

Kinetic Limit for the

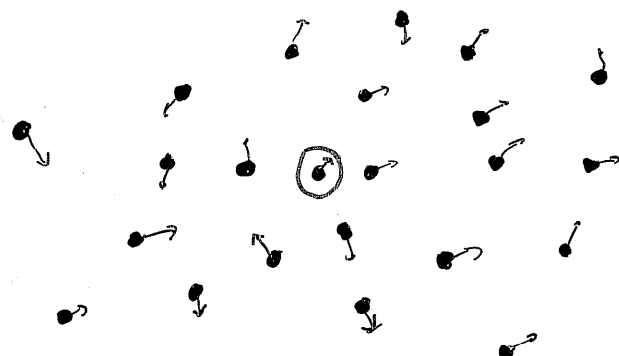
Weakly Nonlinear Schrödinger Equation

(and other wave equations)

joint work with

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motion in random medium



particles

(E. Frey

waves

$$i \frac{d}{dt} \psi = - \Delta \psi + (\psi^* \psi) \psi$$



random potential $V(x, t)$

1.) Kinetic Limit

- classical particles interacting through short range potential

Boltzmann

kinetic limit

range a 3D

$$\| a \rightarrow 0, N \rightarrow \infty, a^2 N \text{ fixed} \|$$

Grad (1951), Lanford (1975)

random initial data!

- What about waves?

Peierls (1929)

obtained kinetic equation

wave turbulence

$$\lambda^2 = \varepsilon$$

limit?

nonlinearity strength λ

time $\lambda^{-2} t$, space $\lambda^{-2} r$,

energy $\cong$ volume = λ^{-6}

$\lambda \rightarrow 0$

random initial data!

Boltzmann distribution function f

$\Rightarrow$ Wigner function of wave field $\psi^* \psi$

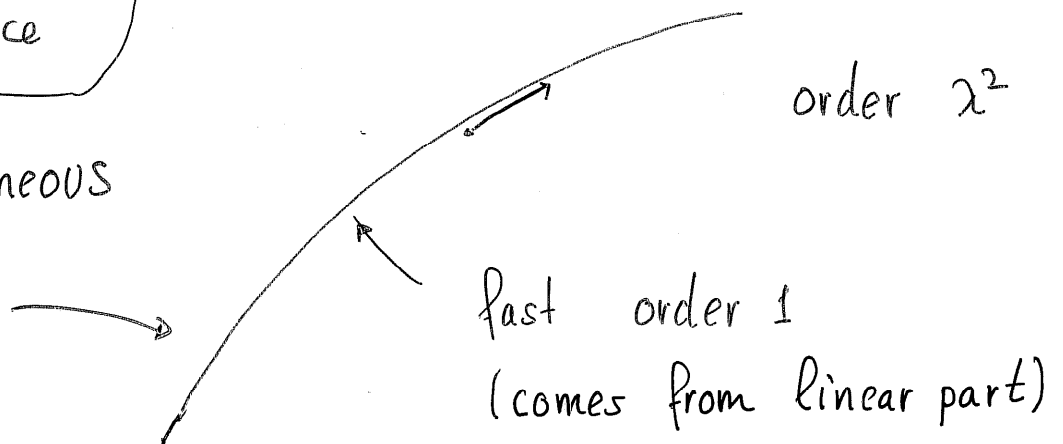
very rough picture

$$i \frac{d}{dt} \psi = -\Delta \psi + \lambda |\psi|^2 \psi$$

space of measures
on phase space

spatially homogeneous

slow
manifold



slow manifold consists of "natural" measures invariant
under the linear dynamics

particles : Poisson measures

waves : subclass of Gaussian measures

(gauge invariant)

Are there any proofs?

- linear wave equation with weak random coefficients

⇒ Schrödinger with random potential.

