservation of energy/momentum to Casimir invariants

Ab initio kinetic plasma simulations and the “ground truth”

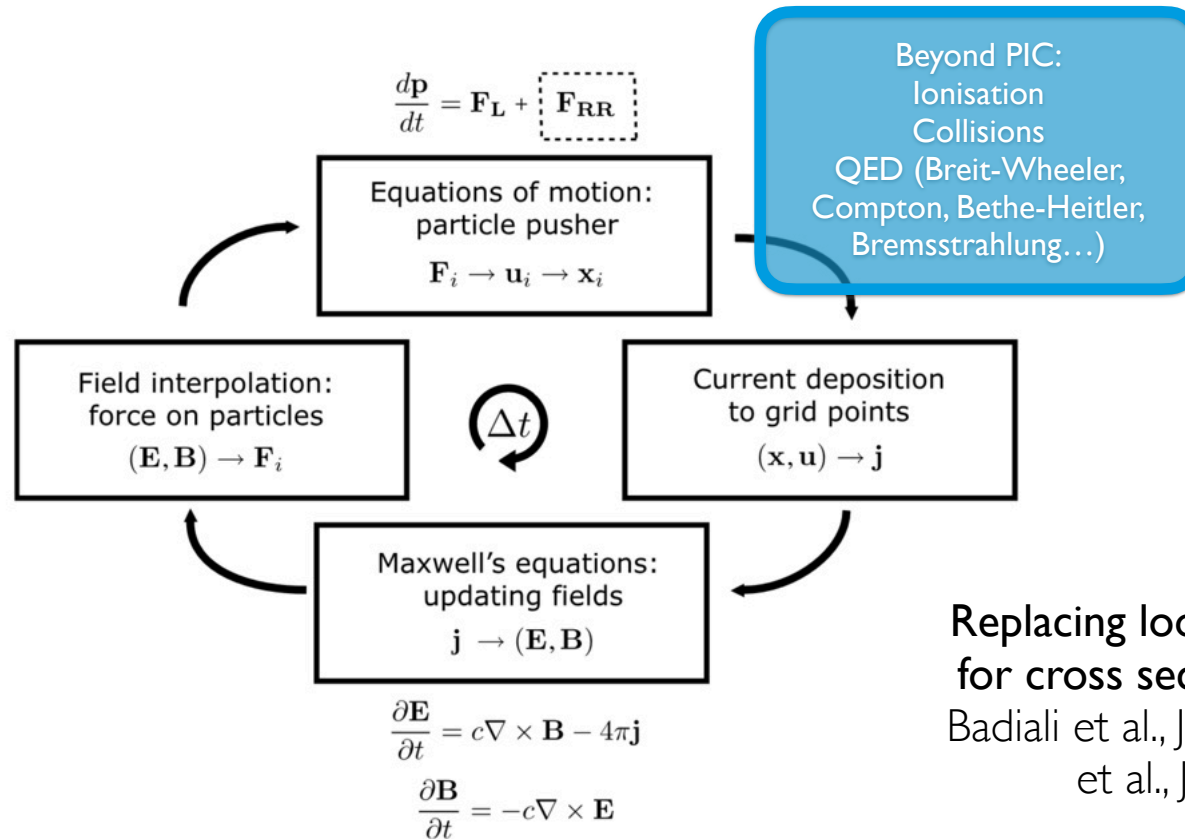
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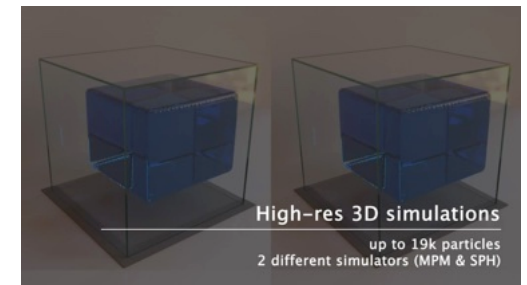
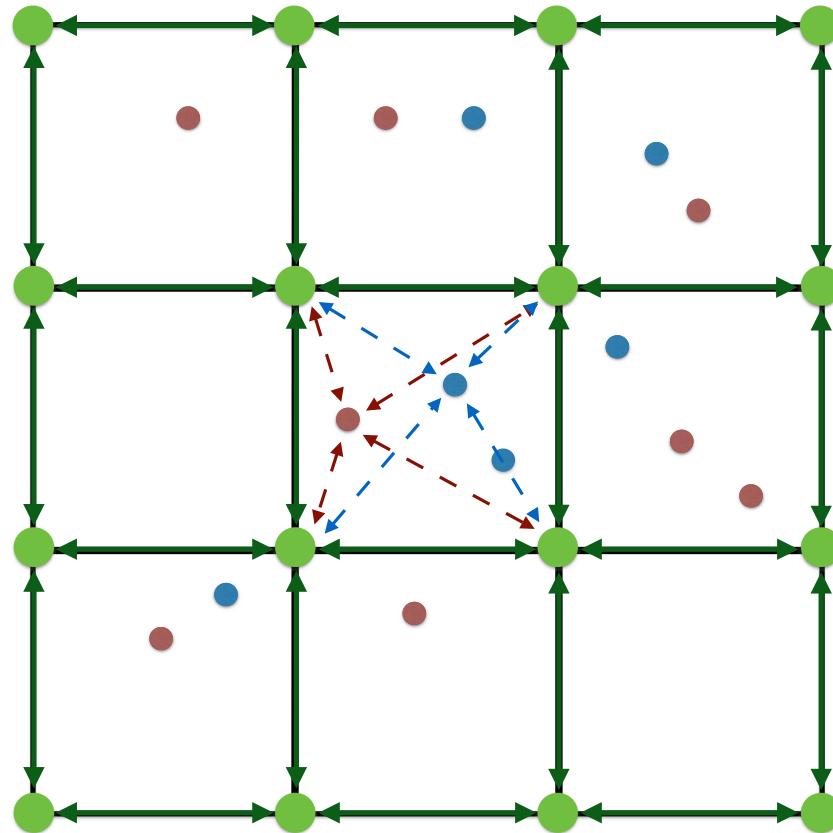
Supplementing the PIC algorithm with additional physics



MC based techniques

Replacing look-up tables e.g. for cross sections with NN
 Badiali et al., JPP 2022; Amaro et al., JCP 2026

PIC codes (and others) can be seen from a graph perspective



A. Sanchez-Gonzalez et al., ICML PMLR 8459–8468 (2020)
R. Lam et al., Science 382.6677 1416–1421 (2023)

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The Breaking of Finite
Amplitude Plasma
Oscillations

by

John Dawson

June 1959

MATT - 4



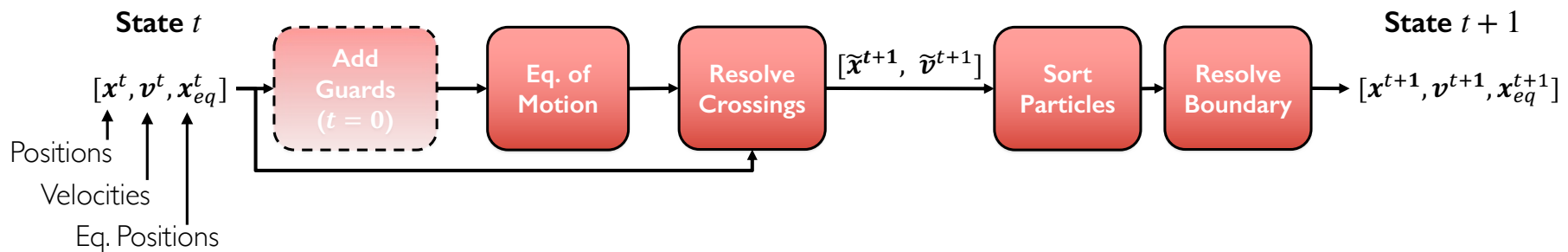
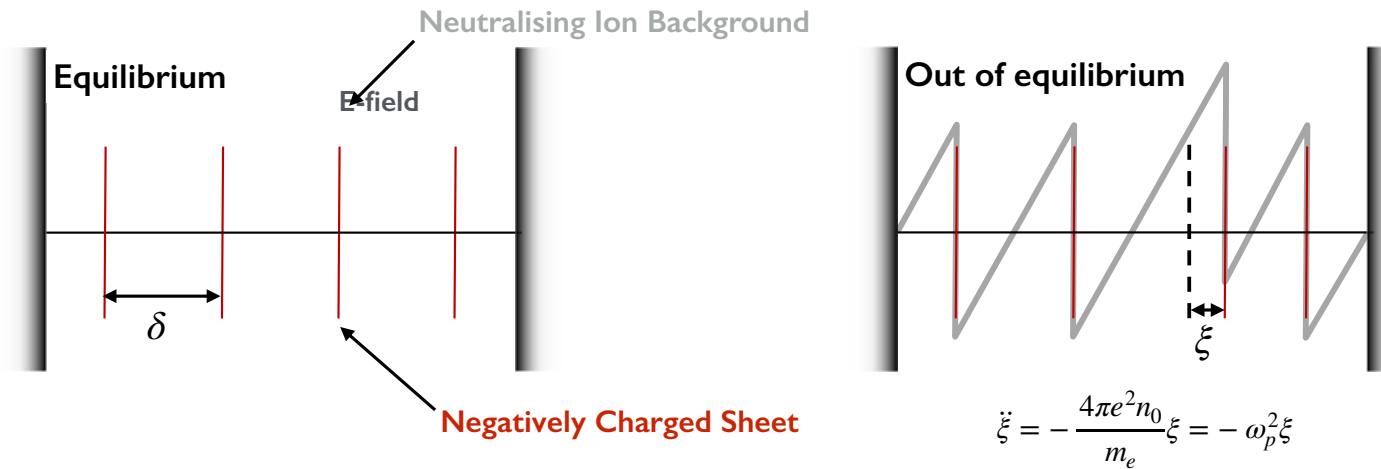
PROJECT MATTERHORN

