

# Unraveling the nature of grain boundary migration by Potts model simulations

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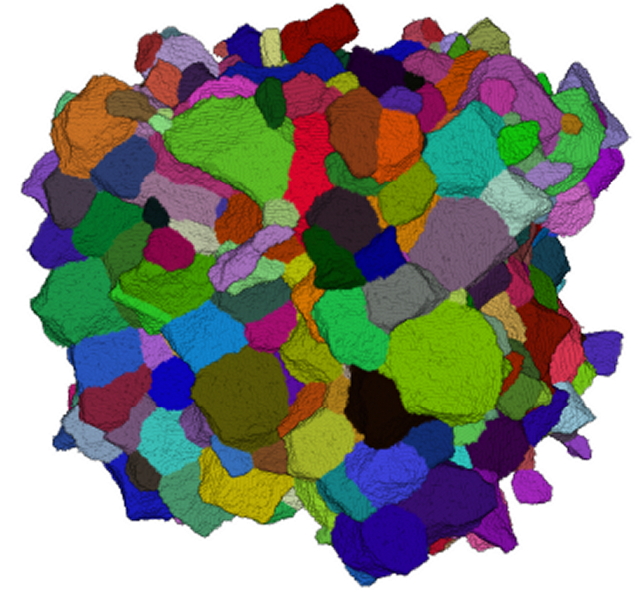
Dana Zöllner

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Department of Materials Physics

Ulm University, Germany  
Institute of Micro and Nanomaterials

# Outline

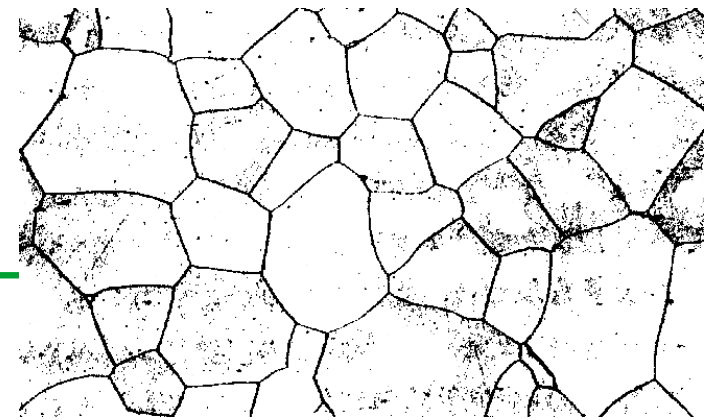
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- Introduction
- Standard Monte Carlo Potts Model
  - History, Algorithms & Problems
  - Simulation of Grain Boundary Motion
- A Potts Model for Nanocrystalline Grain Growth
  - Size Effects in a Size-Invariant Simulation
- Going Back to 2D Grain Growth
  - Grain Boundary Migration in Graphene
- Summary and Outlook

# Introduction

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## Importance of grain growth:

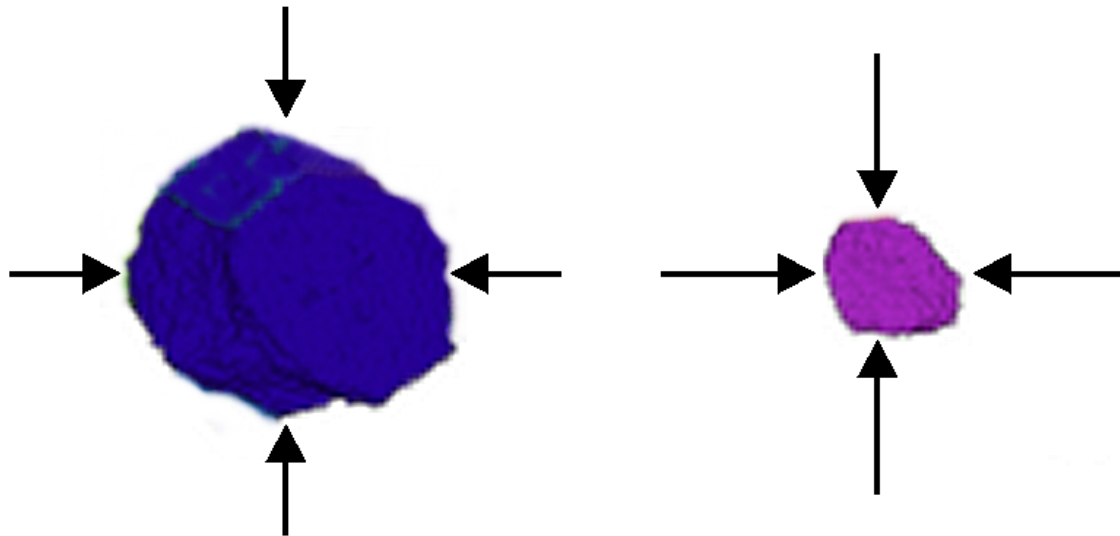
- most materials properties depend on grain size
  - Hall-Petch-effect: small grains = higher yield stress
- grain microstructure = important feature of polycrystalline solids
- understanding & control = key to improve materials properties
  
- Problem: thermally activated grain boundary migration
  - changing grain sizes

# Introduction

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## Basic principles of grain growth:

- All grains of a polycrystalline material try to shrink!  
→ thermodynamically: reduction of surface energy



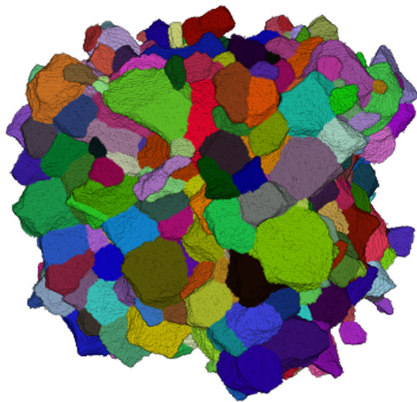
- The smaller a grain is, the stronger it pushes its atoms to the neighboring grains.

# Introduction

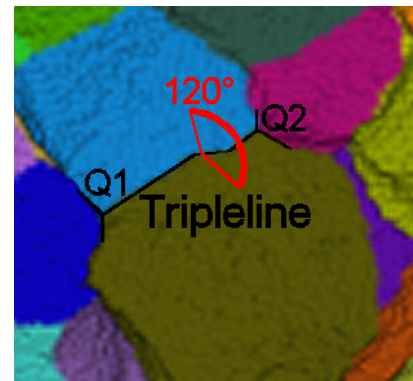
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## Basic principles of grain growth:

- We have a network of grain boundaries.
  - global: volume conservation
  - local: equilibrium at grain boundaries



versus



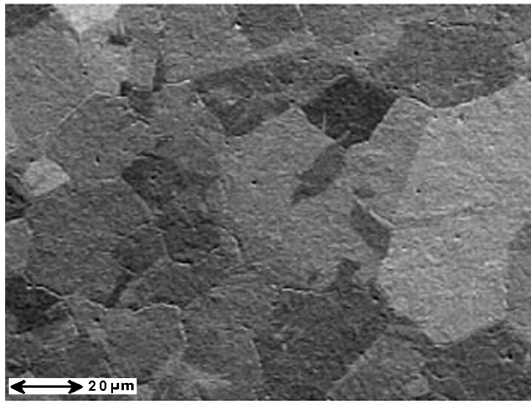
- Therefore, whether a grain will shrink or growth depends on its size and on its neighboring grains.
  - complex 3D problem

# Introduction

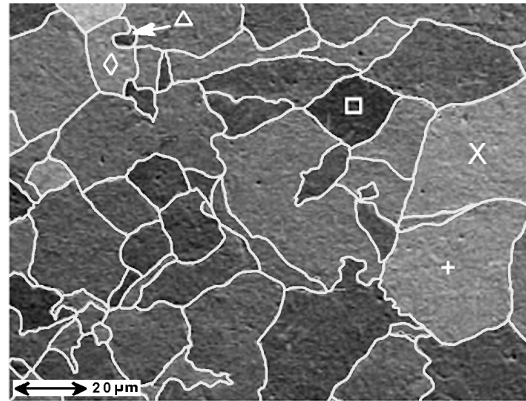
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## Experimental observation of grain growth:

- SEM observation



a



b

[I. Fielden: Mater. Sci. Forum 467-470 (2004) 875]

[D. Zöllner, I. Fielden, P. Streitenberger: Prac. Metallogr. 49 (2012) 428]

# Introduction

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What is known about normal grain growth:

- curvature driven kinetics
- follows **parabolic growth law**

$$\langle R \rangle^{1/n} = b \cdot t + \langle R \rangle_0^{1/n}$$

- shows **self-similar scaling**

$$F(R, t) = g(t) \cdot f(x), \quad g(t) = \frac{N}{\langle R \rangle}, \quad x = \frac{R}{\langle R \rangle}$$

- rate of area resp. volume change

- 2D: von Neumann-Mullins  $\dot{A} = \frac{k \cdot \pi}{3} \cdot (s - 6)$

- 3D: MacPherson-Srolovitz  $\dot{V} = -2\pi \cdot M\gamma \cdot \left( \mathcal{L}(R) - \frac{1}{6} \sum_{i=1}^n E_i(R) \right)$

# Introduction

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## Why use mesoscopic computer simulations:

- bridge the gaps between experiments and theories
- give insight into fundamental physical processes
- give information on morphology and topology
- allow a far more precise manipulation of materials properties than can be realized experimentally
- can be understood as in-situ computer experiments

# Standard Monte Carlo Potts Model

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## History:

- *Monte Carlo method*: all techniques of numerical mathematics using probabilistic methods to solve deterministic problems
  - *first application in physics*: in solid states physics to model ferromagnetic materials → two spins
  - 1952: generalisation by Potts to Q-states model
  - 1984: applied by Anderson et al. to grain growth
- represents *curvature driven grain growth* and takes the *interaction of grains* into consideration *without using complex motion equations*

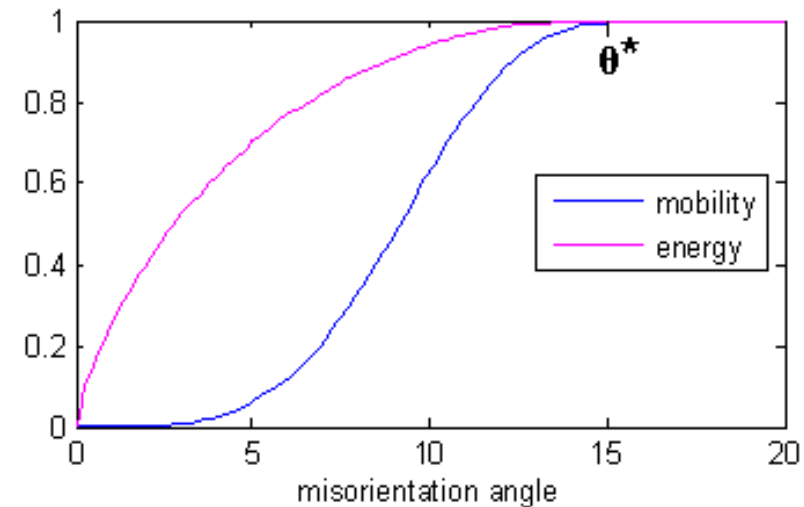
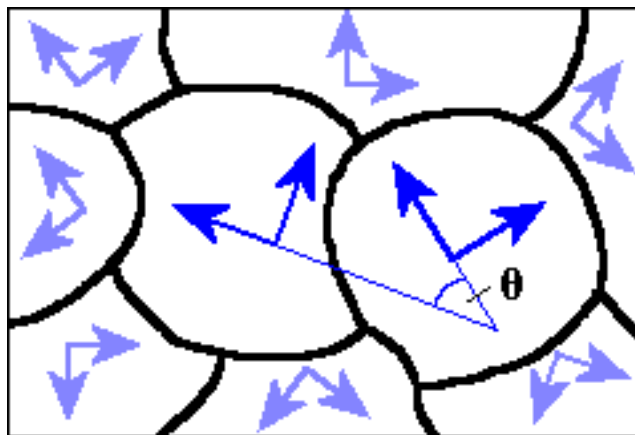
# Standard Monte Carlo Potts Model

## Algorithm - Parameters:

- grain boundary energy  $\gamma$  and mobility  $m$  depend only on misorientation angle  $\theta$  between neighbouring grains

$$\gamma = \begin{cases} \frac{\theta}{\theta^*} \left( 1 - \ln \frac{\theta}{\theta^*} \right) & \text{if } \theta \leq \theta^* \\ 1 & \text{else} \end{cases}$$

$$m = \begin{cases} 1 - \exp \left( -B \left( \frac{\theta}{\theta^*} \right)^n \right) & \text{if } \theta \leq \theta^* \\ 1 & \text{else} \end{cases}$$



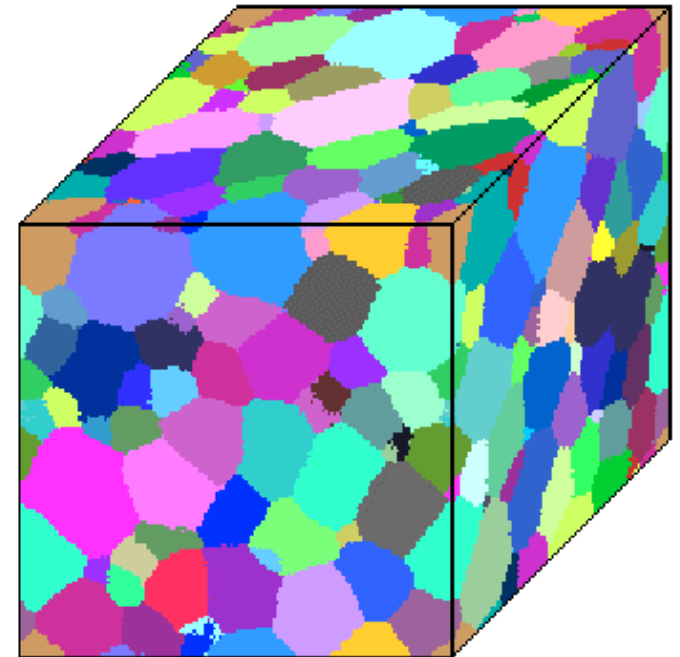
[D. Zöllner: *Monte Carlo Potts Model Simulation...* (2006); ISBN 3-8322-5633-4]

# Standard Monte Carlo Potts Model

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## Algorithm - Parameters:

- cubic 3D lattice
- total  $N > 200^3$  lattice points
- each lattice point has 26 nearest neighbours
- smallest size unit = Monte Carlo Unit (MCU) = one lattice point = represents a cluster of atoms of a grain
- smallest time unit = Monte Carlo Step (MCS) =  $N$  reorientation attempts

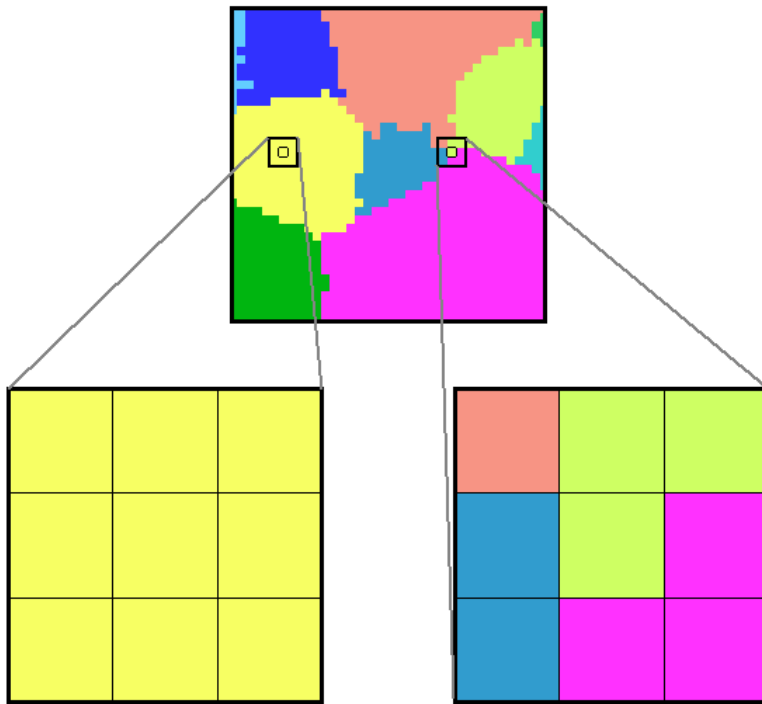


# Standard Monte Carlo Potts Model

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## Algorithm:

1. probabilistic selection of a lattice point
  - lattice point has orientation  $Q_\mu$



! grain boundary migration  
→ atoms change orientation  
from grain to neighboring grain

- lattice point on boundary → proceed with step 2
- else → terminate loop

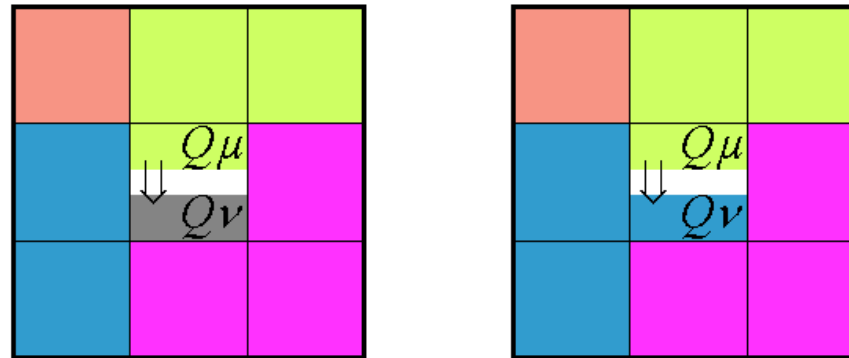
# Standard Monte Carlo Potts Model

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Algorithm:

2. selection of new orientation  $Q_\nu$

- $Q_\nu \neq Q_\mu$



- assigned on probation
- consider only orientations from neighboring lattice points → no nucleation events!

# Standard Monte Carlo Potts Model

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## Algorithm:

### 3. calculation of energy

- given by Hamiltonian

$$H = \gamma \cdot \sum_{i=1}^N \sum_{j=1}^{nn} (1 - \delta_{Q_i Q_j})$$

! only one lattice point is considered at each trial

- hence: difference in energy simplified to

$$\Delta E = \gamma \cdot \sum_{j=1}^{nn} (\delta_{Q_j Q_\nu} - \delta_{Q_j Q_\mu})$$

considering only nearest neighbors

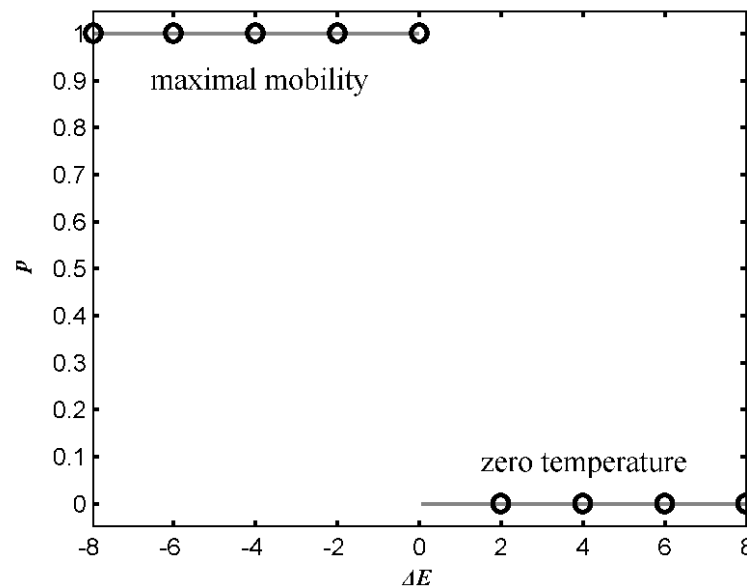
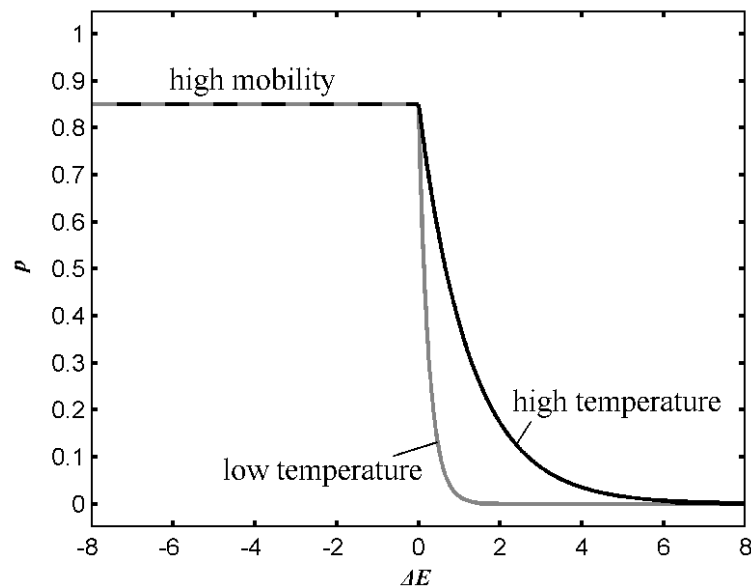
# Standard Monte Carlo Potts Model

## Algorithm:

### 4. selection of final orientation $Q\mu^*$

- chosen with probability

$$p = \begin{cases} m & \text{if } \Delta E \leq 0 \\ m \cdot \exp\left(\frac{-\Delta E}{k_B T}\right) & \text{if } \Delta E > 0 \end{cases}$$

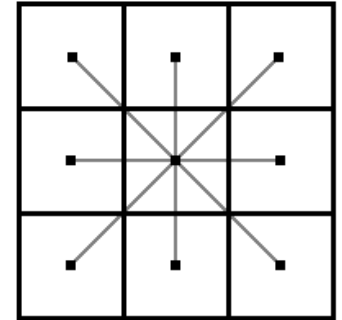


# Standard Monte Carlo Potts Model

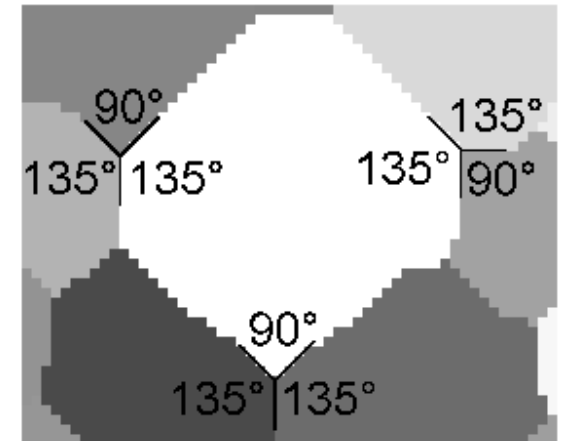
## Problems:

### 1. choice of lattice

- in 2D: quadratic or triangular lattice
- quadratic lattice = easy implementation (\*.bmp)  
→ underlying lattice forces simulation results?



