



Gentlest Ascent Dynamics for Calculating First Excited State and Exploring Energy Landscape of Density Functionals

Chen Li, Jianfeng Lu, Weitao Yang,
Duke University,
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Outline



- Review of Kohn-Sham Eq. (SCF)
- Δ SCF¹
- Gentlest ascent dynamics² (GAD)
- GAD in KS-DFT³

1. T. Ziegler, A. Rauk, and E. J. Baerends, *Theor. Chim. Acta* **43**, 261 (1977).
2. W. E and X. Zhou, *Nonlinearity* 24, 1831 (2011) .
3. C. Li, J. Lu, and W. Yang, *J. Chem. Phys.* 143, 224110 (2015).

Kohn-Sham Equation

Given density functional $E[\rho]$, with ρ coming from $\det(\varphi_1, \varphi_2, \dots, \varphi_N)$,

$$\min_{\{\psi_i\}} E[\rho] - \sum_{ij} \lambda_{ij} (\langle \varphi_i | \varphi_j \rangle - \delta_{ij})$$

gives the ground state density and energy. This leads to Euler-Lagrange equation

$$h_s \varphi_i = \sum_j \lambda_{ij} \varphi_j,$$

where $h_s \varphi_i = \frac{\delta E}{\delta \varphi_i^*} = \frac{\delta E}{\delta \rho_s} \varphi_i$. By a unitary transformation, we can rewrite above Eq. as

$$h_s \psi_i = \varepsilon_i \psi_i,$$

where $\{\psi_i\}$ is a rotation of $\{\varphi_i\}$, and are called canonical orbitals; ε_i is the orbital energy.

One solves the above Kohn-Sham equation to obtain ground state, through iterative procedure, the SCF procedure.

SCF procedure



SCF iteration procedure:

at iteration step k:

$$\text{given } \rho_s^{in} = \sum_{i=1}^{occ} |\psi_i^{in}\rangle\langle\psi_i^{in}|, \text{ compute } h_s^{out} = \frac{\delta E}{\delta \rho_s^{in}}.$$

$$\text{Diagonalize } h_s^{out} = \sum_{i=1}^{\infty} \varepsilon_i^{out} |\psi_i^{out}\rangle\langle\psi_i^{out}|,$$

$$\text{form } \rho_s^{out} = \sum_{i=1}^{occ} |\psi_i^{out}\rangle\langle\psi_i^{out}|.$$

Question: How to choose orbital
to form density matrix?

Check $\rho_s^{out} = \rho_s^{in}$?

Yes: solution found, iteration ended.

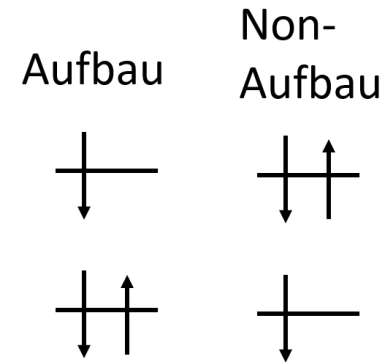
No: set $\rho_s^{in} = \rho_s^{out}$, start next iteration step.



- Aufbau and non-Aufbau solutions

- Δ SCF

- Two SCF calculations:



ground state excited state?

- Both solutions are stationary
- Question: are non-Aufbau solutions meaningful?

Equivalently: if ρ_i is non-ground state stationary density, is $E[\rho_i]$ excited state energy?



- Stationary solutions of the exact density functionals $E[\rho]$ ¹
Stationary state density $\bar{\rho}_i \longrightarrow \bar{\Psi}_i$, excited state wavefunction,
 $E[\bar{\rho}_i]$ is the excited state energy.
 - Implicit assumption: ρ_i is excited state v -representable
 - Numerically give reasonable results, accuracy \sim TDDFT²⁻⁴
 - merit: low computational cost

1. M. Levy and J. P. Perdew, Phys. Rev. A 32, 2010 (1985).
2. C.-L. Cheng, Q. Wu, and T. Van Voorhis, J. Chem. Phys. **129**, 124112 (2008).
3. T. Kowalczyk, S. R. Yost, and T. V. Voorhis, J. Chem. Phys. **134**, 054128 (2011).
4. S. R. Yost, T. Kowalczyk, and T. Van Voorhis, J. Chem. Phys. **139**, 174104 (2013).