Contract AT(30-1) - 1238 with the
US Atomic Energy Commission

AEC RESEARCH AND DEVELOPMENT REPORT

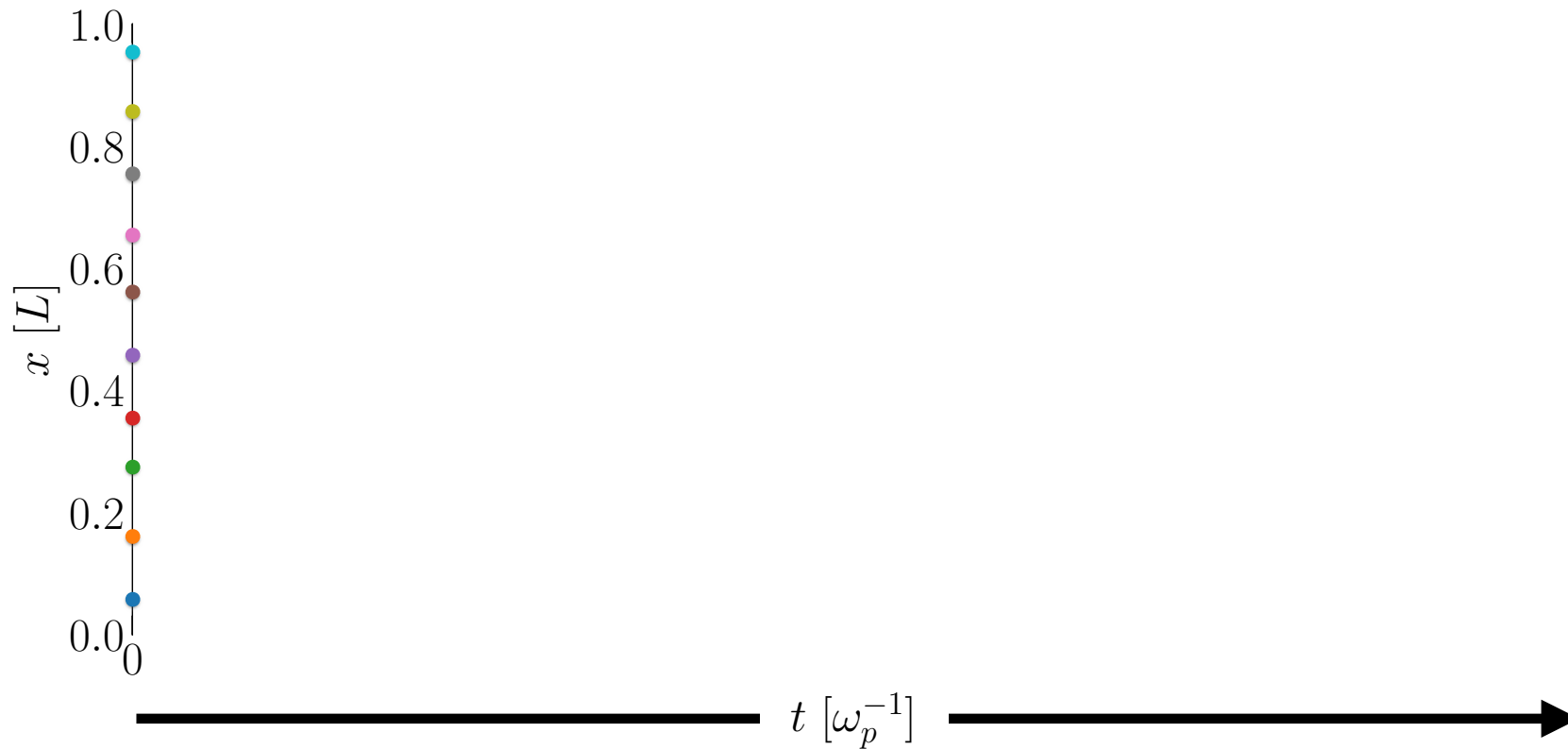
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PRINCETON, NEW JERSEY

1D Plasma Electrostatic Sheet Model



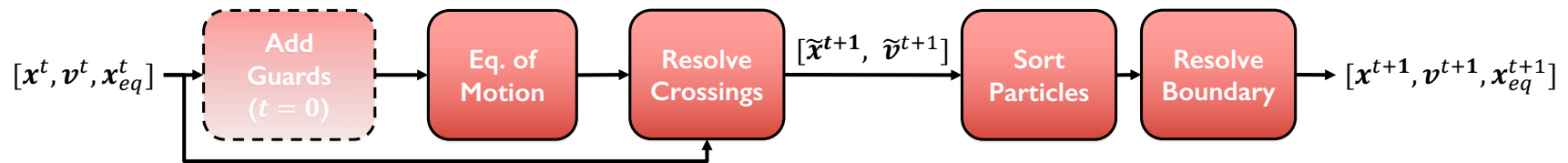
J. Dawson, Phys. Fluids 5.4, 445-459 (1962)
 J. Dawson, Methods in Computational Physics 9, 1-28 (1970)

Example of Simulation

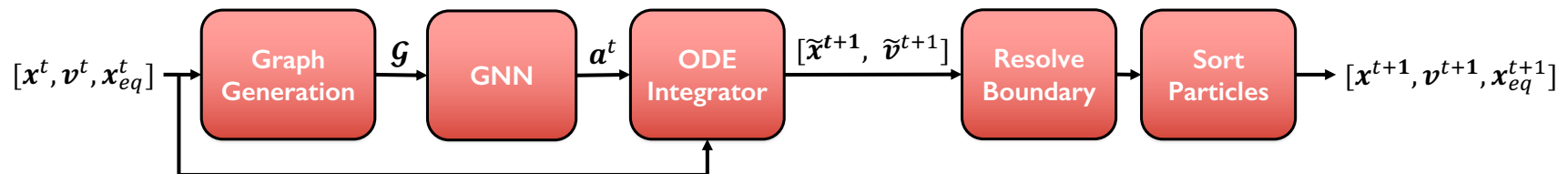


ID Plasma ESM Graph Network Simulator

Sheet Model

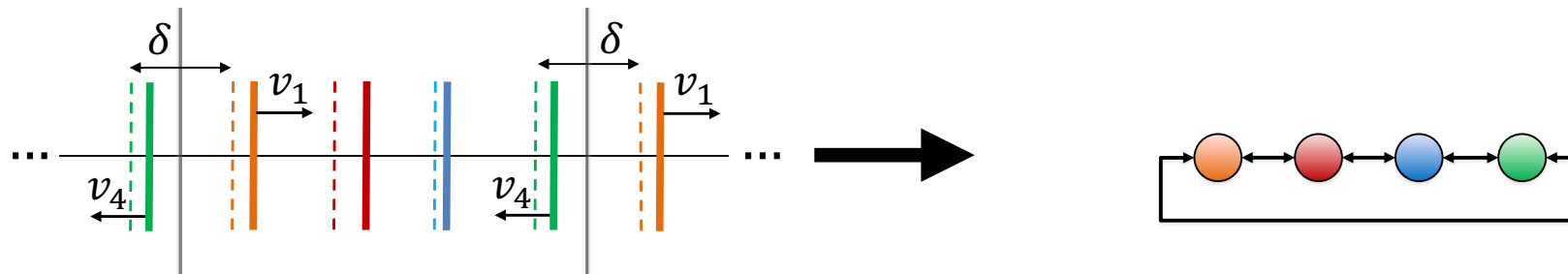


Graph Network Simulator



How do we represent the plasma as a graph?