- Potts model = curvature driven kinetics
- energy minimized by decreasing boundary length
- driving force places boundaries along lattice facets



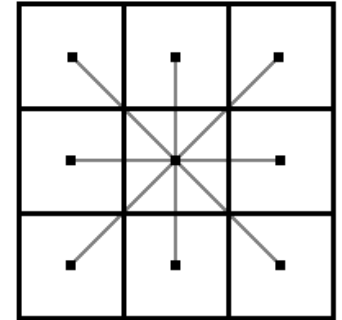
# Standard Monte Carlo Potts Model

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### ○ solution:

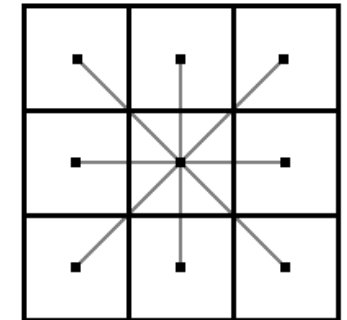
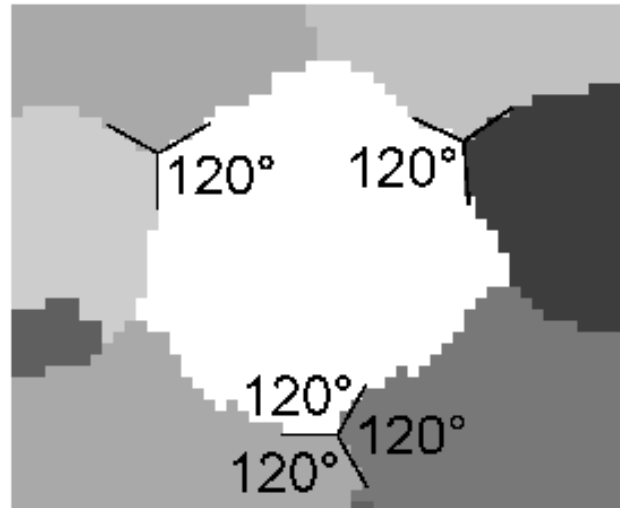
1. increasing number of neighbor lattice points decreases energetic anisotropy of the lattice
2. introduction of simulation temperature  $k_B T$  activates thermal fluctuations that roughens grain boundaries

# Standard Monte Carlo Potts Model

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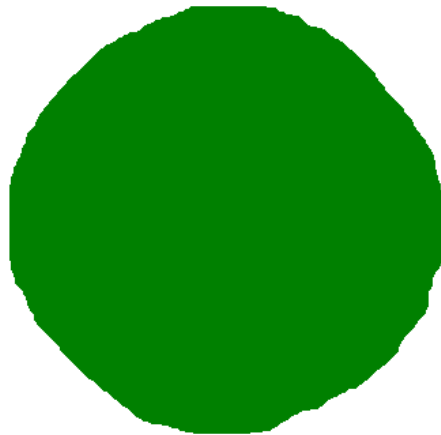
# Standard Monte Carlo Potts Model

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Problems:

## 2. choice of simulation temperature

- simulation temperature not related to real temperature
- usually found by trial and error



$$k_B T = 0.0$$

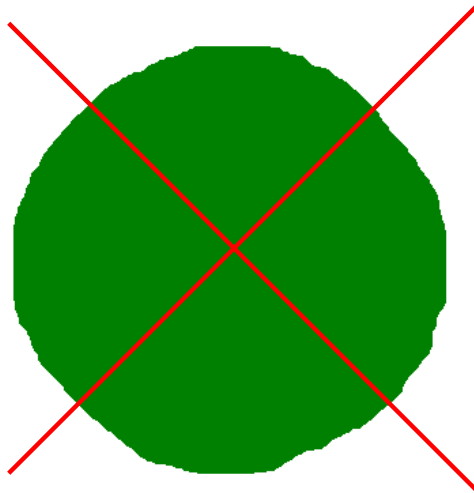
# Standard Monte Carlo Potts Model

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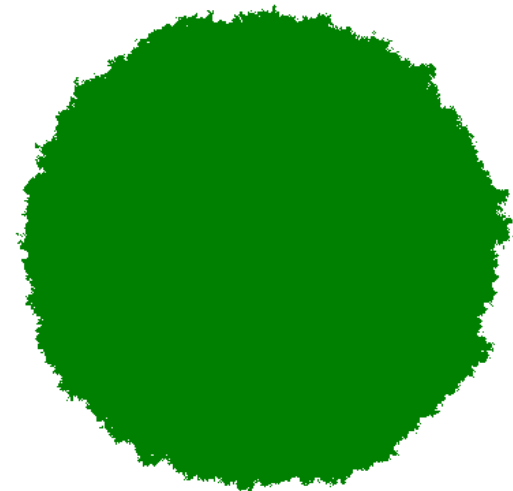
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$$k_B T = 0.0$$



$$2.4$$

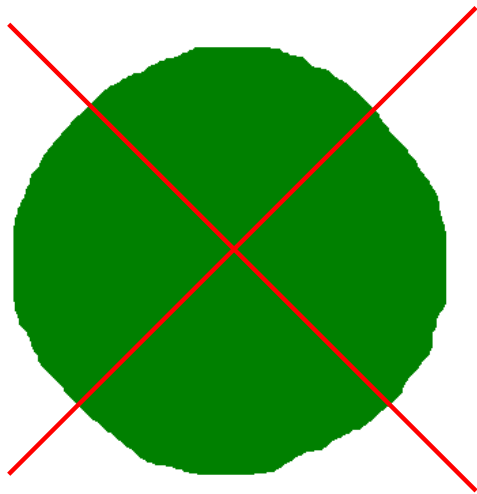
# Standard Monte Carlo Potts Model

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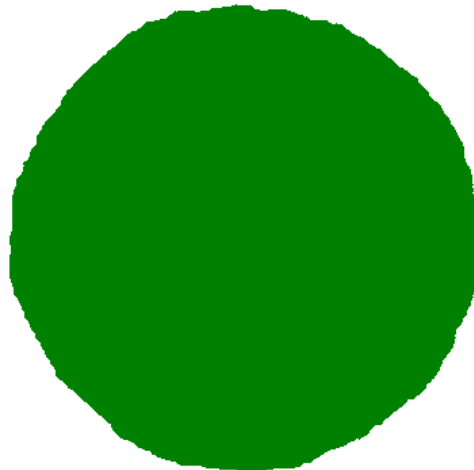
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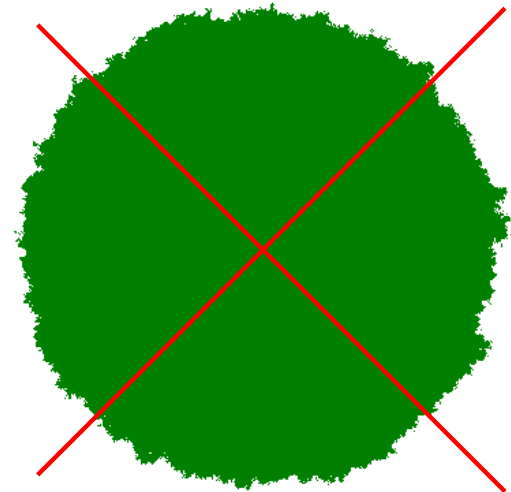
- simulation temperature not related to real temperature
- usually found by trial and error



$k_B T =$  0.0



0.6



2.4

# Standard Monte Carlo Potts Model

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## Problems:

### 3. simulation runtime

- on standard PC (Intel® Core™2 Quad Q9650 CPU @ 3GHz with 3GB RAM)
  - three dimensional grain microstructure
  - $250^3$  lattice points in size
  - 10,000 simulation time steps
- needs nearly 94 hours!

# Standard Monte Carlo Potts Model

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## Problems:

### 3. simulation runtime

- solution:
- in recent years NVIDIA developed parallel computing architecture: Compute Unified Device Architecture
- CUDA allows programming of graphics processing units (GPUs) through standard programming languages (C)
  - GPUs have a parallel architecture
  - GPUs can compute many coexisting threads at the same time, rather than computing only a single thread

[S. Schäfer, D. Zöllner: CD-ROM Proceedings of the 6th Europ. Congr. on Comput. Meth. in Appl. Sci. and Eng. 2012; ISBN: 978-3-9502481-9-7]

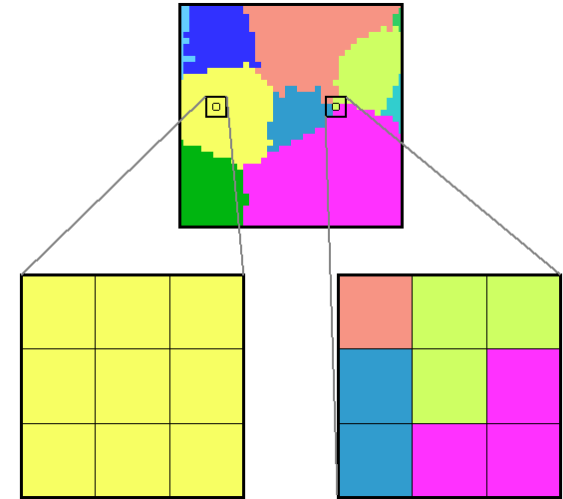
# Standard Monte Carlo Potts Model

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## Problems:

### 3. simulation runtime

- solution:



- drawback = lattice points not completely independent of each other!
- read from and write to the same lattice point?
  1. use two integer arrays
  2. introduce a checkerboard algorithm

# Standard Monte Carlo Potts Model

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## Problems:

### 3. simulation runtime

- solution:
- on a server with
  - two Intel® Xeon™ X5650 CPUs with 2.67GHz (12 cores) and
  - a Tesla™ C2050 graphic card from NVIDIA® with 448 cores) and 24GB RAM
- reduction of runtime down to approx. 18 minutes!  
<< 94 hours

# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

## 1. curvature driven kinetics

- velocity of grain boundary

$$v_n = -k \cdot K$$

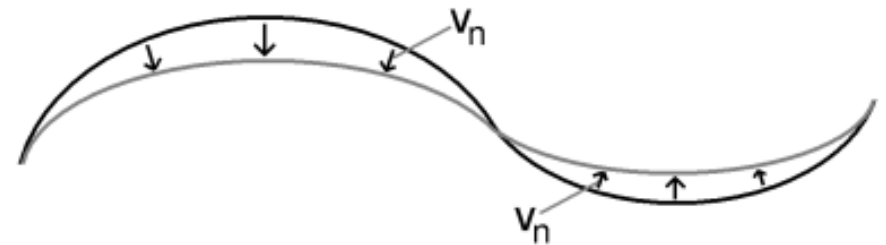
$K$  = curvature;  $k$  = kinetic constant

- temporal change of area

$$\frac{dA}{dt} = \oint_U v_n ds$$

- for any closed plane curve (it can be shown)

$$\oint_U K \cdot ds = 2\pi$$



# Standard Monte Carlo Potts Model

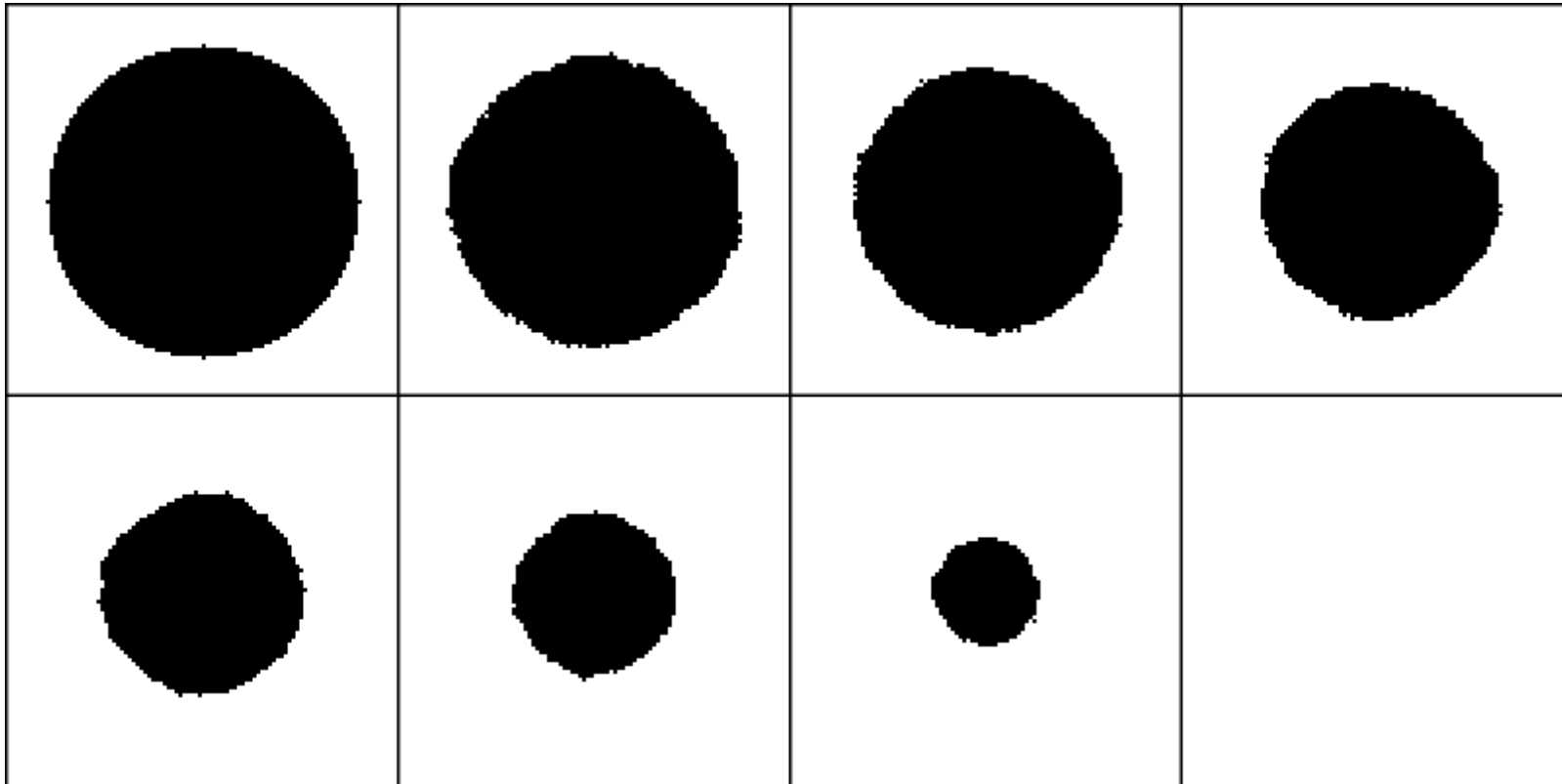
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Simulation of grain boundary motion:

1. curvature driven kinetics – simple case

○ temporal change of area

$$A = -2\pi \cdot k \cdot t + A_0$$



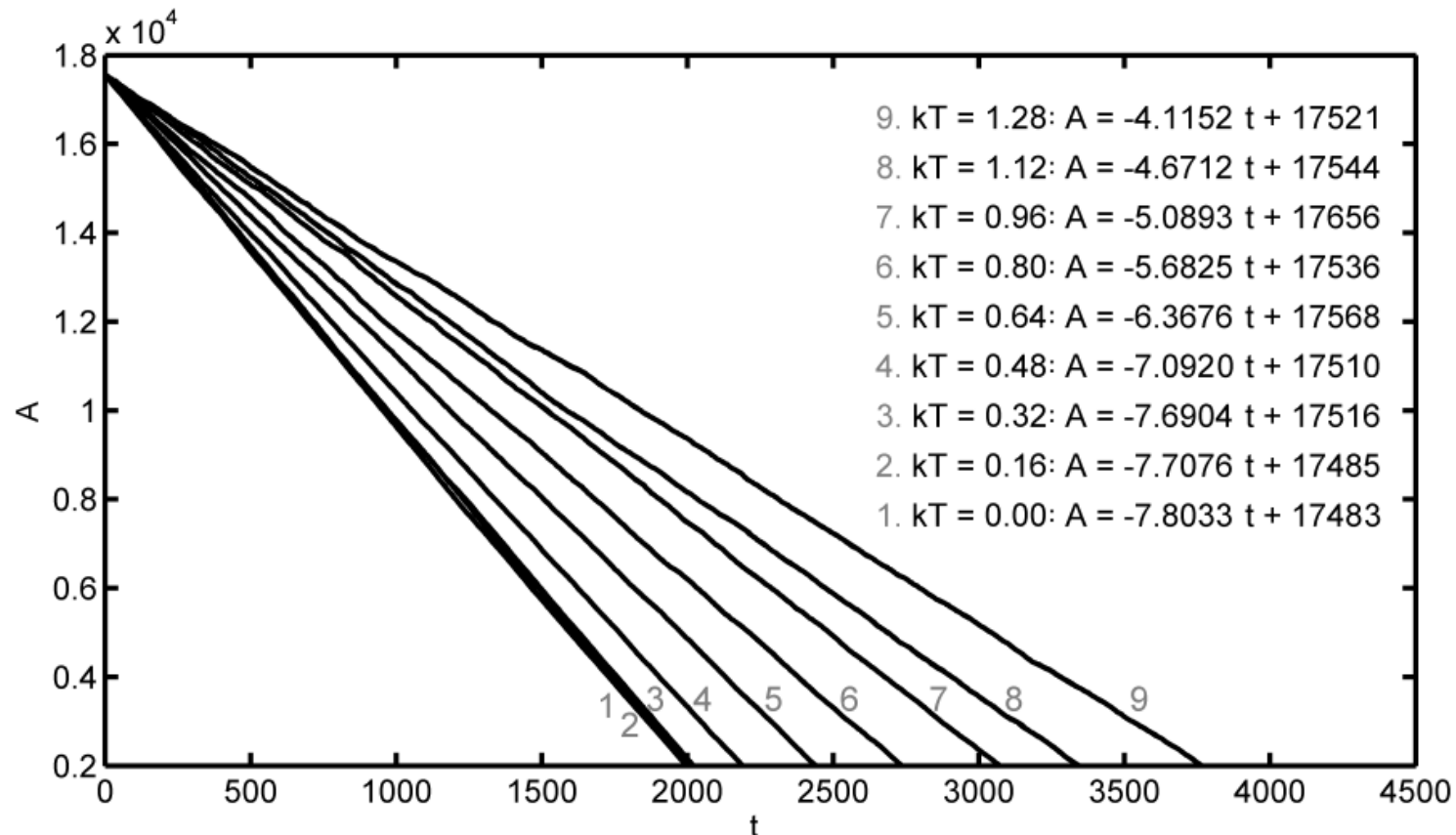
# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

1. curvature driven kinetics – simple case

○ temporal change of area

$$A = -2\pi \cdot k \cdot t + A_0$$



# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

2. ensemble of moving grain boundaries

**remember?**

- parabolic growth law

$$\langle R \rangle^{1/n} = b \cdot t + \langle R \rangle_0^{1/n}$$

- self-similar scaling

$$F(R, t) = g(t) \cdot f(x), \quad g(t) = \frac{N}{\langle R \rangle}, \quad x = \frac{R}{\langle R \rangle}$$

# Standard Monte Carlo Potts Model

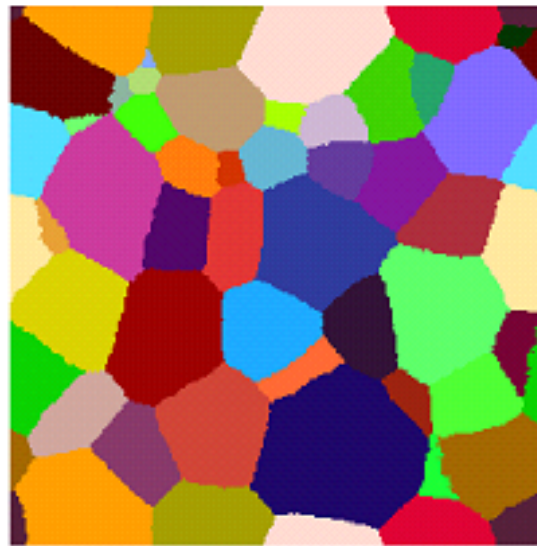
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Simulation of grain boundary motion:

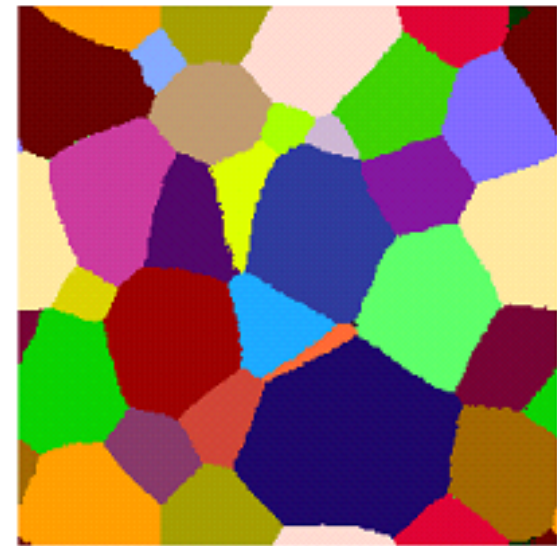
2. ensemble of moving grain boundaries



t = 1000



t = 2000

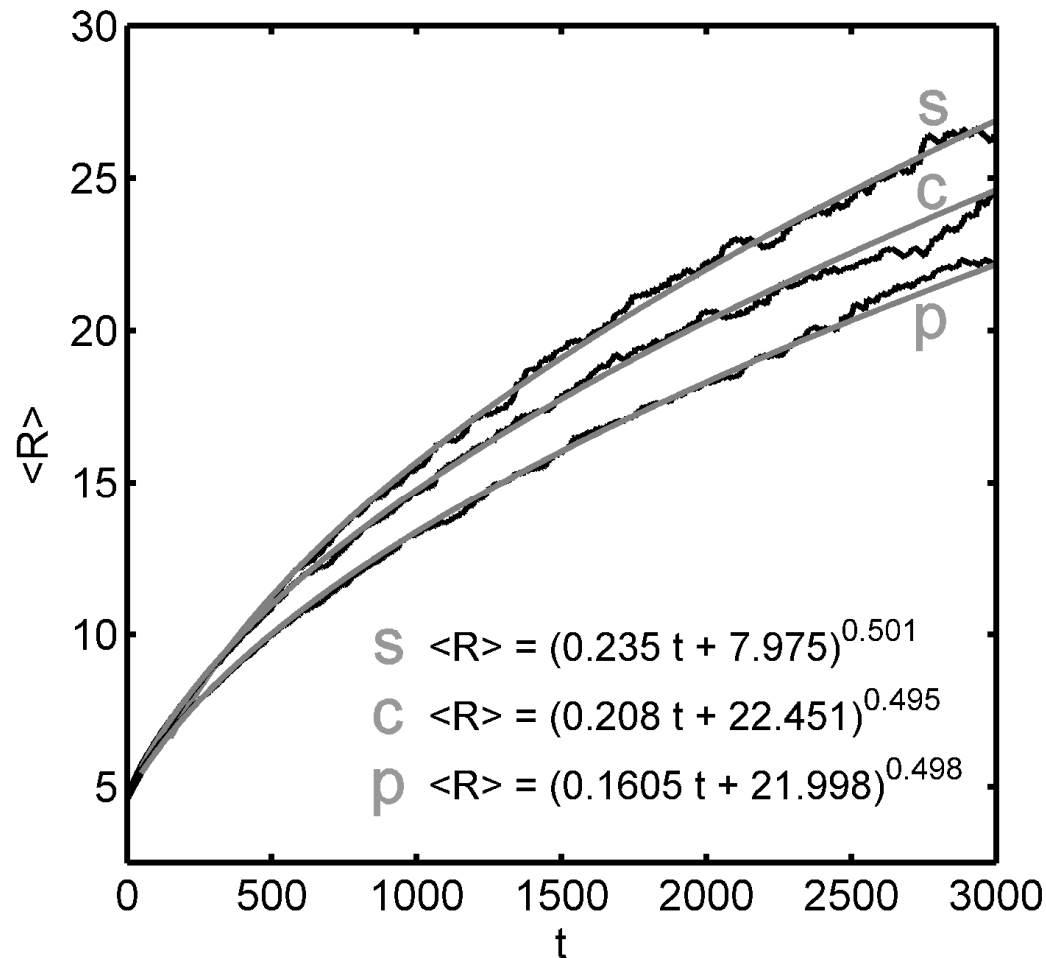


t = 3000

# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

## 2. ensemble of moving grain boundaries



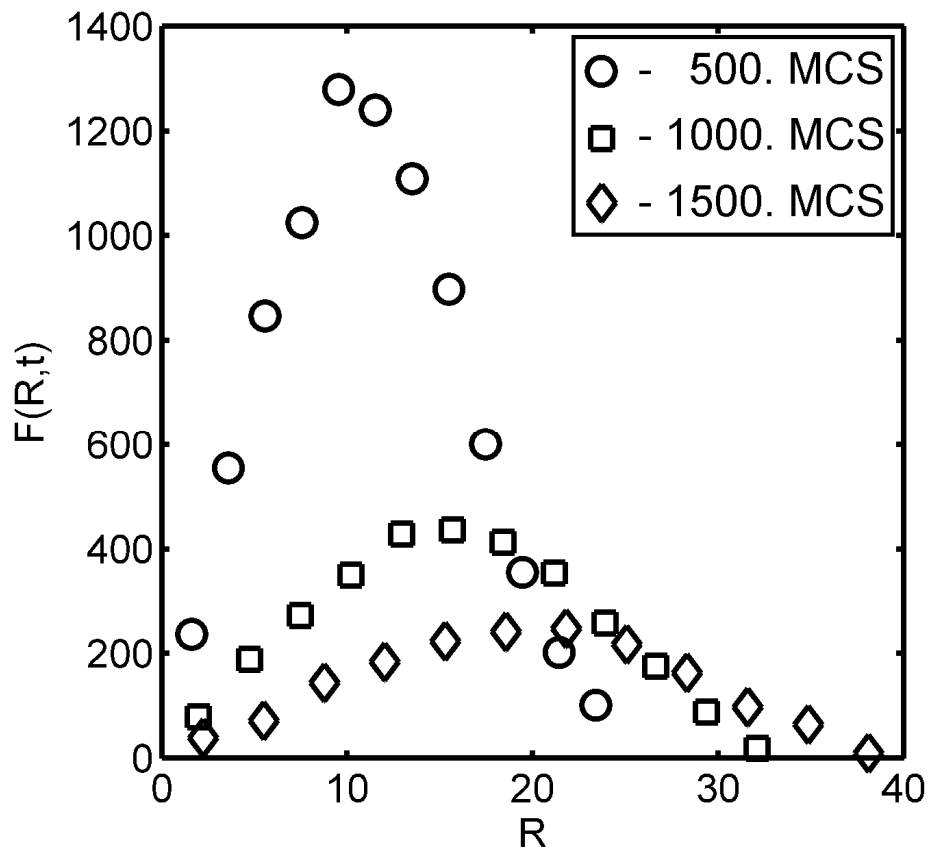
- parabolic growth law
$$\langle R \rangle^{1/n} = b \cdot t + \langle R \rangle_0^{1/n}$$
- exponent  $n \approx 1/2$
- factor  $b$  depends on simulation type
- s = serial
- p = parallel
- c = checkerboard

# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

## 2. ensemble of moving grain boundaries

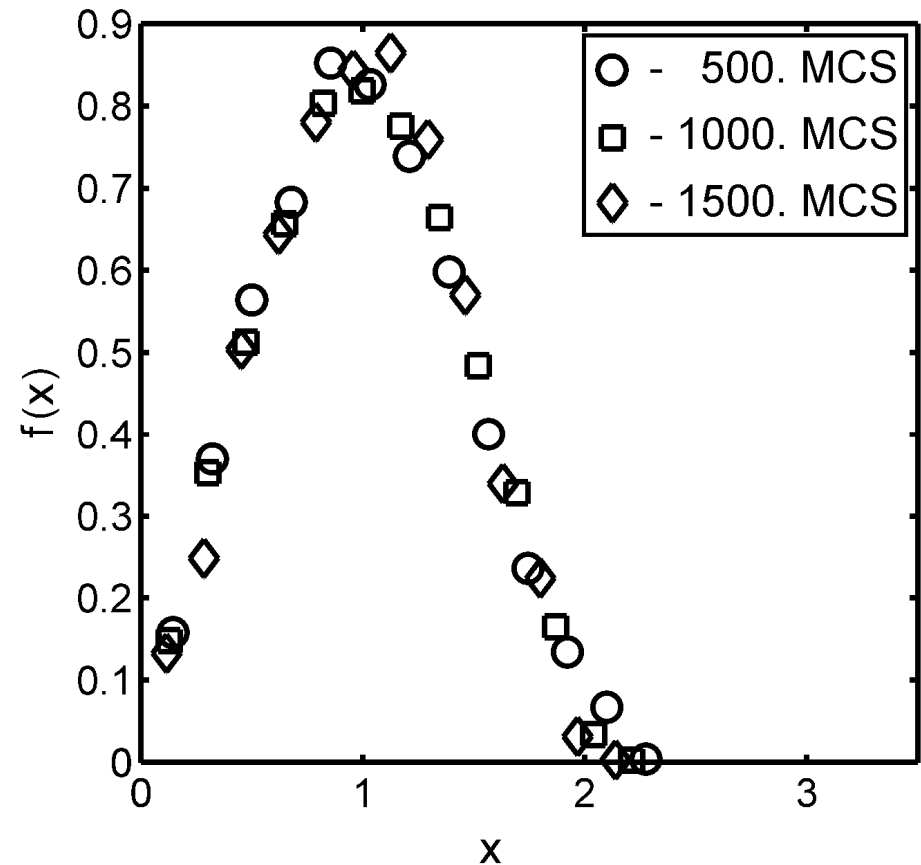
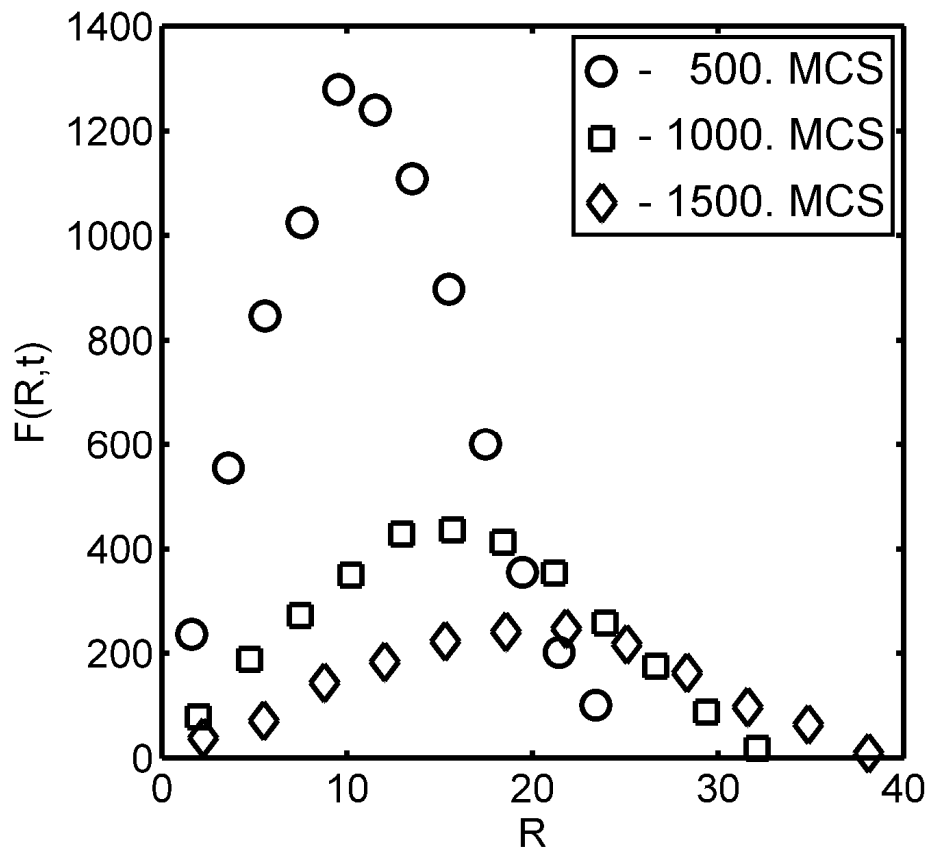


$$F(R,t) = g(t) \cdot f(x)$$

# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

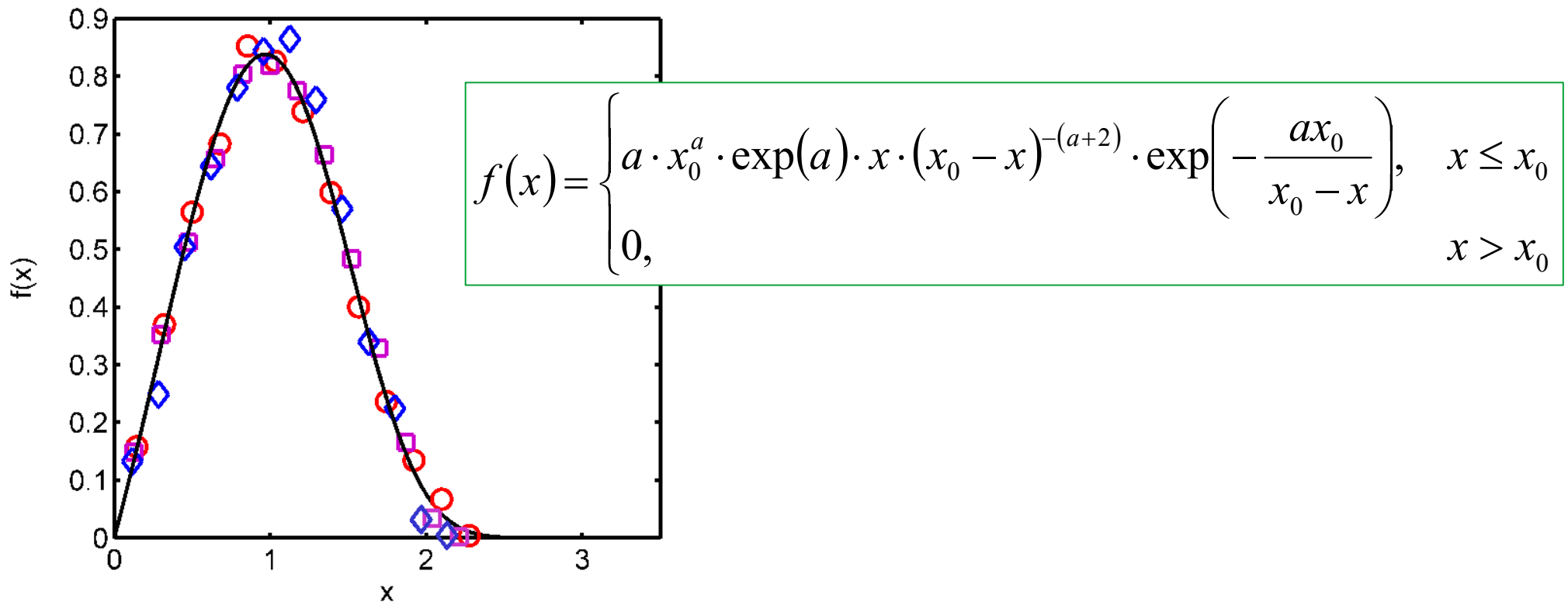
## 2. ensemble of moving grain boundaries



# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

## 2. ensemble of moving grain boundaries



[D. Zöllner, P. Streitenberger: Prac. Metallogr. 47 (2010) 618]

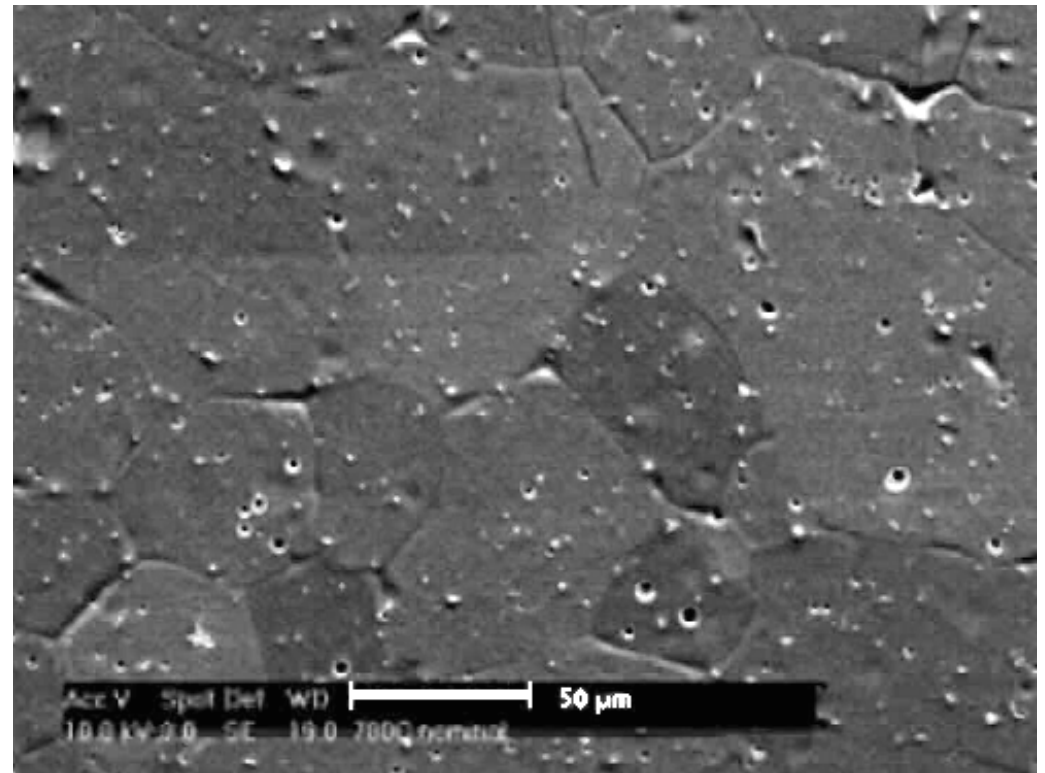
# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

3. individual growth kinetics

Aluminium at approx. 650°C  
from in-situ real-time SEM

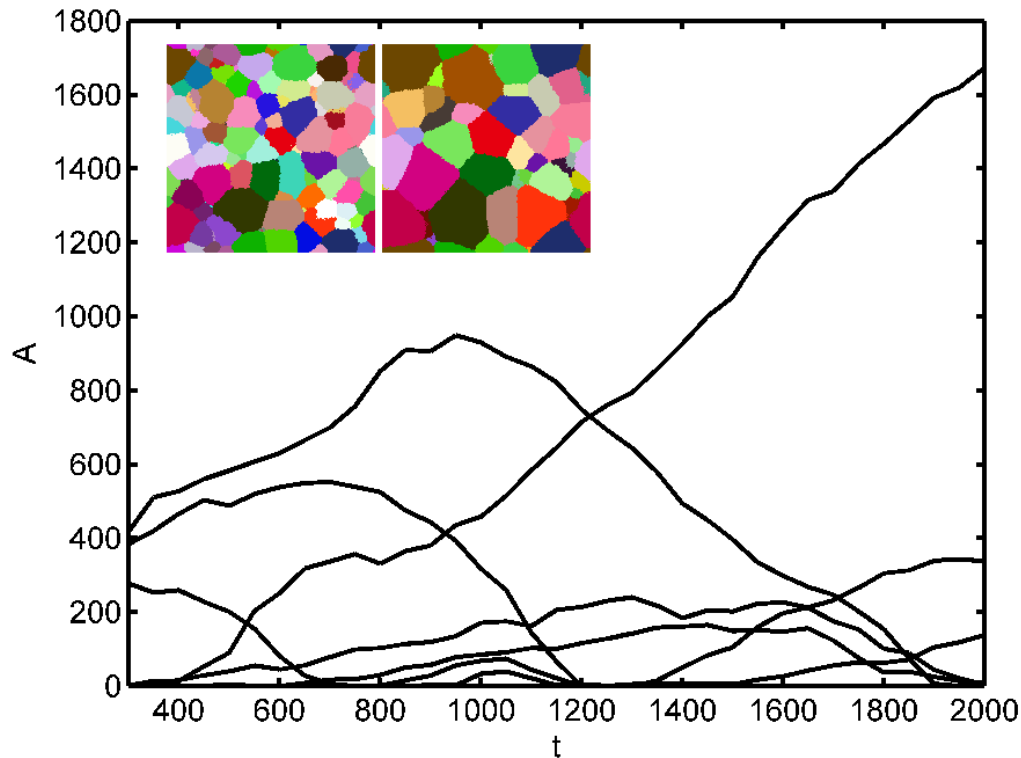


[D. Zöllner, P. Streitenberger, I. Fielden: *Prac. Metallogr.* 49 (2012) 428]

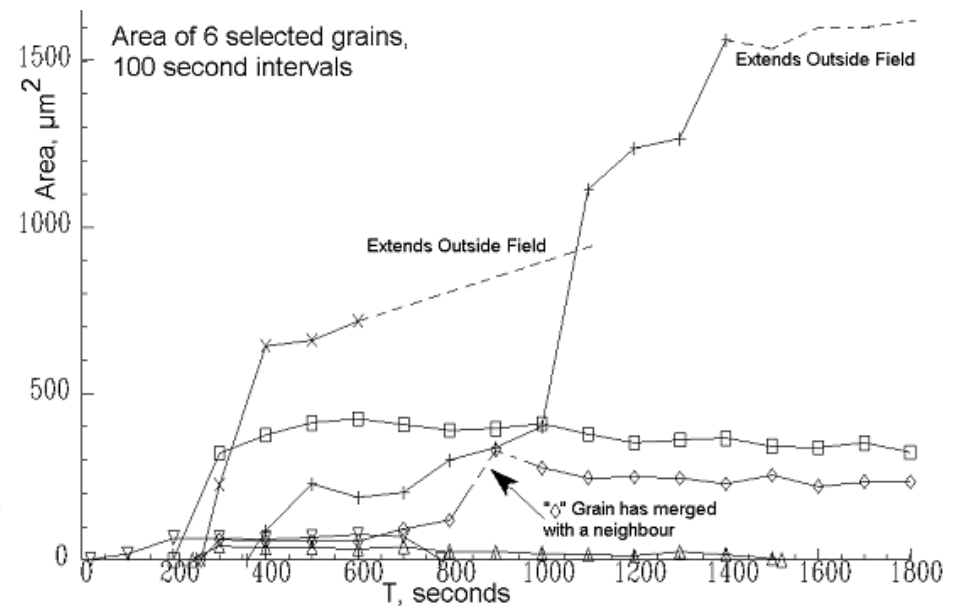
# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

## 3. individual growth kinetics



simulation (left) and real-time, in-situ image sequences (right) of evolving microstructures



# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

3. individual growth kinetics

$$\dot{R}(R, t) = \frac{dR}{dt} = \langle \dot{R} \rangle \cdot G(x)$$

# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

## 3. individual growth kinetics

$$\dot{R}(R, t) = \frac{dR}{dt} = \langle \dot{R} \rangle \cdot G(x)$$

$$R(z, x_i) = \langle R \rangle \cdot z \cdot \sqrt{\tau(z, x_i) + 1}$$

$$t(z, x_i) = \frac{\langle R \rangle^2}{2\gamma} \tau(z, x_i)$$

$$\tau(z, x_i) = \left\{ (x_0 - x_i)^{2a/D} \exp\left(\frac{2ax_0}{D(x_0 - x_i)}\right) \right\} \cdot \left\{ (x_0 - z)^{-2a/D} \exp\left(-\frac{2ax_0}{D(x_0 - z)}\right) \right\}^{-1}$$

# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

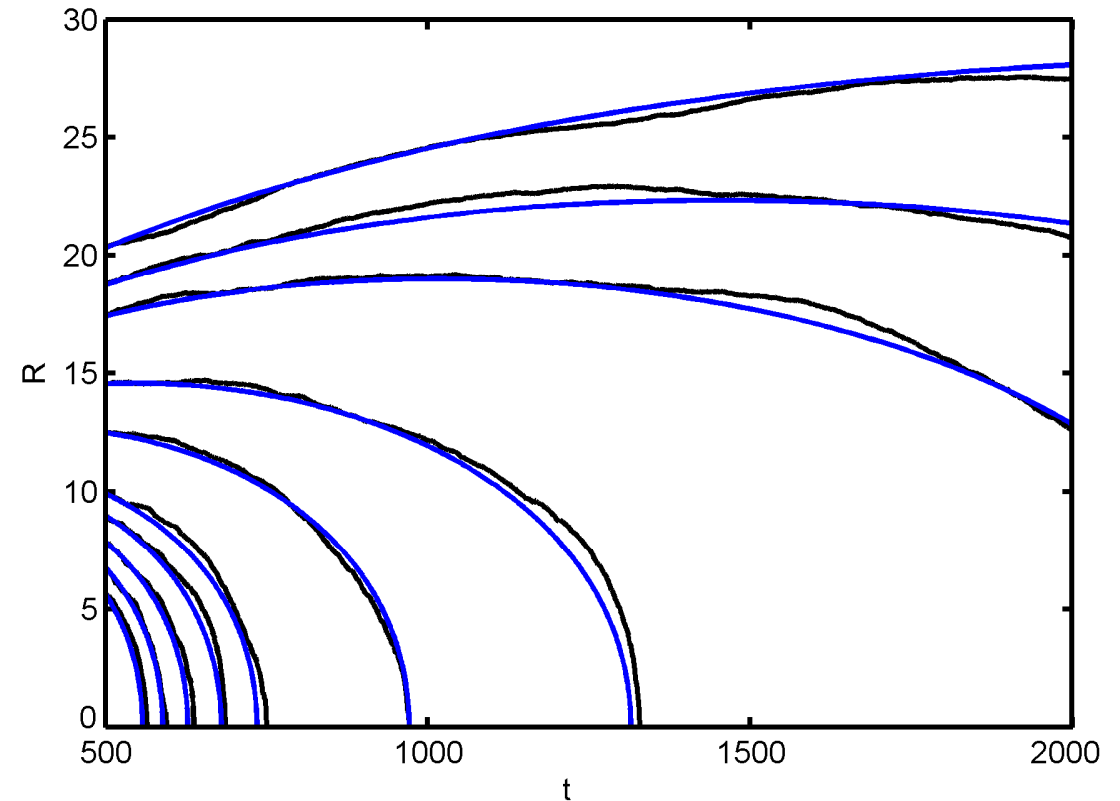
## 3. individual growth kinetics

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# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