- Problem:
 - No guarantee to converge
 - Possible scenario: frontier orbitals switch back and forth
- Other approach
 - Constricted DFT¹
 - Maximum overlap method (MOM)²
 - Meta-dynamics related methods³
- Our approach: an alternative perspective
 - Use stationary nature
 - Target saddle point of energy functional

1. T. Ziegler, M. Seth, M. Krykunov, J. Autschbach, and F. Wang, J. Chem. Phys. **130**, 154102 (2009).

2. N. A. Besley, A. T. B. Gilbert, and P. M. W. Gill, J. Chem. Phys. **130**, 124308 (2009).

3. A. J. W. Thom and M. Head-Gordon, Phys. Rev. Lett. **101**, 193001 (2008).

GAD formalism



- Gentlest ascent dynamics¹ Index-1 saddle point dynamics

Given function $F(x)$ $x = (x^1, x^2, \dots, x^n)$.

Recall steepest descent

$$\frac{d}{dt} x(t) = -\frac{\partial F}{\partial x}$$

$$\frac{d}{dt} v(t) = -\frac{\partial^2 F}{\partial x^2} v, \quad \|v\| = 1$$

$$\frac{d}{dt} x(t) = -\frac{\partial F}{\partial x} + 2 \left\langle \frac{\partial F}{\partial x}, v \right\rangle v$$

1. W. E and X. Zhou, Nonlinearity 24, 1831 (2011)

GAD formalism



- Gentlest ascent dynamics in DFT

- Energy functional $E = E[\Psi]$ $\Psi = (\psi_1, \psi_2, \dots, \psi_N)^T \in \mathcal{W}$
- admissible set $\mathcal{W} = \{\Psi \mid (\psi_i, \psi_j) = \delta_{ij}\}$

- Define inner product

$$\langle \Psi, \Phi \rangle = \sum_{i=1}^N (\psi_i, \phi_i)$$

- Define derivatives

$$\frac{\partial E}{\partial \Psi} \Phi = \left\langle \frac{\partial E}{\partial \Psi}, \Phi \right\rangle = \sum_{i=1}^N \left(\frac{\delta E}{\delta \psi_i}, \phi_i \right)$$

$$\left(\frac{\partial^2 E}{\partial \Psi^2} \Phi \right)_i = \sum_{j=1}^N \left(\frac{\delta^2 E}{\delta \psi_i \delta \psi_j}, \phi_j \right)$$

GAD formalism

- Energy functional expansion

note if $\Psi \in \mathcal{W}$, $\Psi + t\Phi \notin \mathcal{W}$! define $\hat{\mathbf{F}}$, s.t. $\hat{\mathbf{F}}(\Psi + t\Phi) \in \mathcal{W}$.

$$\begin{aligned} L[\Psi + t\Phi] &= E[\hat{\mathbf{F}}(\Psi + t\Phi)] \\ &= L[\Psi] + t \frac{\partial L}{\partial \Psi}[\Phi] + \frac{1}{2} t^2 \langle \Phi, \hat{\mathbf{H}}\Phi \rangle \end{aligned}$$

$$\frac{\partial L}{\partial \Psi} : \text{effective gradient}, \quad \hat{\mathbf{H}} = \frac{\partial^2 L}{\partial \Psi^2} : \text{effective Hessian.}$$

By some calculation,

$$\frac{\partial L}{\partial \Psi} = \frac{\partial E}{\partial \Psi} \quad \hat{\mathbf{H}} = \frac{\partial^2 E}{\partial \Psi^2} - 2\hat{\Gamma}$$

$$\left(\hat{\Gamma}\Phi \right)_i = \sum_j (\hat{h}_s \psi_i, \psi_j) \varphi_j = \sum_j \gamma_{ij} \varphi_j$$

GAD formalism



- Gentlest ascent dynamics in DFT

$$\frac{d}{dt} \Psi = -\frac{\partial E}{\partial \Psi} + 2 \left\langle \frac{\partial E}{\partial \Psi}, \Phi \right\rangle \Phi + \Lambda \Psi$$

$$\frac{d}{dt} \Phi = -\hat{H} \Phi + \mu \Phi + \mathbf{K} \Psi$$

Lagrange Multipliers

μ

Λ

\mathbf{K}

Constraints

$$\|\Phi\| = 1$$

$$\Psi \in \mathcal{W}$$

$$(\psi_i, \varphi_j) = 0 \text{ for all } i \text{ and } j$$



GAD formalism

- Gentlest ascent dynamics in KS-orbitals
- Remove Lagrange multipliers

$$\frac{d\psi_i^\sigma(\mathbf{r})}{dt} = -\hat{h}_s^\sigma \psi_i^\sigma(\mathbf{r}) + 2 \sum_{j\sigma'} (\hat{h}_s^{\sigma'} \psi_j^{\sigma'}, \varphi_j^{\sigma'}) \varphi_i^\sigma(\mathbf{r})$$

$$\frac{d\varphi_i^\sigma(\mathbf{r})}{dt} = -\hat{h}_s^\sigma \varphi_i^\sigma(\mathbf{r}) - \sum_{\sigma'} \hat{V}^{\sigma\sigma'} [2 \sum_j \psi_j^{\sigma'} \varphi_j^{\sigma'}] \varphi_i^\sigma(\mathbf{r}) + \sum_j \gamma_{ij}^\sigma \varphi_j^\sigma(\mathbf{r})$$