Periodic Boundaries



i Node

$$\mathbf{n}_i^t = [\xi_i^t, v_i^t]$$

$i \rightarrow j$ Edge

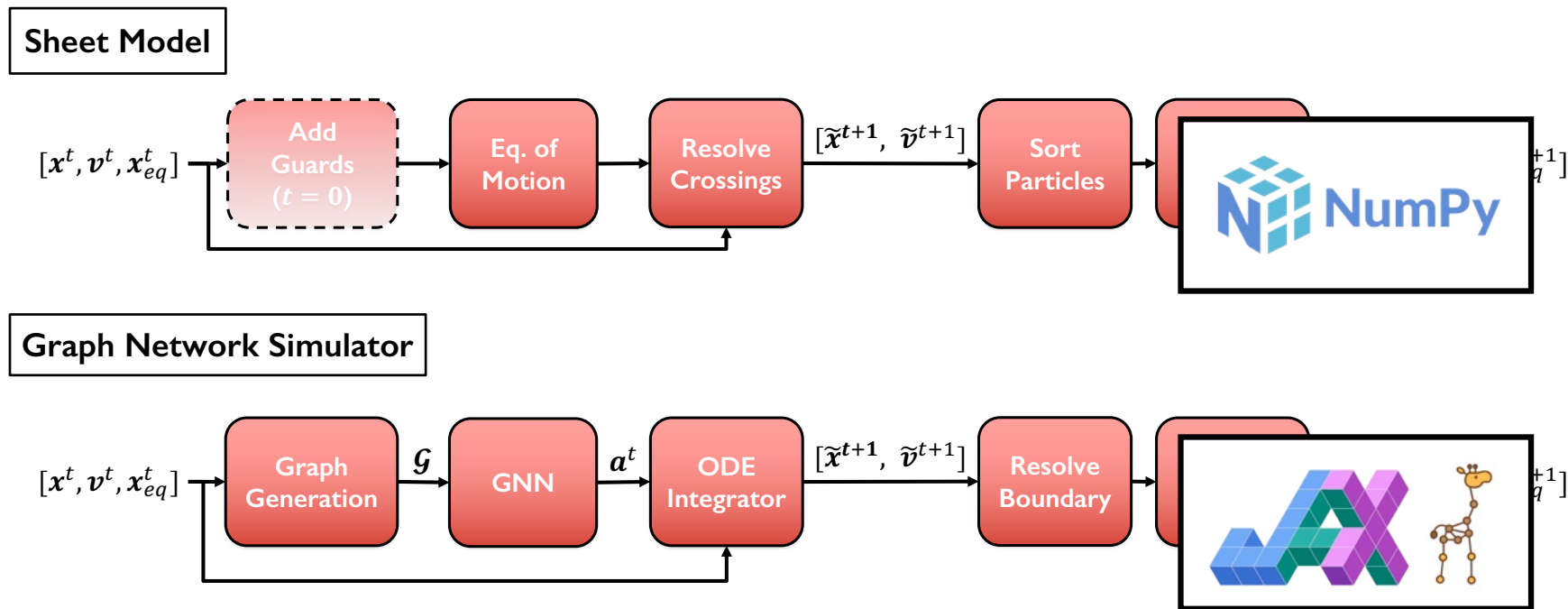
$$\mathbf{r}_{ji}^t = [x_i^t - x_j^t]$$

Target

$$a_i^t = \frac{v_i^{t+1} - v_i^t}{\Delta t}$$

All values are normalised to the intersheet spacing δ

ID Plasma ESM Graph Network Simulator



Code: <https://github.com/diogodcarvalho/gns-sheet-model>

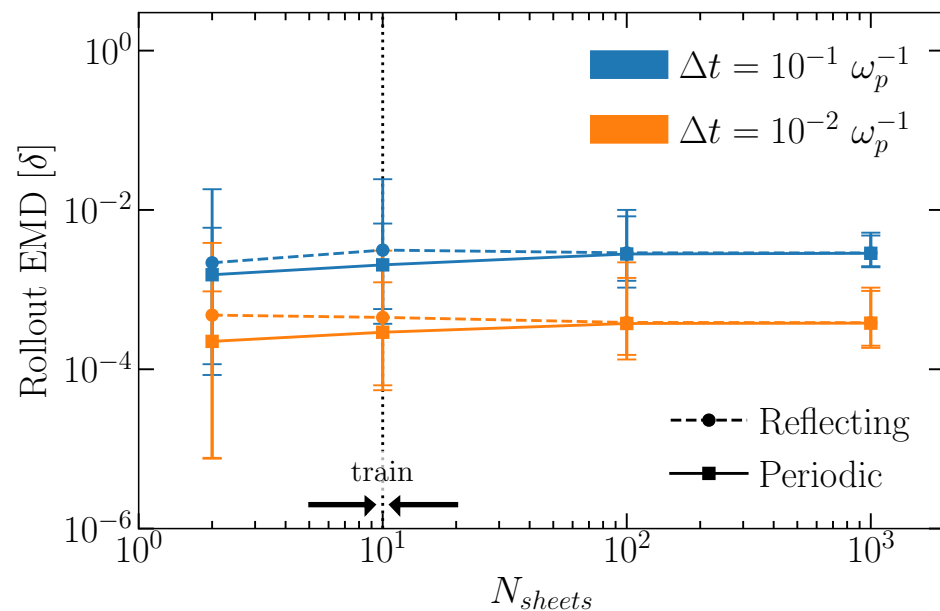
<https://github.com/google/jax>

<https://github.com/deepmind/jraph>

GNS generalizes to different number of sheets and boundary conditions

Trained on **subsampled high temporal resolution data** ($\Delta t_{orig} = 10^{-4} \omega_p^{-1}$) of **10 sheets**
moving inside a **periodic box** ($t_{sim} = 10 \omega_p^{-1}$)

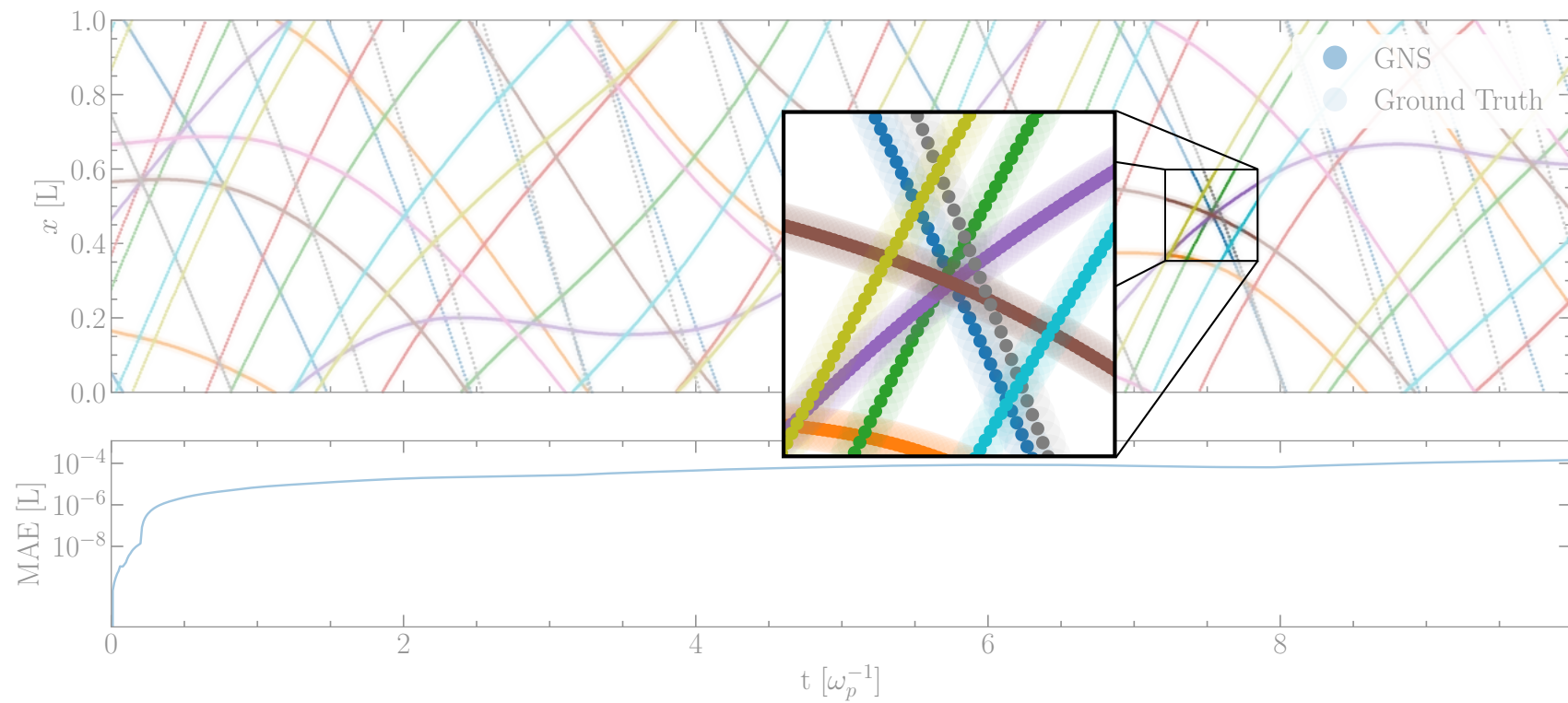
Initial positions and velocities are randomly sampled from a uniform distribution



GNS rollout errors are very small

$$\Delta t = 10^{-2} \omega_p^{-1}$$

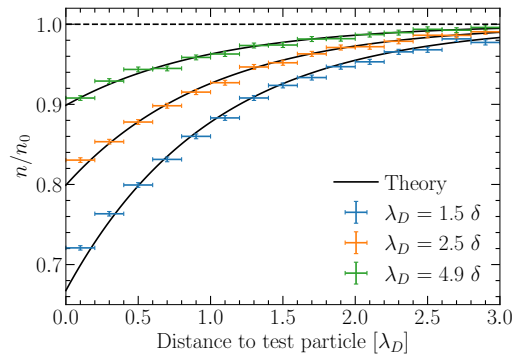
$$\text{Rollout MAE} = 5.6 \times 10^{-4} \delta$$



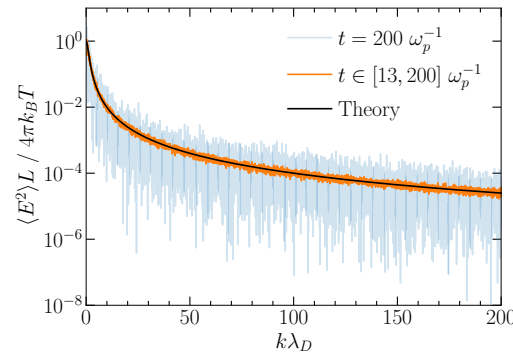
Example shown corresponds to the worst rollout error observed in the test set