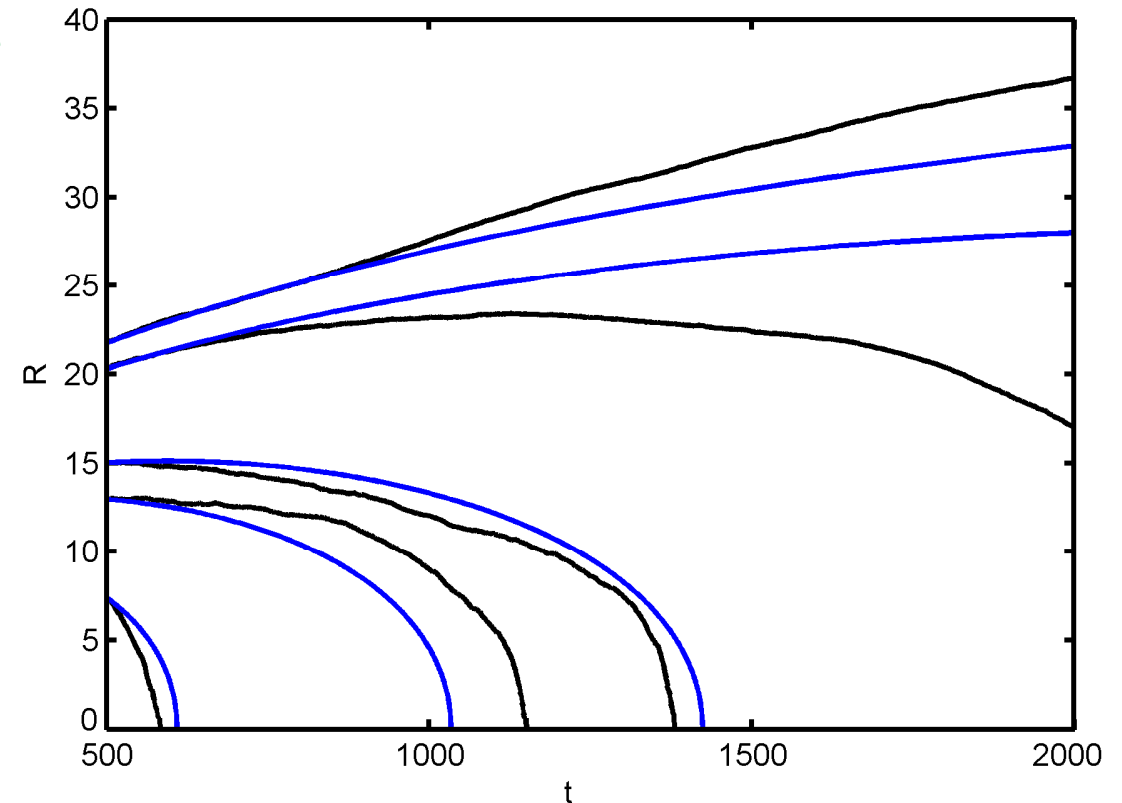
## 3. individual growth kinetics

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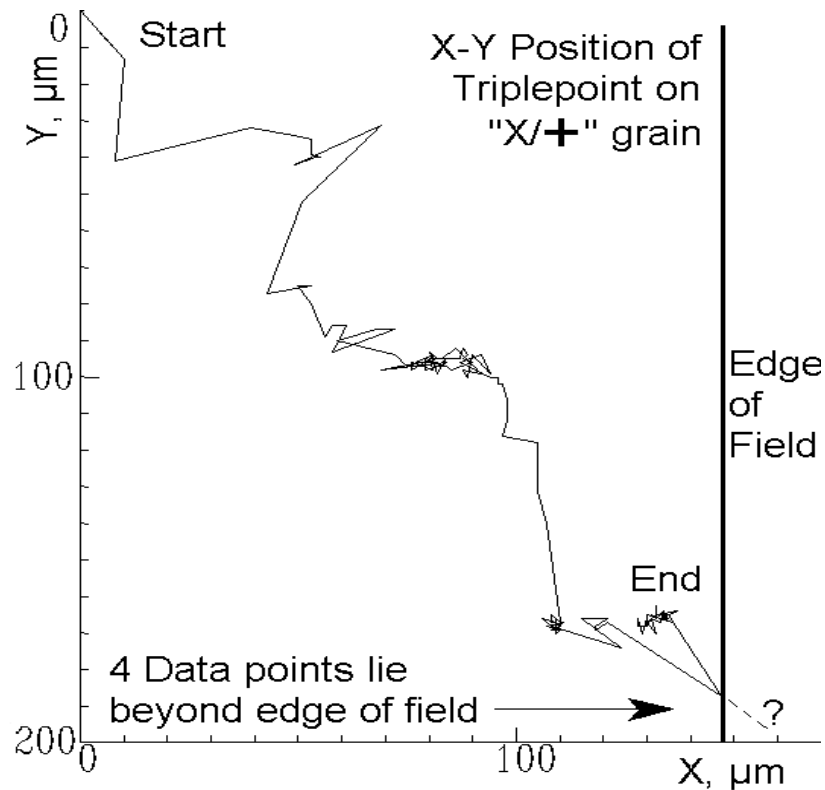
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# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

## 3. individual growth kinetics

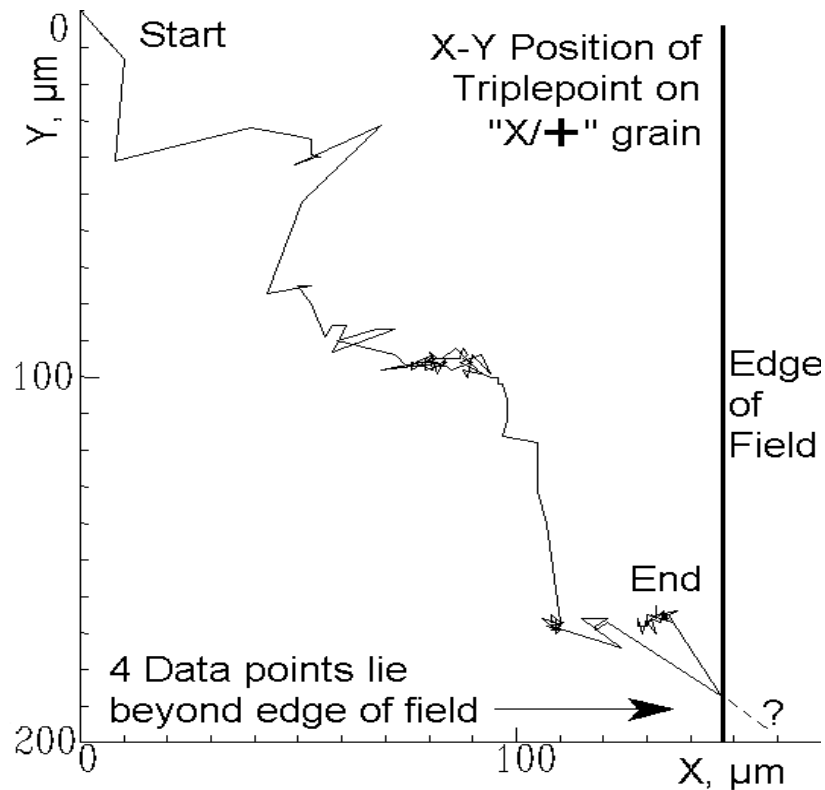


Track of triple point  
position at 10s intervals

# Standard Monte Carlo Potts Model

Simulation of grain boundary motion:

## 3. individual growth kinetics



Track of triple point  
position at 10s intervals

$$\tilde{D} = \frac{msd}{6 \cdot t^*} = \frac{1}{3} \sum_{i=1}^3 \frac{\langle (x_i - x_i(t^* = 0))^2 \rangle}{2 \cdot t^*}$$

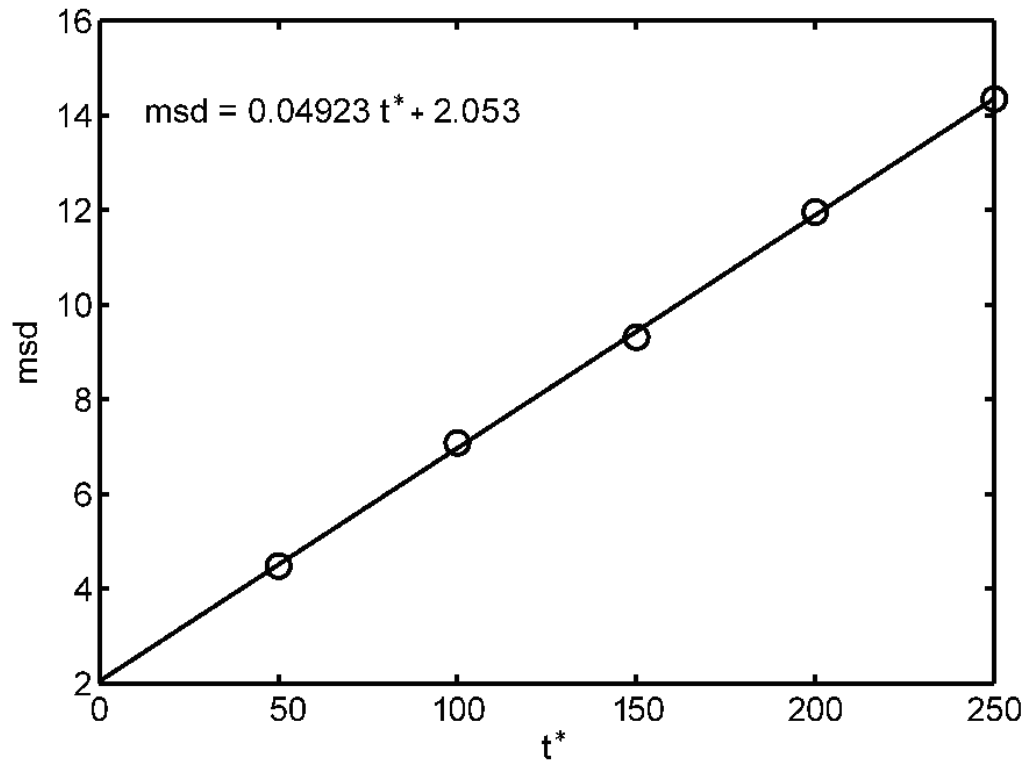
$$\langle R \rangle^{1/n} = \frac{6\tilde{D}}{sd^2} \cdot t + \langle R \rangle_0^{1/n}$$

# Standard Monte Carlo Potts Model

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Simulation of grain boundary motion:

## 3. individual growth kinetics



- $\tilde{D} = 0.0082 \text{ MCU}^2/\text{MCS}$
- $sd = 0.429$ 
  - it follows:  $b = 0.2672 \text{ MCU}^2/\text{MCS}$
- average growth law:  $b = 0.2255 \text{ MCU}^2/\text{MCS}$

# A Potts Model for Nanocrystalline Grain Growth

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# A Potts Model for Nanocrystalline Grain Growth

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What is different?

*experiments on nanocrystalline materials*

- high values of hardness and yield/fracture strength
- linear or exponential growth kinetics in contradiction to parabolic normal grain growth

*experiment, theory and molecular dynamics simulation*

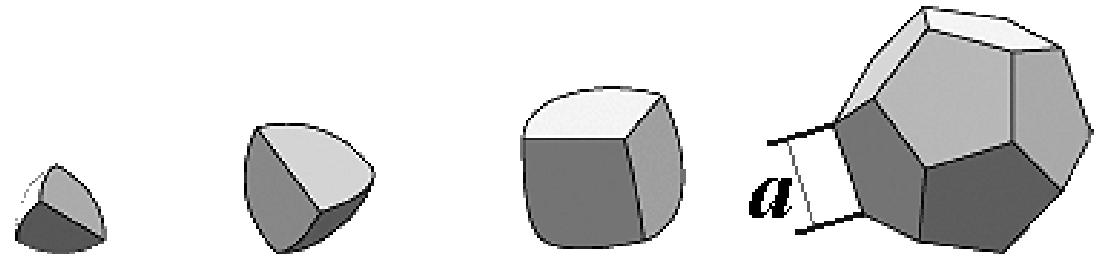
- triple lines and quadruple points have **finite mobilities** different from adjoining boundaries
- failure of standard Potts model

# A Potts Model for Nanocrystalline Grain Growth

theoretical predictions in nanocrystalline size range:

- principle expression

$$v = \frac{m_{gb}\gamma_{gb}K}{1 + \frac{m_{gb}}{am_{tl}} + \frac{m_{gb}}{a^2m_{qp}}}$$



- from thermodynamic approach

→ 1. normal grain growth

2. triple line control

3. quadruple point control

$$\langle R \rangle \sim \left\{ t^{1/2} \quad t \quad e^{t/\tau} \right\}$$

[G. Gottstein, L. Shvindlerman, Scripta Mater. 54 (2006) 1065]

[P. Streitenberger, D. Zöllner, Acta Mater. 59 (2011) 4235]

# A Potts Model for Nanocrystalline Grain Growth

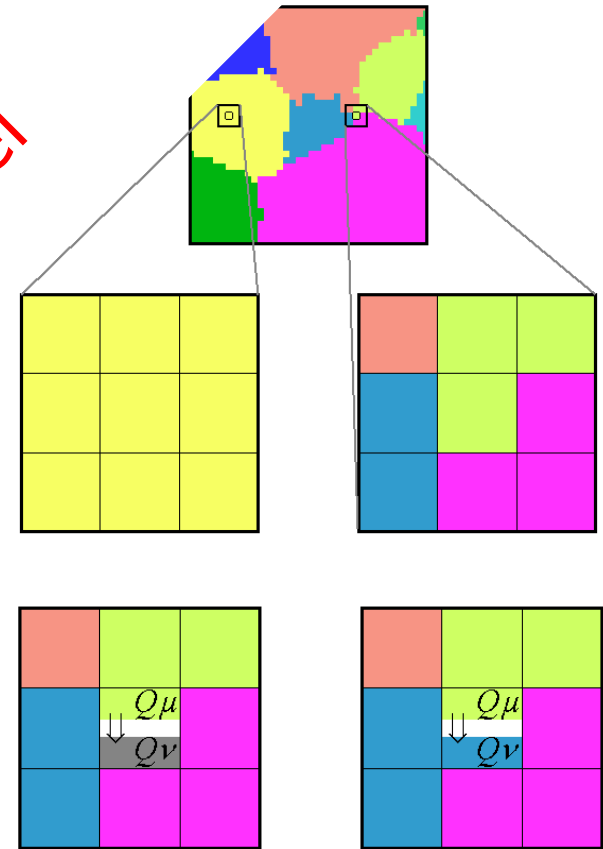
## One Monte Carlo Step (MCS):

1. probabilistic selection of one MCU
2. selecting a new orientation  $Q_v$
3.  $\Delta E = \gamma \cdot \sum_{j=1}^{nn} (\delta_{Q_j Q_v} - \delta_{Q_j Q_\mu})$
4. selecting orientation with

$$p = \begin{cases} m & \text{if } \Delta F \\ m \cdot \exp\left(\frac{-\Delta E}{k_B T}\right) & \text{if } \Delta E \end{cases}$$

repeat  $N$ -times

standard Potts model



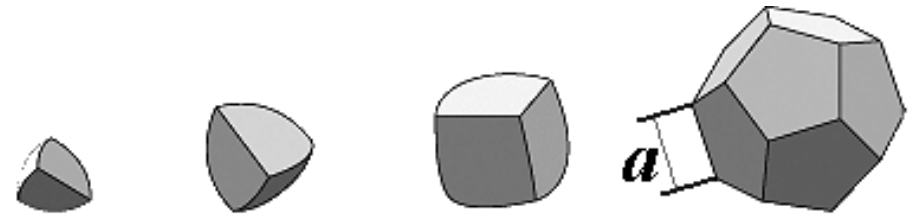
# A Potts Model for Nanocrystalline Grain Growth

## One Monte Carlo Step (MCS):

1. probabilistic selection of one MCU
2. selecting a new orientation  $Q_v$
3.  $\Delta E = \gamma \cdot \sum_{j=1}^{nn} (\delta_{Q_j Q_v} - \delta_{Q_j Q_\mu})$
4. selecting orientation with

$$p = \begin{cases} m & \text{if } \Delta E \leq 0 \\ m \cdot \exp\left(\frac{-\Delta E}{k_B T}\right) & \text{if } \Delta E > 0 \end{cases}$$

$$v = \frac{m_{gb} \gamma_{gb} K}{1 + \frac{m_{gb}}{a m_{tl}} + \frac{m_{gb}}{a^2 m_{qp}}}$$

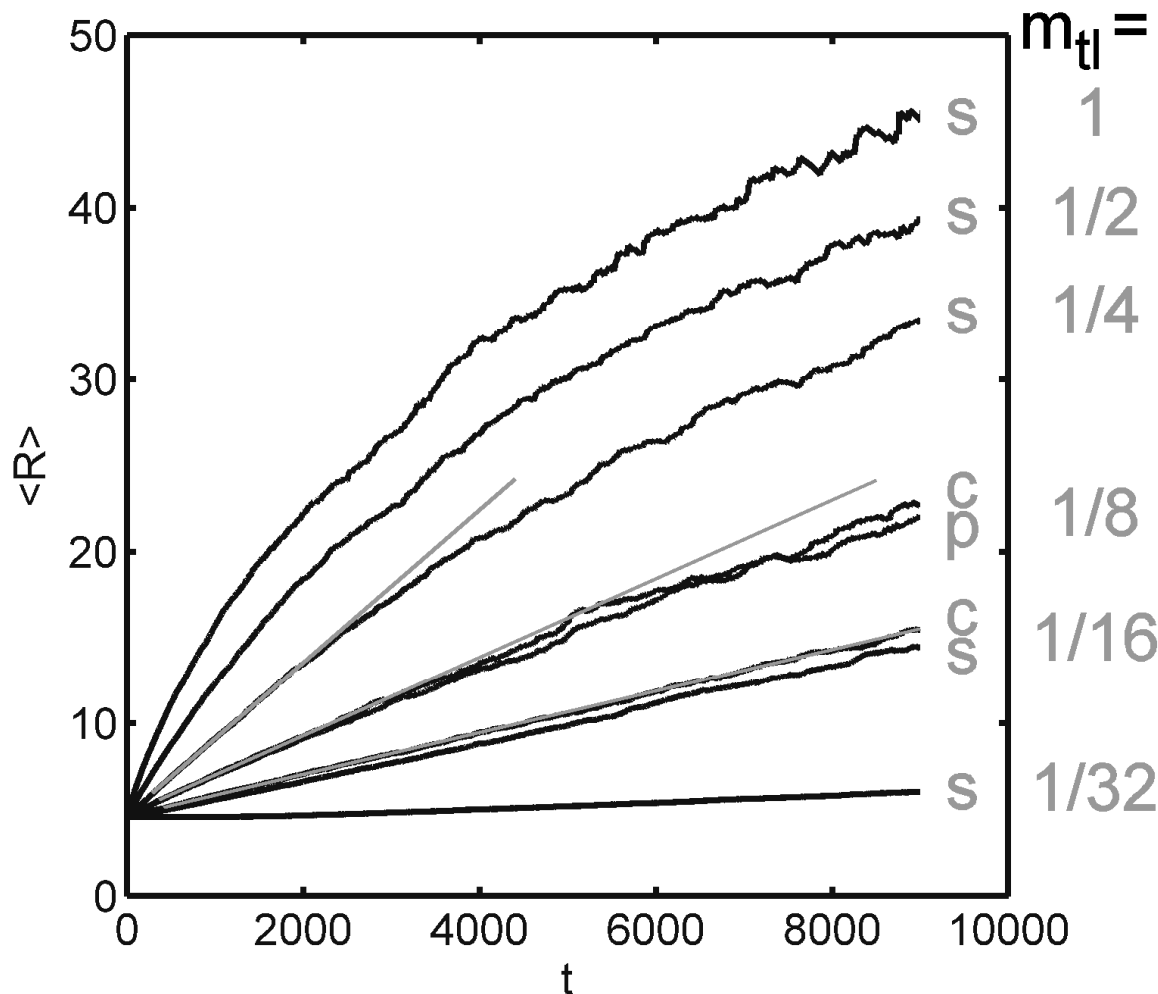


include separate mobilities for all grain features

repeat  $N$ -times

# A Potts Model for Nanocrystalline Grain Growth

triple junction control:



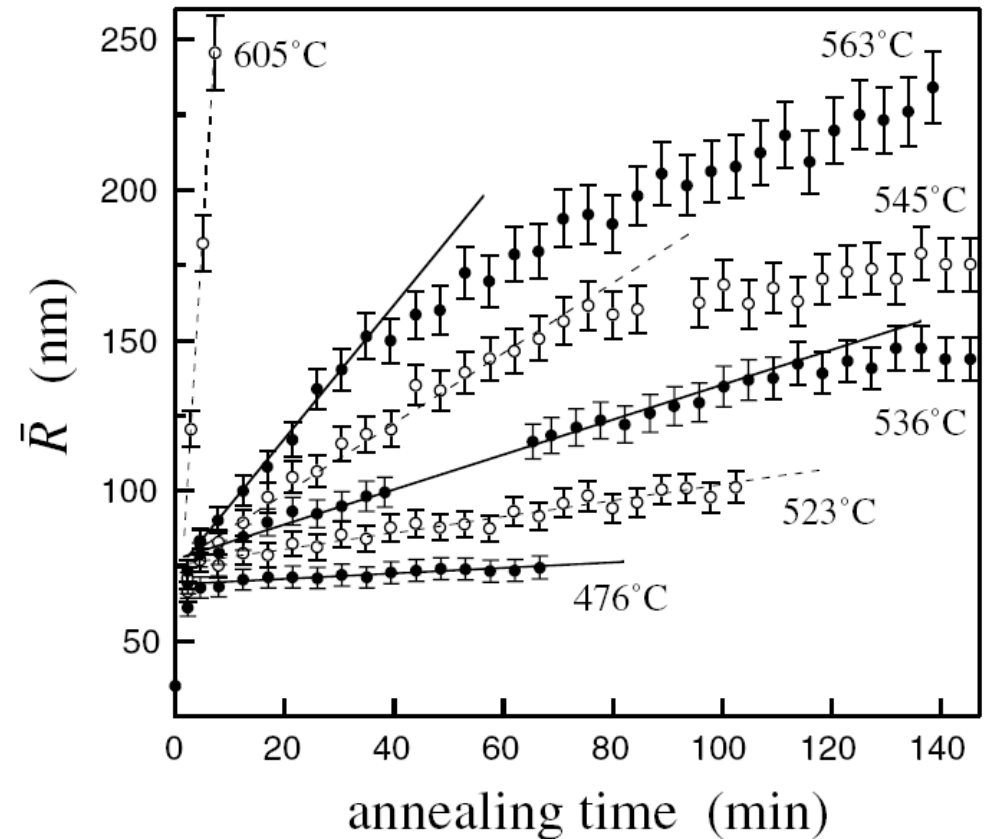
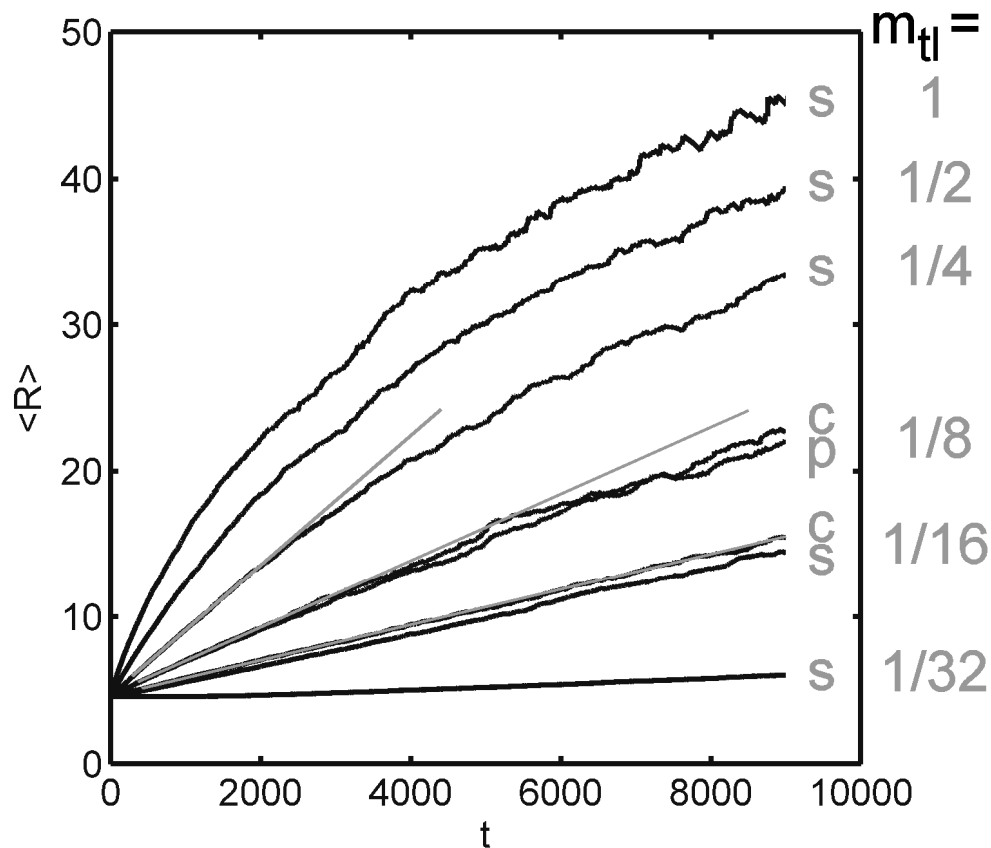
$$m_{gb} = 1 \text{ and } m_{tl} = m_{qp} = \{1, 1/2, 1/4, 1/8, 1/16\}$$

$$\langle R \rangle \sim \left( t^{1/2} \quad t \quad e^{t/\tau} \right)$$

- kinetics changes from parabolic to linear
- kinetics does not depend on the simulation type

# A Potts Model for Nanocrystalline Grain Growth

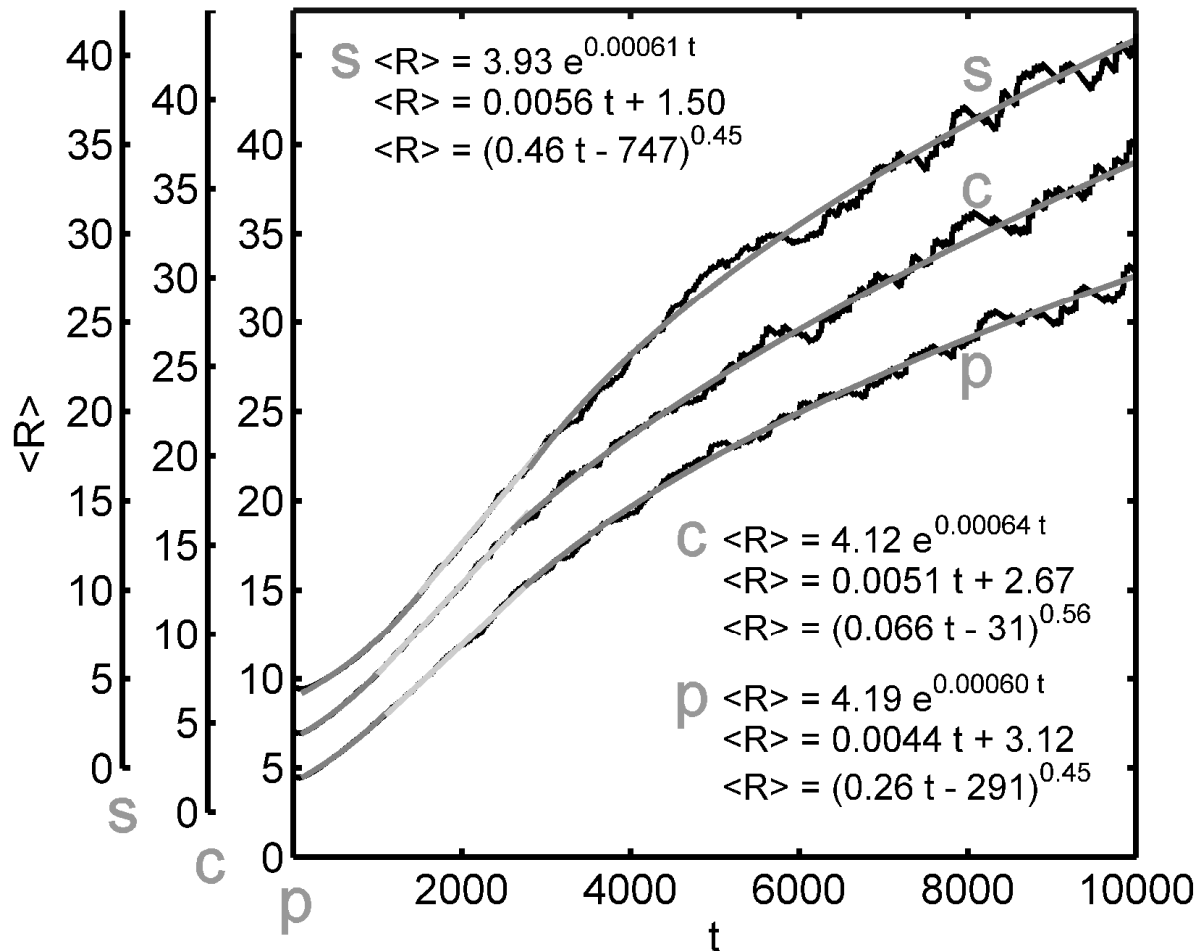
triple junction control:



[C.E. Krill III, L. Helfen, D. Michels, H. Natter, A. Fitch, O. Masson, R. Birringer:  
Phys. Rev. Lett. 86 (2001) 842]

# A Potts Model for Nanocrystalline Grain Growth

quadruple junction control:



$$m_{gb} = m_{tl} = 1 \text{ and}$$

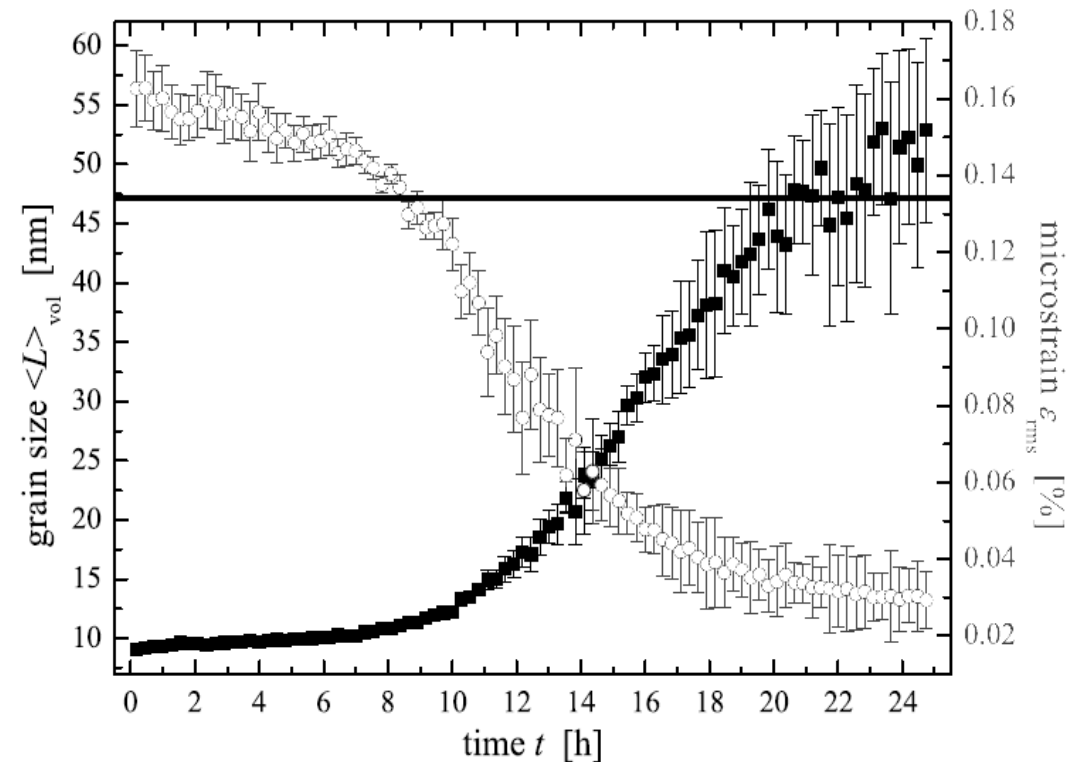
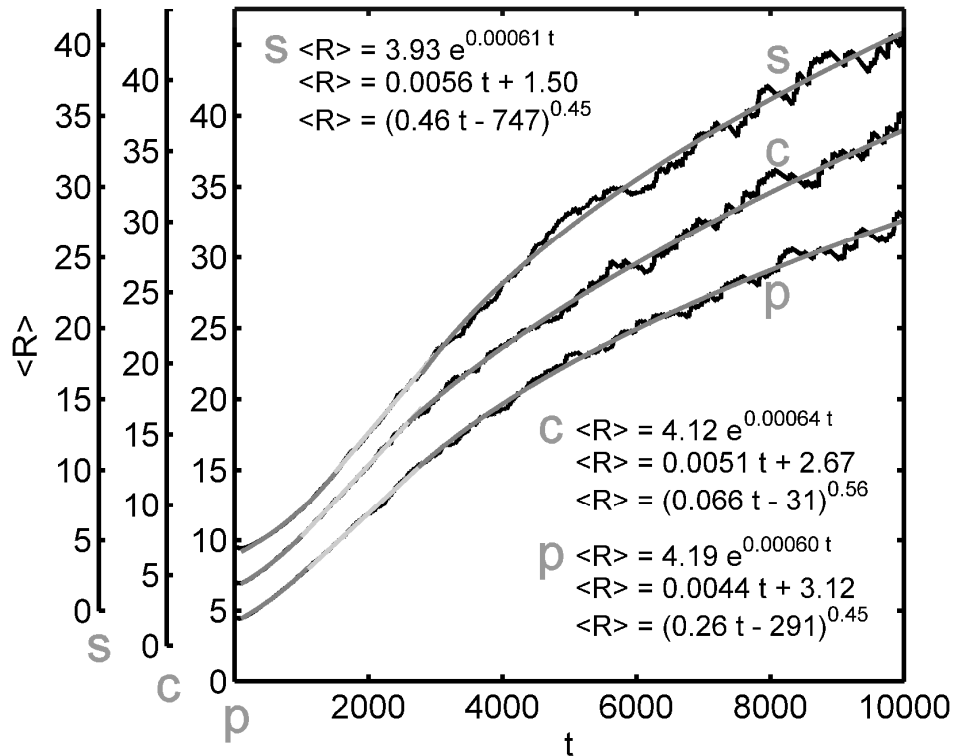
$$m_{qp} < 1$$

$$\langle R \rangle \sim \left( t^{1/2} \quad t \quad e^{t/\tau} \right)$$

- exponential growth behavior for early annealing times
- long-time annealing → linear & parabolic

# A Potts Model for Nanocrystalline Grain Growth

quadruple junction control:



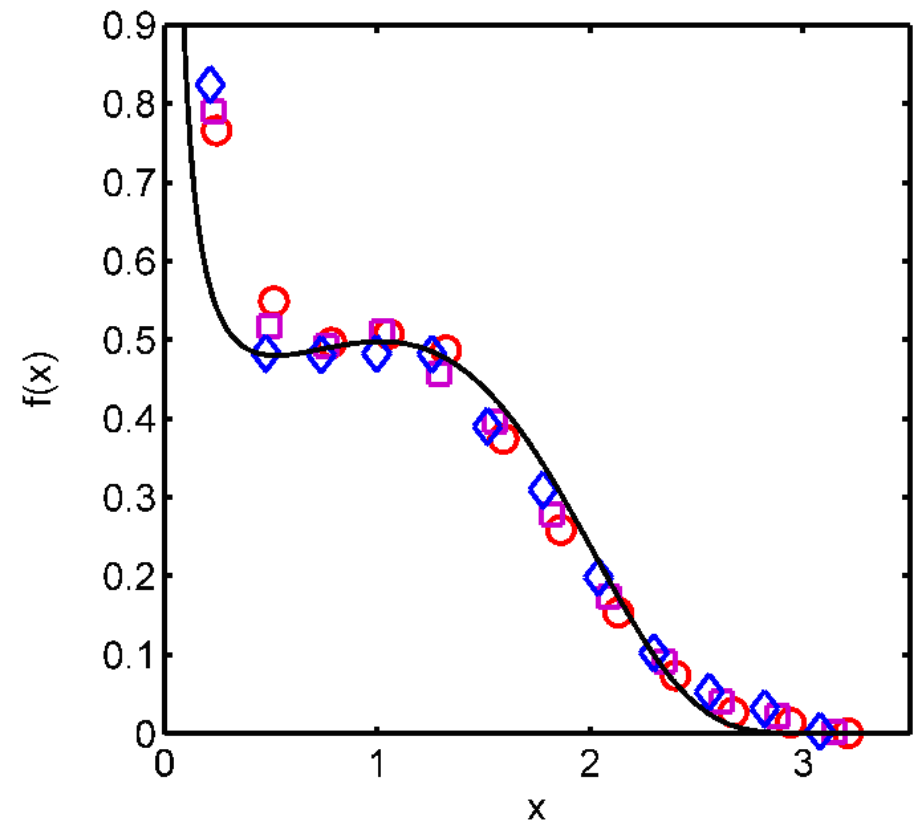
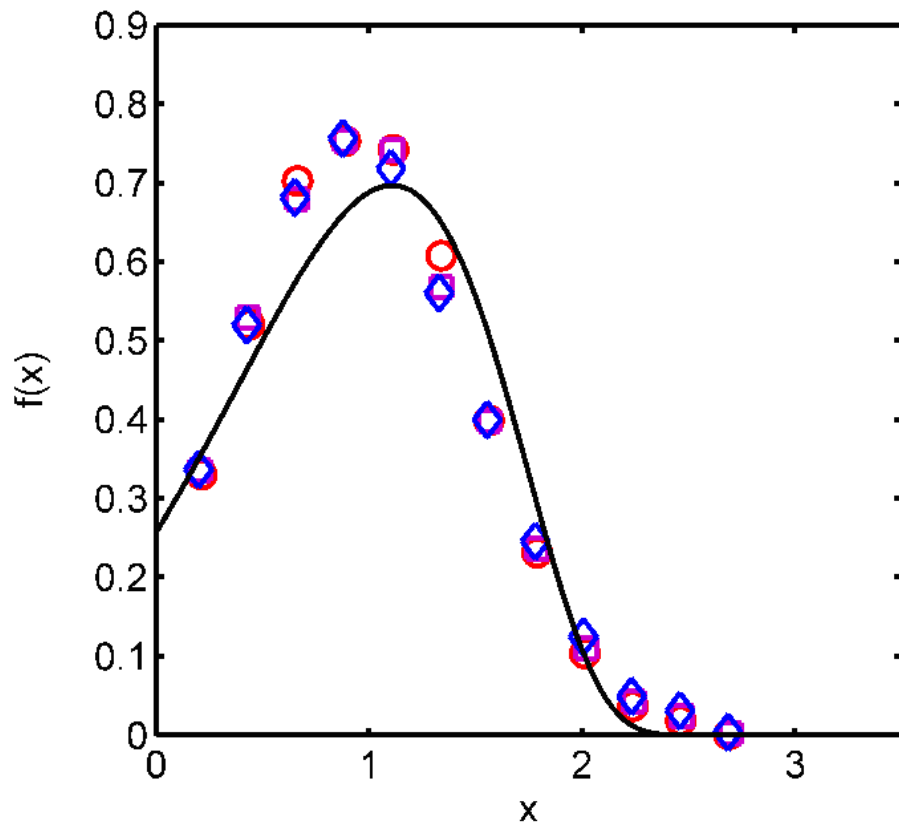
[M. Ames, J. Markmann, R. Karos, A. Michels, A. Tschöpe, R. Birringer: Acta Mater. 56 (2008), p. 4255]

# A Potts Model for Nanocrystalline Grain Growth

triple line control

vs.

quadruple point control:



[P. Streitenberger, D. Zöllner, Acta Mater. 59 (2011) 4235]

# Going back to 2D Grain Growth

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# Going back to 2D Grain Growth

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## What for?

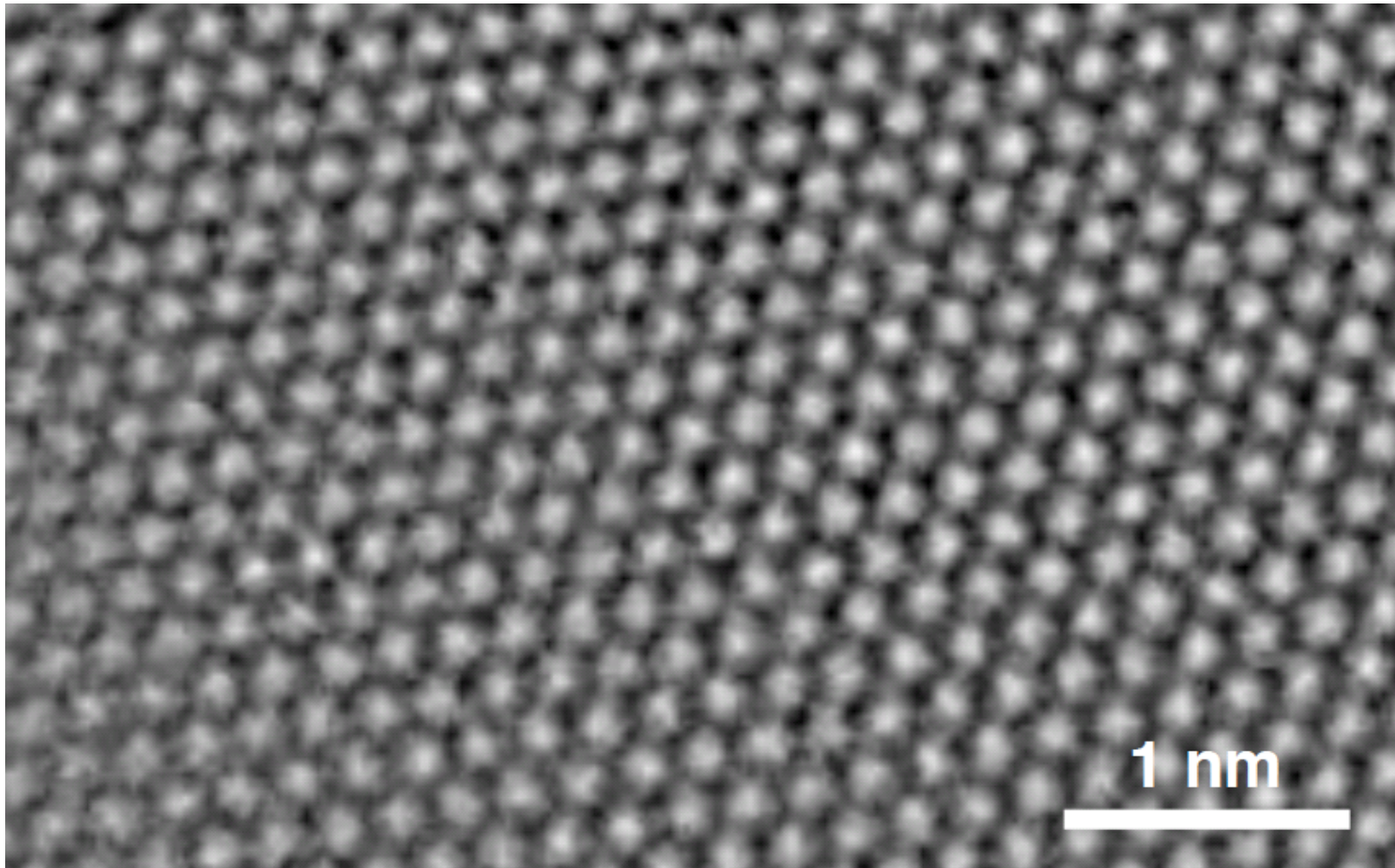
1. Independent of all the new knowledge that has been gained on grain boundary migration and grain growth in the end the motion of grain boundaries is still a complex 3D problem.
2. The dynamics of the individual atoms at the grain boundaries can not be captured (neither temporally nor spatially) by current experimental characterization techniques.

# Going back to 2D Grain Growth

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Where?