Erdős, Yau, ..., 2000 onwards

⇒ harmonic crystal with random masses

Lukkarinen, U.S. (2007)

⇒ wave equation with random index of refraction

Bal, Komorowski, Ryzhik (2003)

(high frequency)

- nonlinear wave equation THIS TALK

nonlinear Schrödinger equation

symmetry between q, p

2.) NLSwave field $\psi(x, t) \in \mathbb{C}$, $x \in \mathbb{R}^D$

$$i \frac{\partial}{\partial t} \psi = -\Delta \psi + \lambda |\psi|^2 \psi$$

$$\lambda \geq 0$$

 $\Rightarrow$ random initial data

singular at short distances

$$\mathbb{R}^D \leadsto \mathbb{Z}^D$$

- $\Delta \leadsto$ convolution with $\alpha(x)$,

|| α real, $\alpha(x) = \alpha(-x)$
 || compact support

 $\leadsto$ dispersion relation

$$\omega(k) = \hat{\alpha}(k), \quad k \in \mathbb{T}^D = [0, 1]^D$$

$$\omega(k) = \omega(-k), \quad \text{real analytic}$$

D-torus

$$\psi: \mathbb{Z}^D \rightarrow \mathbb{C}$$

$$i \frac{\partial}{\partial t} \psi(x, t) = \sum_y \alpha(x-y) \psi(y, t) + \lambda |\psi(x, t)|^2 \psi(x, t) \quad (*)$$

• on-site

$$\text{ALSO} \quad \sum_y |\psi(y, t)|^2 V(y-x) \psi(x, t)$$

• Hamiltonian system

$$\psi(x) = \frac{1}{\sqrt{2}} (p_x + i q_x)$$

energy

$$H = \sum_{x,y} \alpha(x-y) \psi(x)^* \psi(y) + \frac{1}{2} \lambda \sum_{x,y} |\psi(x)|^2 V(x-y) |\psi(y)|^2$$

(*) has solutions global in time of
infinite energy

Lanford, Lieb, Lebowitz 1974

Botta, Caglioti, DiRozza, Marchioro 2007

(propagation estimates)

number $\| \psi \|^2$

energy $H(\psi)$

are locally conserved

Fourier (momentum) space

$$i \frac{\partial}{\partial t} \hat{\psi}(k, t) = \omega(k) \hat{\psi}(k, t)$$

$$+ 2 \int_{(\mathbb{T}^D)^3} dk_2 dk_3 dk_4 \delta(k + k_2 - k_3 - k_4) \\ \times \hat{V}(k_2 - k_3) \hat{\psi}(k_2, t)^* \hat{\psi}(k_3, t) \hat{\psi}(k_4, t)$$

3. The Boltzmann-Peierls equation

random initial data Gaussian, translation
invariant

$$\langle \psi \rangle = 0, \quad \langle \psi \psi \rangle = 0$$

$$\langle \psi(x)^* \psi(y) \rangle = \int dk e^{i2\pi(y-x)k} W(k)$$

Wigner function

$$t > 0 \quad \langle \psi(x,t)^* \psi(x,t) \rangle = \int dk e^{i2\pi(y-x)k} W_\lambda(k,t)$$

Conjecture

$$\lim_{\lambda \rightarrow 0} W_\lambda(k, \lambda^{-2}t) = W(k,t)$$

and

W is solution of

(1D)

$$\frac{d}{dt} W(t) = \mathcal{L} W(t)$$

momentum energy

$$\mathcal{L} W(k_1) = 4\pi \int dk_2 dk_3 dk_4 \delta(k_1 + k_2 - k_3 - k_4) \delta(\omega_1 + \omega_2 - \omega_3 - \omega_4)$$

$$[W_1 W_3 W_4 + W_2 W_3 W_4 - W_1 W_2 W_3 - W_1 W_2 W_4]$$

formal derivation

$$\begin{aligned} \frac{d}{dt} \langle 44 \rangle_t &= \langle 4444 \rangle_t \\ &= \langle 4444 \rangle_0 + \underbrace{\int_0^t ds \langle 444444 \rangle_s} \end{aligned}$$

Gaussian truncation

|| problem remains OPEN ||

slow spatial variation

$$\Rightarrow \text{transport} \quad \frac{1}{2\pi} \nabla_k \omega \cdot \nabla_r$$

law of large numbers
B-P holds almost surely

4. Equilibrium time correlations

$\langle \cdot \rangle_\beta$ equilibrium measure small λ

$$\frac{1}{Z} e^{-\beta H(\psi) + \beta \mu \|\psi\|^2} \prod_{x \in \Lambda} d\psi(x) d\psi(x)^*$$