GNS recovers a broad range of kinetic plasma processes

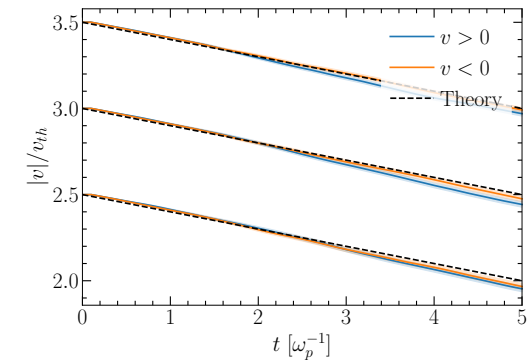
Debye Shielding



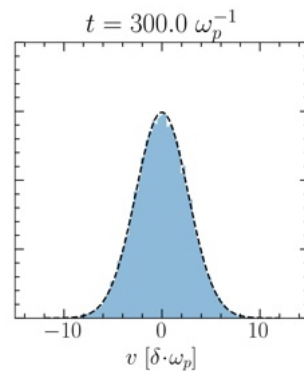
Electrostatic Fluctuations



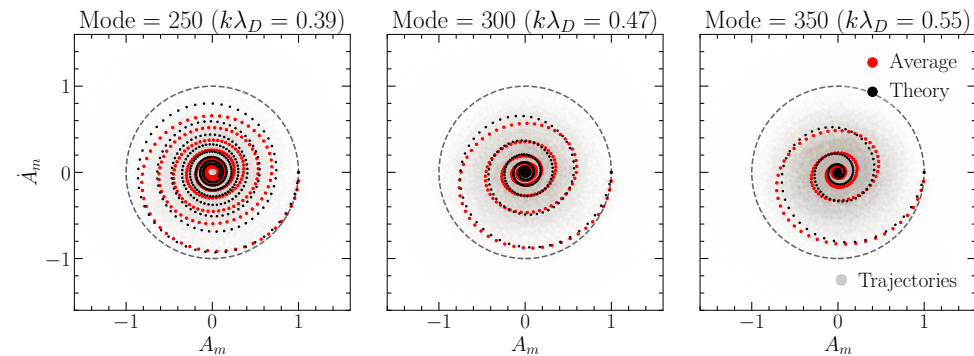
Drag on a Fast Sheet



Plasma Thermalization

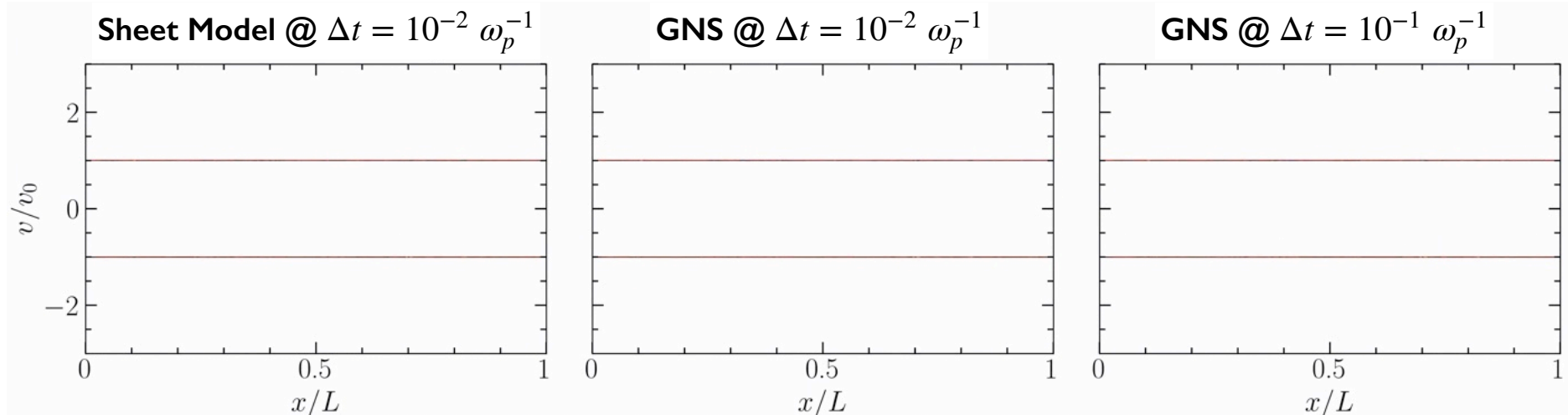


Landau Damping



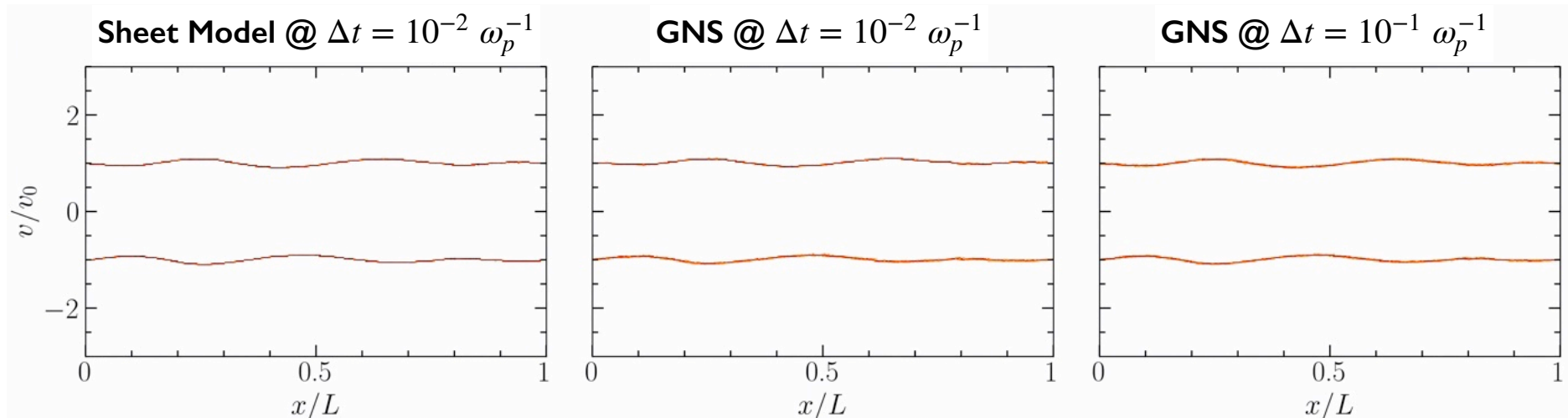
GNS recovers the two stream instability

Parameters: $N_{sheets} = 10,000$ (vs $N_{sheets}^{train} = 10$) $v_0 \approx 500 \delta \cdot \omega_p$ (vs $v_{max}^{train} = 20 \delta \cdot \omega_p$)



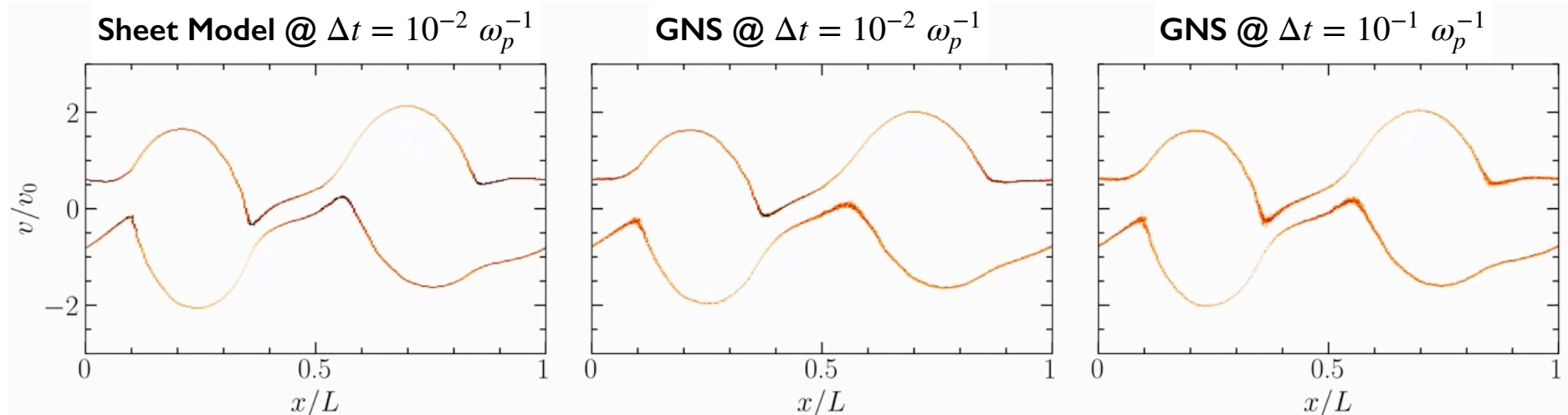
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$$\Delta\epsilon/\epsilon_0 \approx 10^{-6}$$

Run-Time $\approx 1\text{h}$

$$\Delta\epsilon/\epsilon_0 \approx 10^{-2}$$

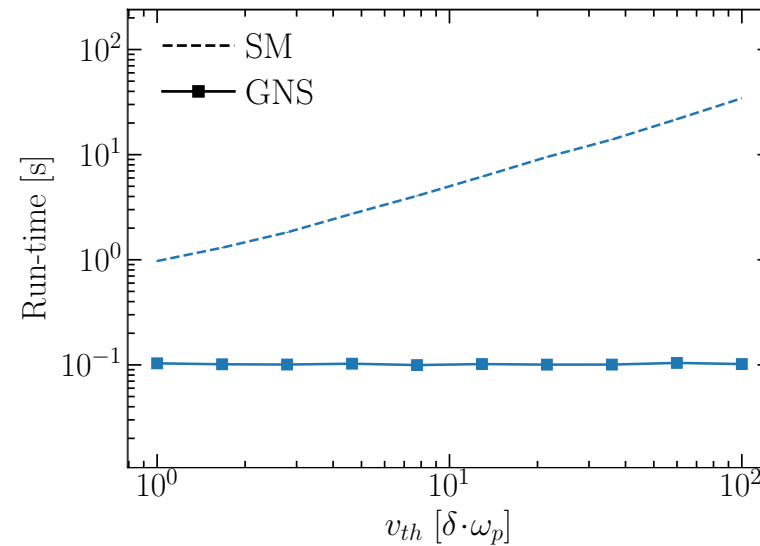
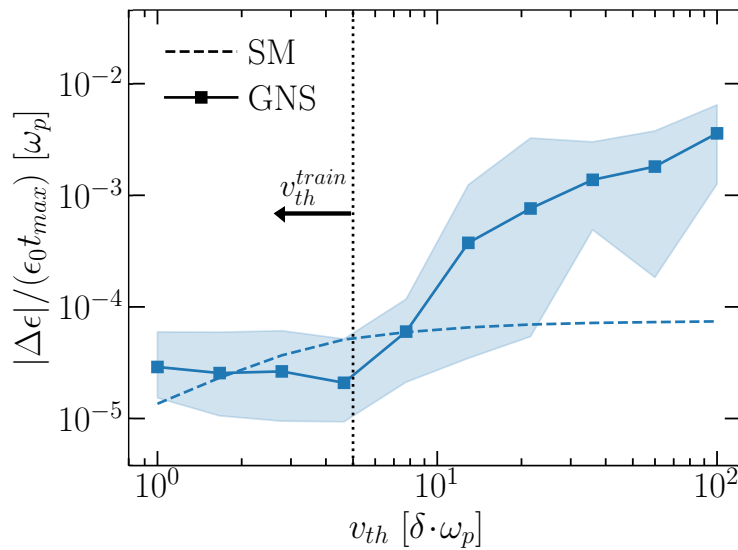
Run-Time $\approx 1\text{min}$

$$\Delta\epsilon/\epsilon_0 \approx 10^{-2}$$

Run-Time $\approx 10\text{s}$

GNS conserves energy similarly to Sheet Model while being significantly faster*

Parameters: $N_{sheets} = 1000$, velocities sampled from thermal distribution, $\Delta t_{GNS} = 10^{-1} \omega_p^{-1}$