- solution: graphene



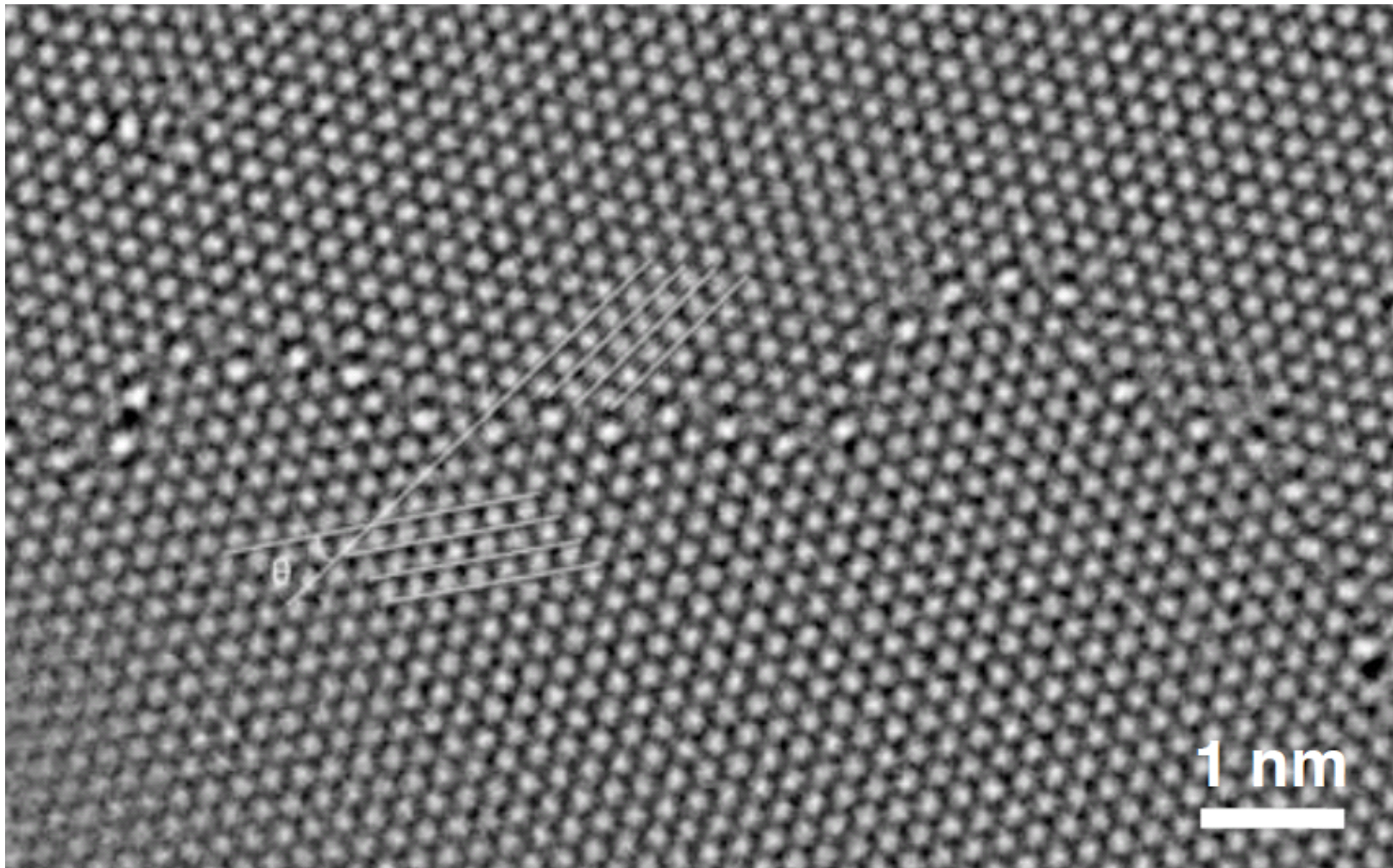
[Kurasch, Kotakoski, Lehtinen, Skákalová, Smet, Krill, Krasheninnikov, Kaiser: Nano Letters 12 (2012) 3168]

# Going back to 2D Grain Growth

---

Where?

- **solution:** polycrystalline graphene

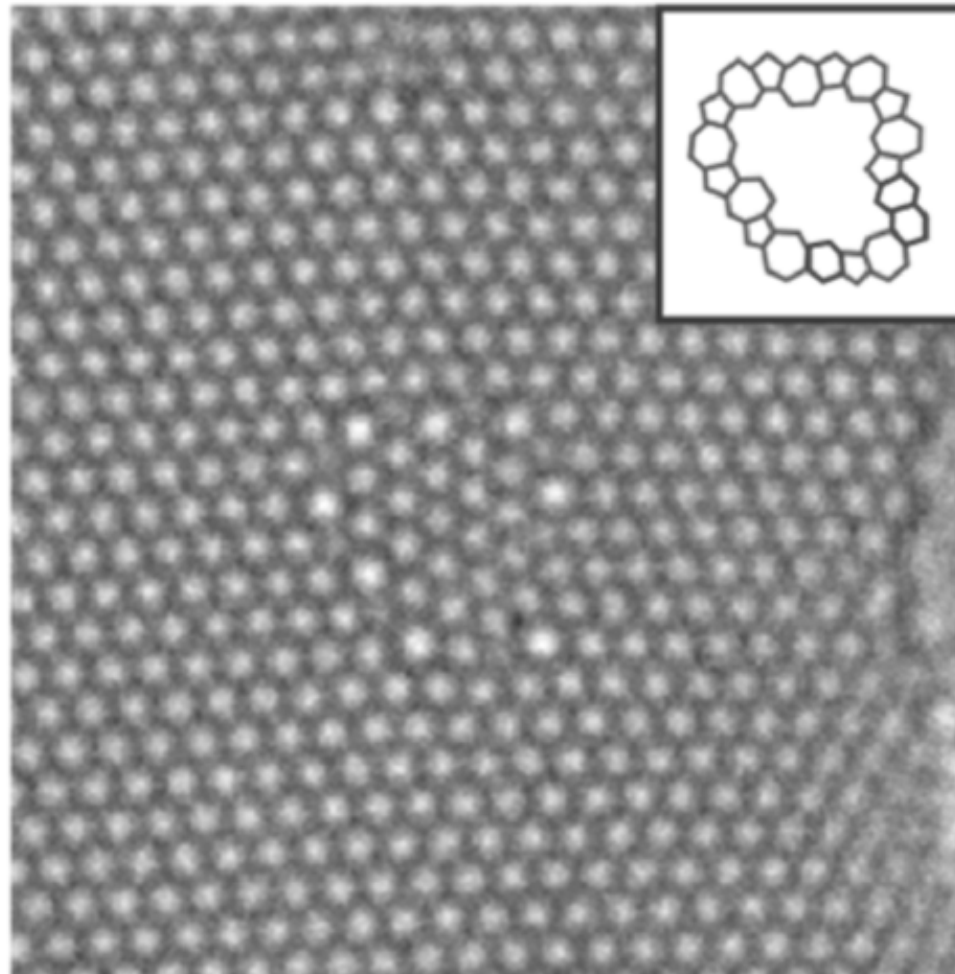


# Going back to 2D Grain Growth

---

Where?

- **solution:** polycrystalline graphene

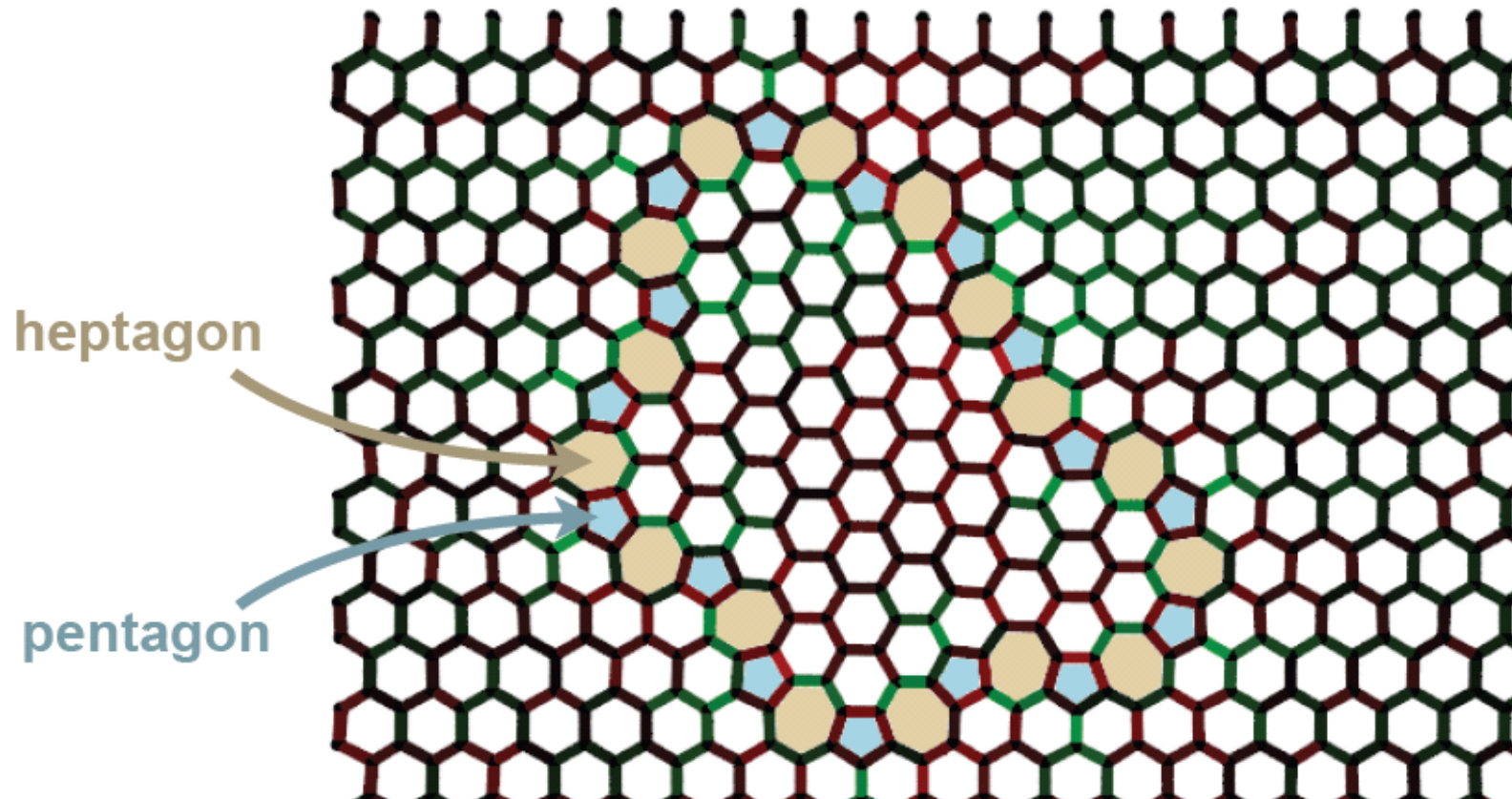


# Going back to 2D Grain Growth

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How?

- irradiation with 80 keV electrons



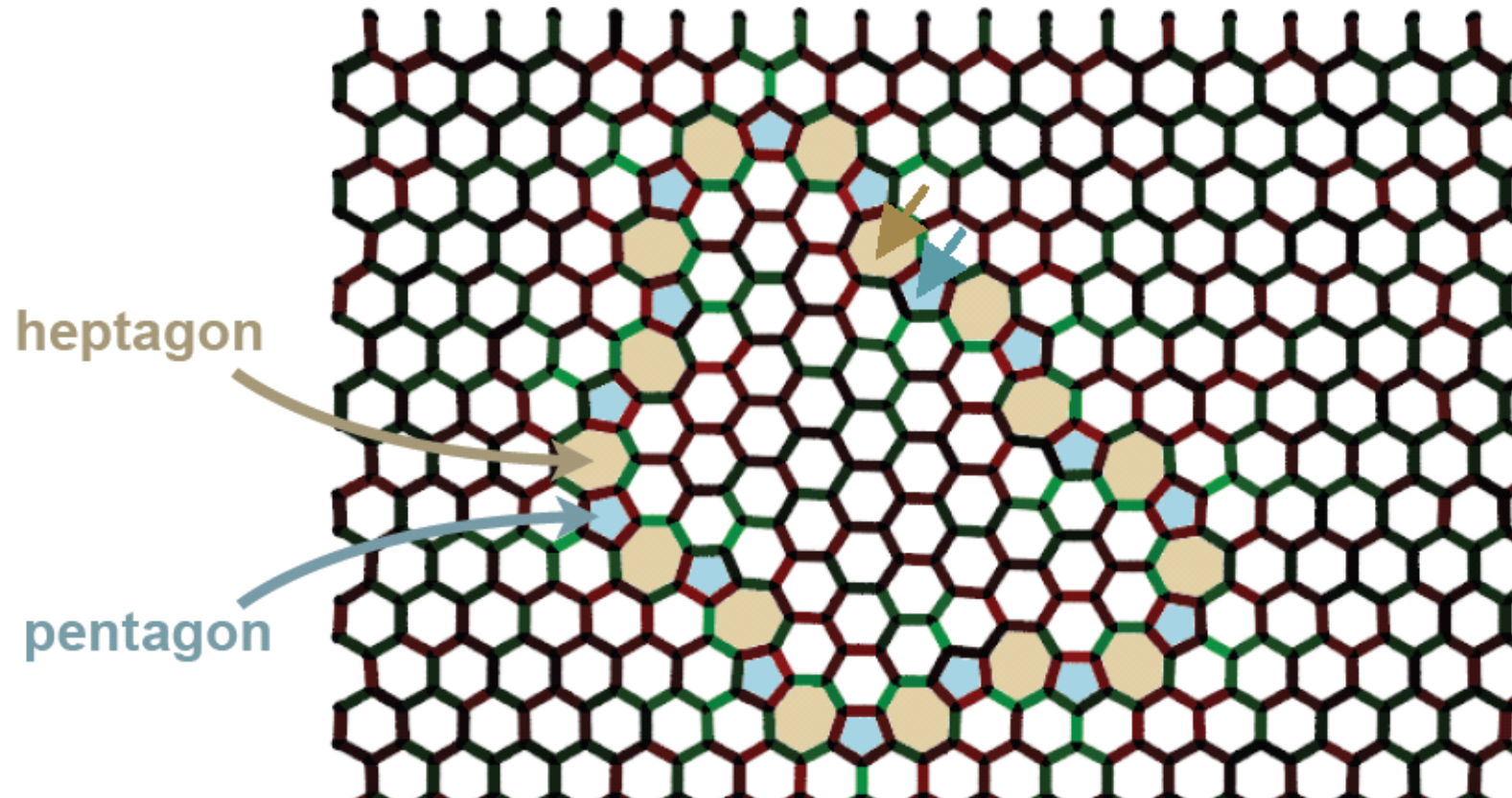
[J. Kotakoski, A.V. Krasheninnikov, U. Kaiser, J. C. Meyer: Phys. Rev. Lett. 106 (2011) 105505]

# Going back to 2D Grain Growth

---

How?

- irradiation with 80 keV electrons

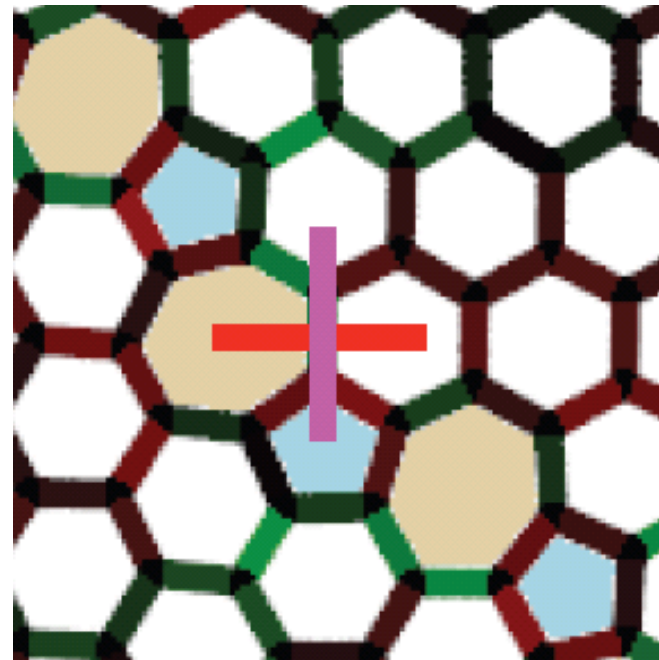
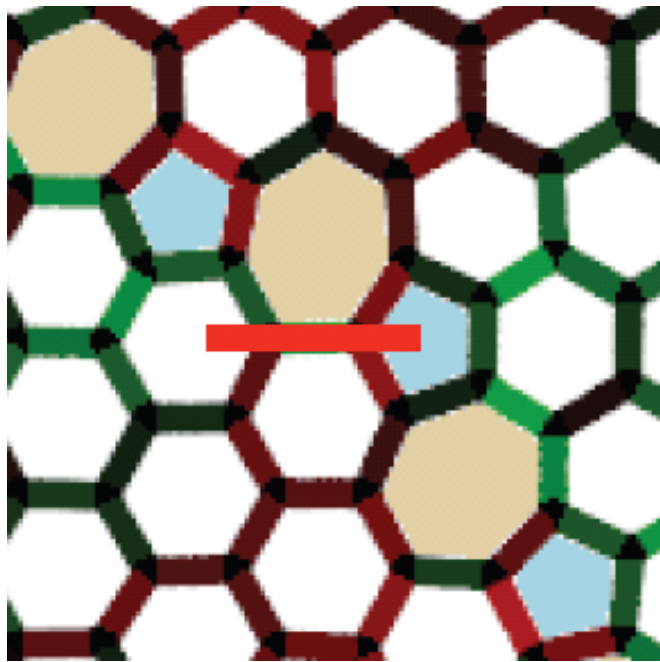


# Going back to 2D Grain Growth

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How?

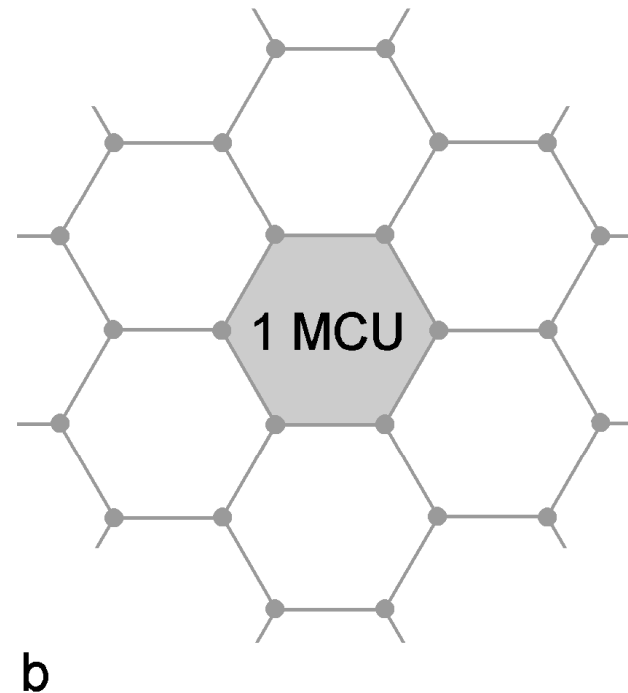
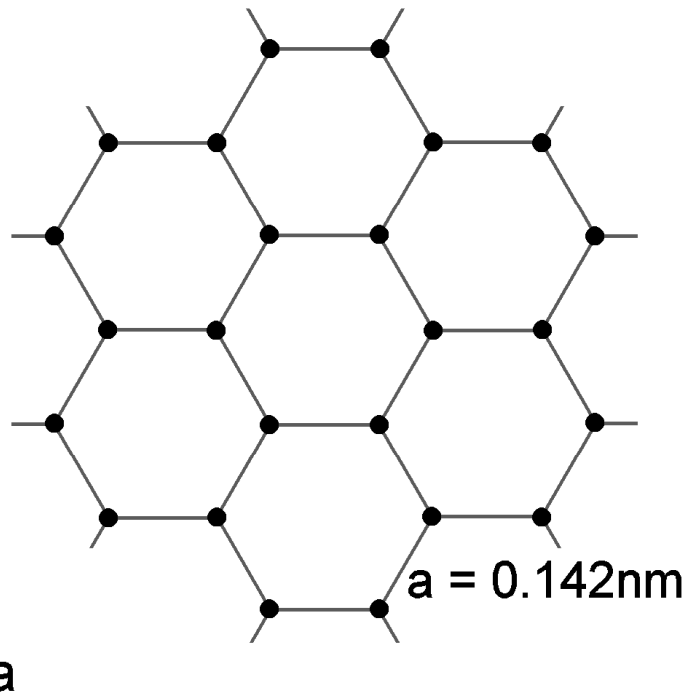
- irradiation with 80 keV electrons



# Going back to 2D Grain Growth

---

2D Potts model simulation:

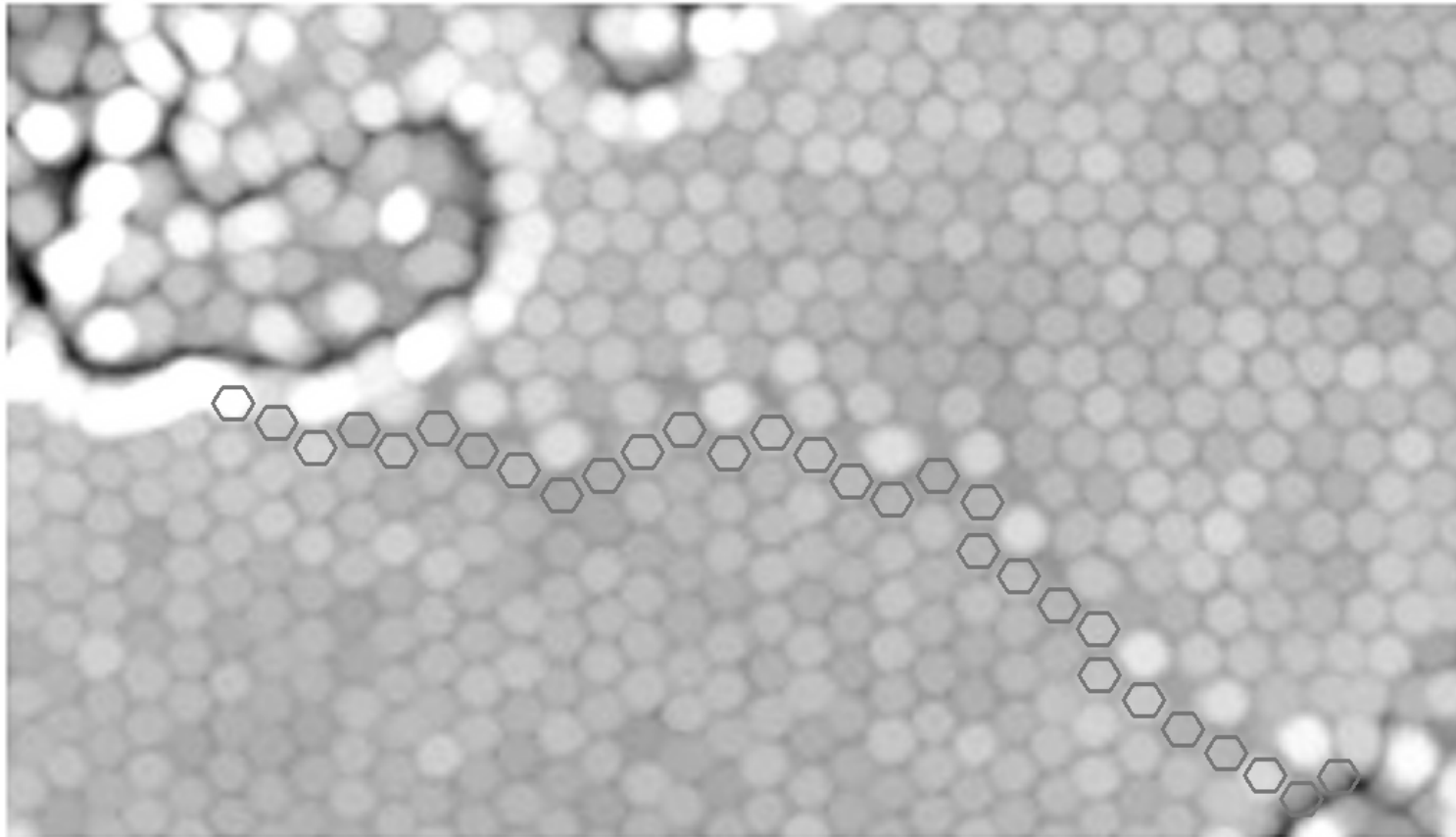


- one problem:
  - experiment: **two** grains = **two** lattice orientations
  - simulation: **two** grains = **one** underlying lattice