$\left\| \min_k \omega(k) > \mu \right\| \rightarrow$ exponential clustering

$\psi(x, t)$ is stationary random field

$$\langle \psi(x, t) \psi(y)^* \rangle_\beta = \int_{\mathbb{T}^D} dk C_\lambda(k, t) e^{i \cdot k(x-y)}$$

$$t=0 \quad \lim_{\lambda \rightarrow 0} C_\lambda(k, 0) = \frac{1}{\beta(\omega(k) - \mu)}$$

$$\lambda=0 \quad C_0(k, t) = e^{-i \cdot \omega(k)t} \frac{1}{\beta(\omega(k) - \mu)}$$

RESULT $t \geq 0$

$$\lim_{\lambda \rightarrow 0} e^{i\omega^\lambda(k)t/\lambda^2} \underline{C_\lambda(k, \lambda^{-2}t)} = \frac{1}{\beta(\omega - \mu)} e^{-\gamma(k)t}$$

• $\text{Re } \gamma(k) > 0$ DECA Y

• renormalized dispersion

$$\omega^\lambda(k) = \omega(k) + \lambda \underline{R_0(k)} + \lambda^2 R_1(k) + O(\lambda^3)$$

• γ, R_0, R_1 are explicit dangerous

There must be some conditions!

conditions

A) ℓ_1 -clustering (static)

$n \geq 4$

fully truncated

$$\sum_{x \in \mathbb{Z}^{nD}} |\delta_{x,0}| < \frac{n}{\prod_{j=1}^n 4(x_j)^{\beta_j}} \leq \lambda(c_0)^n n!$$

Abdessalam, Proccacia, Scoppola
(2009)

B) ℓ_3 -bound

$$p_t(x) = \int_{\pi^D} dk e^{-i\omega(k)t} e^{i k x}$$

$$\sum_x |p_t(x)|^3 \leq c(1+|t|)^{-1-\delta}$$

$$\delta > 0$$

• Stationary phase

$$D \geq 3$$

C) constructive interference

bad set $\mathcal{M}$

$$\left| \int_{\mathbb{T}^D} dk e^{-it(\omega(k) \pm \omega(k-k_0))} \right| \leq \frac{c}{1+|t|} \frac{1}{d(k_0, \mathcal{M})}$$

$\uparrow$
distance

- $0 \in \mathcal{M}$
- $\mathcal{M}$ discrete ? not clear?
requires $D \geq 3$ from other parts of proof
- $\mathcal{M}$ finite collection of smooth curves
o.k. for $-\Delta$ (n.n.)

requires $D \geq 4$

$$D \geq 4$$

D crossing bound

o.k. for $D \gg 3$, nearest neighbor

in addition

finite kinetic time

$$|t| < t_0$$

5. Strategy

- Duhamel expansion

$$\frac{d}{dt} a_t = \lambda (a_t)^3 \quad a_t \in \mathbb{R}$$

$$\frac{d}{dt} (a_t)^n = \lambda \underline{n} (a_t)^{n+2}$$

$$a_t = a_0 + \lambda \int_0^t ds (a_s)^3$$

iterate

$$\Rightarrow a_t = \sum_{n=0}^{\infty} \lambda^n c_n(t) (a_0)^{2n+1}$$

- convergence

$$c_n(t) = n! \frac{1}{n!} t^n$$

times Gaussian pairings for $\psi(t=0)$

$n!$

zero radius

- truncate series

$$\lambda^\alpha N! = 1$$

$$\psi(t) = \psi_{\text{main}}(t) + \underbrace{\psi_{\text{error}}(t)}$$

↗
 Duhamel expanded
 at order N
 only $\psi(t=0)$

full time evolution

- stationarity

$$|\langle \psi_{\text{error}}(t) * \psi(0) \rangle|^2$$

$$\leq \left\langle \left| \int_0^t ds \underbrace{A(t-s)}_{\psi(0)} \psi(s) \right|^{2N+1} \right\rangle \leq \langle |\psi(0)|^2 \rangle$$

// large, but finite N analysis //

Main task

high dimensional oscillatory integrals

with constraints

interaction: $\int \delta(k_1 + k_2 - k_3 - k_4) \hat{\psi}^* \hat{\psi} \hat{\psi}$

initial conditions:

$$\langle (\hat{\psi}^\#)^{2n} \rangle = \sum \text{pairings}$$

$$\langle \hat{\psi}(k)^* \hat{\psi}(k') \rangle = \delta(k - k') \frac{1}{\omega(k) - \mu}$$