***Note: GNS is implemented in JAX** (GPU), Sheet Model is implemented in NumPy (CPU)

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Klimontovich + Maxwell's equations

$$\frac{\partial N}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} N - \frac{q}{m} (\mathbf{E}^m + \mathbf{v} \times \mathbf{B}^m) \cdot \nabla_{\mathbf{v}} N = 0$$

This is the particle-in-cell algorithm (with finite-size particles):

statistical mechanics is well-known (e.g. H. Okuda and C. Birdsall, (1970), A. Langdon and C. Birdsall, (1970), R. Hockney (1971), M. Touati et al. (2022))

Born-Infeld electrodynamics

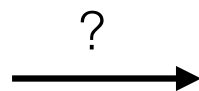
Numerical collision operator has been derived in previous works (but not tested):

Can this be learned from the simulation data in the weakly collisional regime? Can we then use this as a template to study conditions/regimes beyond existing theories?

Can we describe phase-space dynamics using a Fokker-Planck operator?

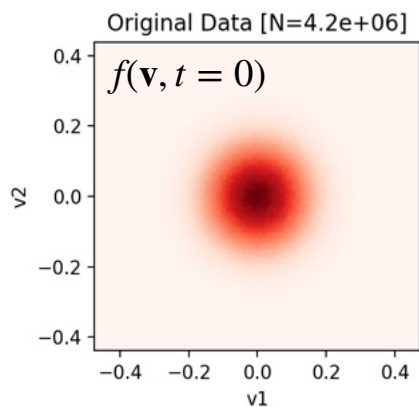
PIC (Klimontovich)

$$\frac{\partial N}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} N - \frac{q}{m} (\mathbf{E}^m + \mathbf{v} \times \mathbf{B}^m) \cdot \nabla_{\mathbf{v}} N = 0$$

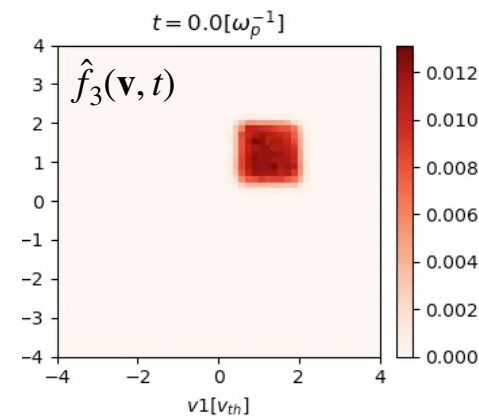
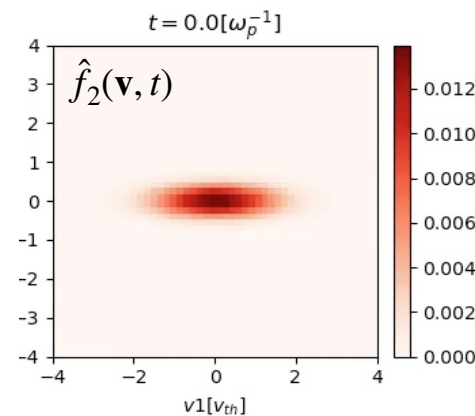
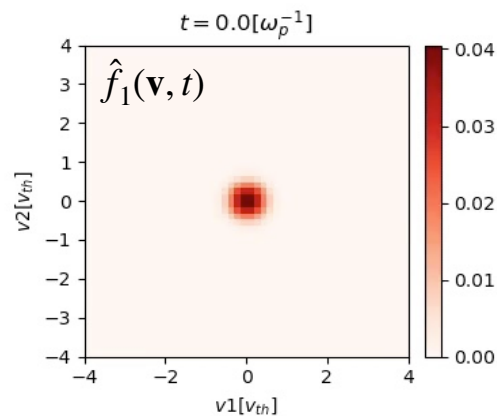


What if we want (Fokker-Planck)?

$$\frac{\partial \hat{f}(\mathbf{v}, t)}{\partial t} = -\nabla_{\mathbf{v}} \cdot (\mathbf{A} \hat{f}) + \frac{1}{2} \nabla_{\mathbf{v}} \nabla_{\mathbf{v}} \cdot (\vec{\mathbf{D}} \hat{f})$$



Thermal Plasma



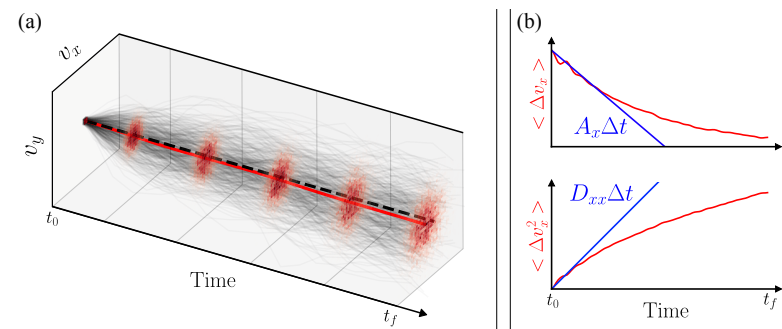
How do we estimate **A** (advection) and $\vec{\mathbf{D}}$ (diffusion)?

How do we estimate \mathbf{A} (advection) and $\overleftrightarrow{\mathbf{D}}$ (diffusion)?

Option 1: From raw particle data

The “correct” approach if possible

Not feasible for larger systems (memory-wise) unless it is done at run-time

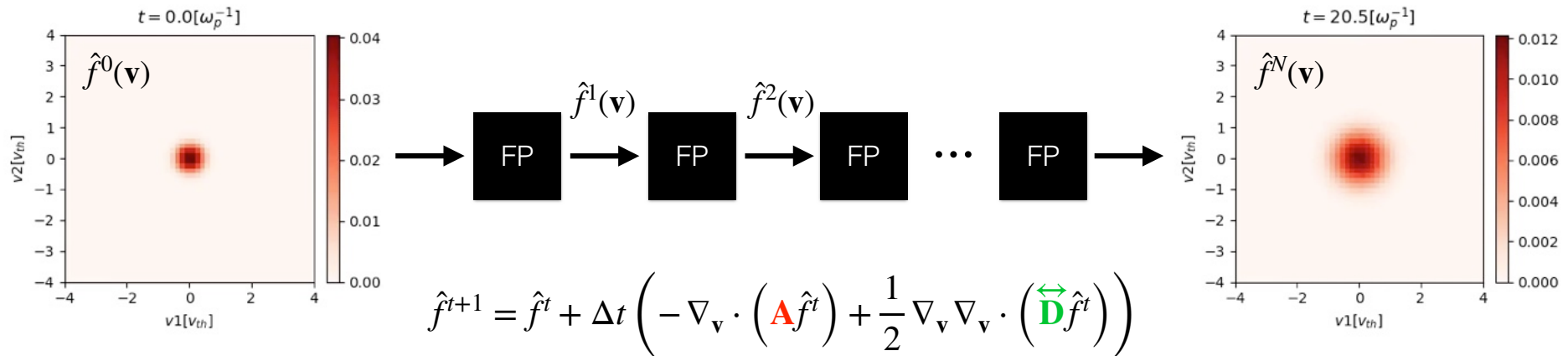


Option 2: From the phase-space evolution of sub-populations

Can be done in post-processing with a differentiable solver

Ill-posed problem: non-unique solution for coefficients

Learning advection / diffusion from evolution of sub-populations

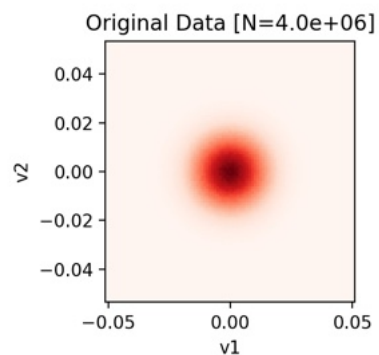


We can make the **Fokker-Planck solver differentiable** and frame this as an **optimisation task**