With constraints

$$\left\{ \begin{array}{l} (\psi_i^\sigma, \psi_j^\sigma) = \delta_{ij}, \\ \sum_{i\sigma} (\varphi_i^\sigma, \varphi_i^\sigma) = 1, \\ (\psi_i^\sigma, \varphi_j^\sigma) = 0. \end{array} \right.$$

$$\hat{V}^{\sigma\sigma'}[f] = \int \frac{\delta v_{\text{eff}}^\sigma(\mathbf{r})}{\delta \rho^{\sigma'}(\mathbf{r}')} f(\mathbf{r}') d\mathbf{r}'$$

- Direct implementation cost $O(N^4)$
- Using density fitting, reduce cost to $O(N^3)$
- Can be modified to target index-2 or higher saddle points

Results

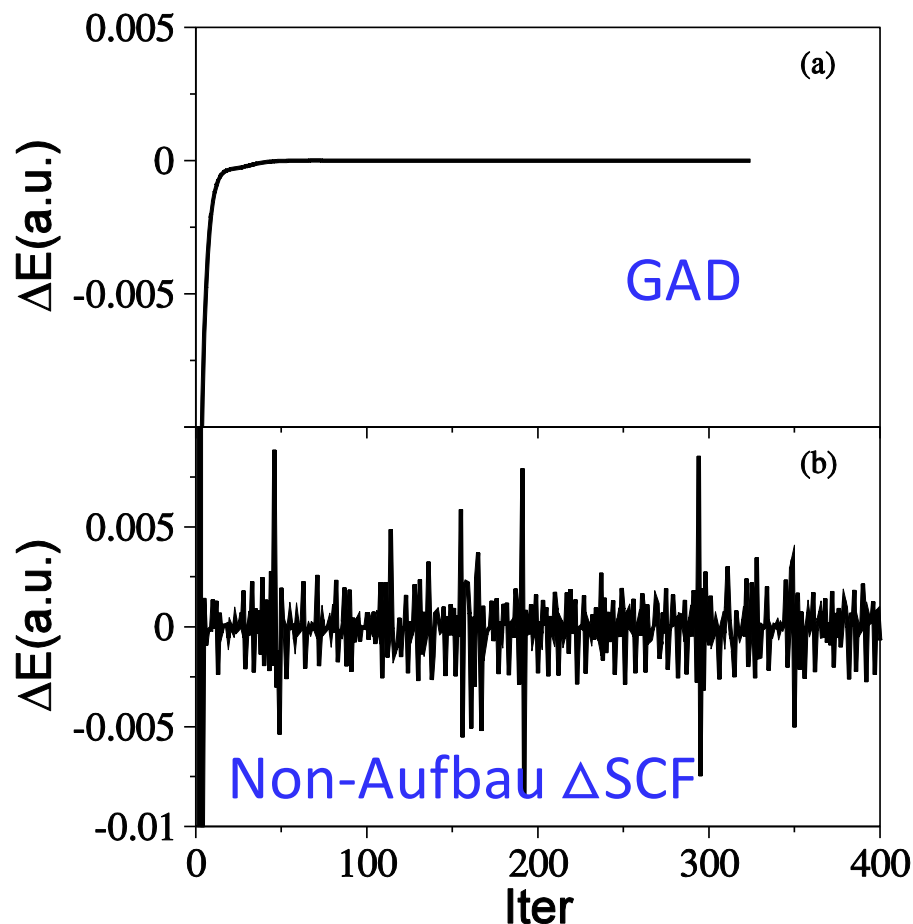
- Comparison between non-Aufbau Δ SCF and GAD