# Going back to 2D Grain Growth

---

2D Potts model simulation: grain boundary migration

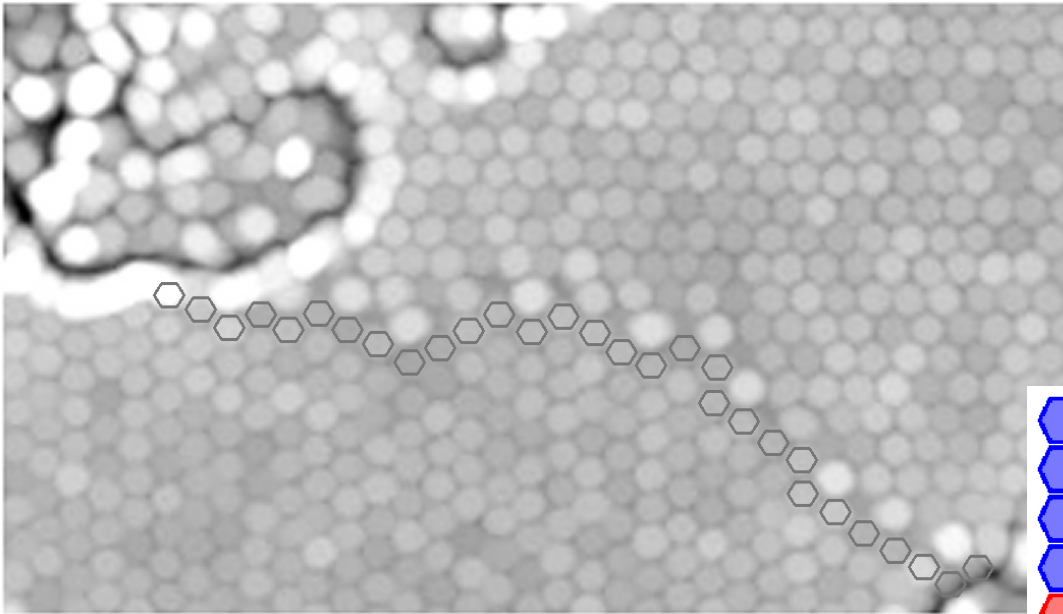


[S. Kurasch, J. Kotakoski, O. Lehtinen, V. Skákalová, J. Smet, C.E. Krill III, A.V. Krasheninnikov, U. Kaiser: Nano Letters 12 (2012) 3168]

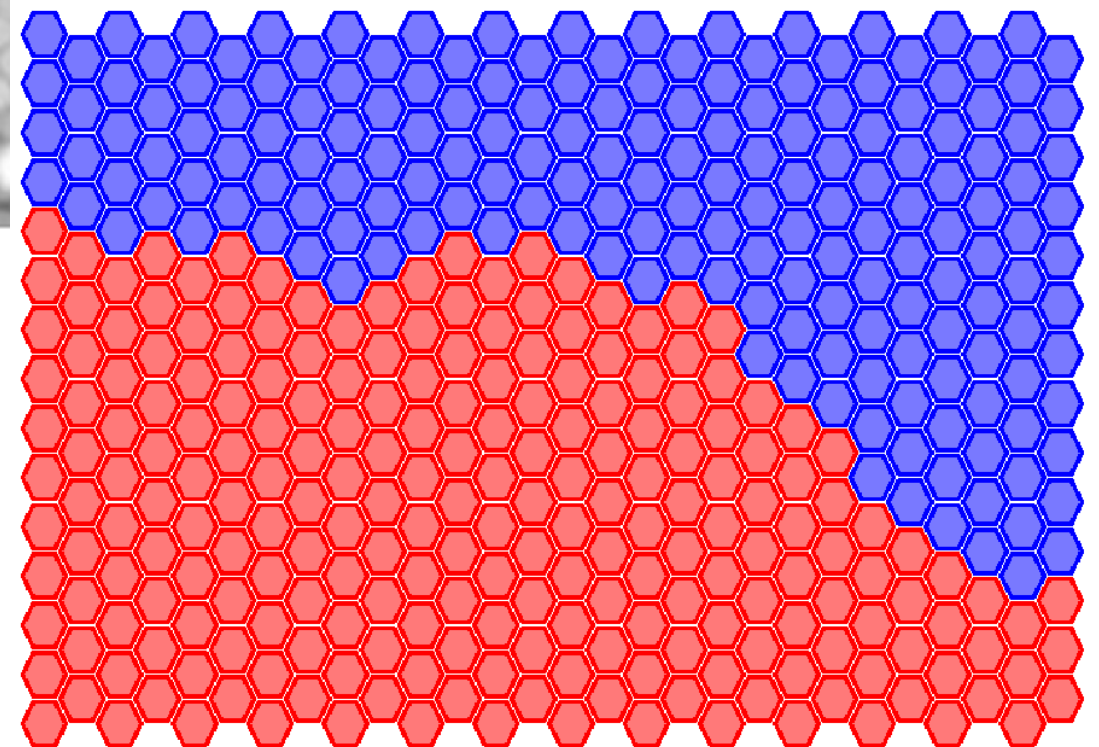
# Going back to 2D Grain Growth

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2D Potts model simulation:



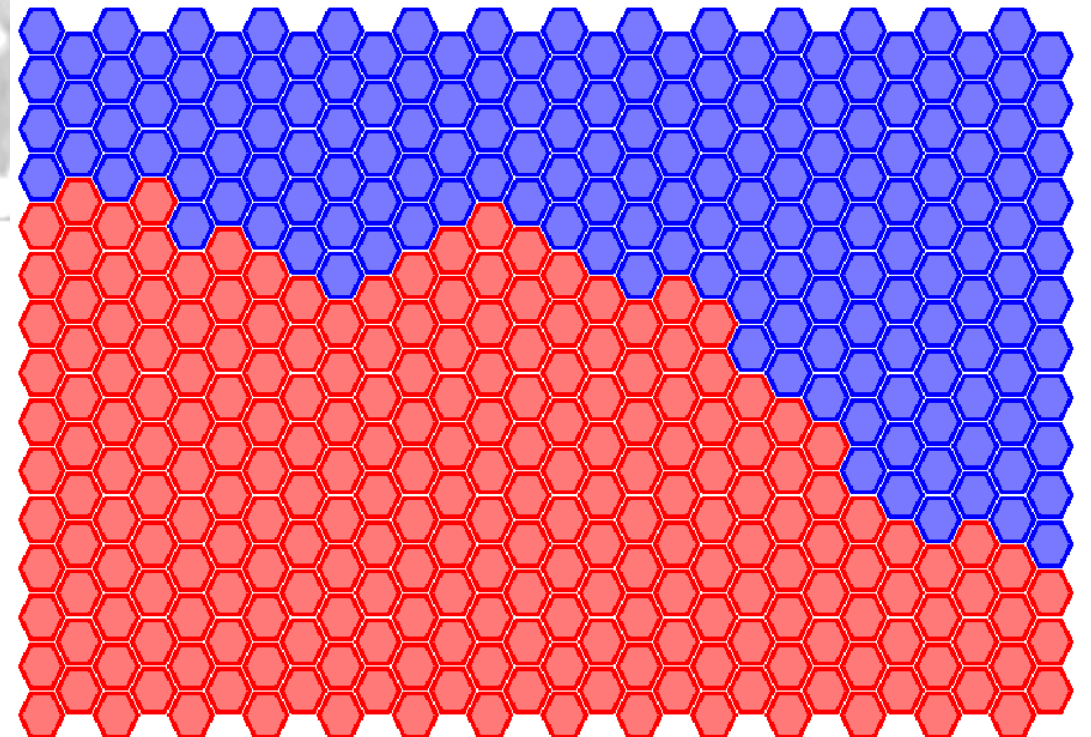
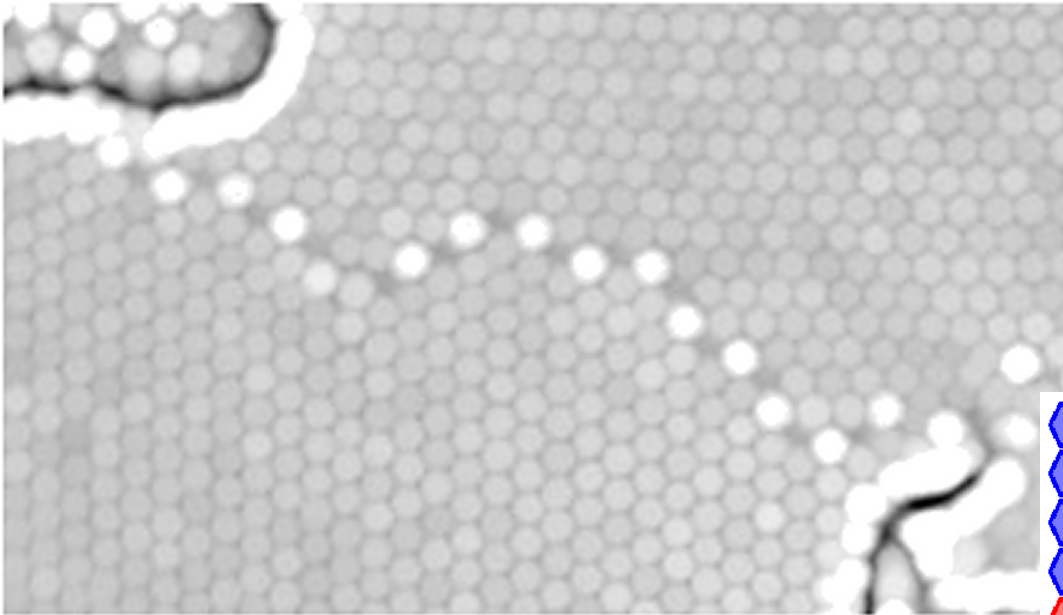
$t = 0.0 \text{ min}$



# Going back to 2D Grain Growth

---

2D Potts model simulation:

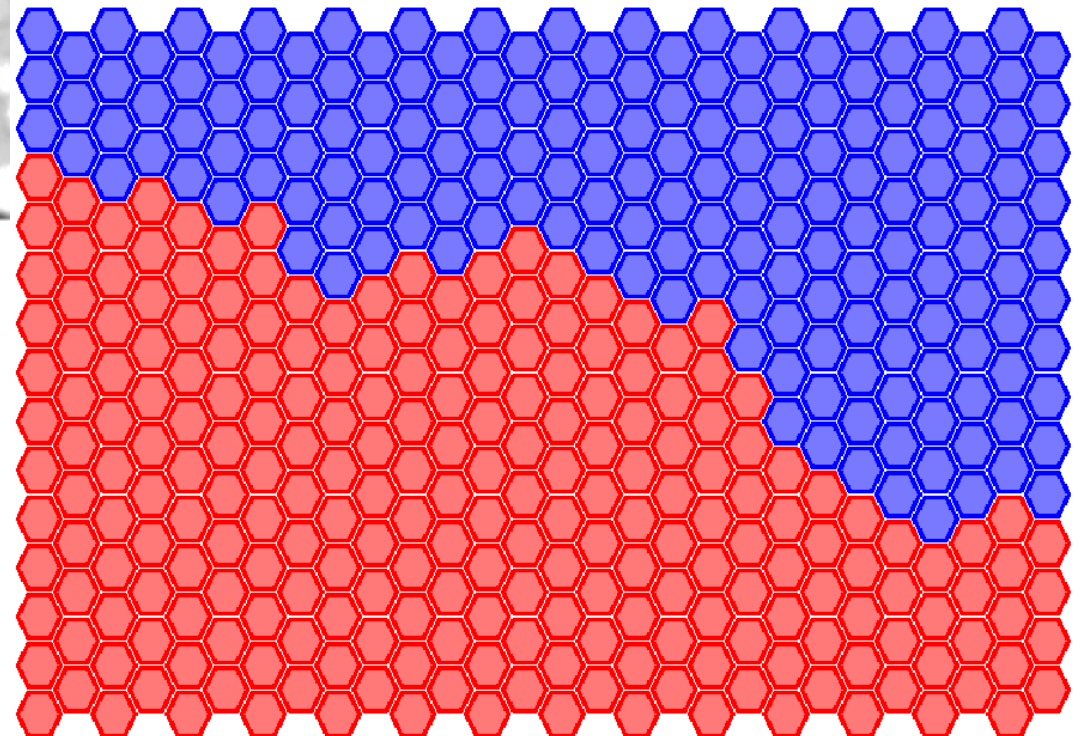
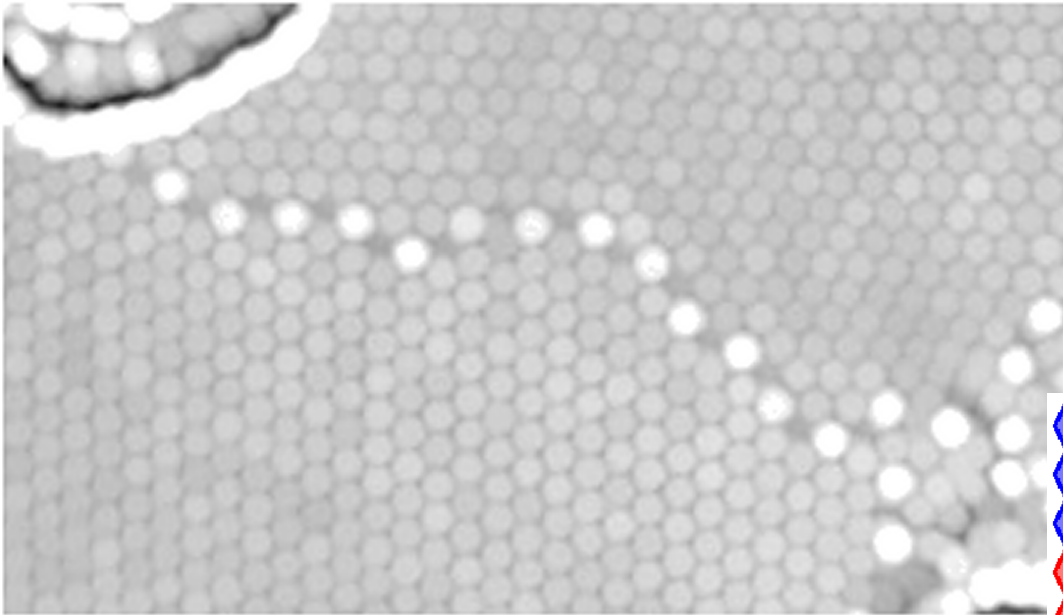


$t = 2.5 \text{ min}$

# Going back to 2D Grain Growth

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2D Potts model simulation:

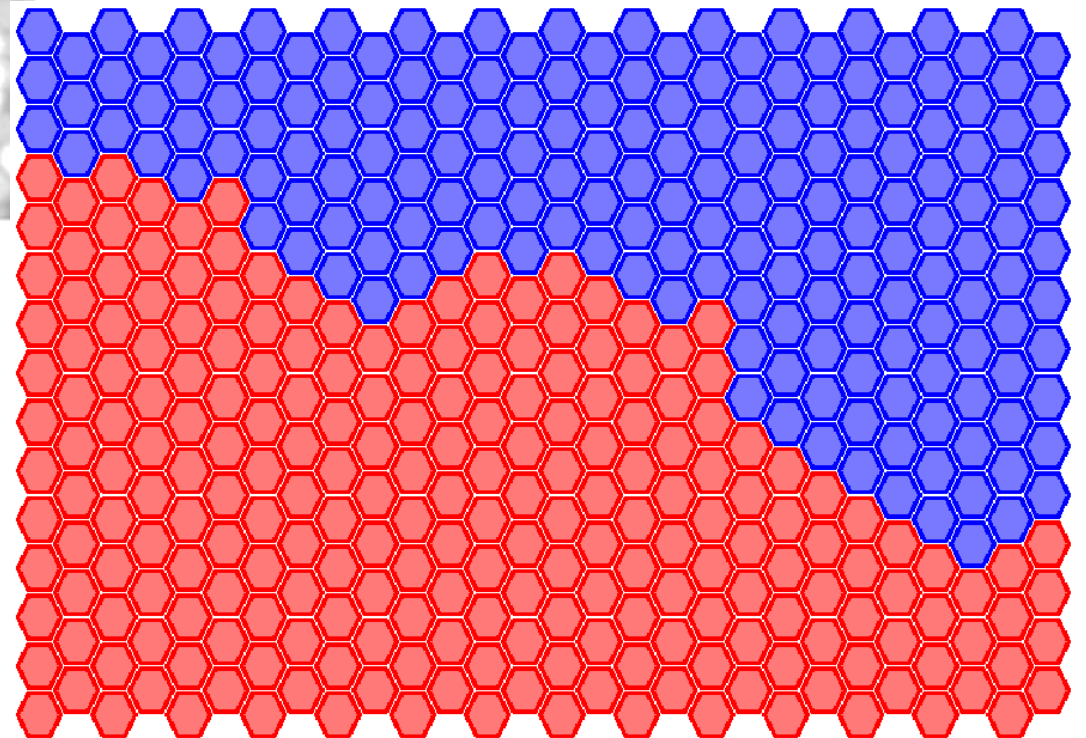
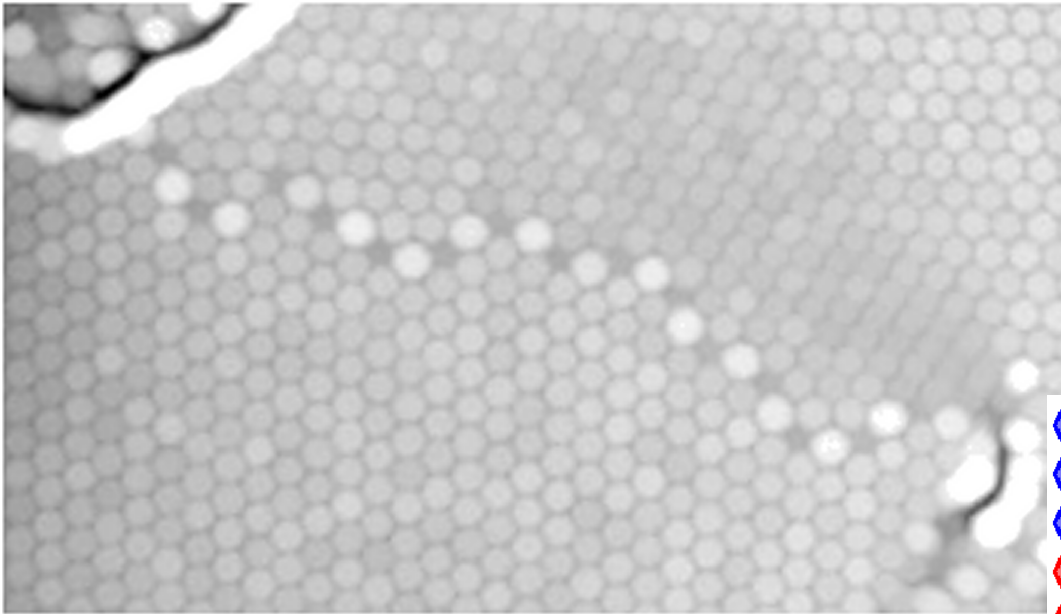


$t = 5.0 \text{ min}$

# Going back to 2D Grain Growth

---

2D Potts model simulation:

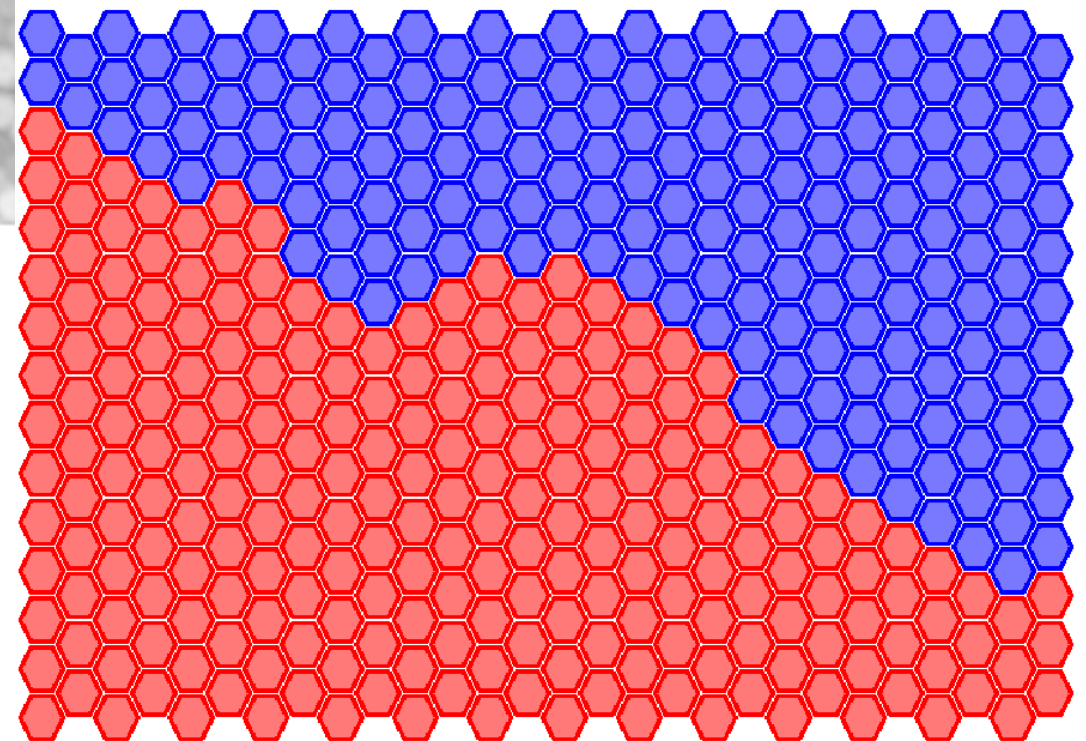
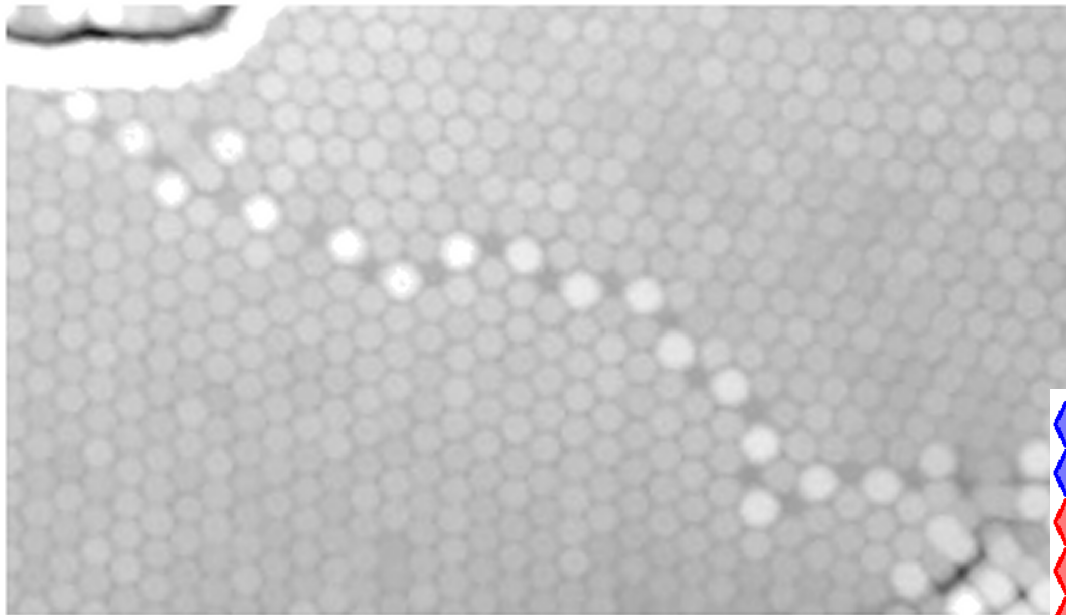