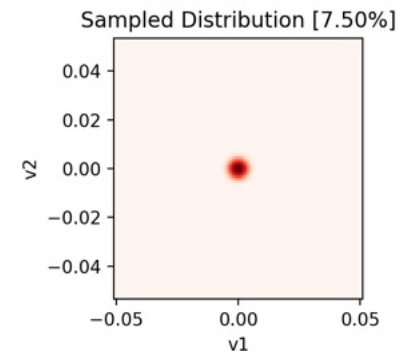
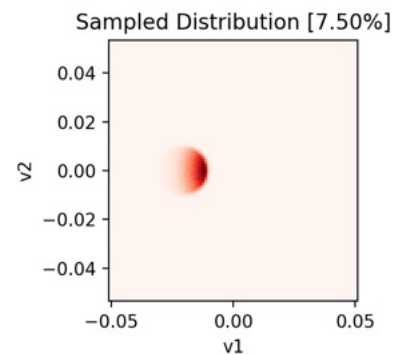
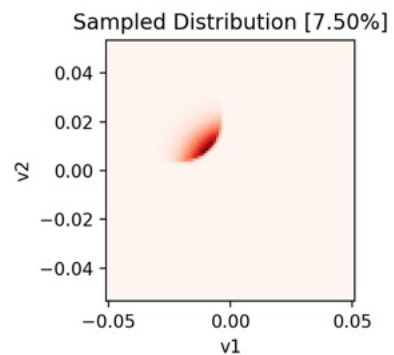
$$\min_{\mathbf{A}, \mathbf{D}} \left\| \hat{f}_{predicted}^N - \hat{f}_{true}^N \right\|$$

This is an ill-posed problem (there exists a family of solutions) \longrightarrow Train with multiple sub-populations

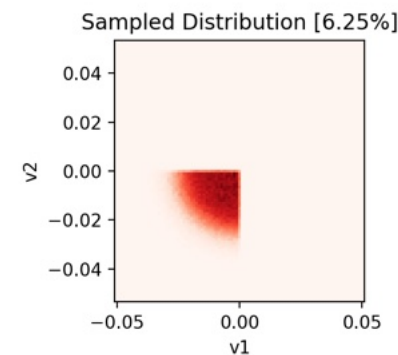
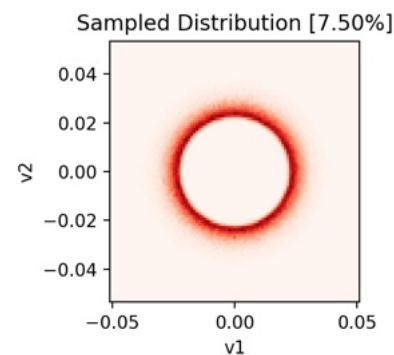
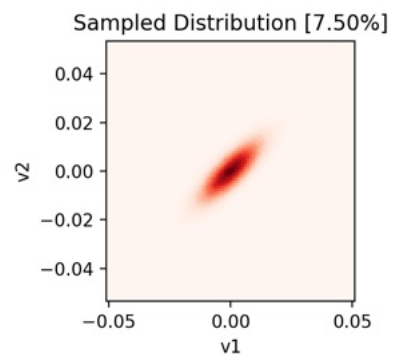
We use different sets of training and test sub-populations



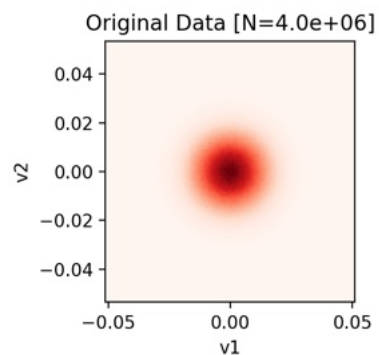
Train (9x)



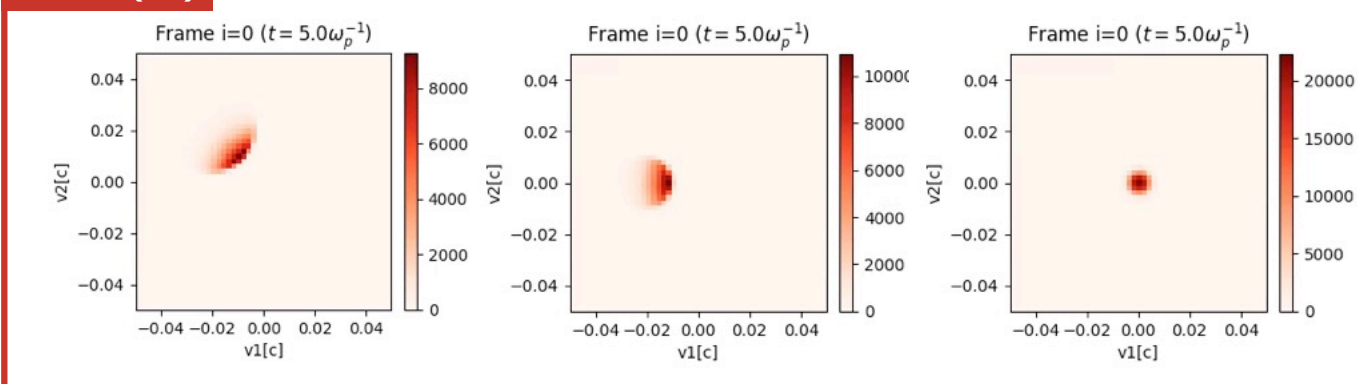
Test (20x)



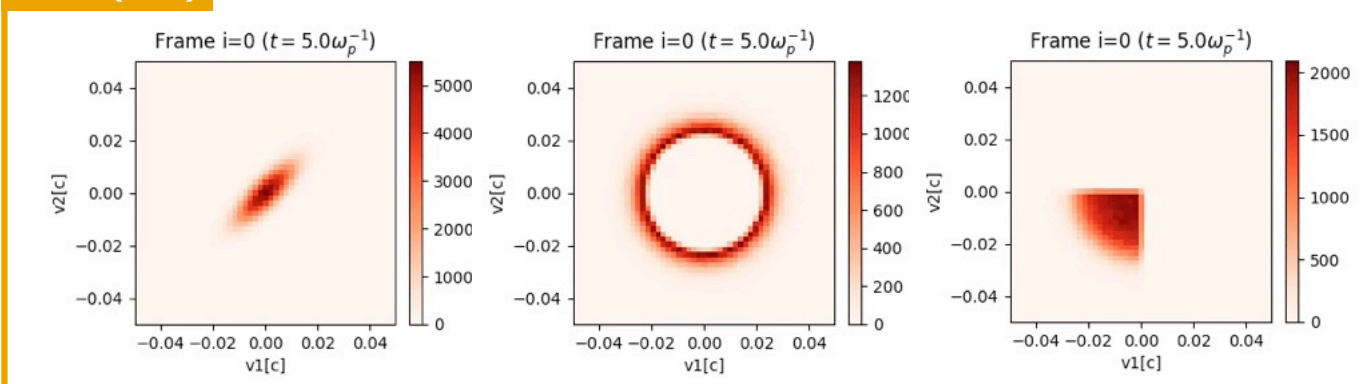
We use different sets of training and test sub-populations



Train (9x)

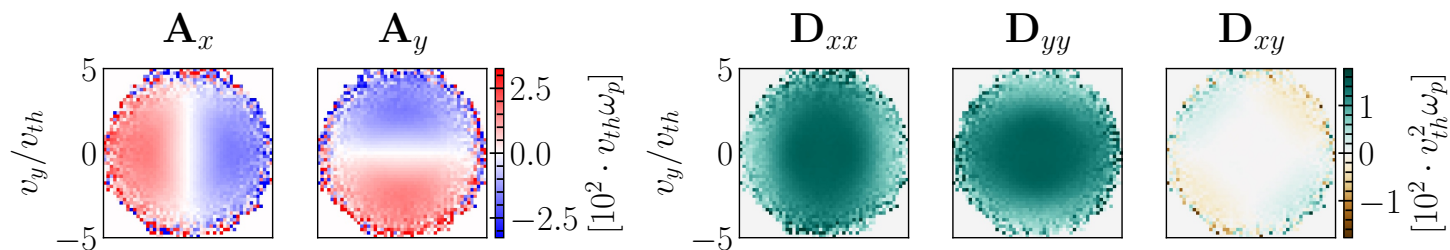


Test (20x)



We can parameterise A/D using a Tensor (discrete) or a NN (continuous)