	Configuration	$E_{\Delta\text{SCF}}$ (a.u.)	E_{diff} (10^{-8} a.u.)
H	2s	-0.12530541	0.4
He	$(1s)^1(2s)^1$	-1.96025338	0.0
Li	$(1s)^2(2p)^1$	-7.27816628	-108
	$(1s)^1(2s)^2$	-5.22918853	4.9
Be	$(1s)^2(2s)^1(2p)^1$	-14.31663987	-1.1
H ₂	$(1\sigma)^1(1\sigma^*)^1$	-0.77725246	-1.9
Li ₂	1 st .	-14.67364711	0.5
CO	1 st .	-112.15621160	1.4
OH	1 st .	Does not converge	converged
HF	1 st .	Does not converge	Converged
H ₂ O	1 st .	-75.55256992	-6.8

Results

- Comparison between non-Aufbau Δ SCF and GAD

first excited state
calculation of HF



Results



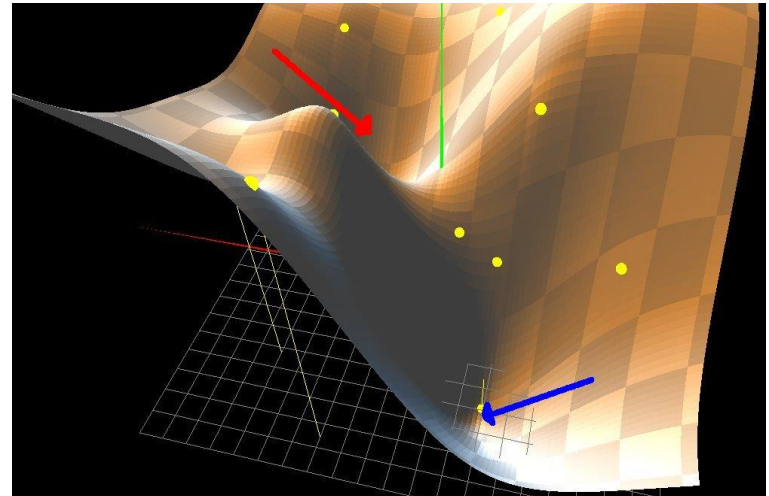
- Eigenvalues of Hessian by direct diagonalization

	Configuration	Eigenvalue (a.u.)				
		1	2	3	4	5
H	1s	0.0000	0.3120	0.4773	0.4773	0.4773
	2s	-0.4417	0.0000	0.1340	0.1340	0.1340
	3s	-0.4315	-0.1084	0.0000	0.0008	0.0027
He	$(1s)^2$	0×2	0.7594	0.7594	0.7594	0.8340
	$(1s)^1(2s)^1$	-0.8715	0×2	0.0033	0.0035	0.6250
Li	$(1s)^2(2s)^1$	0×5	0.0831	0.0831	0.0831	0.1802
	$(1s)^2(2p)^1$	-0.0285	0×5	0.0004	0.0010	0.1279
	$(1s)^1(2s)^2$	-2.402	0×5	0.0341	0.0341	0.0341
Be	$(1s)^2(2s)^2$	0×8	0.0574	0.0574	0.0574	0.2779
	$(1s)^2(2s)^1(2p)^1$	-0.1074	0×8	0.0002	0.0010	0.1536

Results



- 1st excited state
 - Index-1 saddle point
 - Due to the manifold constraint
- Contrast to free dynamics
 - 2nd smallest stationary state given by local minimum



Results

- Eigenvalues of Hessian by direct diagonalization

	Configuration	Eigenvalue (a.u.)				
		1	2	3	4	5
H ₂	(1σ) ²	0 × 2	0.2215	0.5593	0.8388	1.033
	(1σ)(1σ*)	-0.3054	0 × 2	0.5188	0.5993	0.7674
	(1σ*) ²	-0.5135	-0.1142	0 × 2	0.5161	0.6166
Li ₂	gs.	0 × 18	0.0224	0.0369	0.0369	0.0516
	1 st .	-0.0292	0 × 18	0.0211	0.0211	0.0534
CO	gs.	0 × 98	0.1889	0.1889	0.2628	0.3167
	1 st .	-0.2439	0 × 98	0.000005	0.0896	0.0899
HF	gs.	0 × 50	0.3231	0.3231	0.3717	0.3717
	1 st .	-0.4834	0 × 50	0.00001	0.1263	0.5764
H ₂ O	gs.	0 × 50	0.2490	0.2982	0.3166	0.3339
	1 st .	-0.3402	0 × 50	0.0858	0.1373	0.2026

Summary



- Gentlest ascent captures the index-1 saddle point of KS density functionals
- First excited state is an index-1 saddle point
- An alternative to non-Aufbau occupation approach, convergence guaranteed
- GAD useful for scanning energy landscapes

Thank you!