$t = 7.5 \text{ min}$

# Going back to 2D Grain Growth

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2D Potts model simulation:

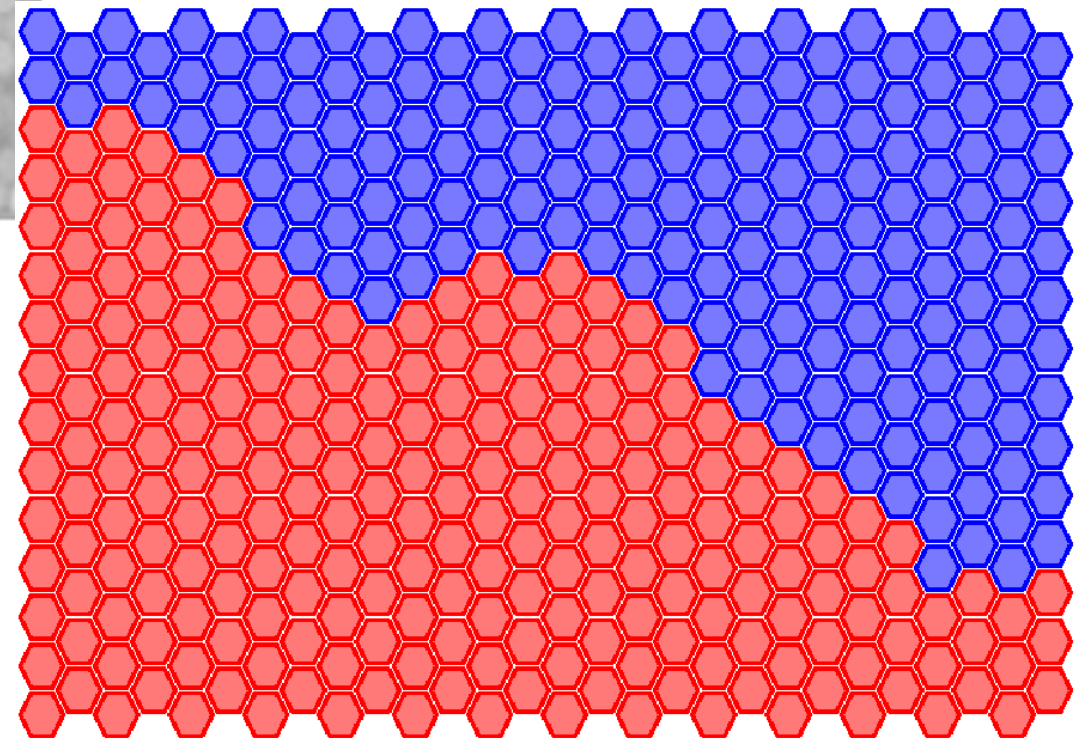
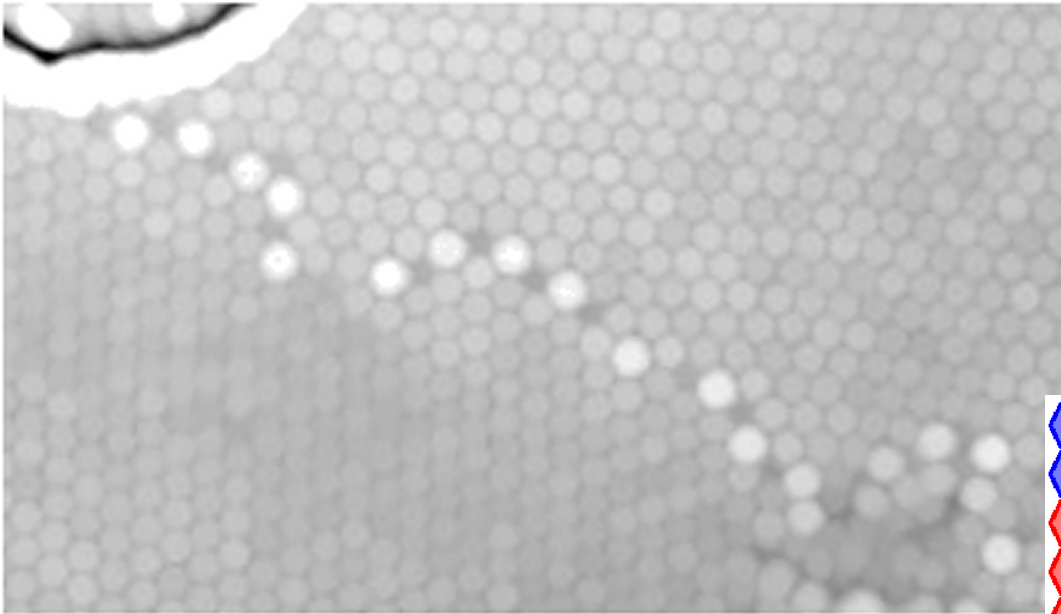


$t = 10.0 \text{ min}$

# Going back to 2D Grain Growth

---

2D Potts model simulation:

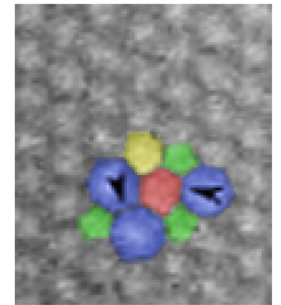
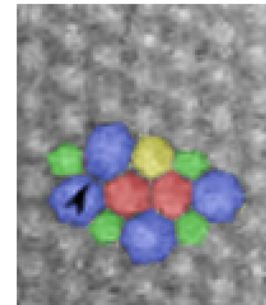
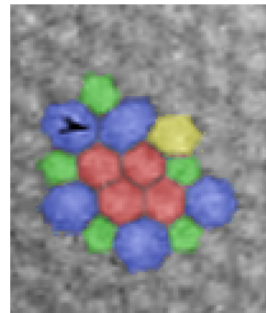
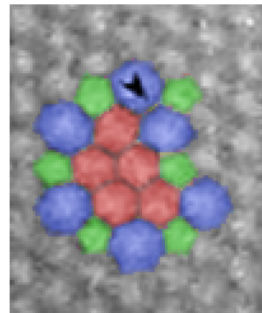
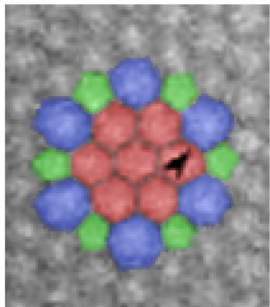


$t = 12.5 \text{ min}$

# Going back to 2D Grain Growth

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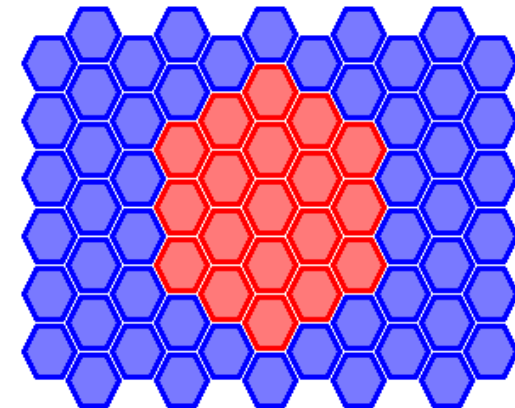
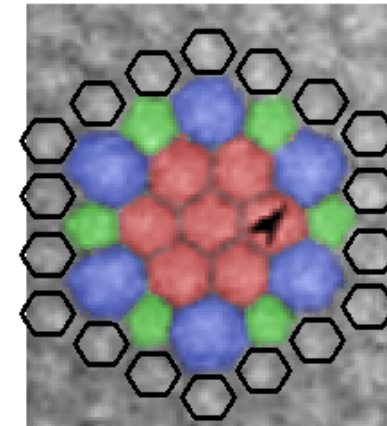
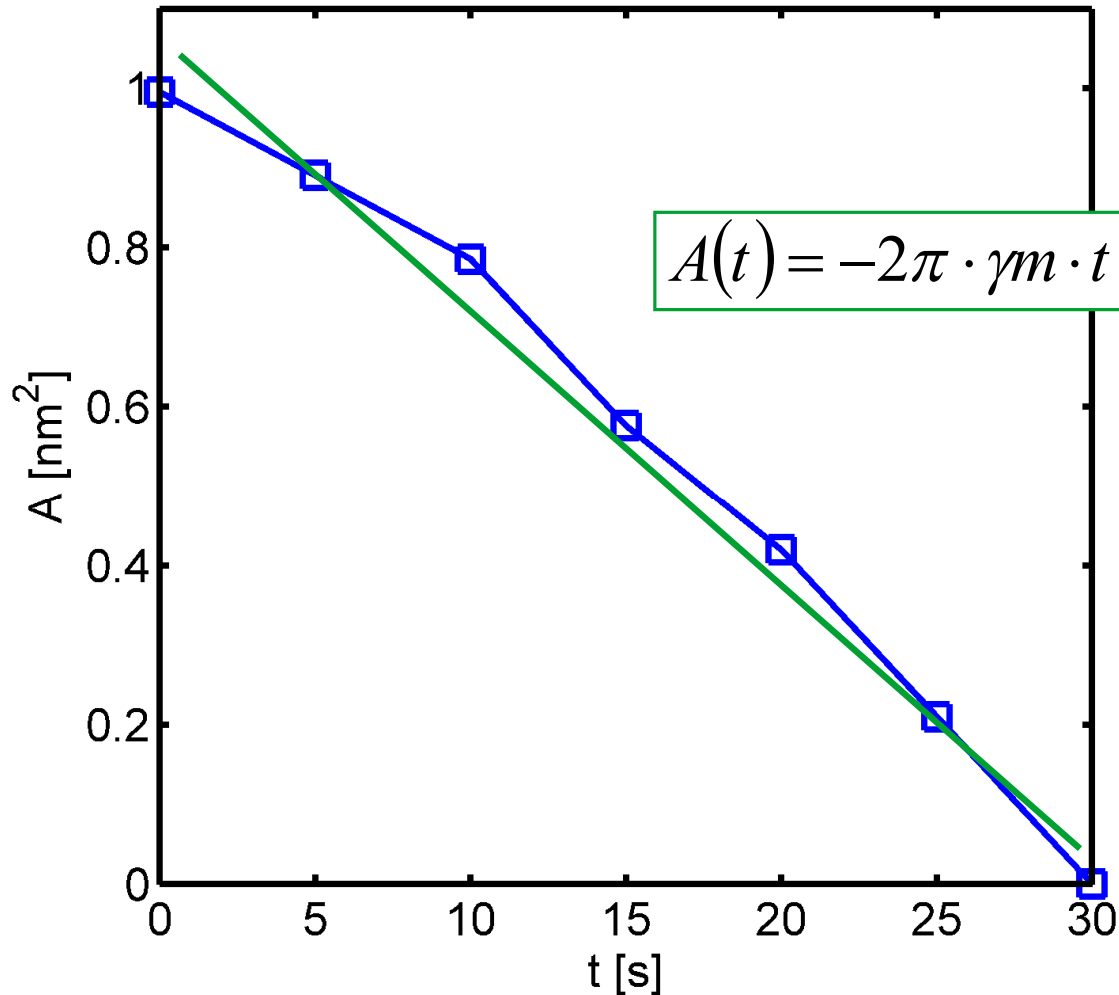
## 2D Potts model simulation: flower defect



[S. Kurasch, J. Kotakoski, O. Lehtinen, V. Skákalová, J. Smet, C.E. Krill III, A.V. Krasheninnikov, U. Kaiser: Nano Letters 12 (2012) 3168]

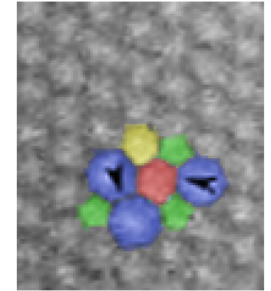
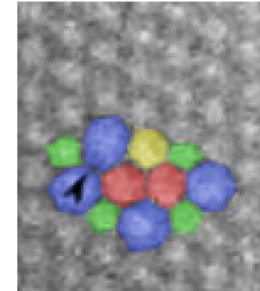
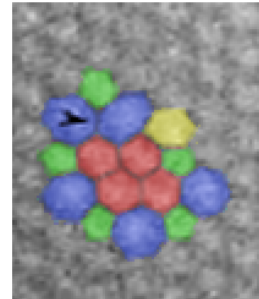
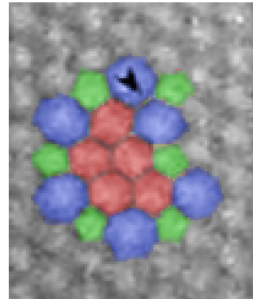
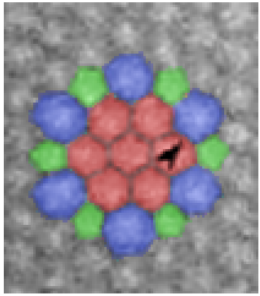
# Going back to 2D Grain Growth

## 2D Potts model simulation: flower defect

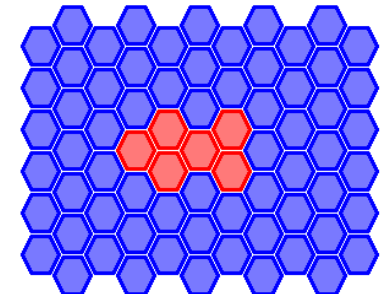
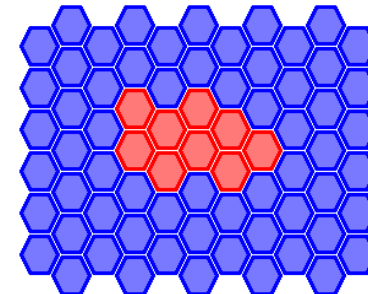
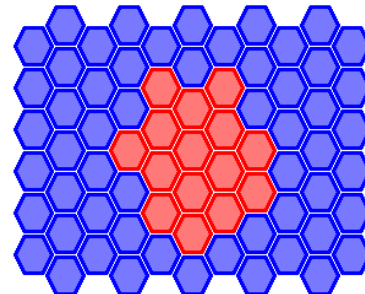
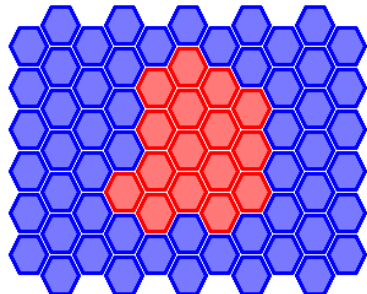
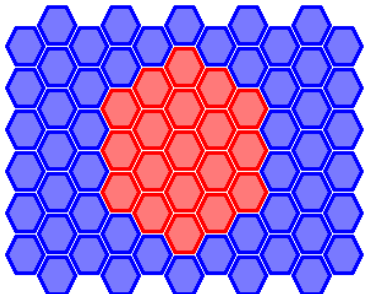
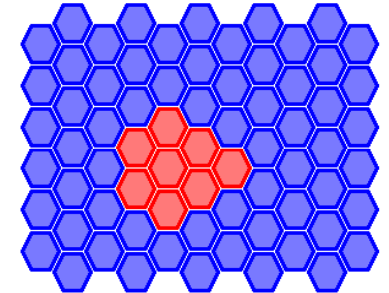
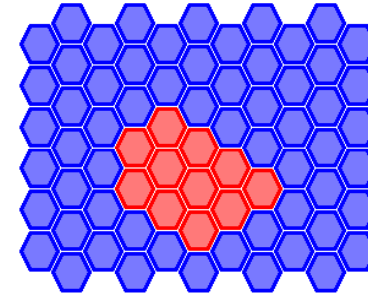
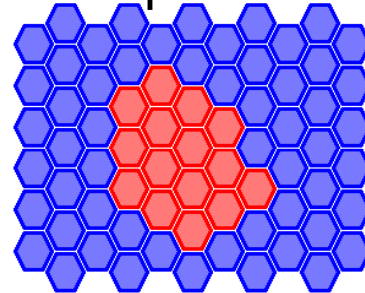
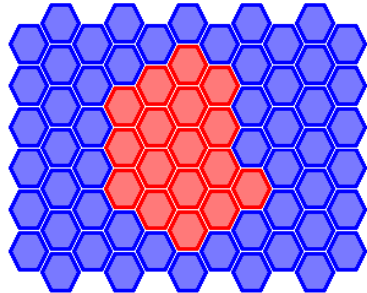
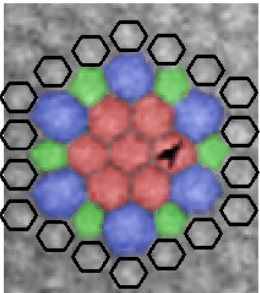


# Going back to 2D Grain Growth

## 2D Potts model simulation: flower defect



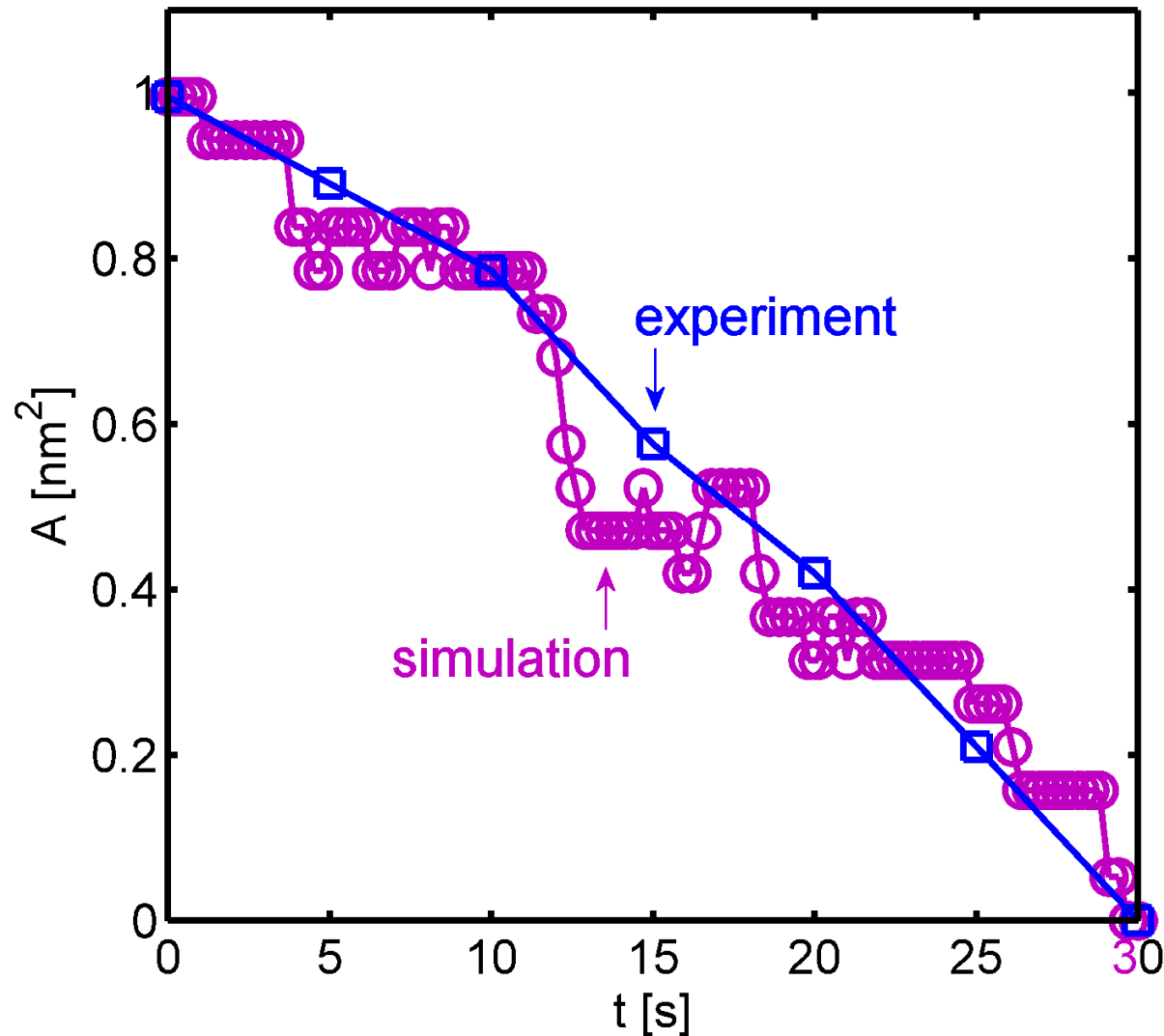
Experiment



Simulation

# Going back to 2D Grain Growth

## 2D Potts model simulation: flower defect



# Summary and Outlook

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## Summary:

- simulation of grain boundary migration and normal as well as nanocrystalline grain growth by Monte Carlo Potts model
- expectations (e.g. parabolic growth law, self-similarity) are fulfilled
- individual growth kinetics allows prediction of average growth law solely from kinetics of quadruple points
- mesoscopic 2D Potts model represents grain boundary migration in graphene quite well

# Summary and Outlook

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## Outlook:

- anomalous growth kinetics in nanocrystalline materials explained by junction control or abnormal grain growth?
- experimental determination of grain size distribution and topology in nanocrystalline materials
- junction drag in nanocrystalline thin films
- grain boundary migration in polycrystalline graphene

# Acknowledgements

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My special thanks goes to

- P. Streitenberger, S. Schäfer (University Magdeburg)
- C.E. Krill, J. Dake (Ulm University)
- T. Stöter (Leibnitz-Institute Magdeburg)
- I. Fielden (Sheffield Hallam University)

and to the

- German Science Foundation (DFG) for financial support.

# Thank you for your attention!

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