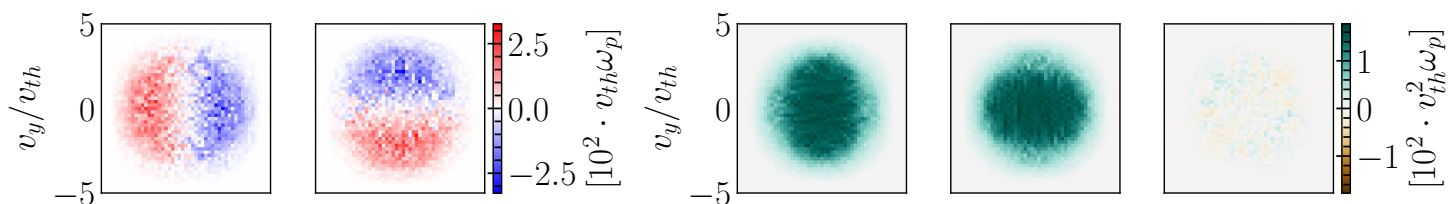
Tracks



Tensor

$$A_i[v_x, v_y]$$

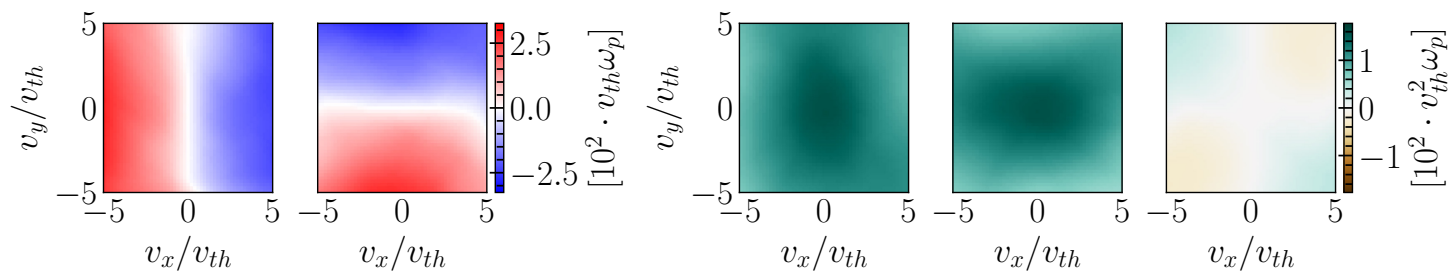
$$D_{ij}[v_x, v_y]$$



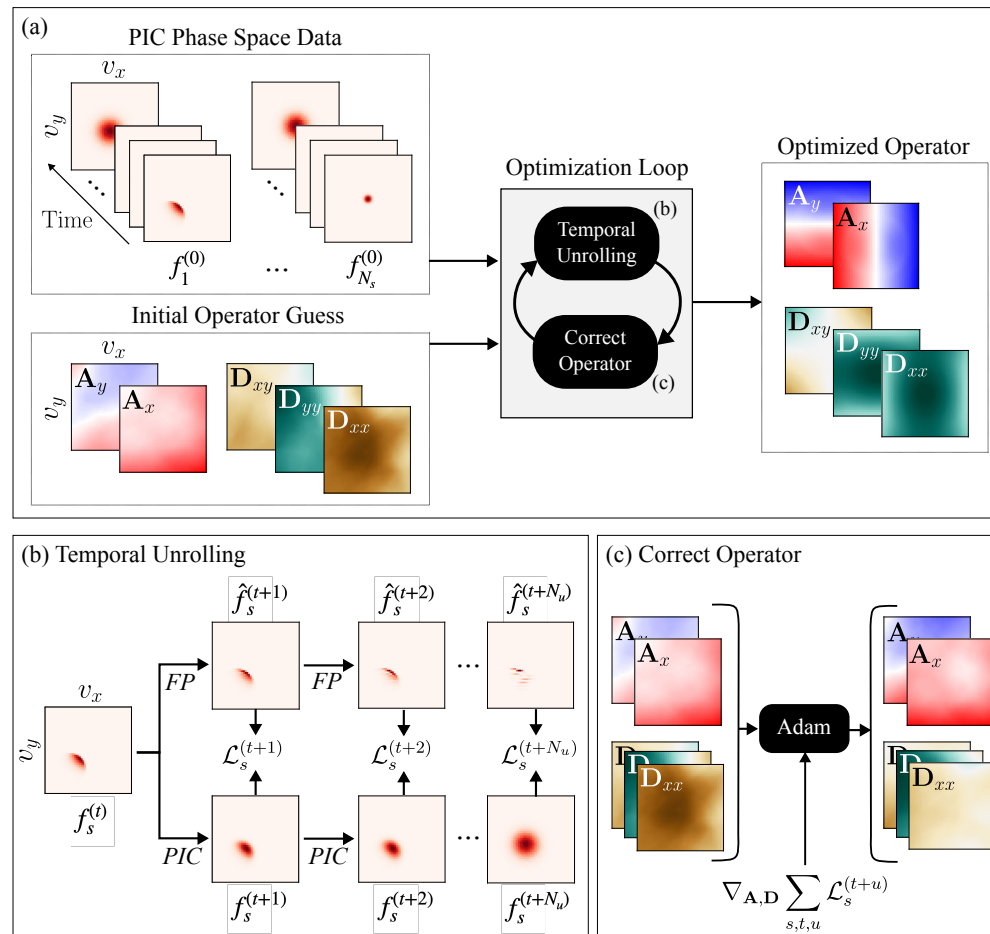
NN

$$A_i = MLP(v_x, v_y)$$

$$D_{ij} = MLP(v_x, v_y)$$



Inference of advection/diffusion from differentiable Fokker-Planck solver

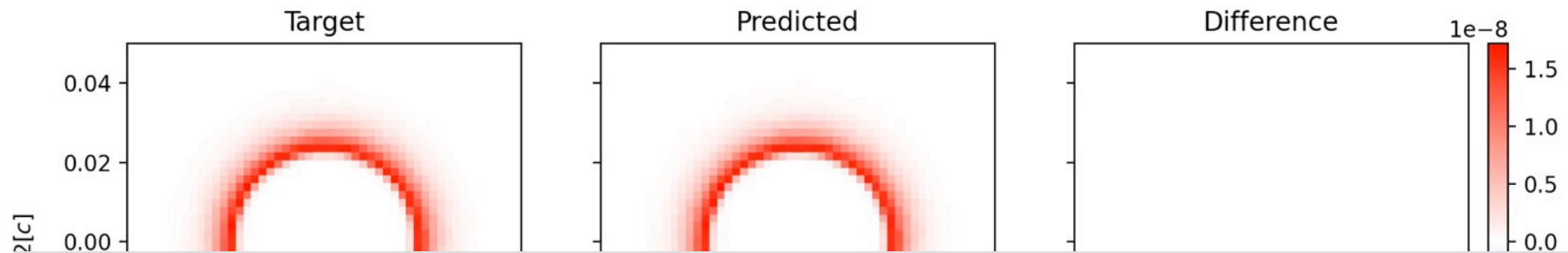


Can these operators reproduce dynamics?

$$f^{t+1} = f^t + \Delta t \left(-\nabla_{\mathbf{v}} \cdot (\mathbf{A}f^t) + \frac{1}{2} \nabla_{\mathbf{v}} \nabla_{\mathbf{v}} \cdot (\overleftrightarrow{\mathbf{D}}f^t) \right)$$

Can these operators reproduce dynamics? (Yes)

$$f^{t+1} = f^t + \Delta t \left(-\nabla_{\mathbf{v}} \cdot (\mathbf{A}f^t) + \frac{1}{2} \nabla_{\mathbf{v}} \nabla_{\mathbf{v}} \cdot (\overleftrightarrow{\mathbf{D}}f^t) \right)$$



Next steps

General purpose library: from phase space data, retrieve **A** and **D** (to be inserted on Fokker-Planck codes/other models) for varying plasma conditions, and from different sources of data

Sub-module to capture (PIC or other) collisions/anomalous transport coefficients for mesoscale simulations

Meta analysis: use different A and D for different plasma conditions (n, B, T) to learn more general behaviour e.g. via sparse regression

Ab initio kinetic plasma simulations and the “ground truth”

Some questions

Speeding standard kinetic simulations + novel simulators

Learning advection and diffusion coefficients from PIC simulation data

Pushing PIC simulations towards capturing *ab initio* collisions:
from many particles per cell to *many cells per particle*

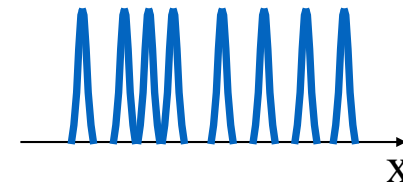
What is the (classical) ground truth?

Klimontovich + Maxwell's equations

$$\frac{\partial N}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} N - \frac{q}{m} (\mathbf{E}^m + \mathbf{v} \times \mathbf{B}^m) \cdot \nabla_{\mathbf{v}} N = 0$$

This is the particle-in-cell algorithm (with finite-size particles): statistical mechanics is well-known (e.g. H. Okuda and C. Birdsall (1970), A. Langdon and C. Birdsall (1970), R. Hockney (1971), M. Touati et al. (2022)) + Born-Infeld electrodynamics

What if the cell/particle size is (shorter than) the classical electron radius r_e ?



What are the challenges of running $\ll 1$ ppc?

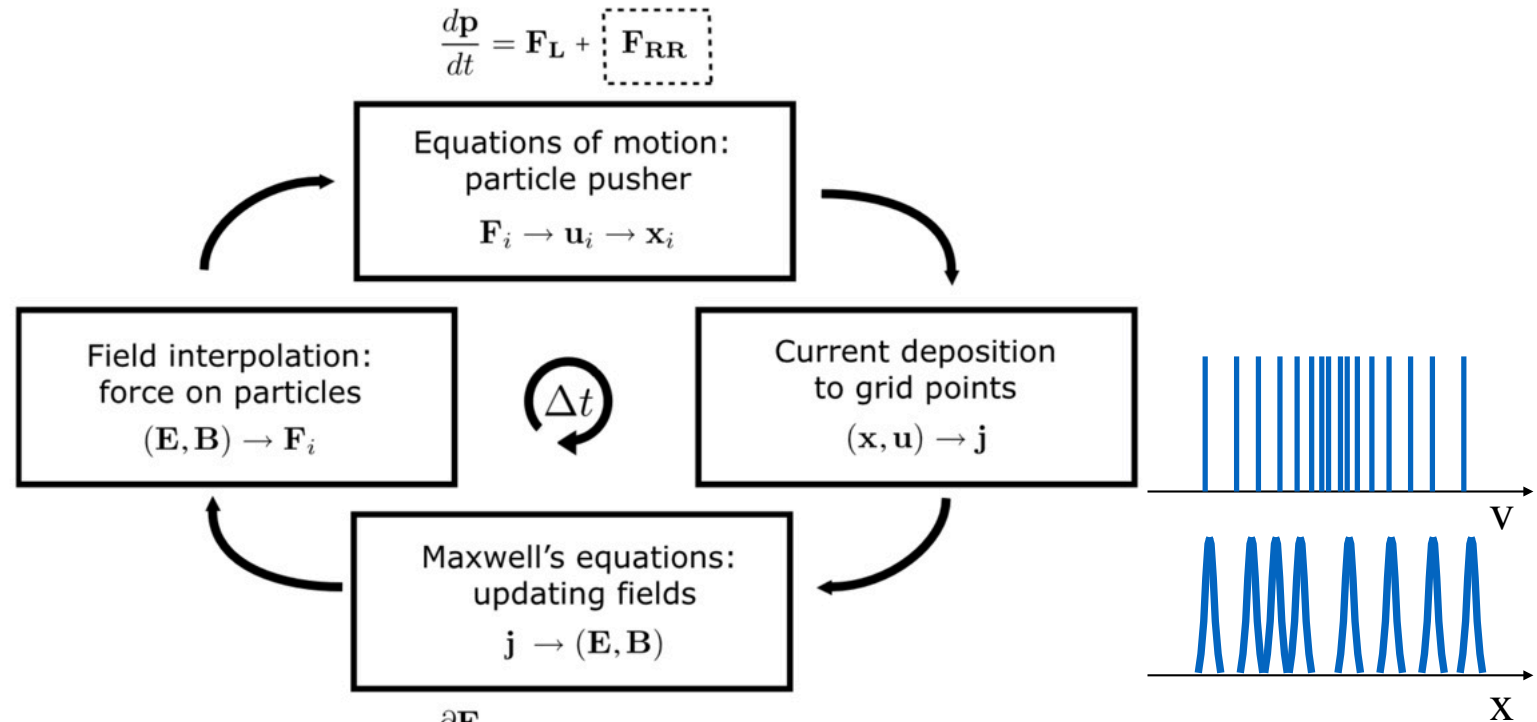
Field initialization becomes critical + computation determined by grid (N^3 or N^2)

Numerical heating (still need to resolve Debye length) + Very small time steps (CFL) + numerical transition radiation

Validation against theory (but theory is very limited) or computational models (MD non relativistic)

Shape functions to capture quantum effects?

High resolution simulations can capture collisional dynamics



Nonlinear collisional absorption in laser-driven plasmas

C. D. Decker *et al.*, *Phys. Plasmas* 1, 4043–4049 (1994)

Disorder induced heating

M. D. Acciarri *et al.*, *Plasma Sources Sci. Technol.* 33 035009 (2024)

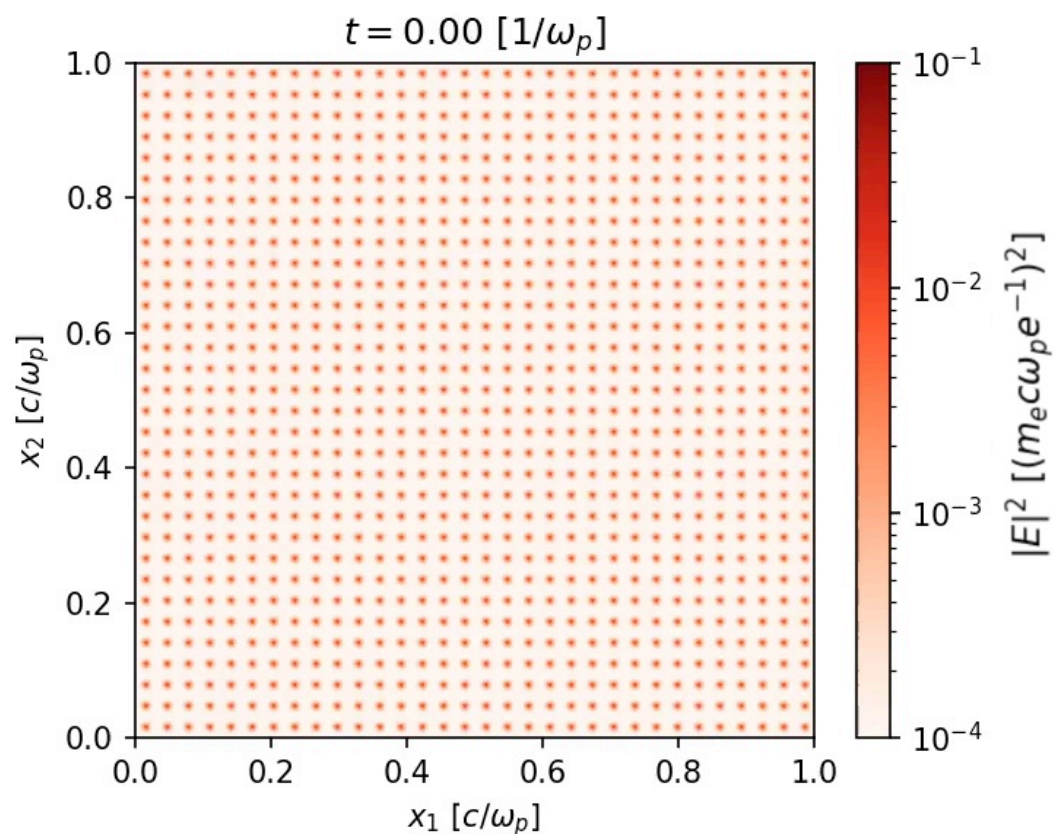
$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi \mathbf{j}$$

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

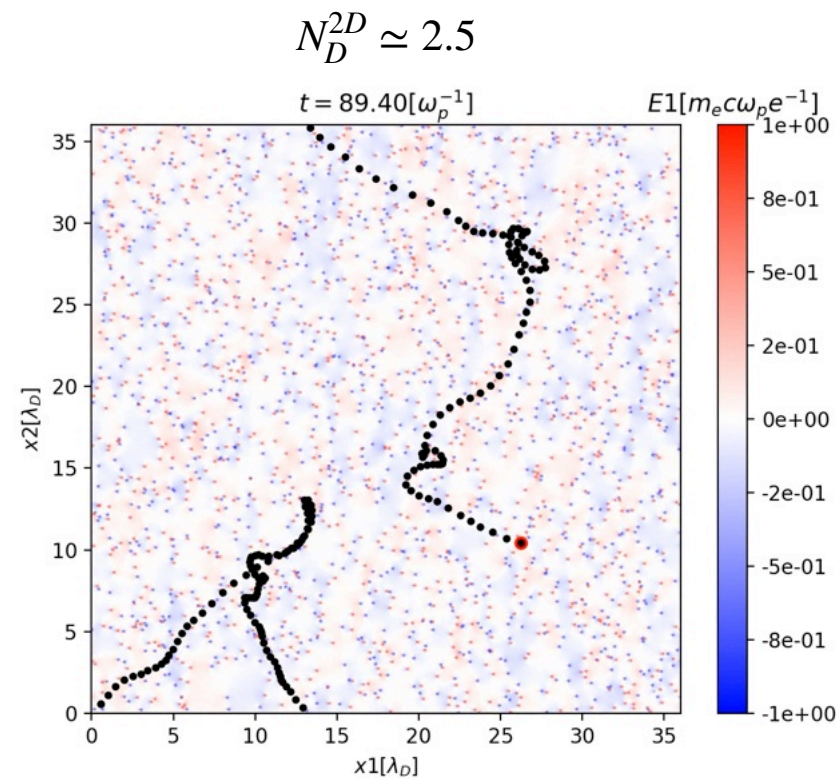
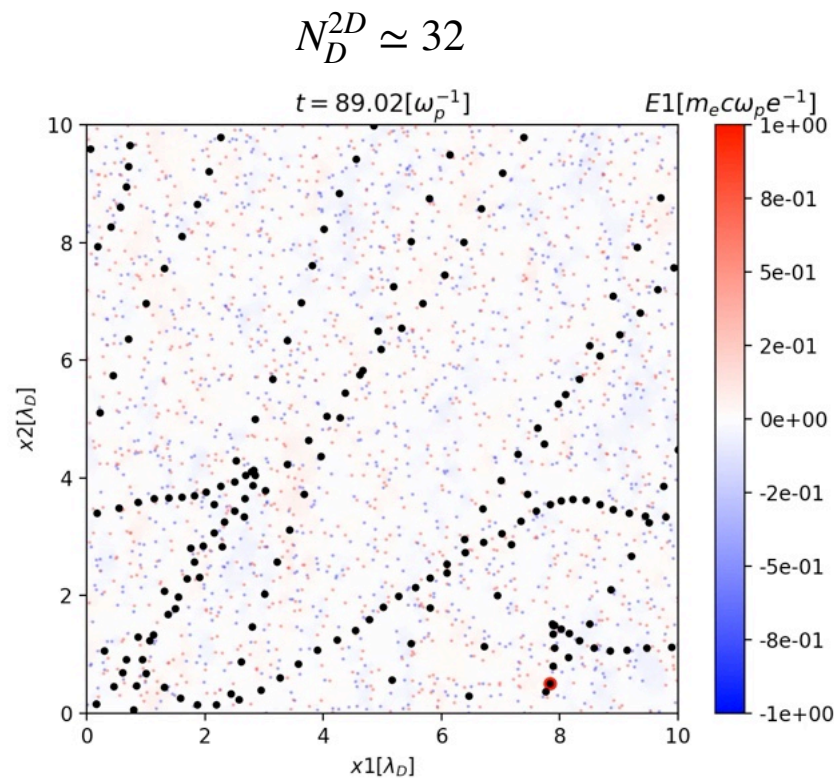
$$N_s(\mathbf{x}, \mathbf{v}, t) = \sum_{i=1}^{N_0} \delta[\mathbf{x} - \mathbf{X}_i(t)] \delta[\mathbf{v} - \mathbf{V}_i(t)]$$

Pushing PIC simulations to capture self-consistent (e.m.) collisions

$$n\lambda_D^2 \simeq 0.1 \quad (\text{Average interparticle distance} \gg \lambda_D)$$



Can we learn operators that describe the collisional dynamics?



Summary

Exciting opportunities from the interplay between *ab initio* PIC simulations and AI

Improvements in standard simulators

Novel approaches to standard algorithms

Tapping on unique properties of synthetic high fidelity PIC data

Interplay between HPC + PIC and AI is just starting: “There are (many) unknown unknowns” (which is great for science!)

